Linear



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Linear Imperfections

equation of motion in an accelerator

→ Hills equation

sine and cosine like solutions + one turn map

dipole perturbations

— closed orbit response

— dispersion orbit

→ integer resonances

quadrupole perturbations

→ tune error

→ beta-beat

half-integer resonances

orbit correction



most effective corrector

Variable Definition

Variables in moving coordinate system:



$$x' = \frac{d}{ds} x$$



Hill's Equation:

 $\frac{d^2x}{ds^2} + K(s) \cdot x = 0; \quad K(s) = K(s + L);$

$$K(s) = \begin{cases} 0 & drift \\ 1/\rho^2 & dipole \\ 0.3 \cdot \frac{B[T/m]}{p[GeV]} & quadrupole \end{cases}$$

Sinelike and Cosinelike Solutions

system of first order linear differential equations:

$$\underline{\mathbf{y}} = \begin{pmatrix} \mathbf{x} \\ \mathbf{x}^{\dagger} \end{pmatrix} \longrightarrow \underline{\mathbf{y}}^{\dagger} + \begin{pmatrix} \mathbf{0} & \mathbf{1} \\ \mathbf{K} & \mathbf{0} \end{pmatrix} \cdot \underline{\mathbf{y}} = \mathbf{0}$$

Floquet theorem:

$$\underline{\mathbf{S}}(\mathbf{s}) = \begin{pmatrix} \sqrt{\beta(\mathbf{s})} \cdot \sin(\phi(\mathbf{s}) + \phi_0) \\ [\cos(\phi(\mathbf{s}) + \phi_0) + \alpha(\mathbf{s}) \cdot \sin(\phi(\mathbf{s}) + \phi_0] / \sqrt{\beta(\mathbf{s})} \end{pmatrix}$$

$$\underline{\mathbf{C}}(\mathbf{s}) = \begin{pmatrix} \sqrt{\beta(\mathbf{s})} \cdot \cos(\phi(\mathbf{s}) + \phi_0) \\ -[\sin(\phi(\mathbf{s}) + \phi_0) + \alpha(\mathbf{s}) \cdot \cos(\phi(\mathbf{s}) + \phi_0] / \sqrt{\beta(\mathbf{s})} \end{pmatrix}$$

$$\left(\beta (s) = \beta (s + L); \phi (s) = \int \frac{1}{\beta} ds; \alpha (s) = -\frac{1}{2} \beta'(s) \right)$$



<u>particles oscillate around an ideal orbit:</u>



additional dipole fields perturb the orbit:

error in dipole field

energy error

$$\alpha = \frac{l}{\rho} = \frac{q \cdot B \cdot l}{p + \Delta p} \approx \left(1 - \frac{\Delta p}{p}\right) \cdot \frac{q \cdot B \cdot l}{p}$$
offset in quadrupole field

$$B_x = -g \cdot y$$

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$$B_{y} = -g \cdot x \qquad x = x_{0} + \tilde{x} \rightarrow B_{y} = -g \cdot x_{0} - g \cdot \tilde{x}$$

dipole component

Sources for Orbit Errors

Quadrupole offset:

- ground motion
 slow drift
 civilisation
 moon
 seasons
 civil engineering
- *Error in dipole strength*
 - *power supplies*
 - **estimation**
- *Energy error of particles*
 - injection energy (RF off)
 - **RF** frequency
 - momentum distribution

inhomogeneous equation:



$$y(s) = a \cdot \underline{S}(s) + b \cdot \underline{C}(s) + \psi(s)$$

we need to find only one solution!

variation of the constant:

 $\underline{\Psi}(s) = \underline{\phi}(s) \cdot u(s); \qquad \underline{\phi}(s) = c \cdot \underline{S}(s) + d \cdot \underline{C}(s);$

variation of the constant in matrix form:

 $\Psi(s) = \Phi(s) \cdot u(s);$ with

$$\underline{\phi(s)} = \left(\frac{\sqrt{\beta(s)} \cdot \sin(\phi(s) + \phi_0)}{\cos(\phi(s) + \phi_0)} \sqrt{\beta(s)} \cdot \cos(\phi(s) + \phi_0) \right) = \left(\frac{\sqrt{\beta(s)}}{\cos(\phi(s) + \phi_0)} \sqrt{\sqrt{\beta(s)}} - \frac{\sqrt{\beta(s)}}{\cos(\phi(s) + \phi_0)} \sqrt{\sqrt{\beta(s)}} \right)$$

periodic boundary conditions:

 $\underline{y(s)} = a \cdot \underline{S(s)} + b \cdot \underline{C(s)} + \underline{\phi(s)} \cdot \int_{s0}^{-1} \underline{G(t)} dt$

with

$$\underline{\mathbf{y}(\mathbf{s})} = \begin{pmatrix} \mathbf{x}(\mathbf{s}) \\ \mathbf{x}^{\dagger}(\mathbf{s}) \end{pmatrix}; \quad \mathbf{x}(\mathbf{s}) = \mathbf{x}(\mathbf{s} + \mathbf{L}); \quad \mathbf{x}^{\dagger}(\mathbf{s}) = \mathbf{x}^{\dagger}(\mathbf{s} + \mathbf{L})$$



periodic boundary conditions determine coefficients a *and* b

$$\mathbf{x}(s) = \frac{\sqrt{\beta(s)}}{2\sin(\pi \cdot Q)} \cdot \int_{s0}^{s0+circ} \sqrt{\beta(t)} \cdot \mathbf{G}(t) \cos[\phi(t) - \phi(s) - \pi Q] dt$$



$\bigcirc \underline{Q}$: number of β -oscillations per turn



 $\mathbf{x}(s) = \frac{\neg \beta(s)}{2\sin(\pi \cdot Q)} \cdot \oint \overline{\beta(t)} \cdot \mathbf{G}(t) \cos[\phi(t) - \phi(s)] - \pi Q] dt$

Example:

 $\mathbf{x}(s) = \frac{\neg \beta(s)}{2\sin(\pi \cdot \mathbf{Q})} \cdot \oint \overline{\beta(t)} \cdot \mathbf{G}(t) \cos[\phi(t) - \phi(s)] - \pi \mathbf{Q}] dt$

with

$$\mathbf{G}(\mathbf{t}) = \frac{-1}{\rho(\mathbf{t})} \cdot \frac{\Delta \mathbf{p}}{\mathbf{p}}$$



$$\mathbf{x}(\mathbf{s}) = \mathbf{D}(\mathbf{s}) \cdot \frac{\Delta \mathbf{p}}{\mathbf{p}}$$

with

$$D(s) = \frac{-\sqrt{\beta(s)}}{2\sin(\pi \cdot Q)} \cdot \oint \frac{-\sqrt{\beta(t)}}{\rho(t)} \cdot \cos[\phi(t) - \phi(s)] - \pi Q] dt$$



β – **Beat**

quadrupole error:





variation of the constant:

$$\Delta\beta(s) = \frac{\beta(s)}{2\sin(2\pi \cdot Q)} \int_{s0}^{s0+\text{circ}} \beta(t) \cdot \Delta k(t) \cos[2[\phi(t)-\phi(s)]-2\pi Q] dt$$

 β – beat oscillates with twice the betatron frequency



Orbit Stability

<u>Quadrupole Error:</u>

• orbit kick proportional to

beam offset in quadrupole

Q = N + 0.5



amplitude increase

2. Turn: x < 0



amplitude increase



Tune Error

one turn map:

cosine- and sine-like solutions to Hill's equation

$$\underline{z}_{n+1} = \mathbf{M} \cdot \underline{z}_n \qquad \underline{z} = \begin{pmatrix} \mathbf{X} \\ \mathbf{X} \end{pmatrix}$$

with

 $\mathbf{M} = \mathbf{I} \cdot \cos(2\pi \mathbf{Q}) + \mathbf{J} \cdot \sin(2\pi \mathbf{Q})$

$$\mathbf{I} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \qquad \mathbf{J} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}; \ \gamma = \begin{bmatrix} 1 + \alpha^2 \end{bmatrix} / \beta$$

remember: $\cos(2\pi Q) = \frac{1}{2}$ trace M

 $\longrightarrow \ \text{the coefficients of:} \quad \frac{M - I \cdot \cos(2\pi Q)}{\sin(2\pi Q)}$

provide the optic functions at s_0

transfer matrix for single quadrupole:

$$\mathbf{m}_{0} = \begin{pmatrix} 1 & \mathbf{0} \\ -\mathbf{k} \cdot \mathbf{I} & 1 \end{pmatrix}$$

matrix for single quadrupole with error:

$$\mathbf{m} = \begin{pmatrix} 1 & 0 \\ [-\mathbf{k} + \Delta \mathbf{k}] \cdot \mathbf{i} & 1 \end{pmatrix}$$

one turn matrix with quadrupole error:

$$\mathbf{M} = \mathbf{m} \bullet \mathbf{m}_0^{-1} \bullet \mathbf{M}_0$$

trace M

 $\cos(2\pi Q) = \cos(2\pi Q_0) - \frac{1}{2} \beta \cdot \Delta k \cdot k \sin(2\pi Q_0)$

Tune Error

distributed perturbation:



Problems Generated by Orbit Errors



orbit correctors



the orbit determines the particle energy!

assume: L > design orbit



->> E depends on orbit and magnetic field!

Orbit Correction

- the orbit error in a storage ring with conventional
 magnets is dominated by the contributions
 from the quadrupole alignment errors
 - orbit perturbation is proportional to the local β -functions at the location of the dipole error
 - alignment errors at QF cause mainly horizontal orbit errors
 - alignment errors at QD causes mainly vertical orbit errors

Orbit Correction

aim at a local correction of the dipole error due to the quadrupole alignment errors

 place orbit corrector and BPM next to the main quadrupoles

horizontal BPM and corrector next to QF
 vertical BPM and corrector next to QD



orbit in the opposite plane?

relative alignment of BPM and quadrupole?

Local Orbit Bumps I





trajectory response:

[no periodic boundary conditions]

$$\longrightarrow x(s) = \neg \beta_i \beta(s) \cdot \theta_i \cdot \sin[\phi(s) - \phi_i]$$

$$\longrightarrow x'(s) = \neg \langle \beta_i / \beta(s) \cdot \theta_i \cdot \cos[\phi(s) - \phi_i] \rangle$$

Local Orbit Bumps II

closed orbit bump:

compensate the trajectory perturbation with

additional corrector kicks further down stream

closure of the perturbation within one turn

local orbit excursion

possibility to correct orbit errors locally

closure with one additional corrector magnet
 π - bump
 closure with two additional corrector magnets
 three corrector bump

Local Orbit Bumps III



requires 90° lattice

sensitive to lattice errors

requires horizontal BPMs at QF and QD

sensitive to BPM errors

requires large number of correctors

Local Orbit Bumps IV

3 corrector bump: (quasi local correction of error)



Harmonic Filtering I

unperturbed solution (smooth approximation):

 $x'' + \left[\frac{2\pi}{C} \cdot Q\right]^2 \cdot x = 0 \longrightarrow x(s) = A \cdot e^{i\frac{2\pi}{C} \cdot Q \cdot s}$

orbit perturbation due to random kicks:

$$\mathbf{x}^{"} + \mathbf{K}(\mathbf{s}) \cdot \mathbf{x} = \sum_{i=1}^{m} \boldsymbol{\theta}_{i} \cdot \boldsymbol{\delta}(\mathbf{s} - \mathbf{s}_{i})$$
$$= \mathbf{F}(\mathbf{s})$$

periodic boundary conditions:

$$\longrightarrow \quad \mathbf{x}(\mathbf{s}) = \sum_{n} \mathbf{d}_{n} \cdot \mathbf{e}^{i \frac{2\pi}{C}} \cdot \mathbf{n} \cdot \mathbf{s}$$

$$\longrightarrow \quad \mathbf{F}(\mathbf{s}) = \sum_{n} \mathbf{f}_{n} \cdot \mathbf{e}^{i \frac{2\pi}{C}} \cdot \mathbf{n} \cdot \mathbf{s}$$



inserting Ansatz into Hill's equation:



small number of correctors are efficient

SVD I

linear relation between BPM and corrector data:

COR: vector of corrector amplitudes

 \longrightarrow <u>COR</u> = [a_1, a_2, \dots, a_m]

BPM: vector of all BPM data

 \longrightarrow <u>BPM</u> = [b₁, b₂,..., b_n]

 \longrightarrow <u>BPM</u> = A · <u>COR</u>; A = n × m matrix

global orbit correction:

ORB: vector of all measured orbit data

 \longrightarrow <u>ORB</u> = [c₁, c₂,..., c_n]

— find a set of corrector settings that satisfies:

 $\underline{ORB} - A \cdot \underline{COR} = \underline{0}$



mathematical solution:

 $\underline{\mathbf{COR}} = \mathbf{A}^{-1} \cdot \underline{\mathbf{ORB}}$

problem:

A is normally not invertible

(A is normally not even a square matrix)

 \longrightarrow minimise the norm: $|| ORB - A \cdot COR ||$

with
$$\left\| \underline{\mathbf{X}} \right\| = \left(\sum_{i=1}^{m} \left| \mathbf{x}_{i} \right|^{p} \right)^{1/p}$$

— find a matrix B such that:

 $|| \underline{ORB} - A \cdot B \cdot \underline{ORB} ||$

attains a minimum

SVD III

singular value decomposition SVD:



D is a diagonal matrix:

$$\mathbf{D} = \begin{pmatrix} \mathbf{\sigma}_{11} & \mathbf{0} & \cdots & \mathbf{0} \\ \mathbf{\sigma}_{22} & & \vdots \\ \mathbf{0} & \mathbf{\sigma}_{kk} & \cdots & \mathbf{0} \end{pmatrix}$$

 $k \leq \min(n,m)$





define the correction matrix:

 $\longrightarrow B = O_2^t \cdot D \cdot O_1^t$

$$\longrightarrow \mathbf{A} \cdot \mathbf{B} = (\mathbf{O}_1 \cdot \mathbf{D} \cdot \mathbf{O}_2) \cdot (\mathbf{O}_2^{\mathsf{t}} \cdot \overset{\wedge}{\mathbf{D}} \cdot \mathbf{O}_1^{\mathsf{t}})$$

SVD allows you to adjust k corrector magnets k = min(m,n)

- if k = m = n one obtains a zero orbit(by using a all possible corrector magnets)
 - if m ≠ n SVD algorithm minimises the norm (by using all possible corrector magnets)

algorithm is not stable if <u>D</u> has small
 eigenvalues

Most Effective Corrector

orbit is perturbed by a few large perturbations:



 $|| \mathbf{BPM} - \mathbf{A} \cdot \mathbf{COR} || \quad \text{with } || \underline{\mathbf{X}} || = \left(\sum_{i=1}^{m} |\mathbf{x}_{i}|^{p}\right)^{1/p}$

with a small set of 'k' corrector magnets

brut force: select all possible combinations



- selective: keep the already selected correctors
 - → much faster!



— can generate orbit bumps

MICADO: selective + cross-correlation between orbit residues and remaining correctors



the orbit determines the particle energy!

assume: L > design orbit



->> E depends on orbit and magnetic field!