

# Plasma Wake Acceleration



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**New Accelerator Concepts for Particle Therapy  
(Laser-electron, Laser-ion.. etc)**

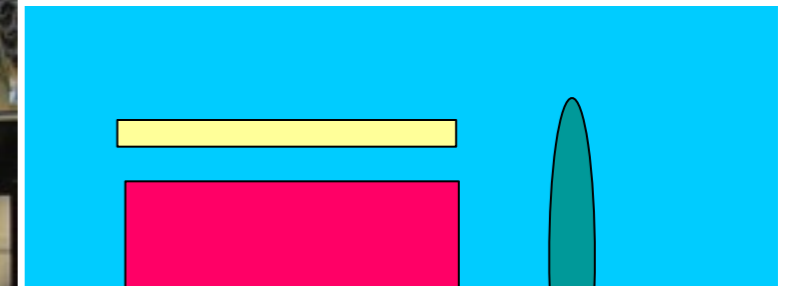
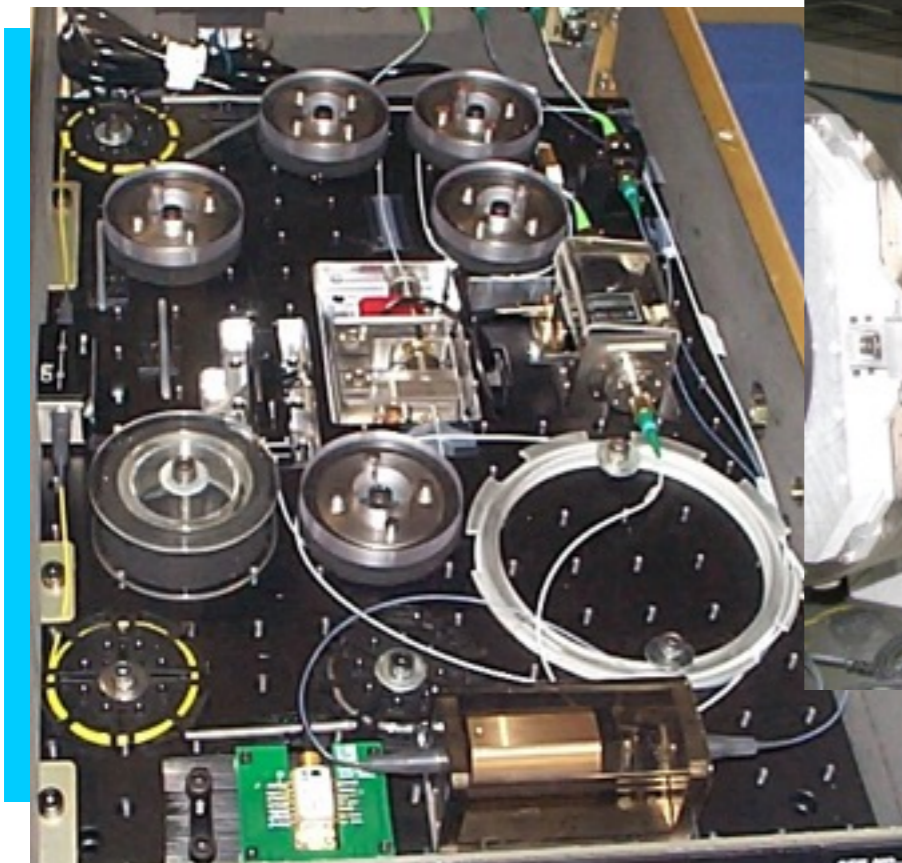
Markus Roth

Institute for Nuclear Physics  
TU Darmstadt

- Ultra-Intense-Lasers
- Laser-driven electrons
  - Potential for Applications in Therapy
  - Use of secondary Radiation
- Laser-driven Ion Beams
  - Properties and Mechanism
  - Experiments
- Applications as Sources for Therapy
  - Transport of laser-accelerated protons
- Summary

# How to do A LOT of laser light

MOPA- scheme:



spatial filter

Limits of the amplification:

Damage threshold of the medium  
non-linear refraction index

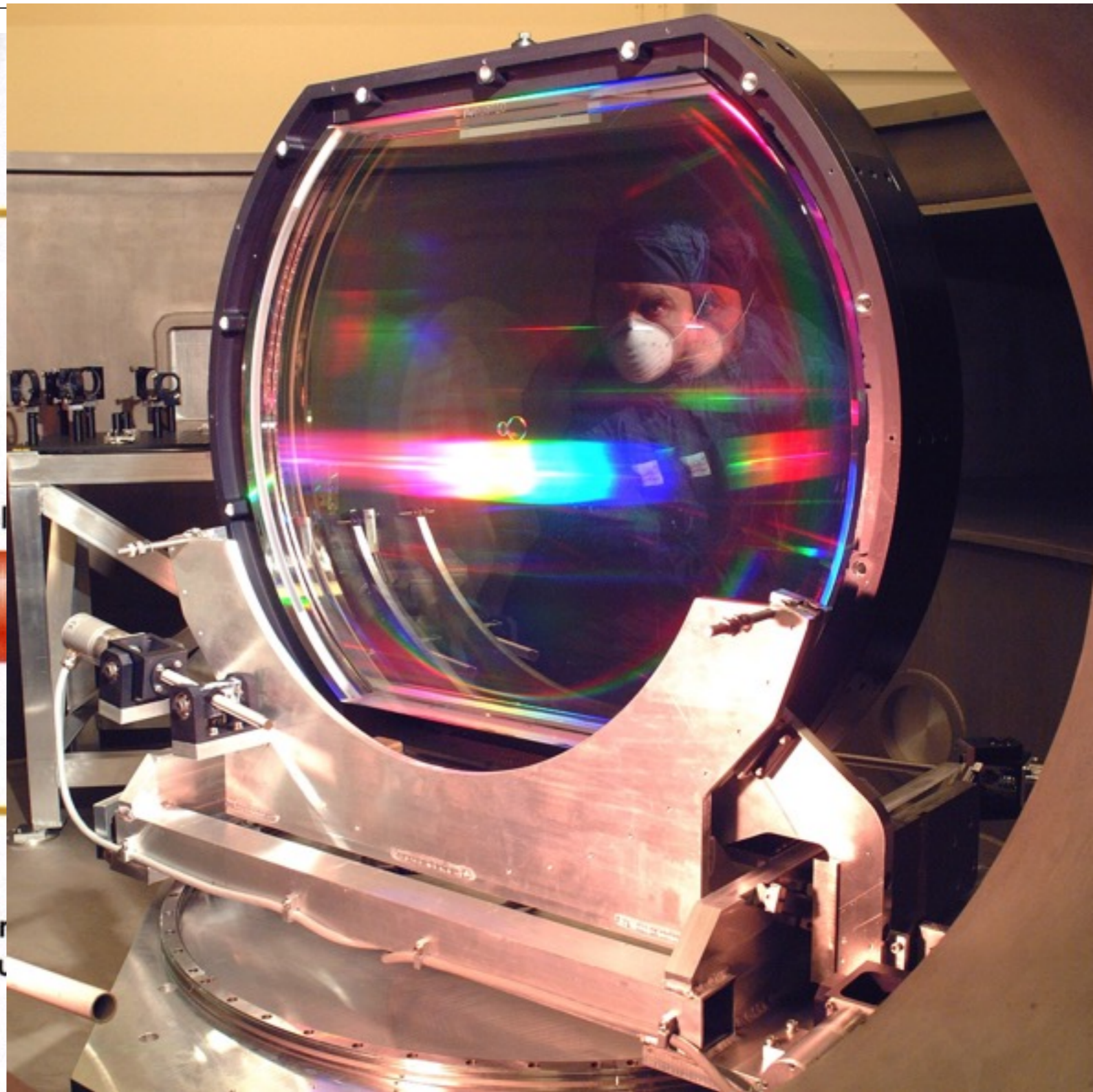
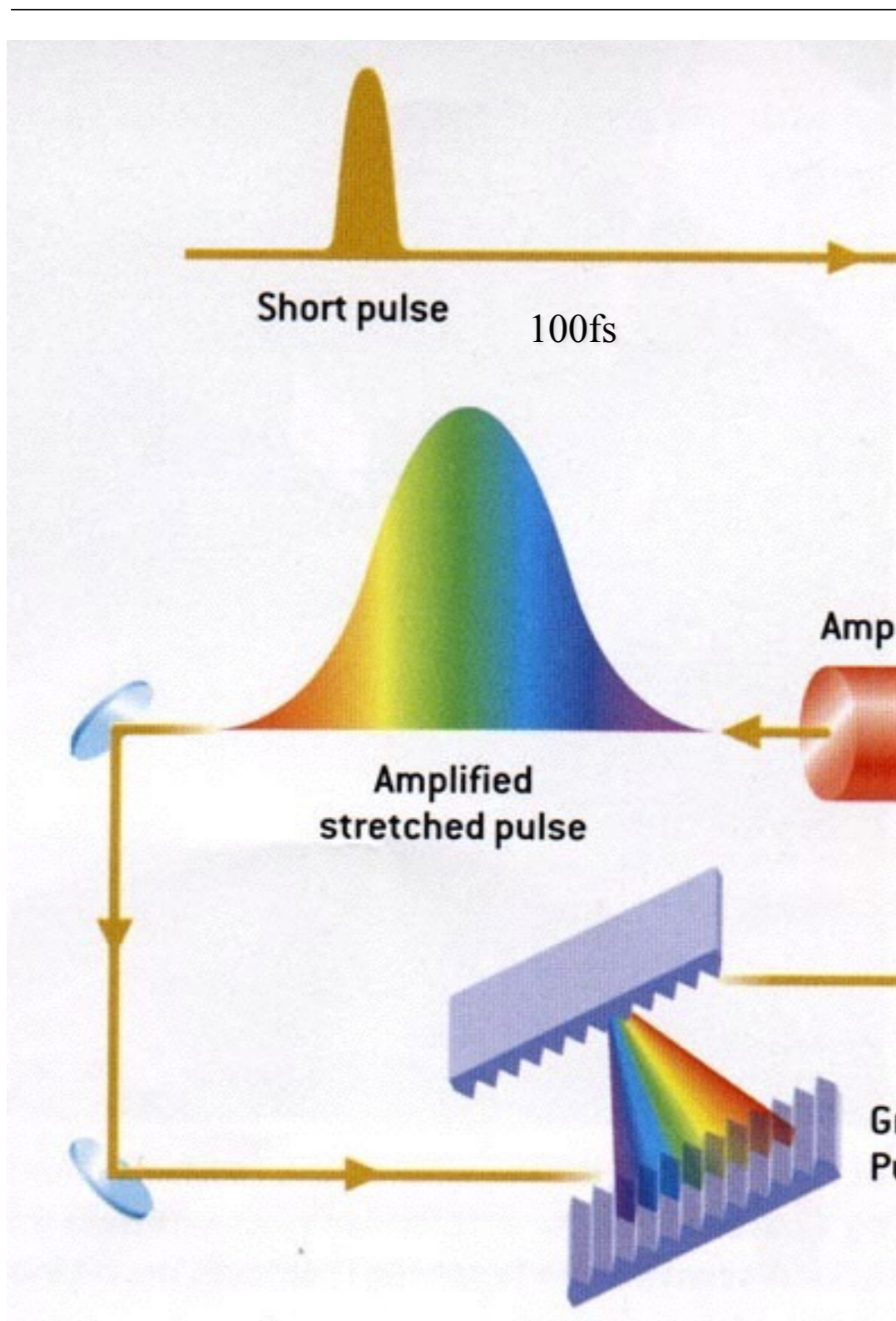
self- filamentation  
beam quality

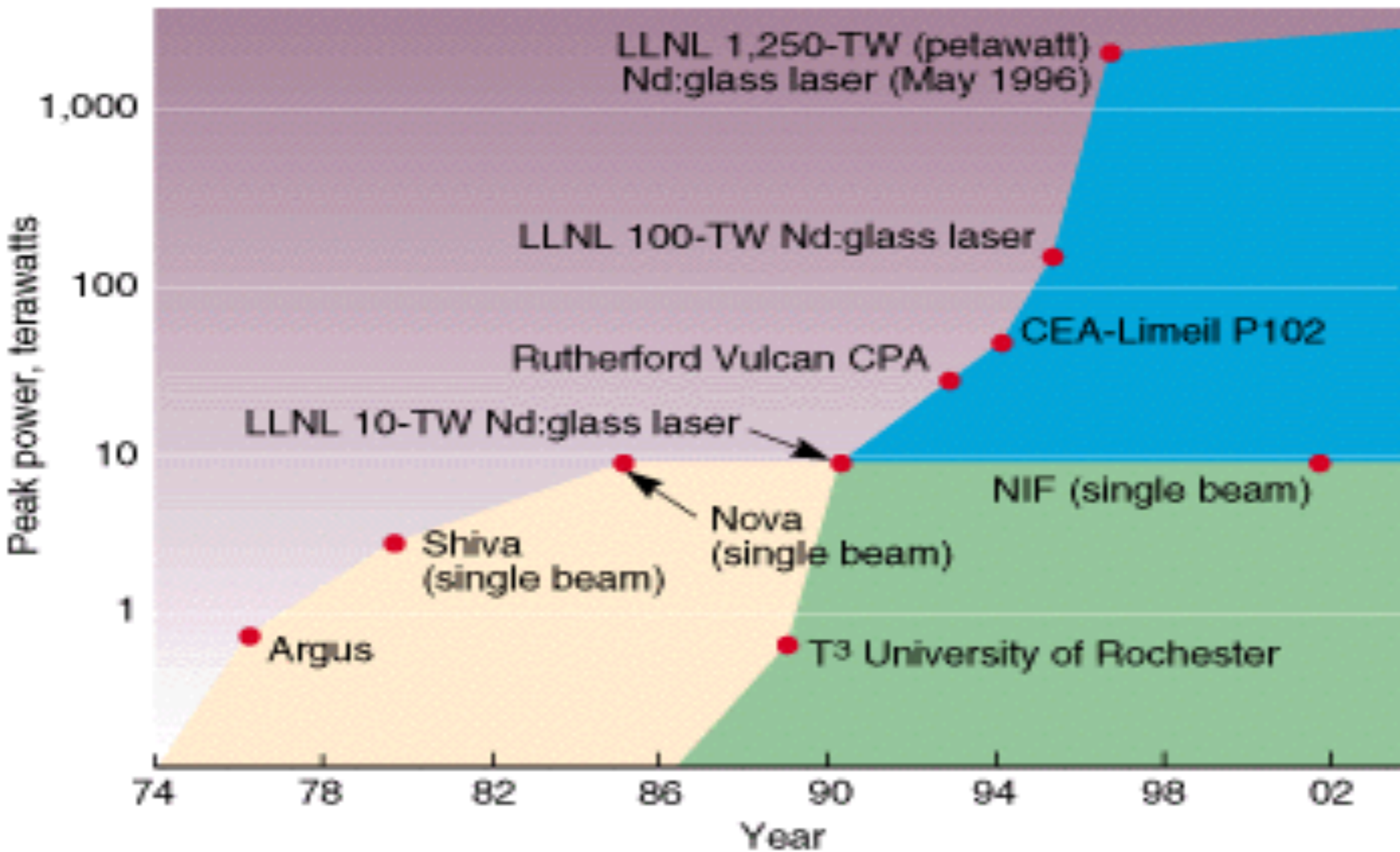
$$B = \Delta\Phi = \frac{2\pi}{\lambda} \int I(l) n_{nl} dl$$

# CPA technique



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# The Petawatt opens a new regime of ultra-relativistic laser-matter interactions

## Enormous EM fields

at  $I = 10^{21}$  W/cm<sup>2</sup> :

$$E \sim I^{1/2} \lambda = 10^{14} \text{ V/m}$$

$$B = E/c = 3 \times 10^5 \text{ Tesla}$$

$$P_{\text{rad}} = I/c = 3 \times 10^{10} \text{ J/cm}^3$$

$$= 300 \text{ GBar}$$

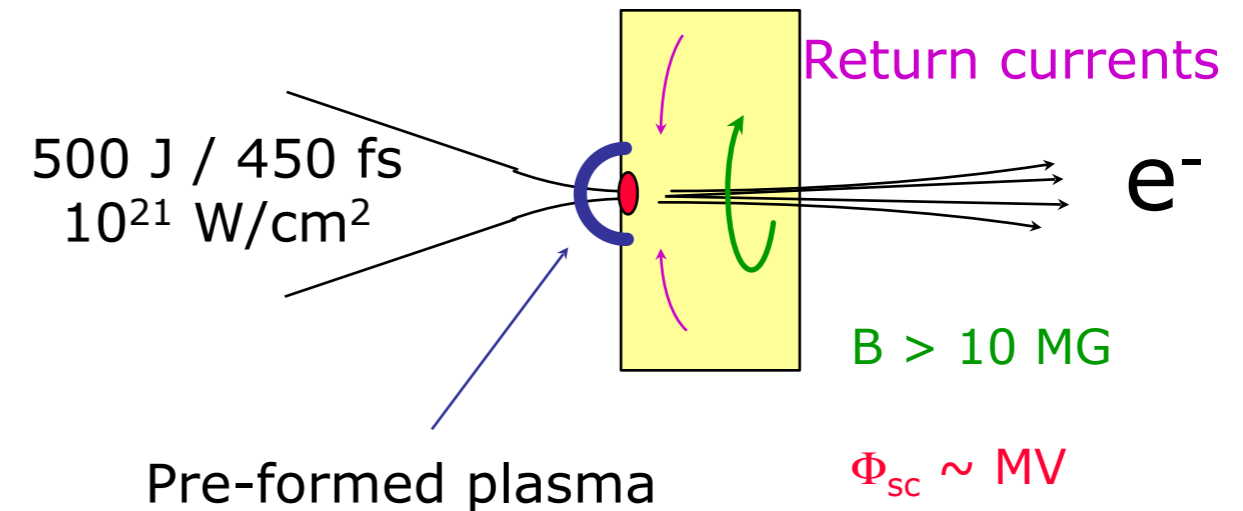
## Relativistic Electron Motion

Cycle-averaged oscillation energy:

$$E_{\text{avg}} = mc^2 [1 + a_0^2/2]^{1/2}$$

$$a_0 = eA/mc^2 = [I/(1.37 \times 10^{18})]^{1/2} \lambda (\mu\text{m})$$

at  $10^{21}$  W/cm<sup>2</sup>,  $a_0 \sim 27$ ,  $E_{\text{avg}} > 10 \text{ MeV}$



## Laser-plasma effects

filamentation

self-focusing

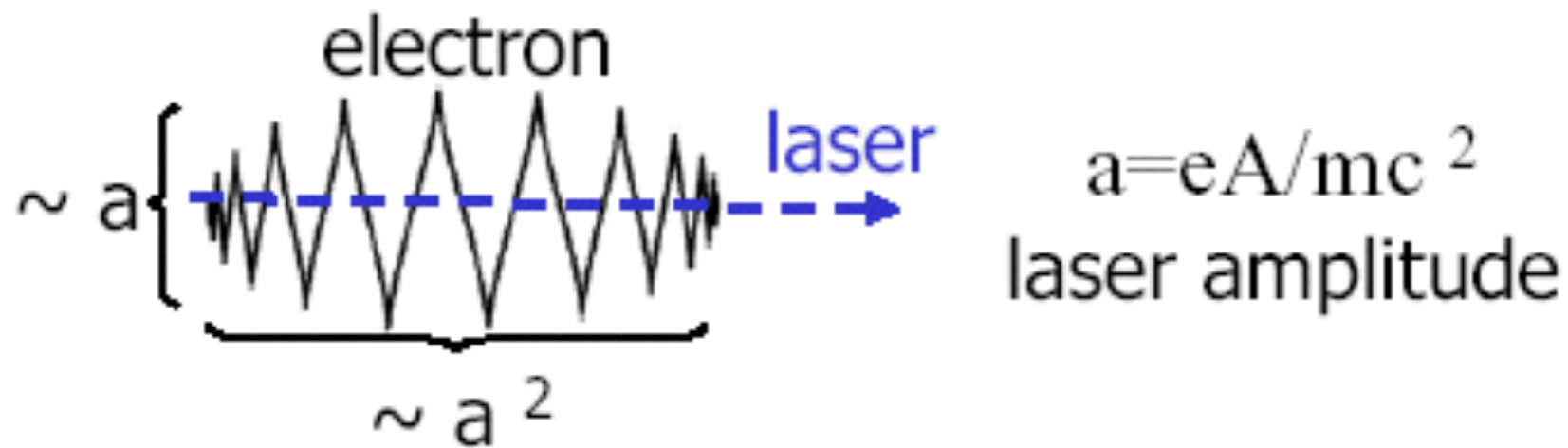
plasma acceleration

» dependence on pre-pulse

*Electron & bremsstrahlung energies can greatly exceed nuclear excitation & pair-creation thresholds*

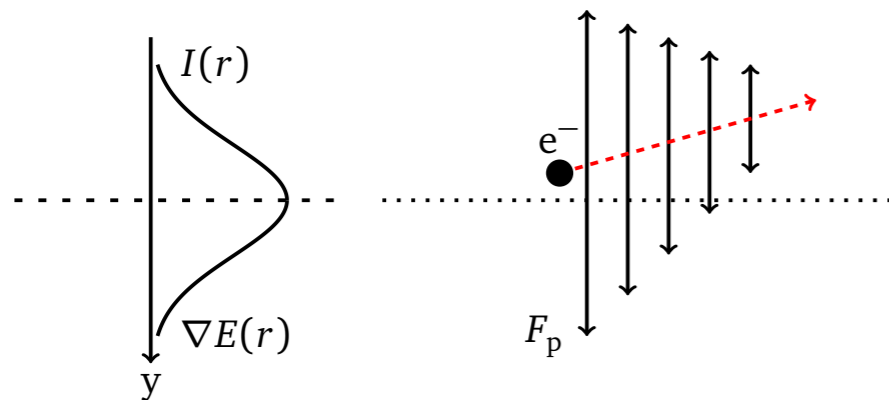
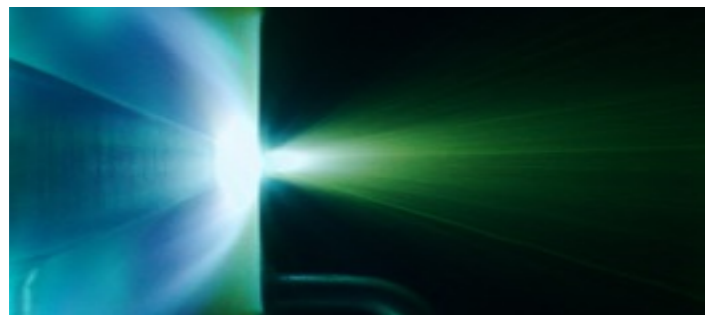
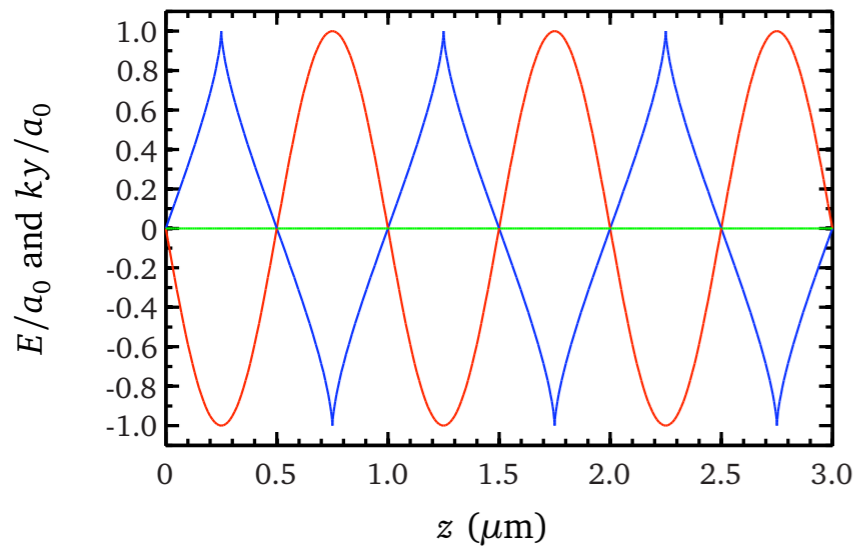
# Electron trajectory in plane laser pulse

$$\vec{F} = -e \left[ \underbrace{\vec{E}}_{\sim a} + \underbrace{(\vec{v} / c) \times \vec{B}}_{\sim a^2} \right]$$

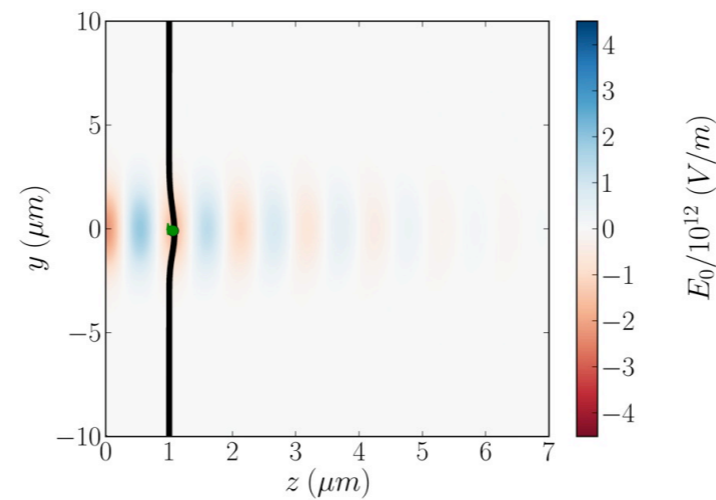


an electron cannot gain energy from a plane wave

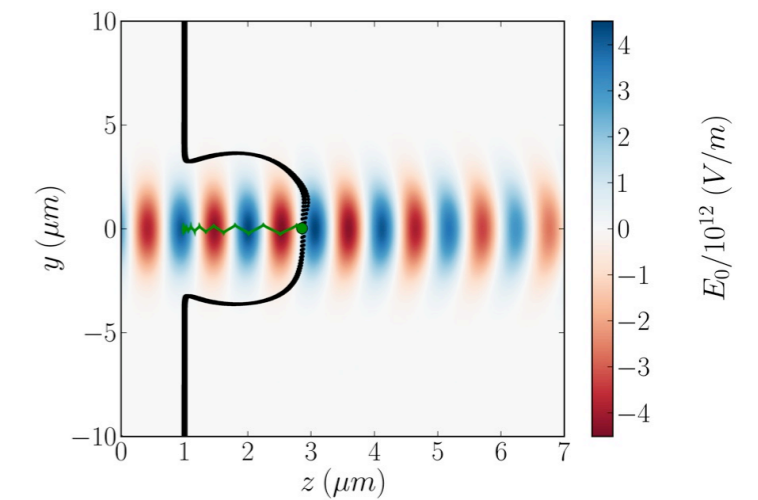
# But in reality.....



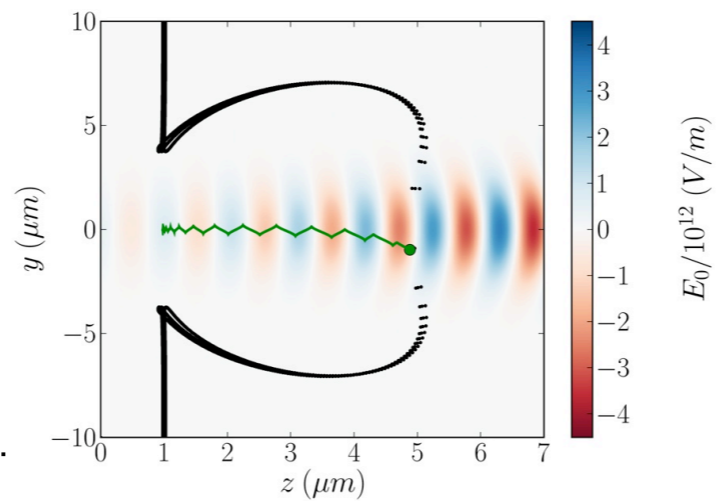
$t = 00043.41 \text{ fs}$



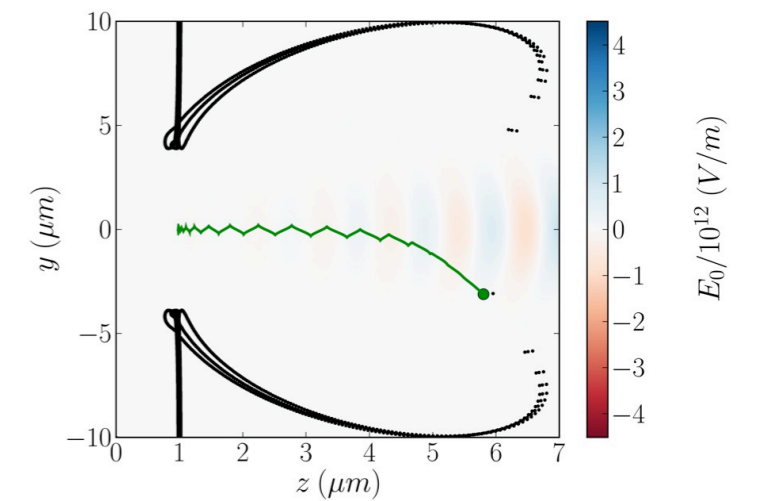
$t = 00069.28 \text{ fs}$



$t = 00090.52 \text{ fs}$



$t = 00106.84 \text{ fs}$

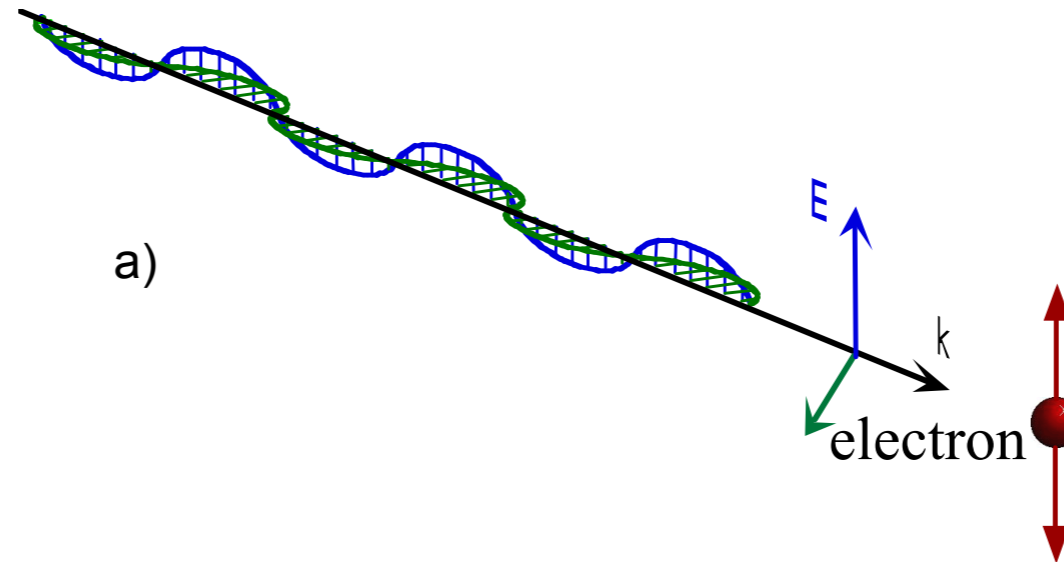


$$f_p = -\frac{e^2}{4m_e\omega_L^2} \nabla(\mathbf{E} \cdot \mathbf{E}^*). \quad f_p = -\frac{c^2}{\gamma} \left[ \nabla m_{\text{eff}} + \frac{\gamma - 1}{v_0^2} (\mathbf{v}_0 \cdot \nabla m_{\text{eff}}) \mathbf{v}_0 \right] \quad m_{\text{eff}} = \left( 1 + \frac{e^2 \mathbf{A} \cdot \mathbf{A}^*}{2m_e^2 c^2} \right)^{1/2} = \bar{\gamma},$$



# Relativistic Effects (e-m rectification)

*classical Optics*

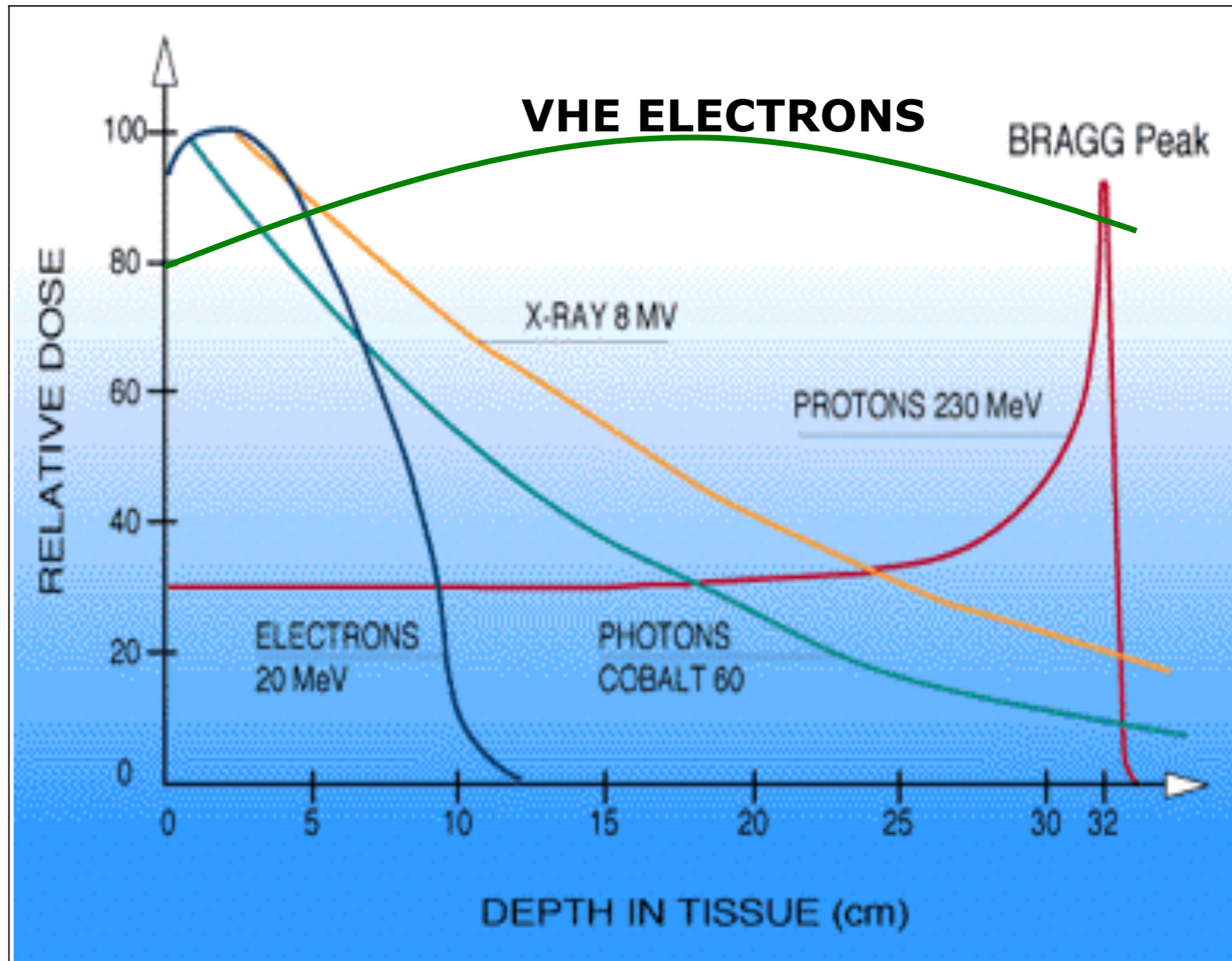


**Ultra-Intense Lasers will push large amounts of electrons in FORWARD direction!**

# Dose deposition : Photon X, electrons VHE (very high energy) electrons, & ions



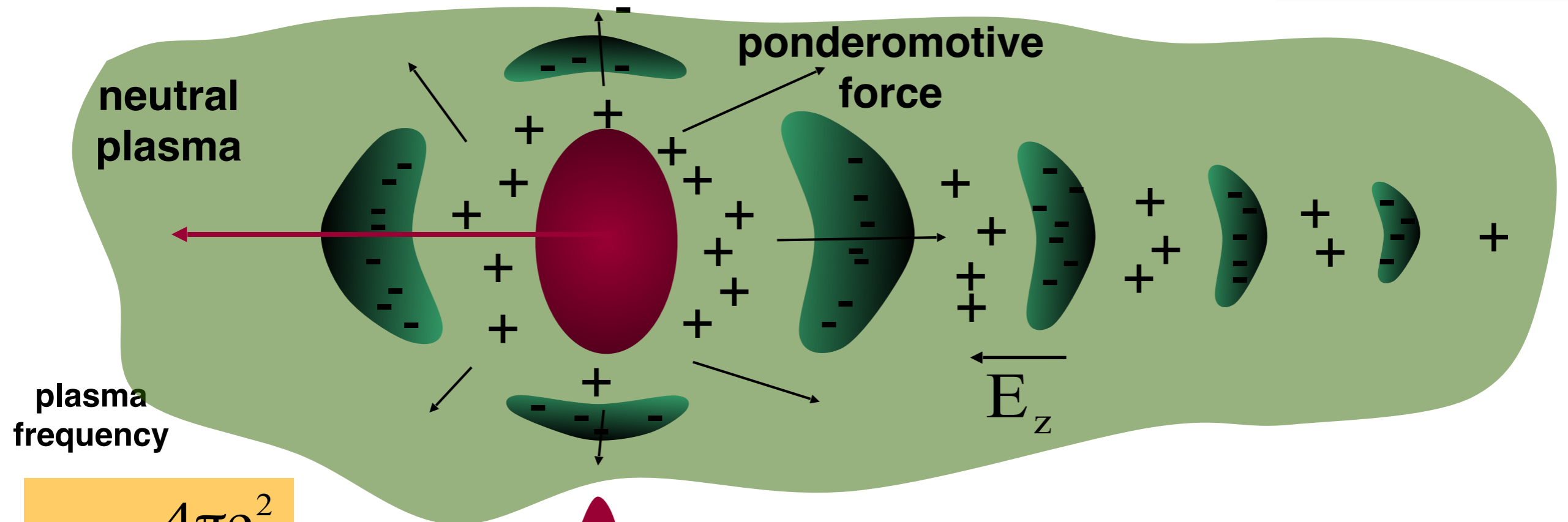
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„Surfing“ on plasma waves – Wake Field Acceleration

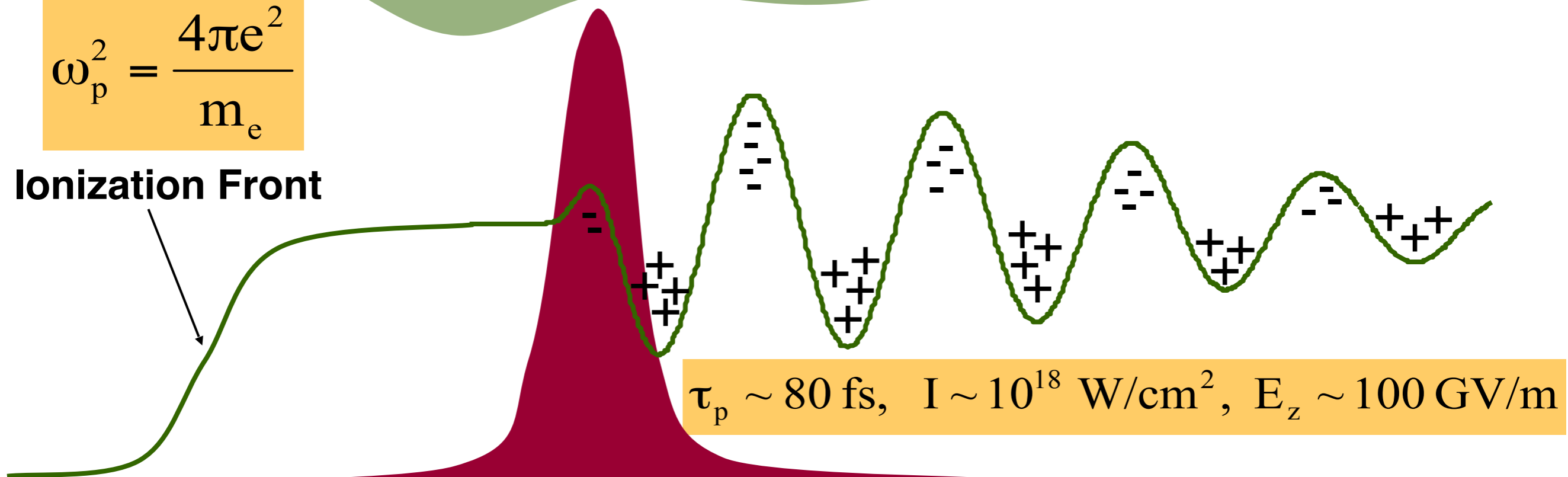


# Resonant Laser Wakefield Generation:



$$\omega_p^2 = \frac{4\pi e^2}{m_e}$$

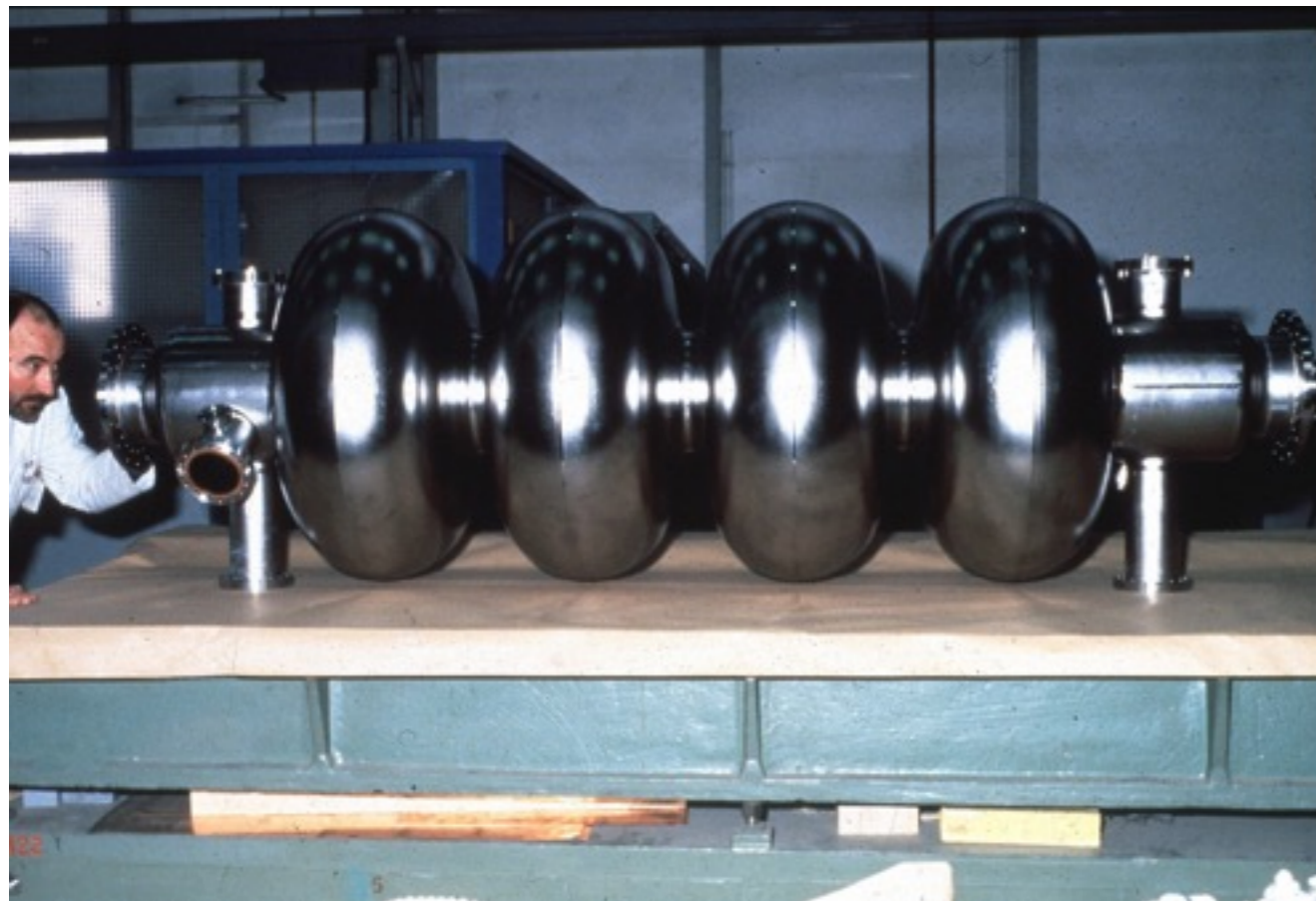
**Ionization Front**



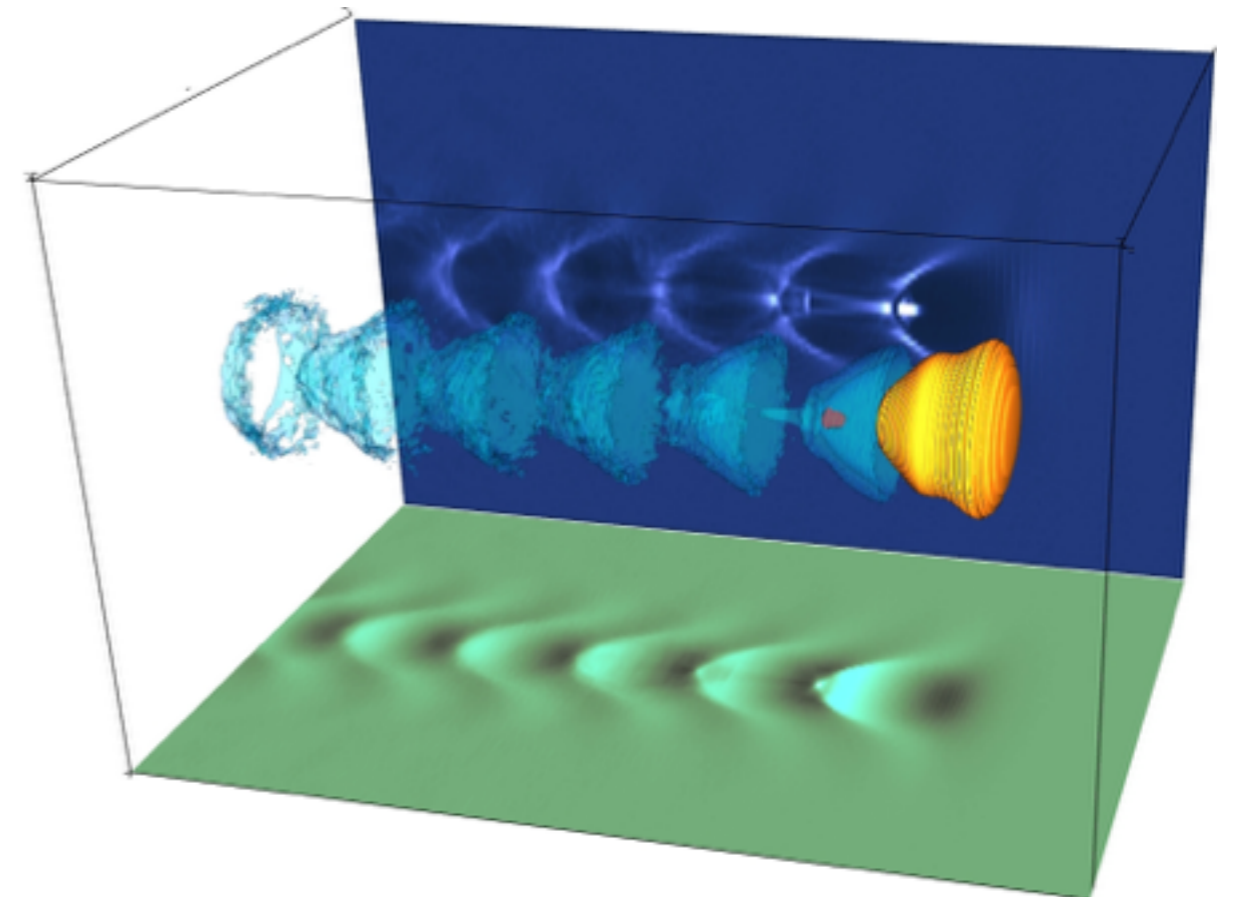
# Efficient acceleration in plasmas = compact accelerator ?

**RF cavity : 1 m**

**Plasma cavity : 1 mm**



**$E = 10-100 \text{ MeV/m}$**



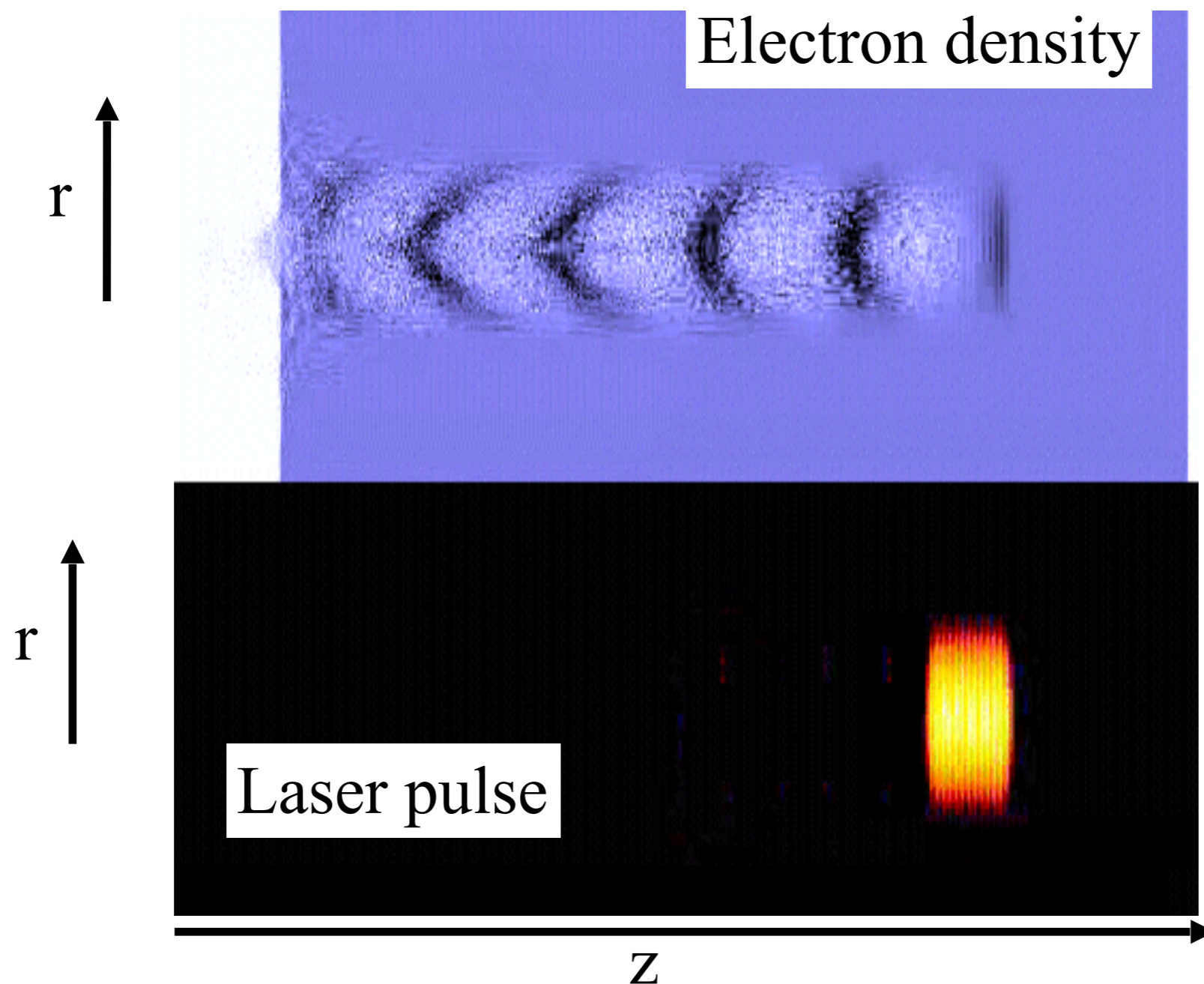
**$E = 100-1000 \text{ GeV/m}$**

courtesy of **V. Malka et al., Science 2002**

# 2D PIC Resonant LWFA Simulation

A. Pukov et al, MPI - Quantenoptic, Germany

$$I=4 \times 10^{18} \text{ W/cm}^2, t=70 \text{ fs}, n_e=3 \times 10^{17} \text{ cm}^{-3}$$



# Limits of the electron acceleration

For the optimum plasma condition,  $\lambda_p = \pi\sigma_z$

**Diffraction limitation:**

$$L_{diff} = \pi Z_R = \frac{\pi^2 r_0^2}{\lambda_0}$$

$$\Delta W_{dif} [\text{GeV}] \cong 0.85 P [\text{TW}] \lambda_0 [\mu\text{m}] / (\gamma_0 \tau_L [\text{fs}])$$

**Dephasing limitation:**

$$L_d [\text{cm}] = 0.18 \times 10^{-4} \tau_L^3 [\text{fs}] \gamma_0 / \lambda_0^2 [\mu\text{m}]$$

$$\Delta W_d [\text{GeV}] \cong 0.01 P [\text{TW}] \tau_L^2 [\text{fs}] / r_0^2 [\mu\text{m}]$$

**Pump depletion:**

$$L_{pd} [\text{cm}] = 1.06 \times 10^{-6} \tau_L^3 [\text{fs}] r_0^2 [\mu\text{m}] \gamma_0^3 / (\lambda_0^4 [\mu\text{m}] P [\text{TW}])$$

$$\Delta W_{pd} [\text{GeV}] = 0.91 \times 10^{-3} \tau_L^2 [\text{fs}] \gamma_0^2 / \lambda_0^2 [\mu\text{m}]$$

e.g.

$$\lambda_0 = 0.8 \mu\text{m},$$

$$P = 2 \text{TW},$$

$$\tau = 100 \text{fs},$$

$$r_0 = 10 \mu\text{m}$$

$$Z_R \cong 0.4 \text{mm}$$

$$a_0 = 0.77 (\gamma_0 = 1.14)$$

$$n_0 \cong 3.5 \times 10^{17} \text{cm}^{-3}$$

$$\lambda_p \cong 60 \mu\text{m}$$

**Diffraction limit:**

$$L_{diff} \cong 1.2 \text{mm}$$

$$\Delta W_{dif} \cong 12 \text{MeV}$$

**Dephasing limit:**

$$L_d \cong 32 \text{cm}$$

$$\Delta W_d \cong 2 \text{GeV}$$

**Pump depletion limit:**

$$L_{pd} \cong 192 \text{cm}$$

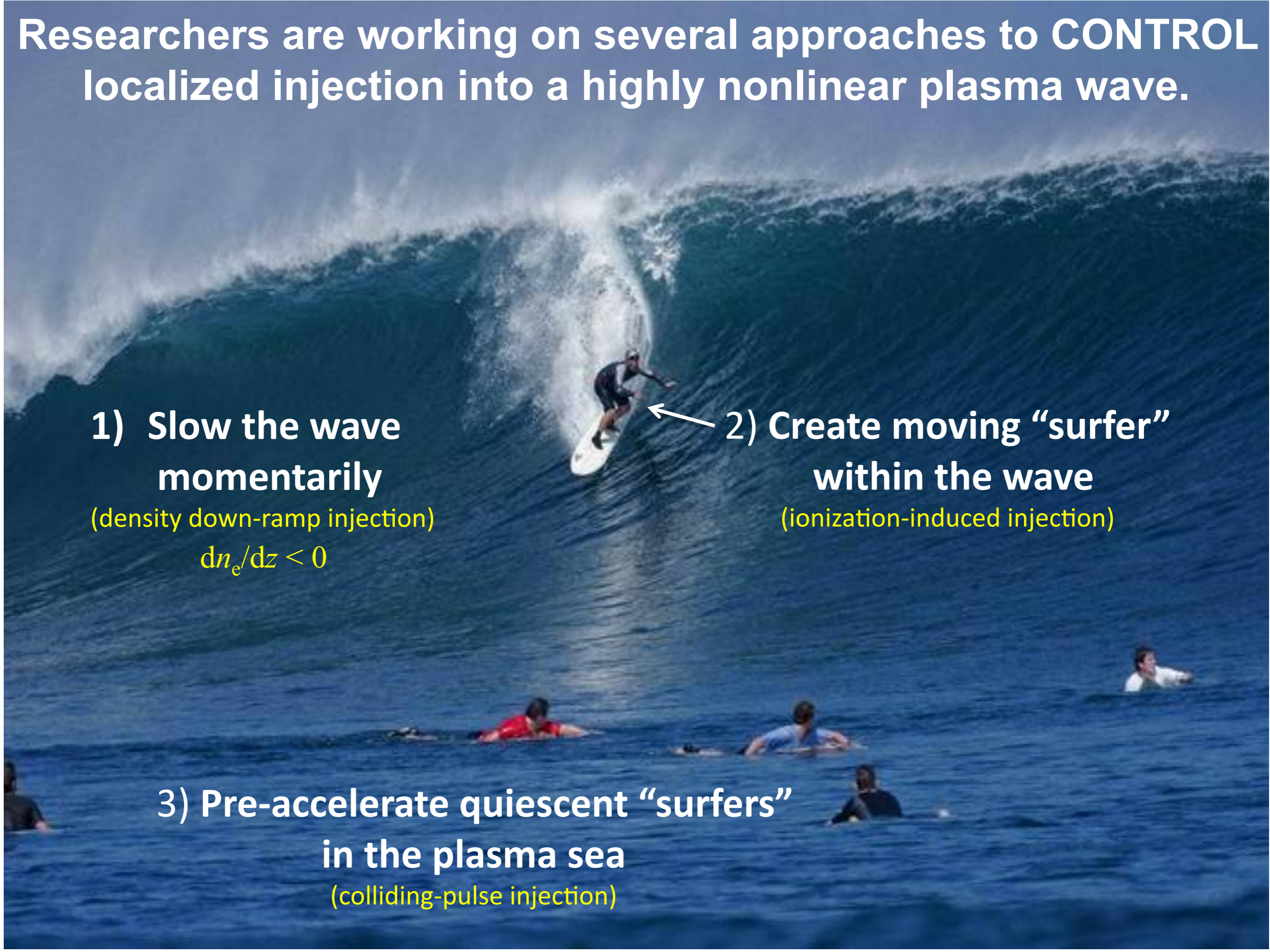
$$\Delta W_{pd} \cong 18.5 \text{GeV}$$

Researchers are working on several approaches to **CONTROL** localized injection into a highly nonlinear plasma wave.

**1) Slow the wave momentarily**  
(density down-ramp injection)  
 $dn_e/dz < 0$

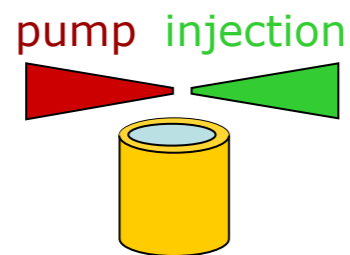
**2) Create moving “surfer” within the wave**  
(ionization-induced injection)

**3) Pre-accelerate quiescent “surfers” in the plasma sea**  
(colliding-pulse injection)

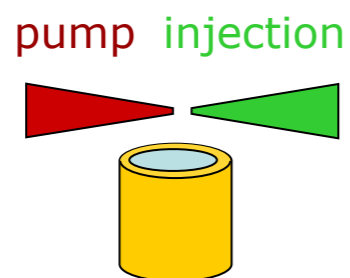




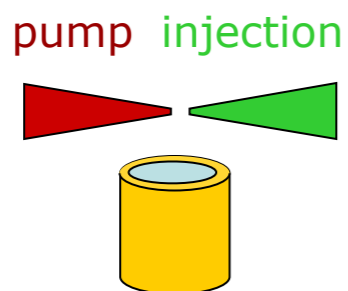
# Use of a two laser scheme: Mono energetic e-beam :1% energy spread



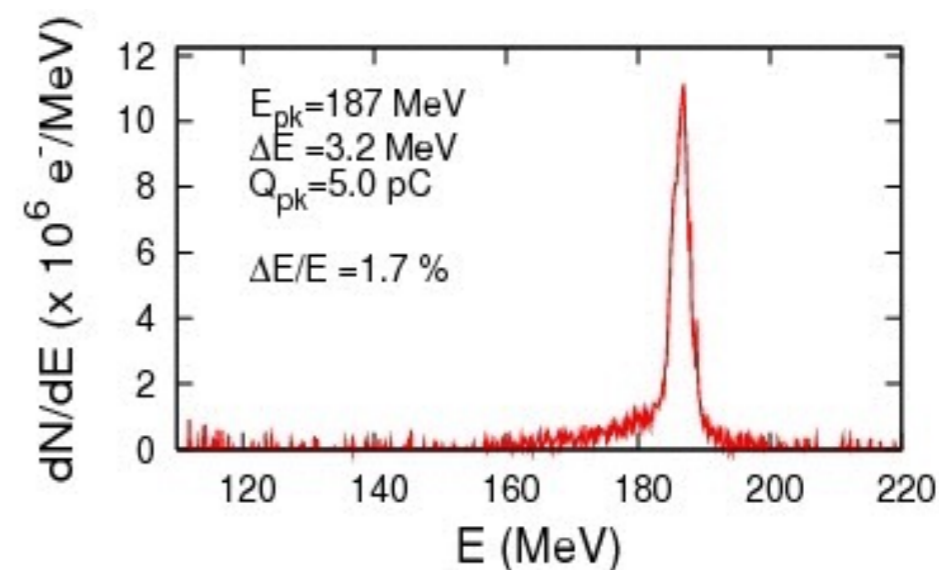
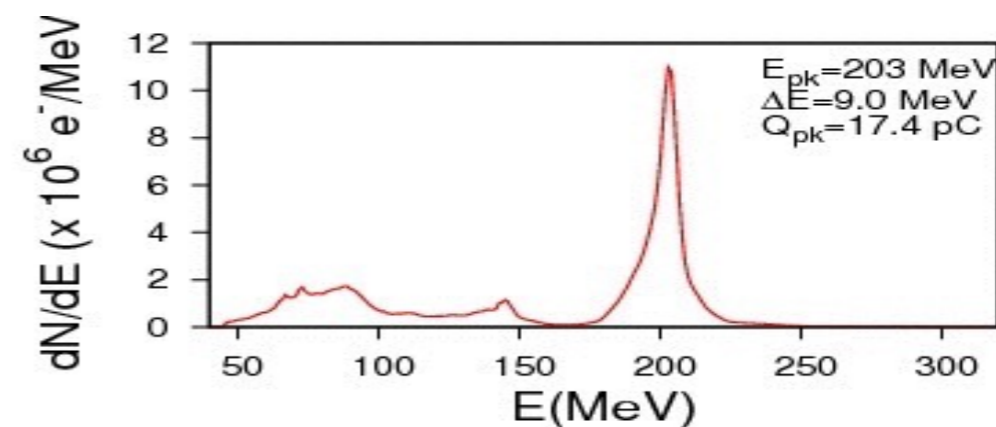
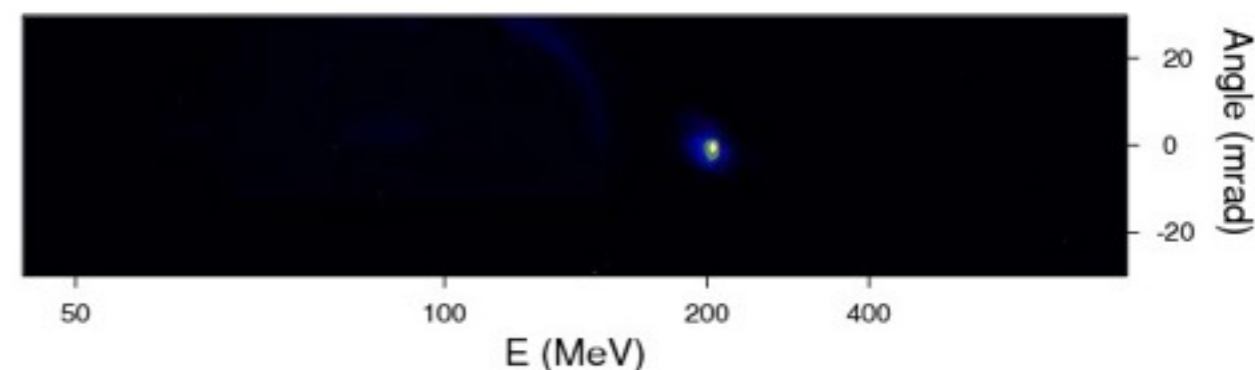
Late Injection



Injection at center



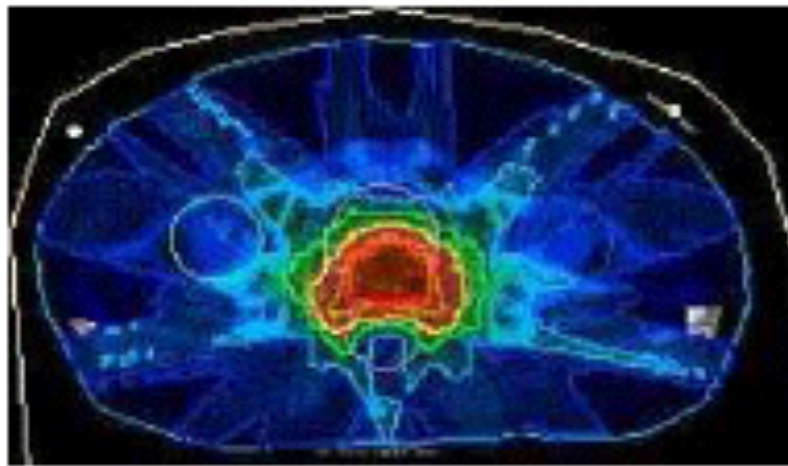
Early Injection



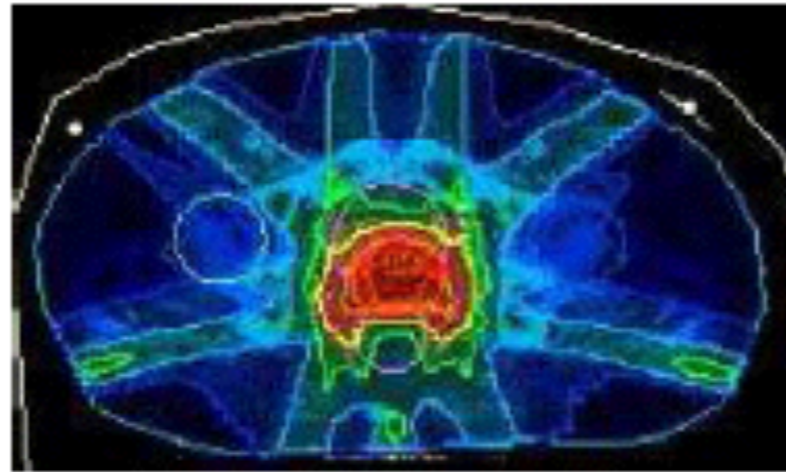
**J. Faure et al., Nature 2006**

# Application to radiotherapy: Improvement of some cancer treatments

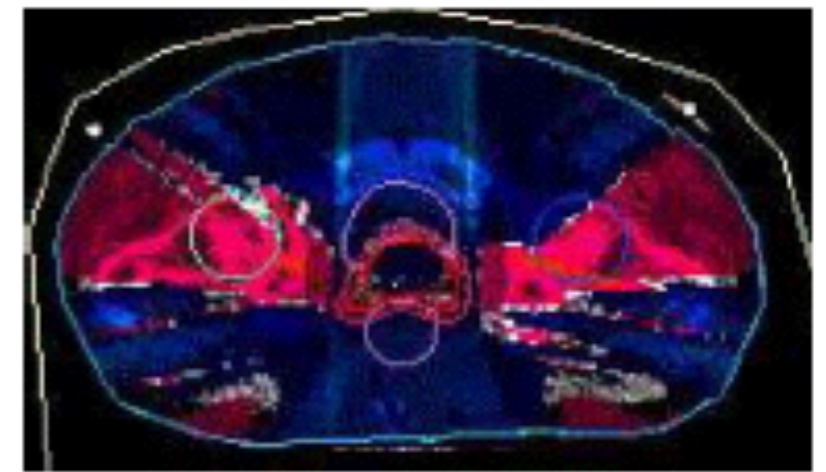
**A typical transversal dose distribution with 7 beams.**



(a)



(b)



(c)

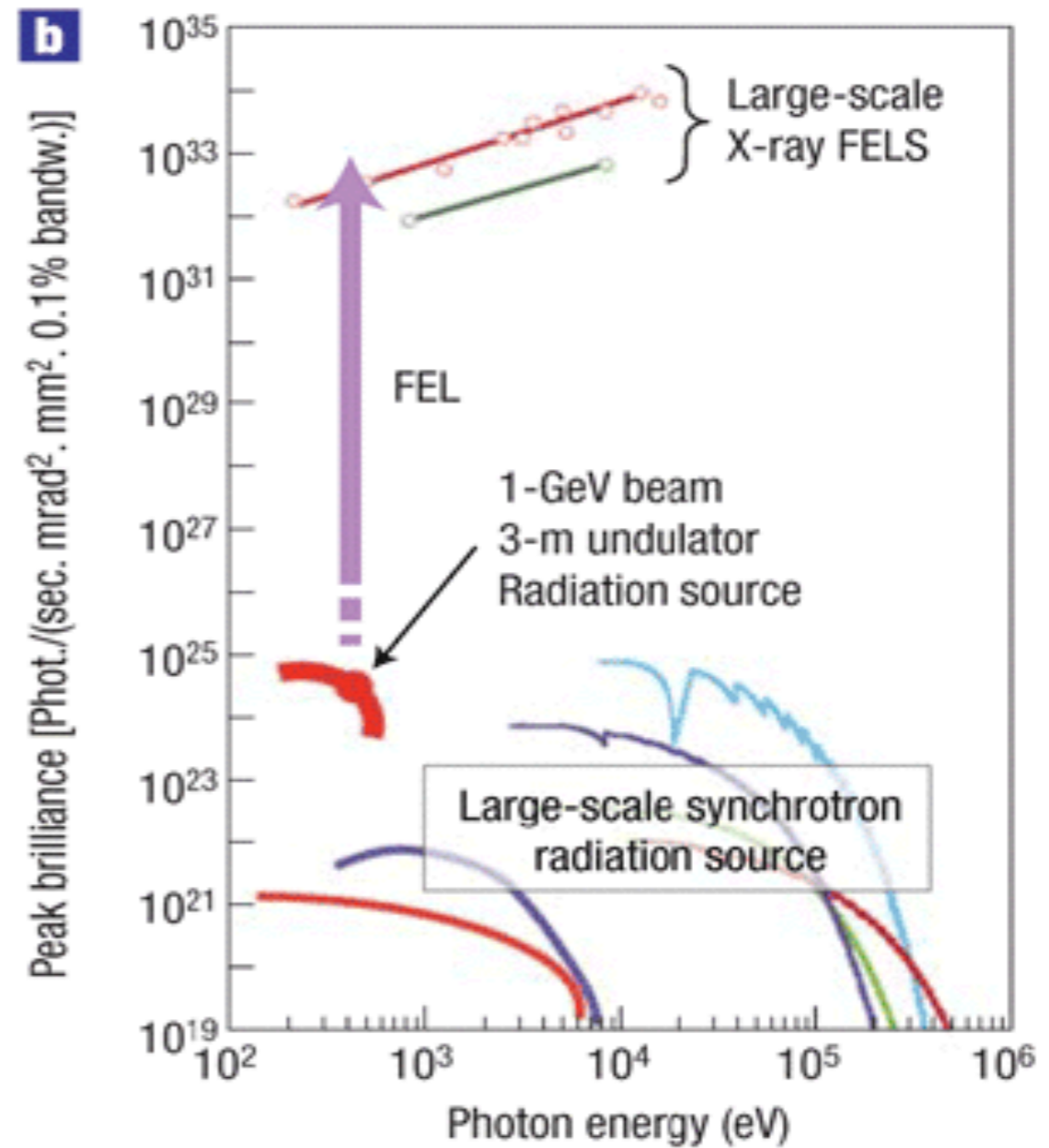
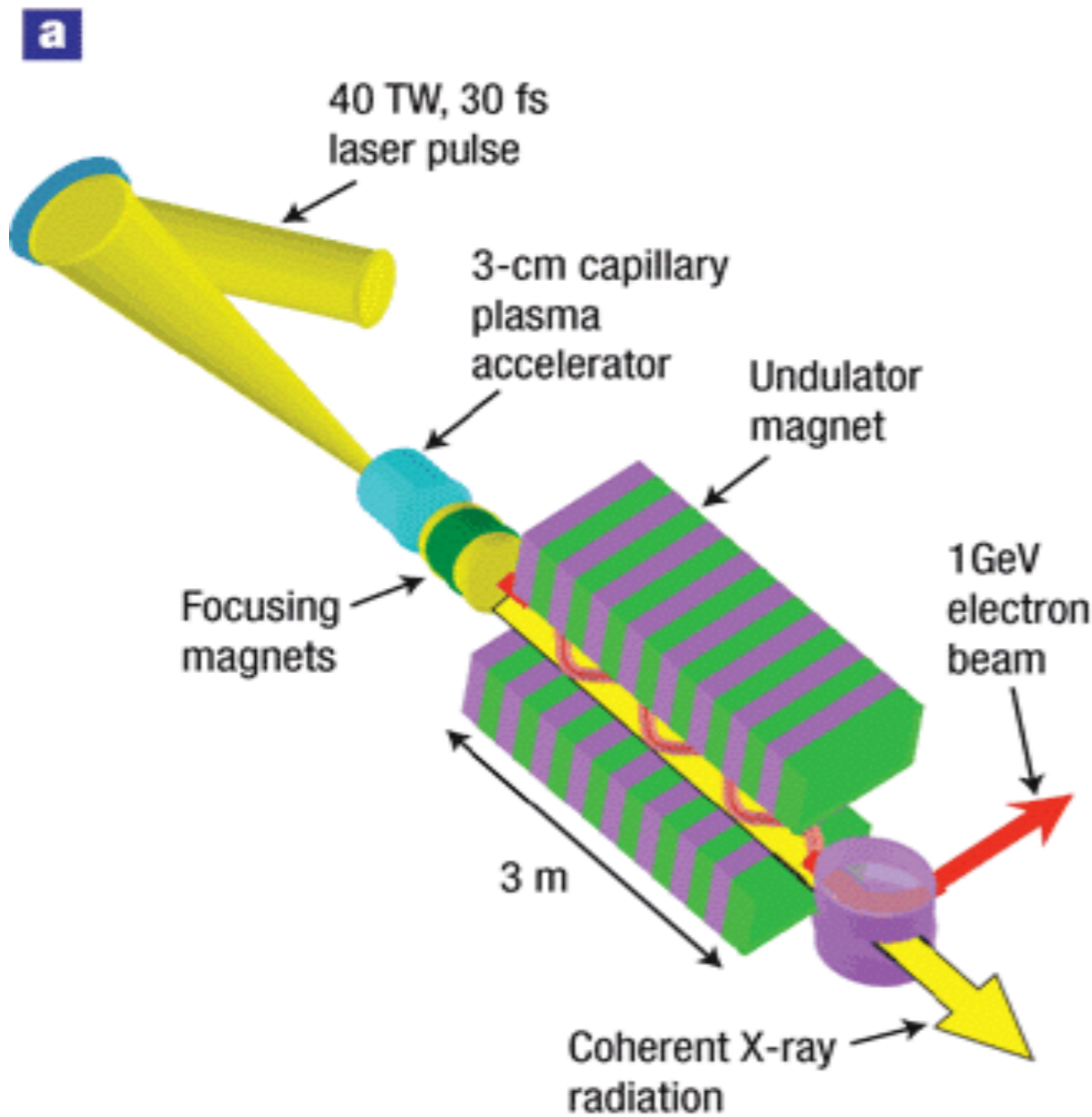
## Electrons

## Photons X

## Différence

A comparison of dose deposition with 6 MeV X ray an improvement of the quality of a clinically approved prostate treatment plan. While the target coverage is the same or even slightly better for 250 MeV electrons compared to photons the dose sparing of sensitive structures is improved **up to 19%**.

# Use of secondary radiation



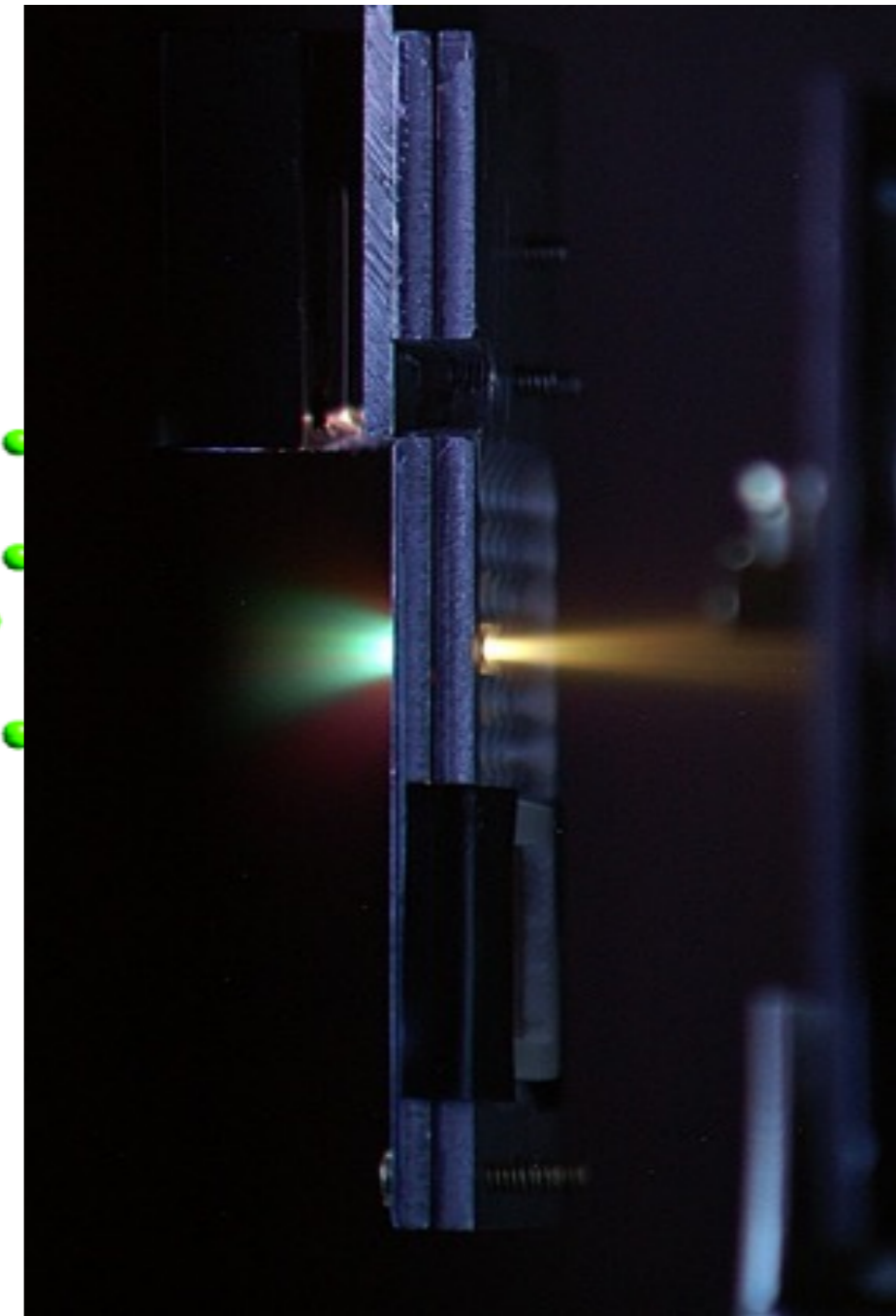
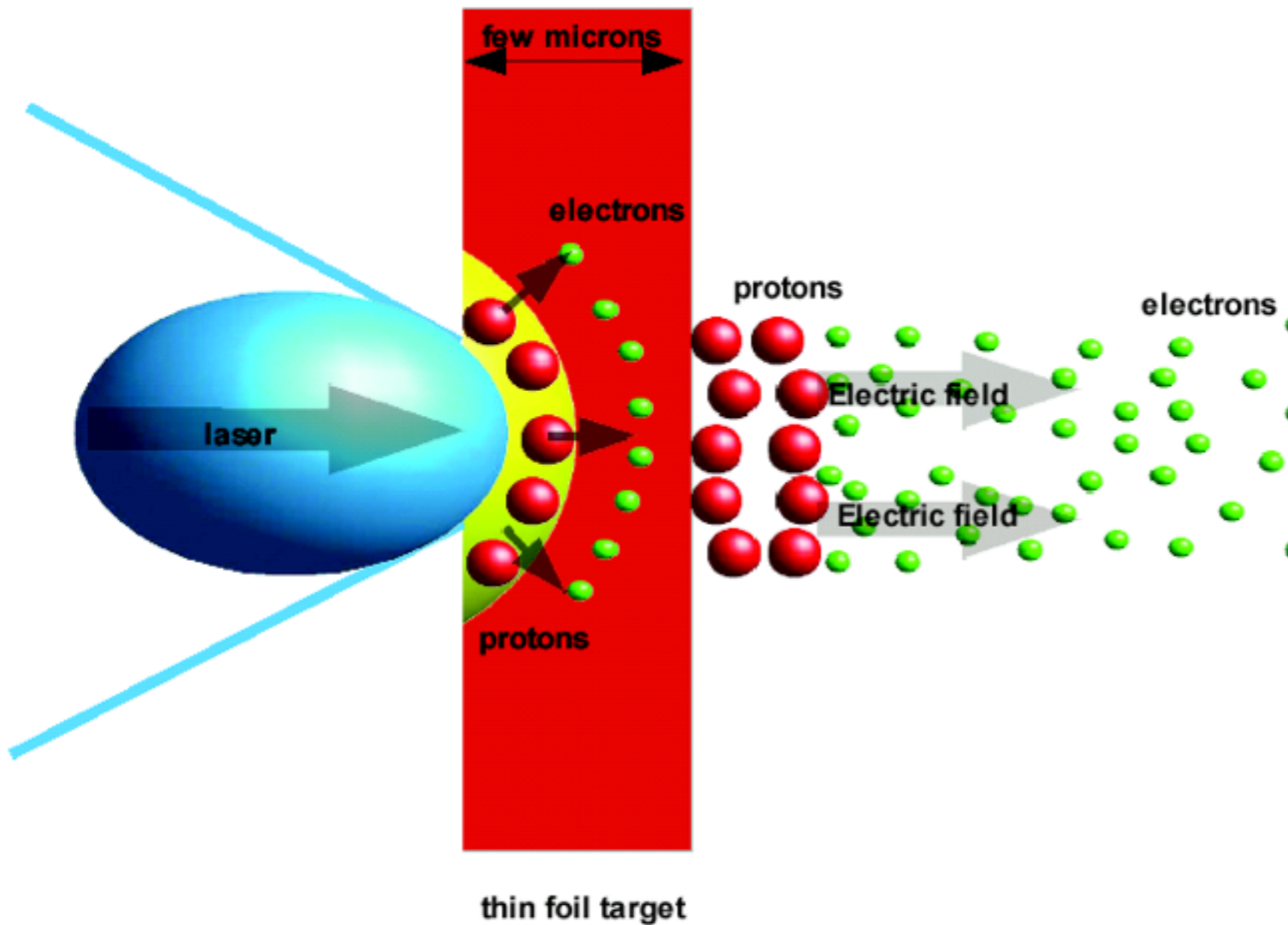
Kazuhisa Nakajima, *Nature Physics* 4, 92 - 93 (2008)

C. B. Schroeder, W. M. Fawley, E. Esarey, W. P. Leemans  
Lawrence Berkeley National Laboratory, Berkeley, CA 94720,  
USA

# Proton acceleration with lasers : Static electric fields



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# Requirements for ion acceleration

The requirements strongly depend on the application: a few examples

## ● Ion source as a new injector:

- Rep rate matched to conventional accelerator structures (e.g. 50 Hz)
- Ion energy a few tens of MeV
- Radial beam shaping for divergence optimization
- Ion species selectable
- Energy matched to particle number acceptable to acc structure

## ● Medical Application:

- Ion energy  $>250$  MeV for protons and  $>400$  MeV/u for e.g. Carbon (prob. no TNSA)
- High contrast
- Rep rate 10 to 30 Hz
- Energy stability better 3%
- Relatively low particle numbers required ( $10^{11}$  or  $10^9$  per patient)
- Uniform ion beam --> Laser beam shaping

# General remark

- For most applications high particle numbers are required
- Energy conservation asks for laser energy, not only intensity
- High contrast is key as most applications cannot rely on TNSA
- Wall plug efficiency becomes important

ICFA/ICUIL Meeting

**Table 1:** Comparison of wall-plug efficiency of various accelerators.

<i>Accelerator</i>	<i>Beam</i>	<i>Beam energy (GeV)</i>	<i>Beam power (MW)</i>	<i>Efficiency AC to beam</i>	<i>Note on AC power</i>
PSI Cyclotron	H <sup>+</sup>	0.59	1.3	0.18	RF + magnets
SNS Linac	H <sup>-</sup>	0.92	1.0	0.07	RF + cryo + cooling
TESLA (23.4 MV/m)	e <sup>+</sup> /e <sup>-</sup>	250 × 2	23	0.24	RF + cryo + cooling
ILC (31.5 MV/m)	e <sup>+</sup> /e <sup>-</sup>	250 × 2	21	0.16	RF + cryo + cooling
CLIC	e <sup>+</sup> /e <sup>-</sup>	1500 × 2	29.4	0.09	RF + cooling
LPA	e <sup>+</sup> /e <sup>-</sup>	500 × 2	8.4	0.10	Laser + plasma

# Example: Medical:

	Experiments	Status	Theory	Relevance to Therapy
<b>TNSA</b>	> 1999	>10 <sup>13</sup> ions, robust, reproducible	Analytical + 2D/3D simulations	+
<b>TNSA/BOA (Break-out- afterburner)</b>	> 2011	experimental evidence	2 D / 3D simulations	++(+)
<b>RPA</b>	>2008	experimental evidence not conclusive	2 D / 3D simulations >GeV	++
<b>Coulomb explosion</b>	-	-	2 D simulation	+
<b>Gas Jet - RPA</b>	> 2009	2 MeV observed	2D	++



International Committee for Future Accelerators  
Sponsored by the Particles and Fields Commission of IUPAP

**Beam Dynamics  
Newsletter**

No. 56

# Direct Treatment

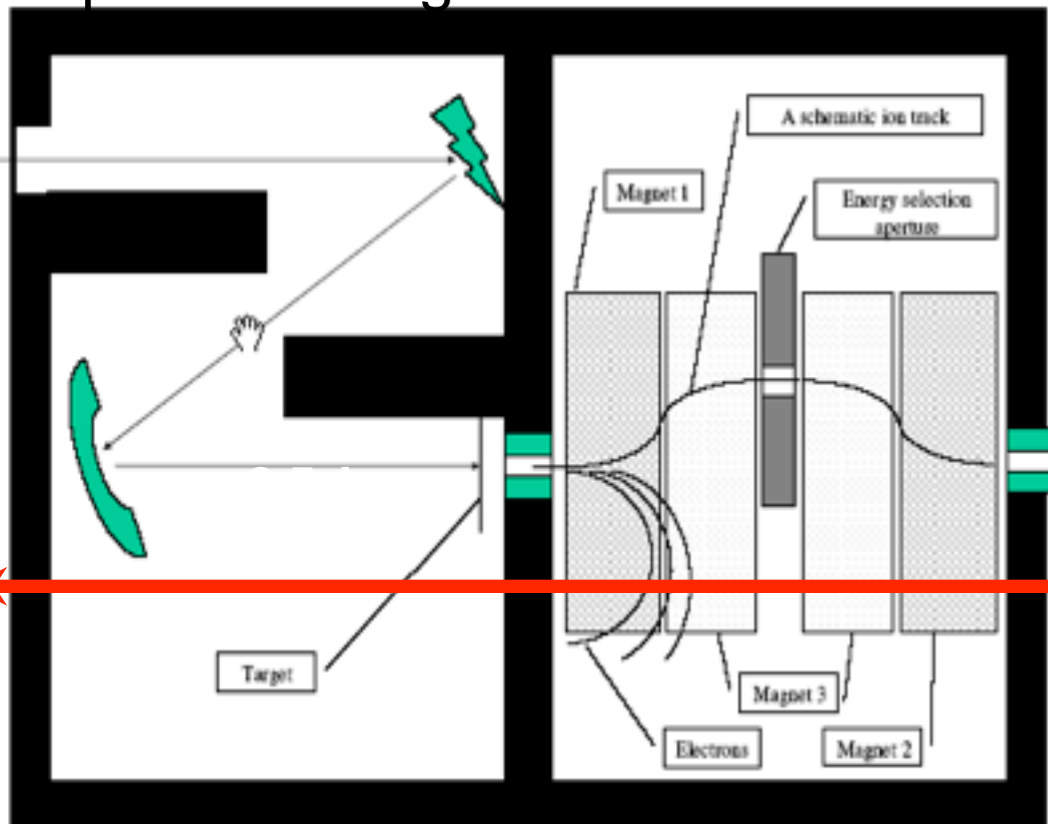
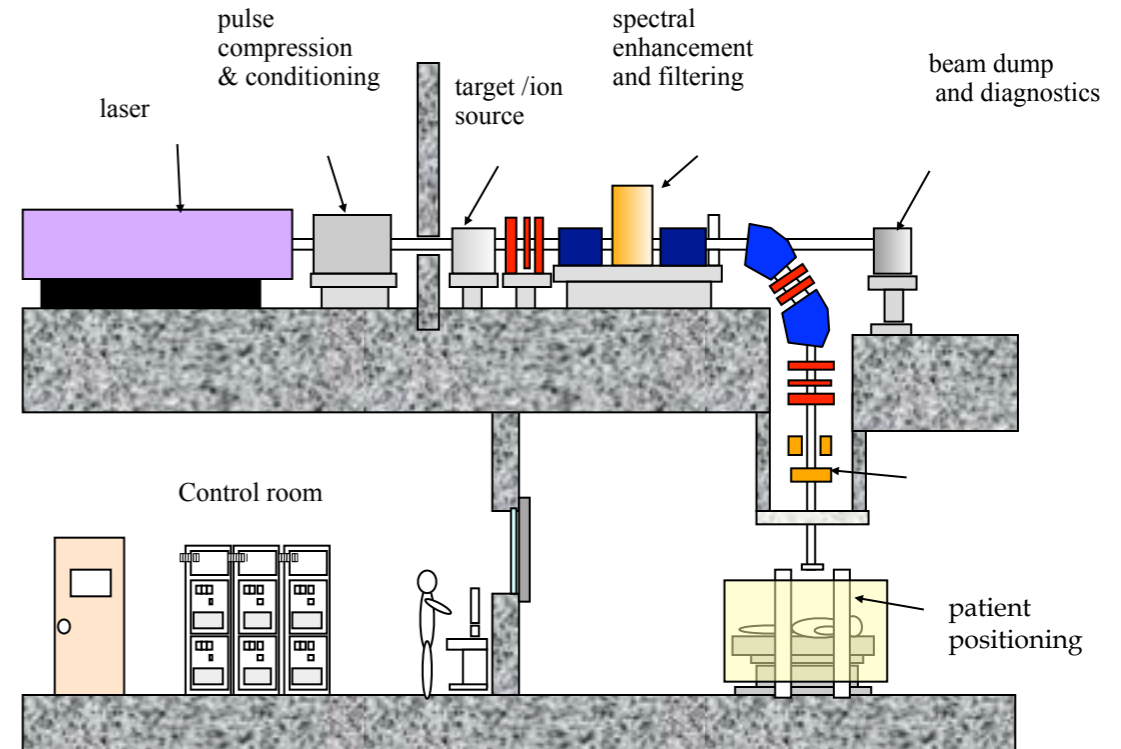
## Particle Selection and Beam Collimation

The Idea of direct laser treatment:

Laser must deliver 250 MeV Protons  
(400 MeV/u Carbon)

Non- Maxwellian Distribution  
Shielding Essential (gamma's, Neutrons)  
(Might require 10's of PW Laser Power  
@ high rep rate)

Spot Scanning???

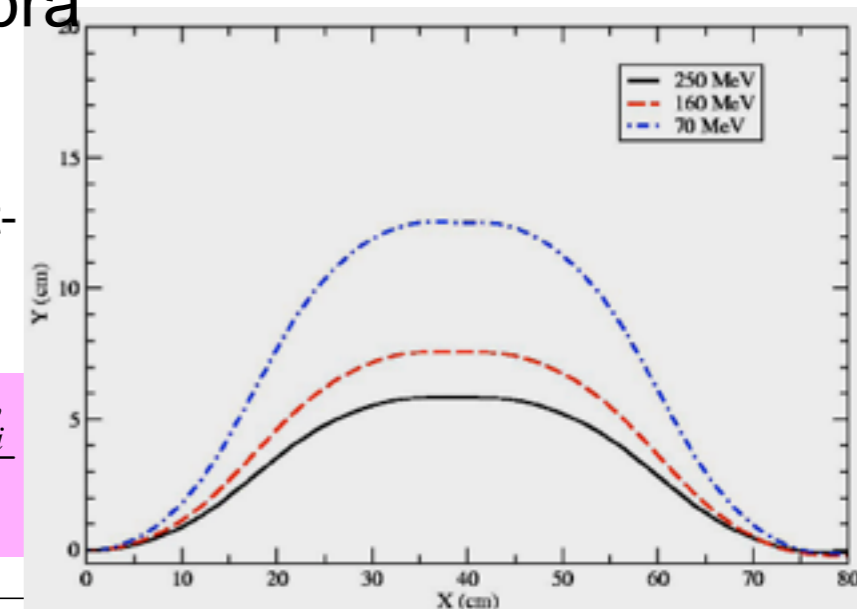


Movable aperture to select protons of desired energy with sharp beam penumbra

Nb<sub>3</sub>Ti superconducting coils can provide I = 85 A per loop with magnetic field to 4.4 Tesla by Biot-Savart law:

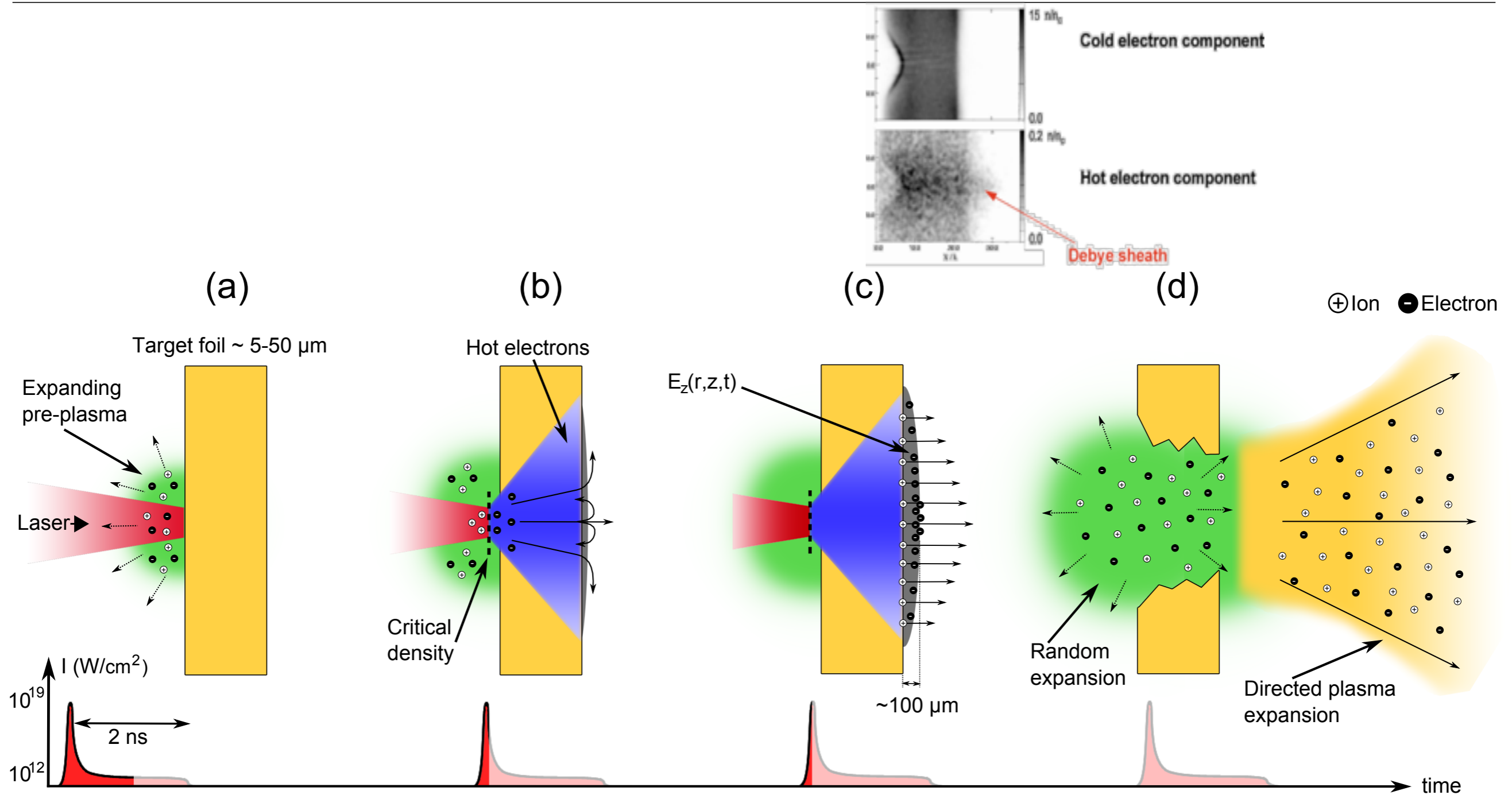
$$B = \oint \frac{\mu_0 I dl \times r}{4\pi r^3} = \sum_i^4 \int \frac{\mu_0 I dl \times r_i}{4\pi r_i^3}$$

C. Ma  
(Fox Chase Cancer Center)



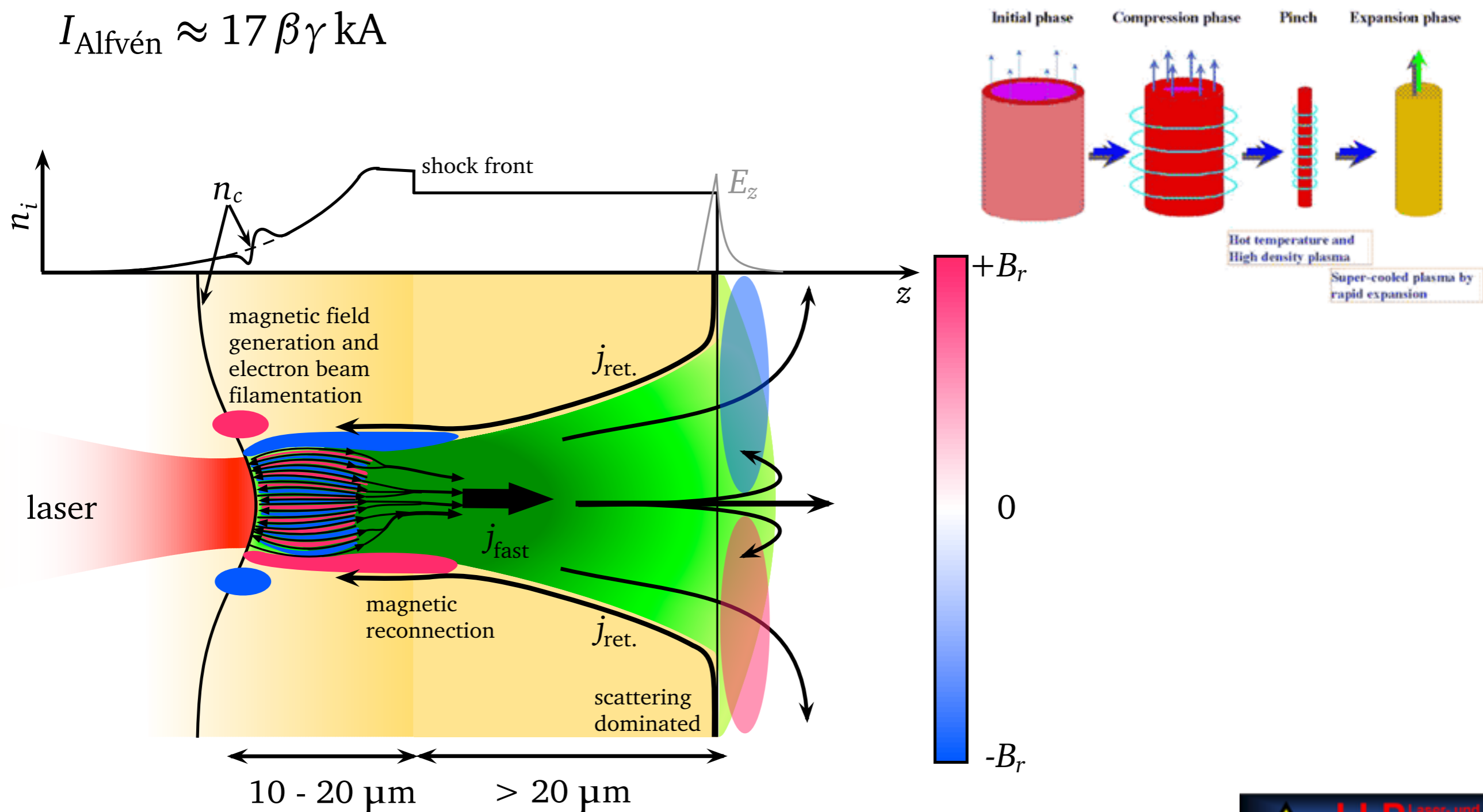


# Overview



# Electron driver III

$$I_{\text{Alfvén}} \approx 17 \beta \gamma \text{ kA}$$



- Electrons escape into vacuum at the rear side

- Charge separation: Potential  $\epsilon_0 \frac{\partial^2 \Phi}{\partial z^2} = en_e.$

- if one half space perfectly compensates the potential (target) and for  $z \rightarrow \infty$  it goes to infinity the field can be obtained analytically

$$n_e = n_{e,0} \exp\left(\frac{e\Phi}{k_B T_{hot}}\right) \quad \text{Ansatz } e\Phi/k_B T_{hot} = 2 \ln(\lambda z + 1)$$

the potential is 
$$\Phi(z) = -\frac{2k_B T_{hot}}{e} \ln\left(1 + \frac{z}{\sqrt{2}\lambda_D}\right)$$

- and the field 
$$E(z) = \frac{2k_B T_{hot}}{e} \frac{1}{z + \sqrt{2}\lambda_D}.$$

- if one takes the electron density and the spreading into account the Debye length is:

$$\lambda_D = \left( \frac{\epsilon_0 k_B T_{hot}}{e^2 n_{e,0}} \right)^{1/2} \approx 1.37 \mu m \frac{r_0 + d \tan \theta / 2}{r_0} \frac{\sqrt{1 + 0.73 I_{18} \lambda^2} - 1}{I_{18}^{7/8}}$$

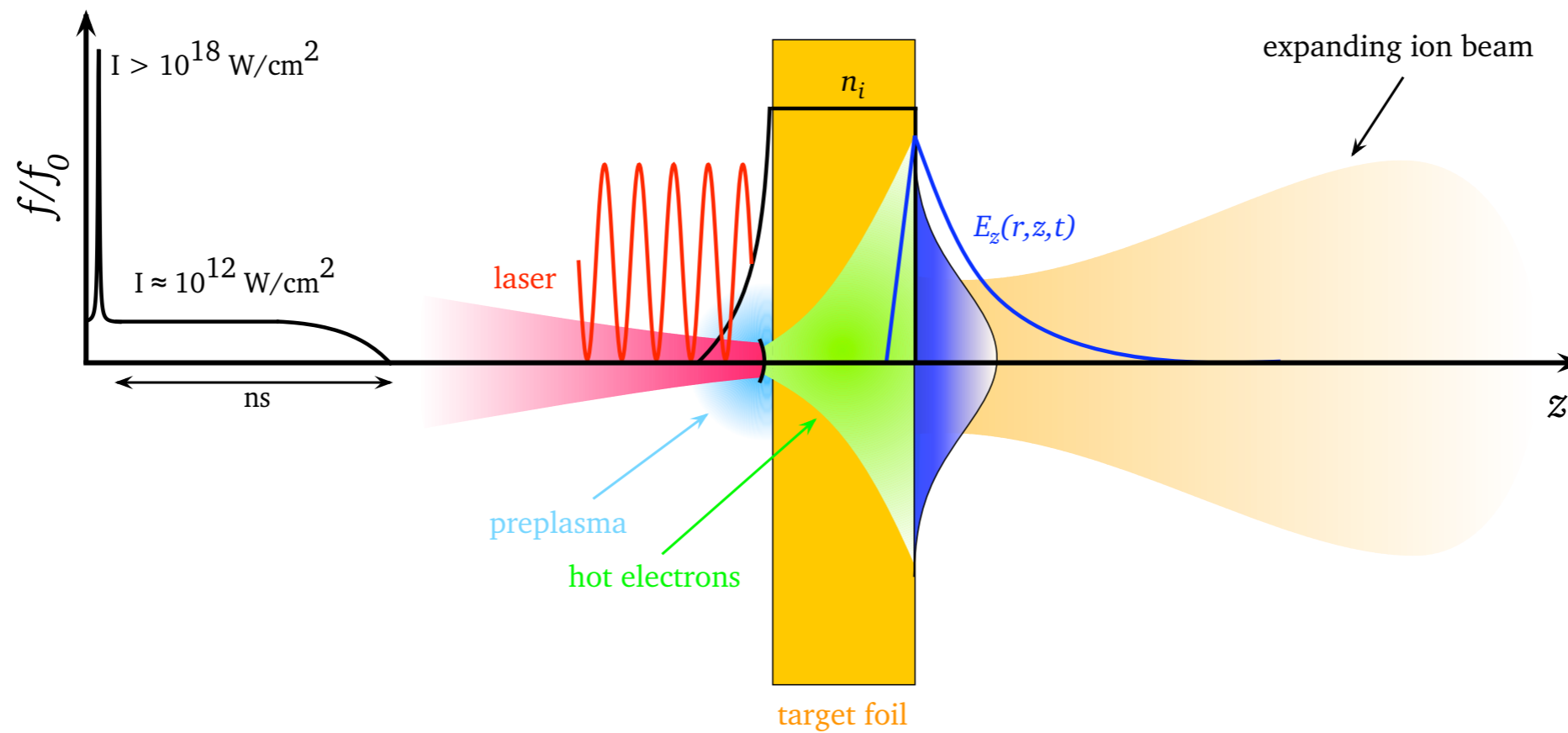
the Debye length: very short (0.6  $\mu m$  in the example), depends on laser intensity and target thickness

$$\begin{aligned} E_{max}(z=0) &= \frac{\sqrt{2} k_B T_{hot}}{e \lambda_D} \\ &\approx 5.2 \times 10^{11} V/m \frac{r_0}{r_0 + d \tan \theta / 2} I_{18}^{7/8} \\ &= 9 \times 10^{10} V/m \frac{r_0}{r_0 + d \tan \theta / 2} E_{12} E_{12}^{3/4}. \end{aligned}$$

Field can easily exceed several  $10^{12}$  V/m

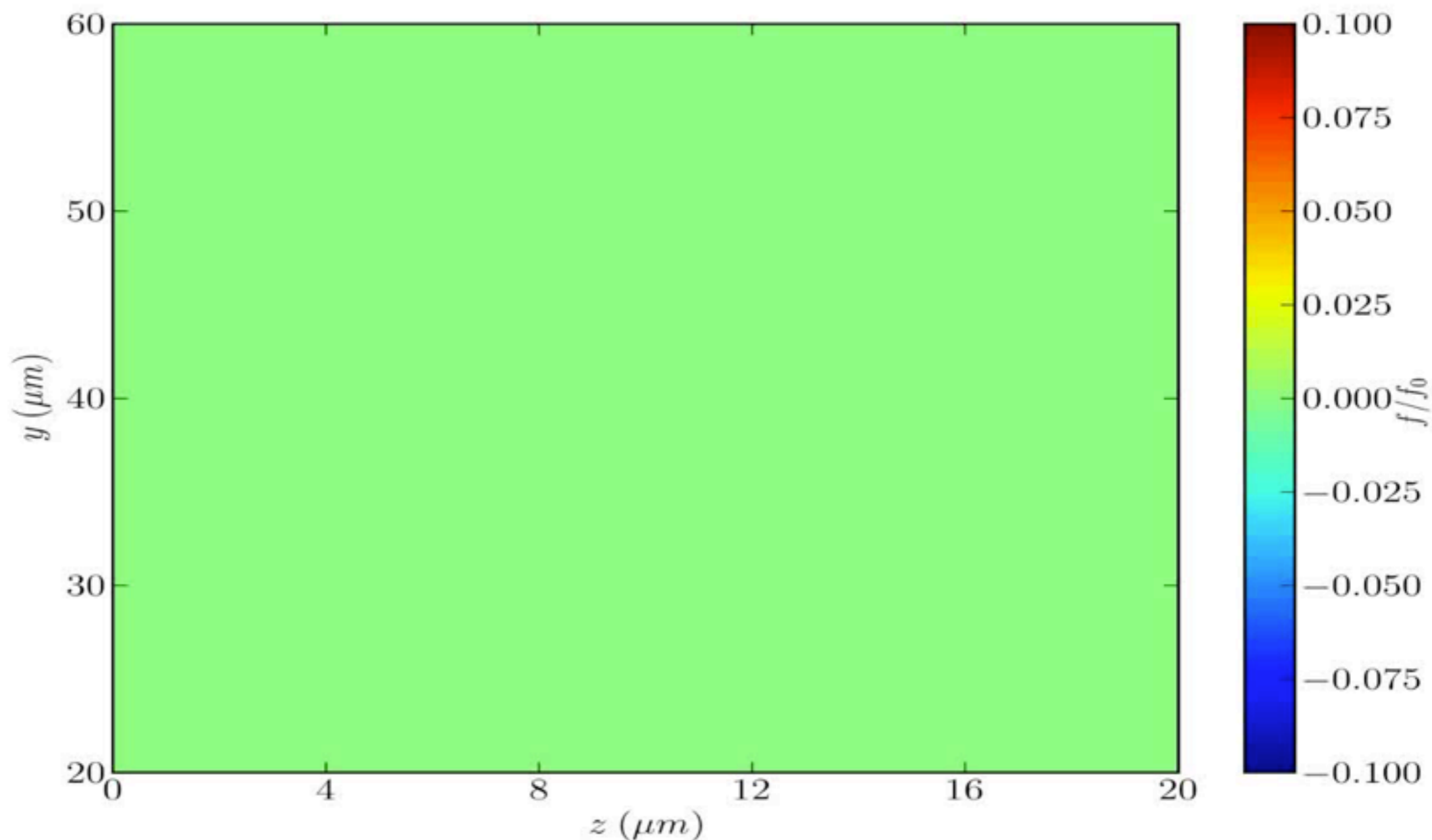
Immediate field ionization of the atoms at the rear surface

# TNSA III

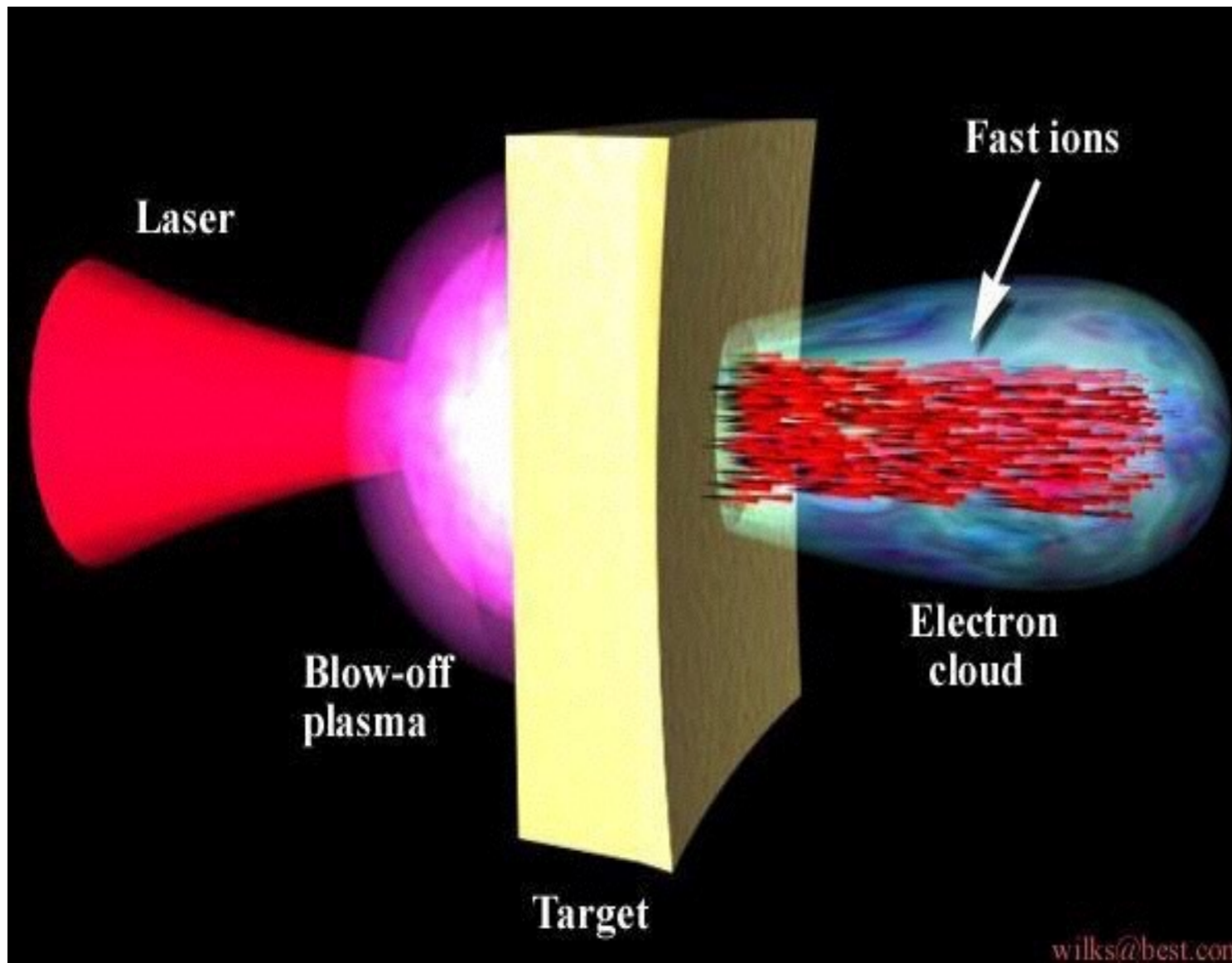


# to summarize... a TNSA simulation

$t = 00000.00 \text{ fs}$   
 $x = 25.00 \mu\text{m}$   
 $min = 0.00$   
 $max = 0.00$



# Ion Beam Properties



Number of Protons:  
 $10^{14}$  (LLNL- Petawatt)

Pulse duration:  
few Picoseconds

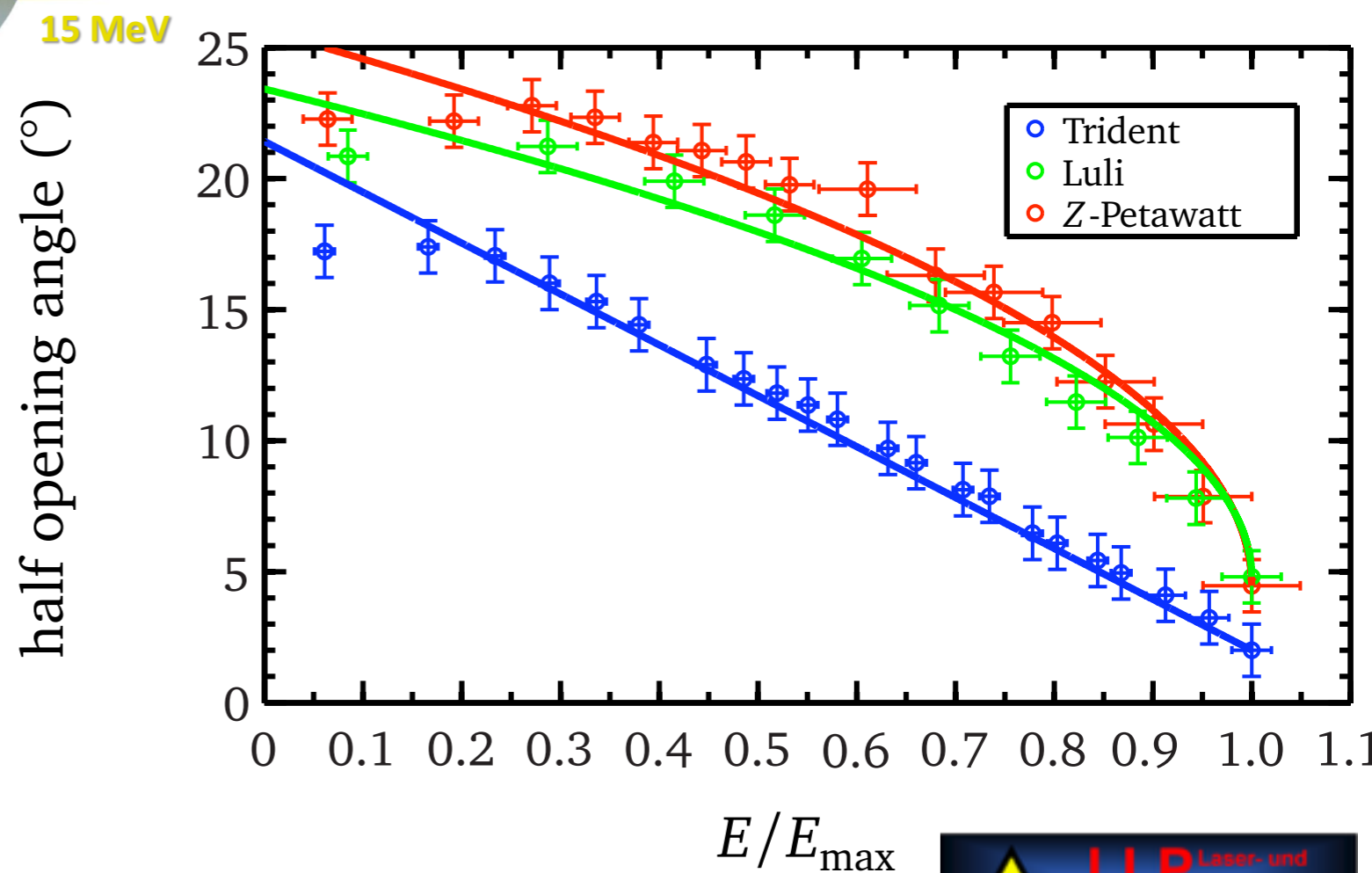
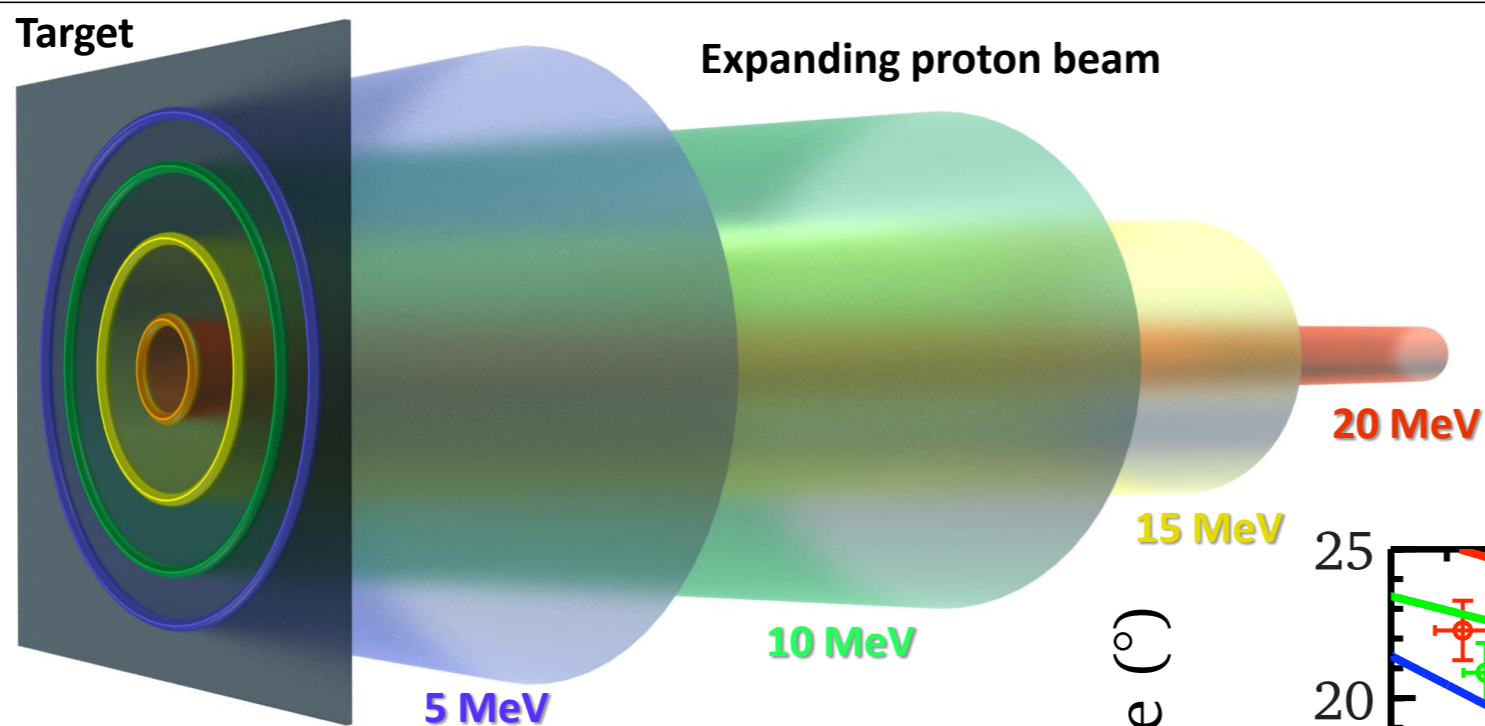
Maximum Energy:  
60 MeV (LLNL Petawatt)

Divergency:  
<math>10^\circ</math> for highest  
energies

Always normal to the  
rear

Origin of Protons: Surface Contaminants

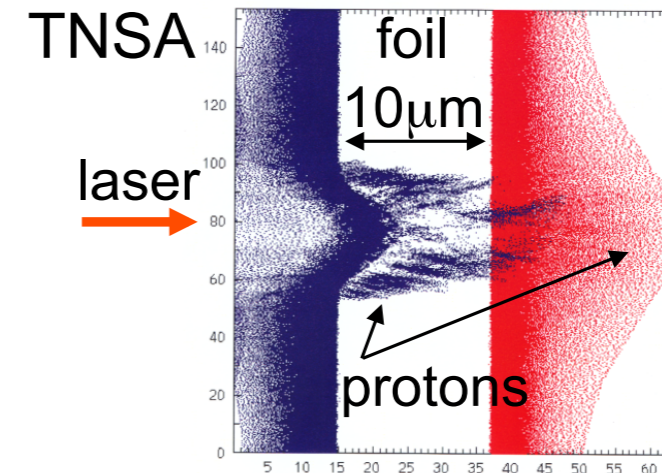
# Spatial Particle Distribution





# Ion Acceleration Mechanisms

Ion acceleration mechanism	Acronym	Ion Accel. process
Target-Normal Sheath Acceleration <i>S. Hatchett et al., Phys. Plas. 7, 2076 (2000)</i>	TNSA	Charge separation  GeV protons? <b>X</b>



# TNSA vs. BOA



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Accessible with moderate contrast lasers  
Micrometer sized targets  
Spectrum limited to 70 MeV  
Surface acceleration

High contrast lasers needed  
Sub-Micrometer sized targets  
Ion energies exceeding 120 MeV/u  
Volume acceleration  
Heavy ions (deuterons) at same speed as protons  
Lower EMP and less debris

# High contrast Lasers (PHELIX)

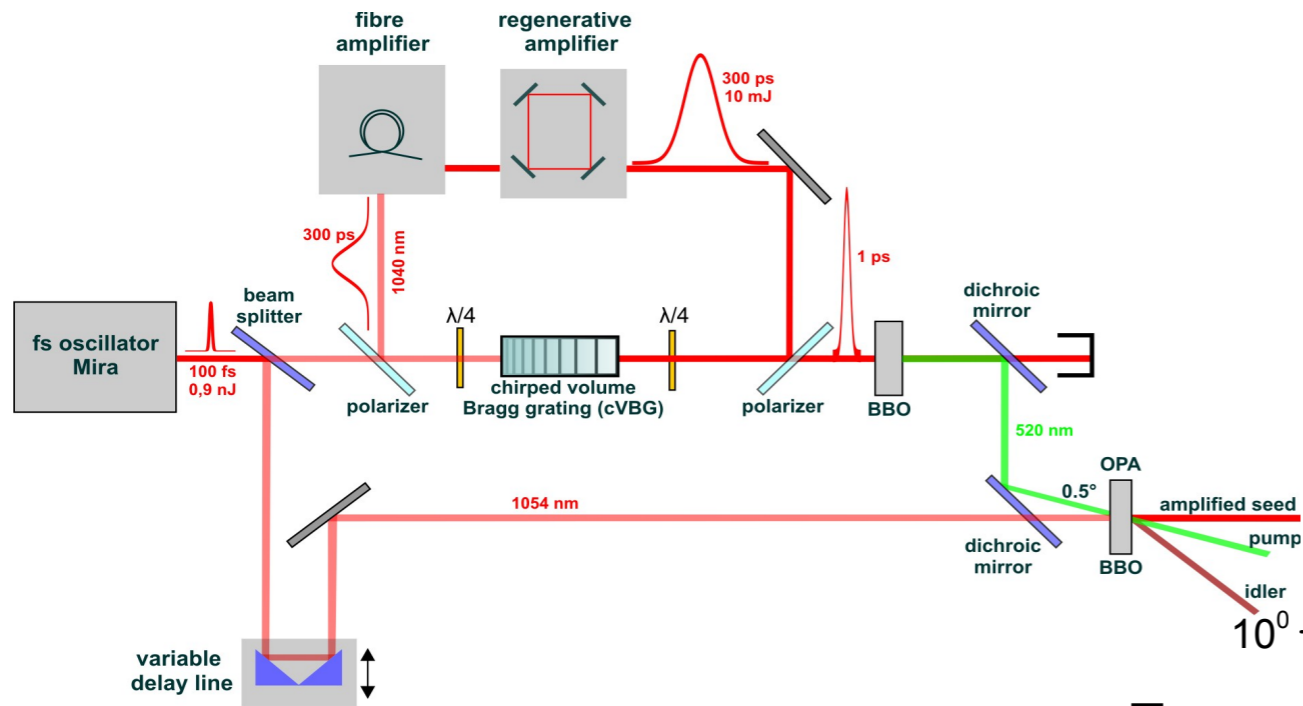
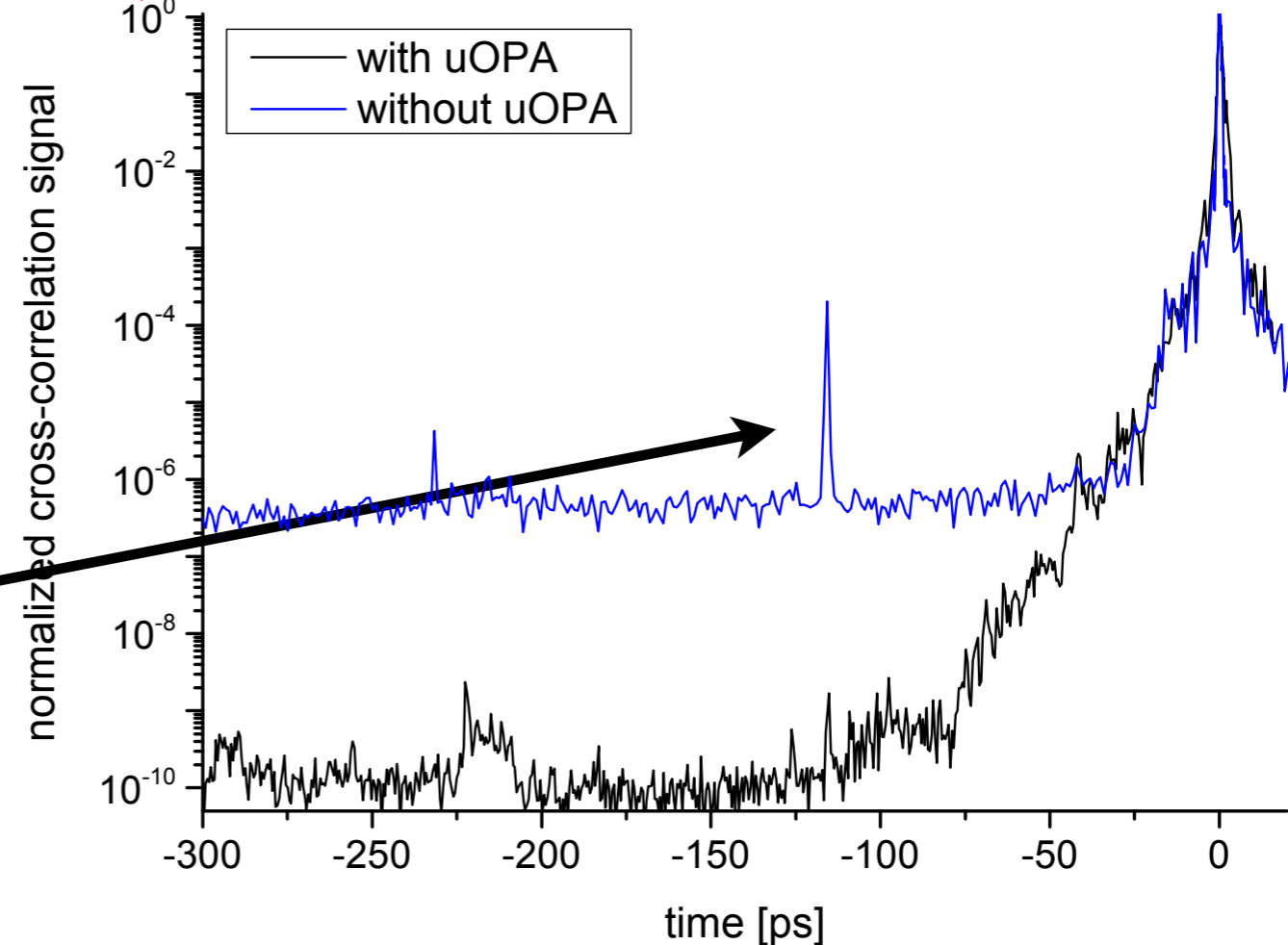
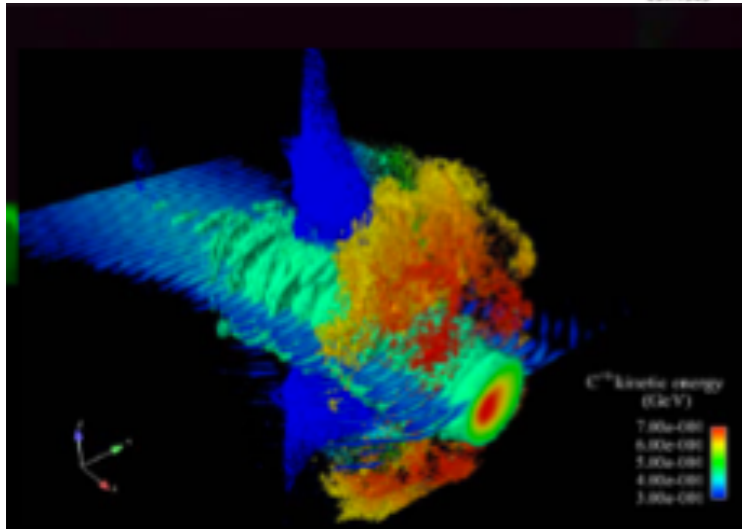


Fig. 2: Setup of the contrast-boosting module

Prepulse from the oscillator



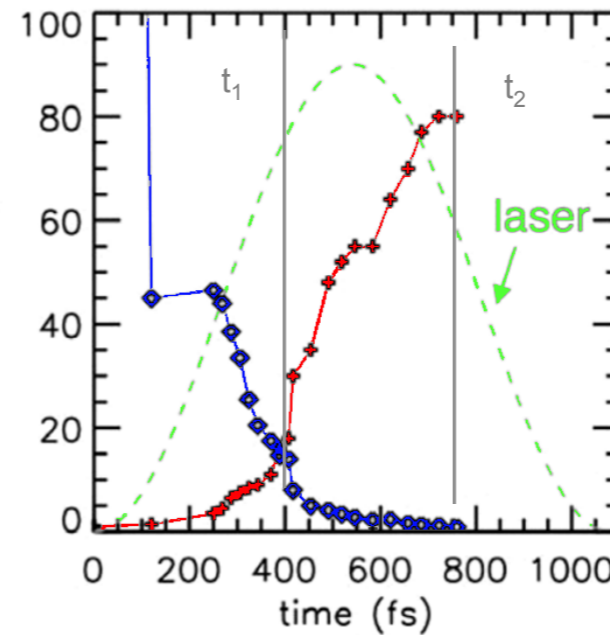
# Break out Afterburner (BOA)



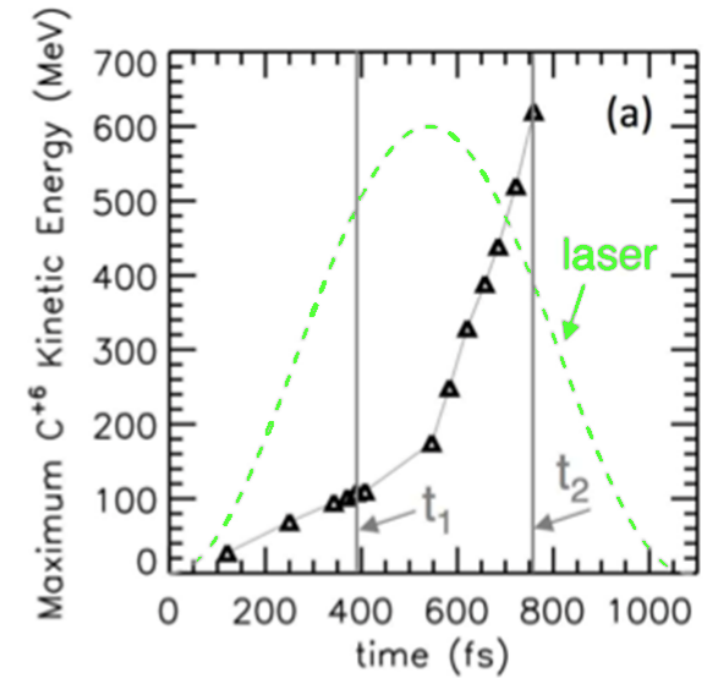
$t_1$ : relativistic transparent  
 $n' > 1 \geq n'/\gamma$

$t_2$ : classically underdense  
 $n' < 1$

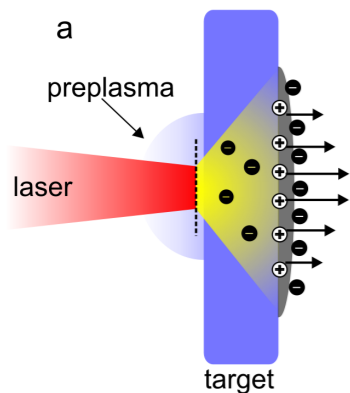
$\gamma, n_e/n_c$



Yin, et al., *Laser and Particle Beams* 24 (2006), 1–8  
Yin, et al., *Phys. Plasmas* 14, 056706, (2007)  
Yin, et al., *Phys. Plasmas* 18, 063103 (2011)



Albright, et al., *Phys. Plasmas* 14, 094502 (2007)  
Yin, et al., *Phys. Rev. Lett.* 107, 045003 (2011)



- a) Target Normal Sheath Acceleration (TNSA) phase
- b) Intermediate phase
- c) Laser Breakout Afterburner (BOA) phase

$$\eta = \sqrt{1 - \frac{n_e}{\gamma n_c}} = \sqrt{1 - \frac{\omega_p^2}{\gamma \omega_L^2}} \quad \omega_p = \sqrt{\frac{e^2 n_e}{\epsilon_0 m_e}}$$

VPIC: 100nm CH<sub>2</sub> target & Trident laser with  $2 \times 10^{20} \text{W/cm}^2$

Max. energy	proton	carbon
Ideal laser	132 MeV	450 MeV
Real laser	121 MeV	447 MeV

# First BOA experimental results TRIDENT/ PHELIX

High contrast beam

thin targets 200-900nm

cryo targets

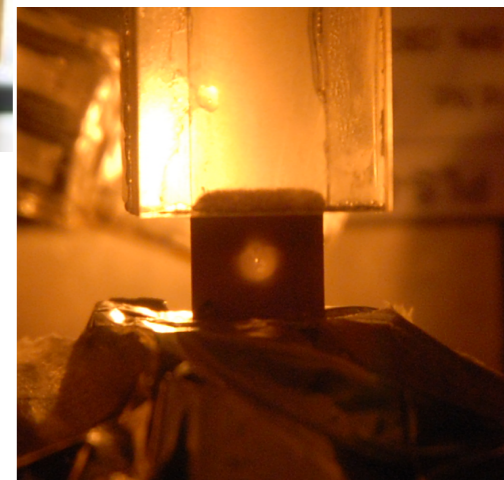
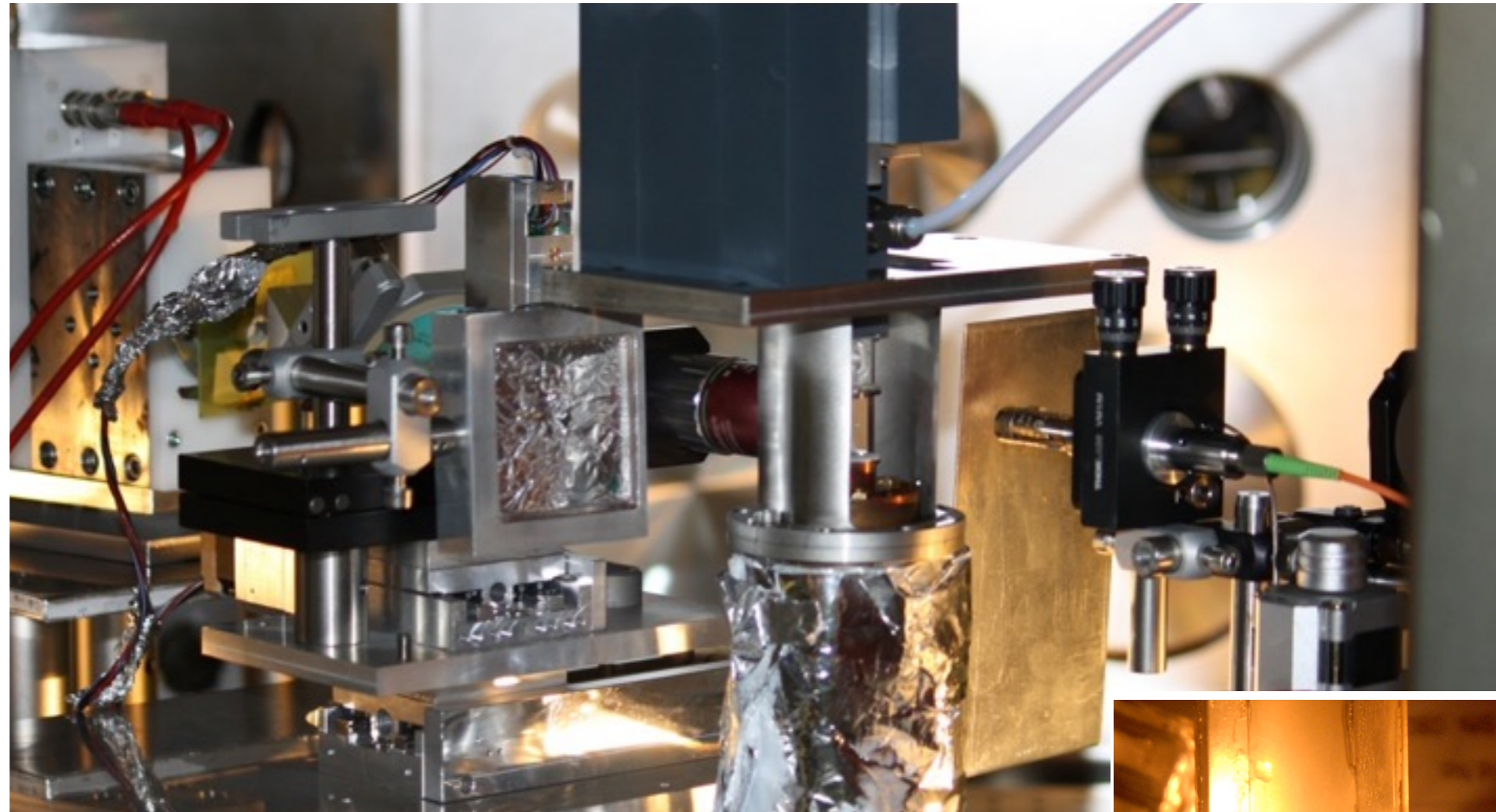
short  $f\#$  parabola

About 80 J on Target

Using CH foils around  
100 MeV protons

First test on Cryo target  
55 MeV protons

Pure deuteron beam

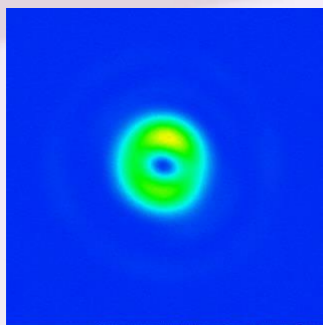


# Tailoring and transport of ion beams - the LIGHT project -



# LIGHT

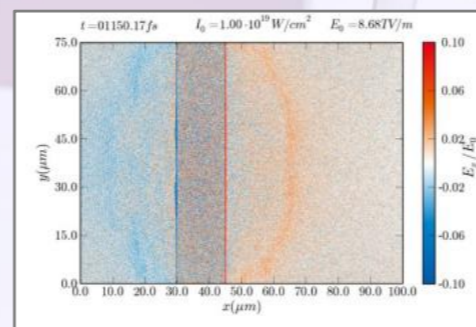
## Laser Ion Generation, Handling and Transport



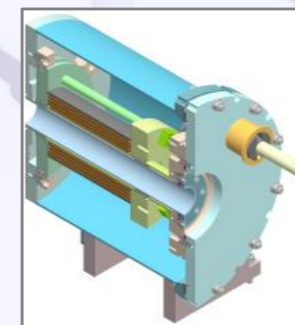
laser shaping



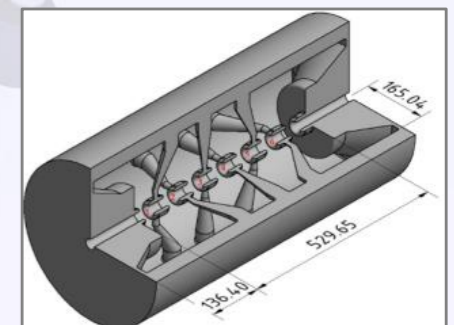
target shaping



simulation

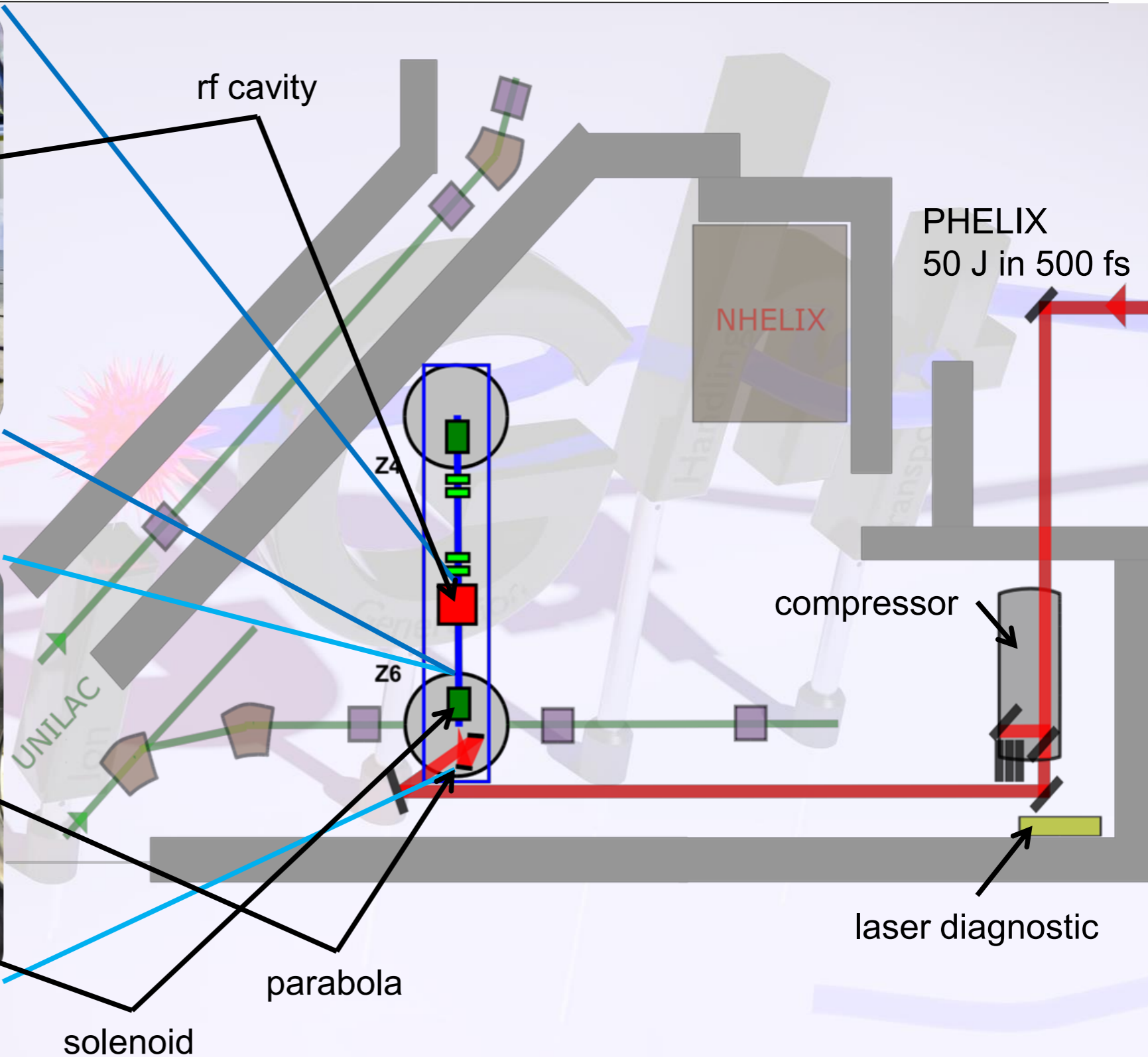
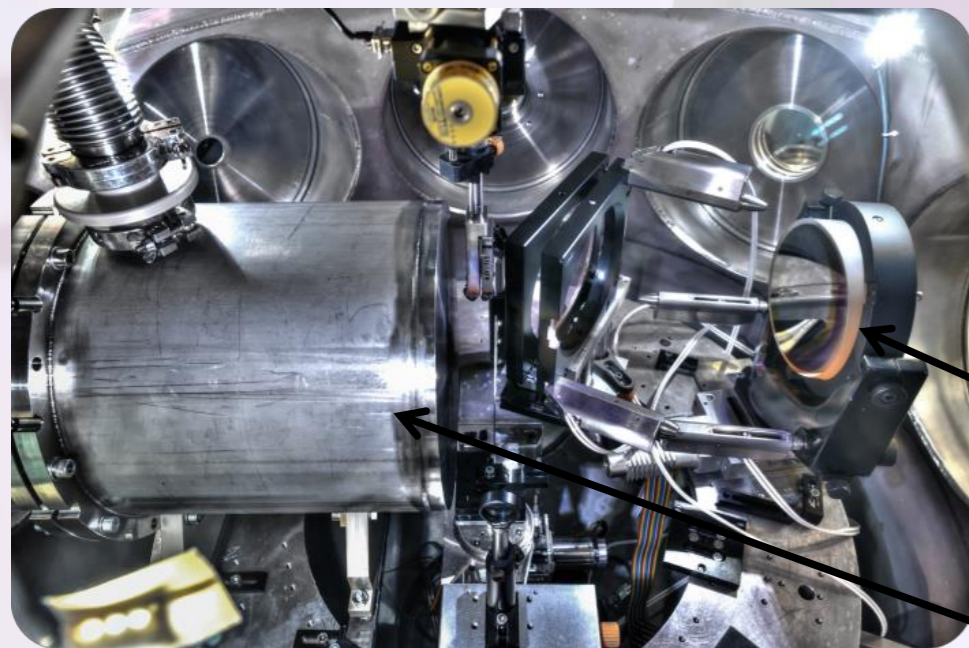
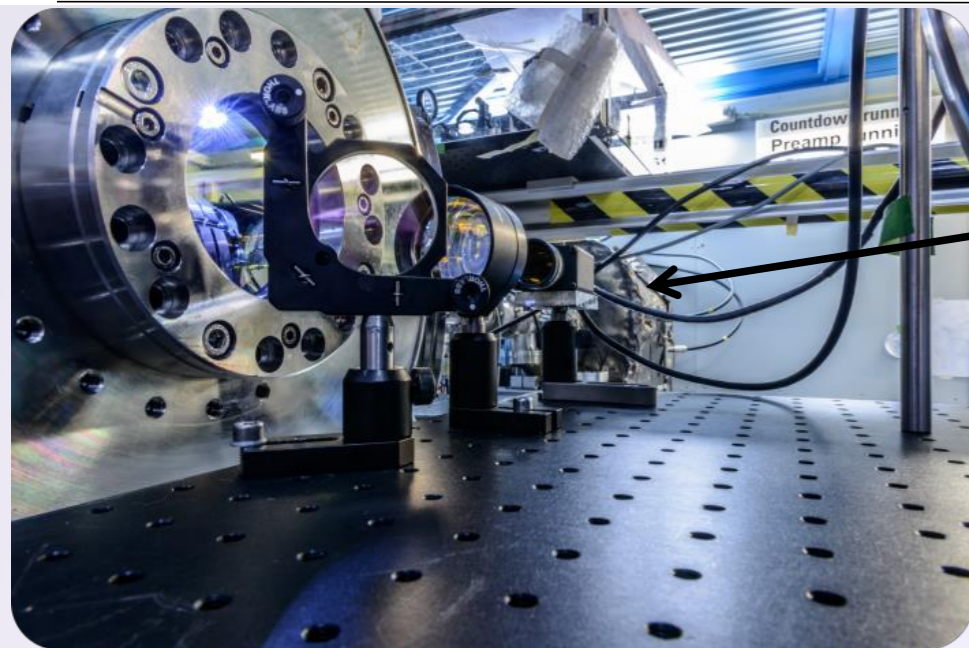


transport



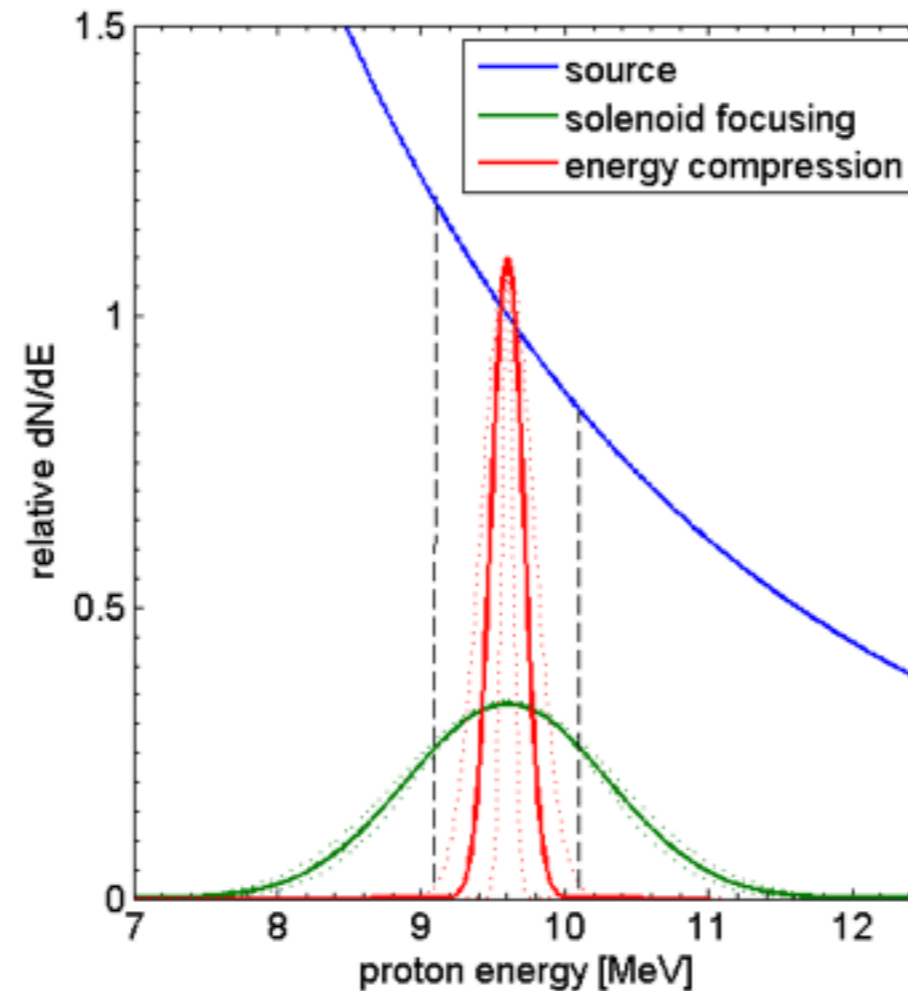
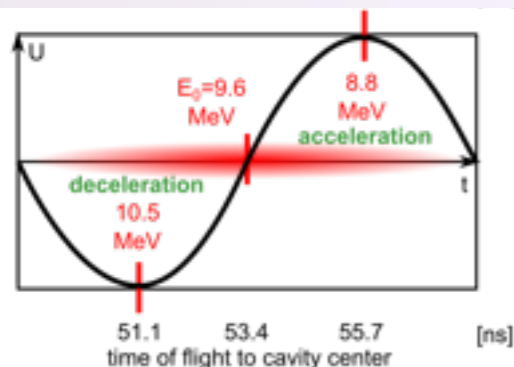
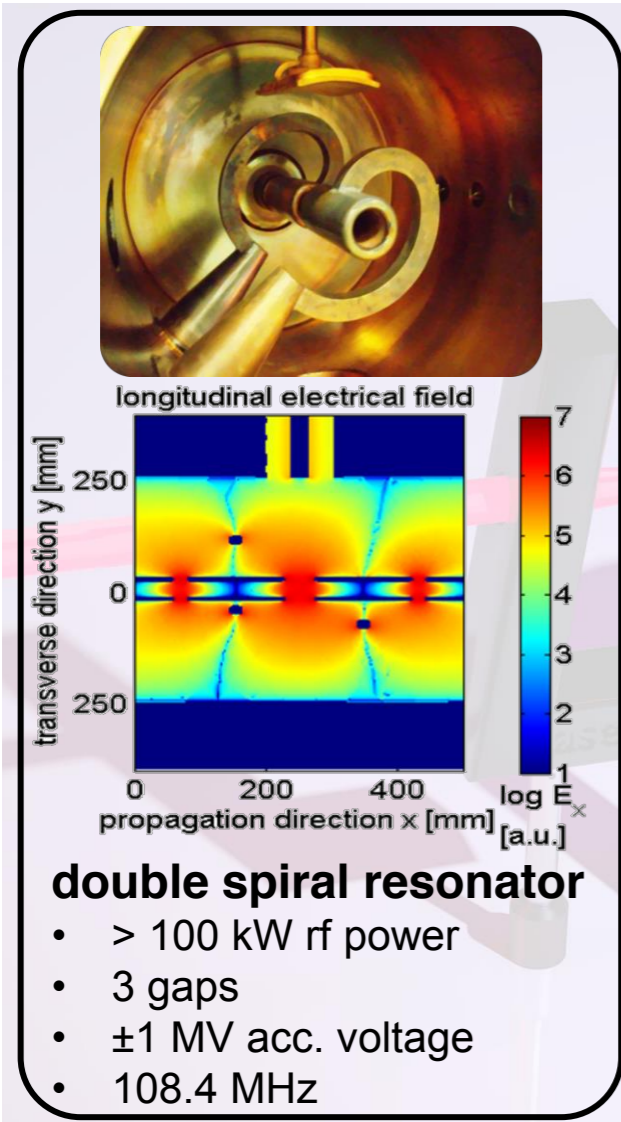
post acceleration

# The LIGHT Beamline



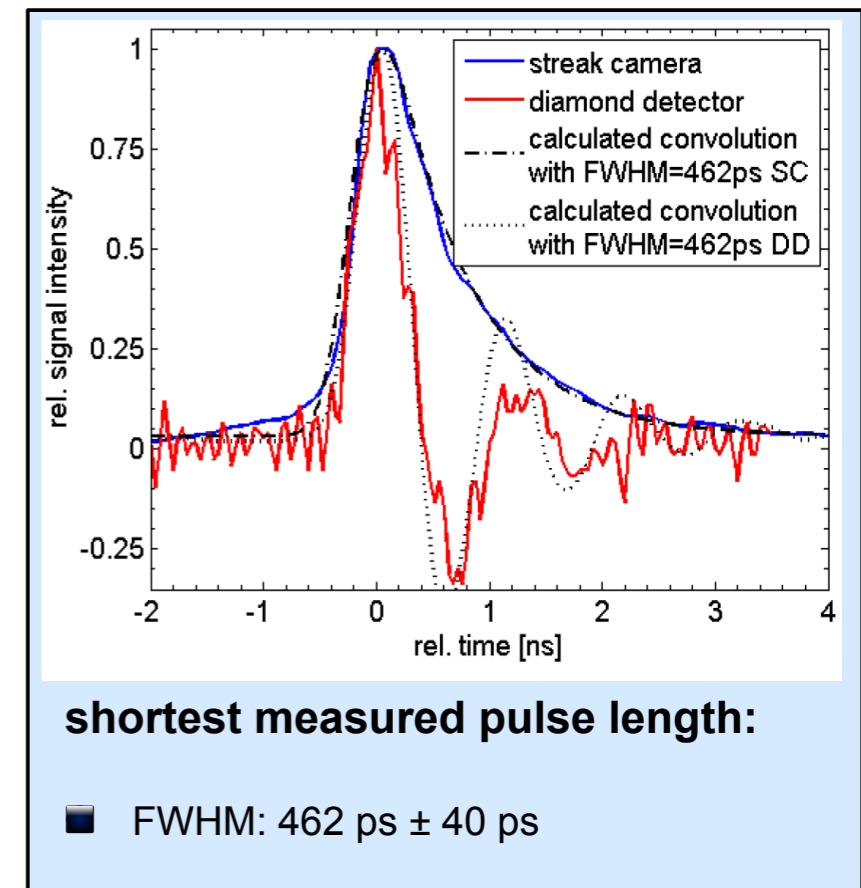
# phase rotation

S. Busold et al., PR-STAB 17, 031302 (2014)



**Energy selection and width for 9.6 MeV :**  
 **$18.0 \pm 3.0$  % due to chromatic focusing of the solenoid**  
 **$2.7 \pm 1.7$  % using the cavity**

$n_p = 5 \times 10^8$   
 $t = (462 \pm 40)$  ps  
 $I = 100$  mA





# Laser Ion Parameter for Treatment



	Spot scanning	Passive formation	Comments
Protons / laser shot	$2 \times 10^7$	$2 \times 10^8$	reach 2 Gy by accumulation
# transverse	10x10 spots	10 reps for lateral uniformity	
Energy steps	10	10	$\Delta E/E = \pm 5\%$
Reps $\rightarrow$ dose spec. (~30% intensity jitter)	40	40	10 reps 4 gantry directions
Total # shots per fraction	10000	1000	$\frac{1}{4}$ applied
Duration of fraction Laser rep rate	5 min 30 Hz	1.5 min 10 Hz	

# Laser requirements for medical applications:

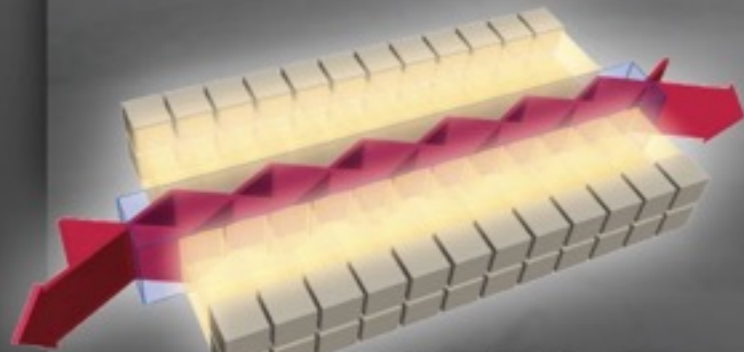
- Based on the need for RPA or BOA acceleration to reach the required particle energy:

	laser proton	laser carbon
Rep rate (spot/passive)	30 Hz / 10 Hz	
Laser intensity (W/cm <sup>2</sup> )	1-3 10 <sup>21</sup>	1-3 10 <sup>22</sup>
Pulse duration (fs)	50-150	
Rise time (fs)	<20	
Contrast (5 ps / 500 ps)	<10 <sup>-8</sup> / 10 <sup>-12</sup>	<10 <sup>-9</sup> / 10 <sup>-13</sup>
Laser energy stability	1-5%	
Spot radius (μm)	5	
Peak power (PW)	1-3	10-30
Pulse energy (J)	50-150	500-1500
Average power (kW) 10 Hz (30 Hz)	0.5-1.5 (1.5-4.5)	5-15 (15-45)

Laser cost assumption	<10 M€	~15 M€
Laser wavelength (nm)	800-1054	
Efficiency	1-10%	
Polarization	lp/cp	
Laser beam quality	diffraction limit	
Pulse stability	0.01	
Laser pointing (μrad)	1-10	
Laser availability	12 h/day (50% duty factor)	
Failure rate	<2%	

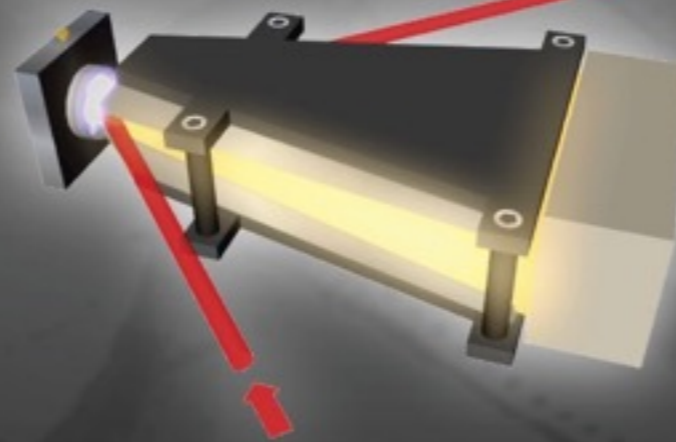
# Architectures for rep-rated 100 J systems

**HALNA – Japan**  
Dr. Toshiyuki Kawashima



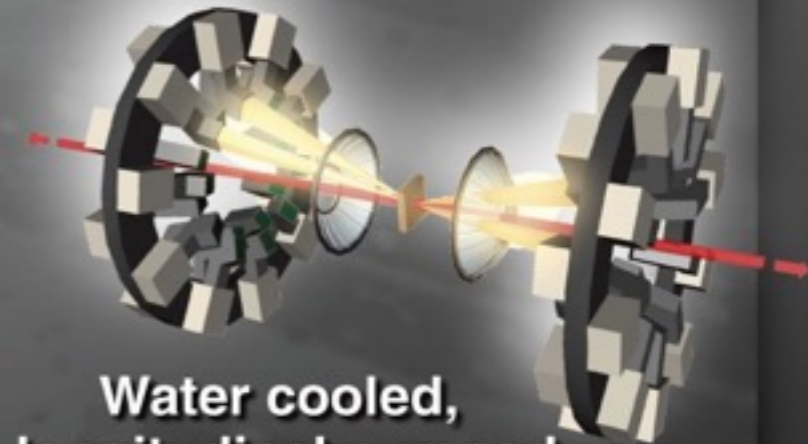
Water cooled,  
side pumped  
Nd:Phosphate slab

**Lucia – France**  
Dr. Jean-Christophe Chanteloup



Water cooled,  
longitudinal pumped  
Yb:YAG disk

**Polaris – Germany**  
Dr. Joachim Hein



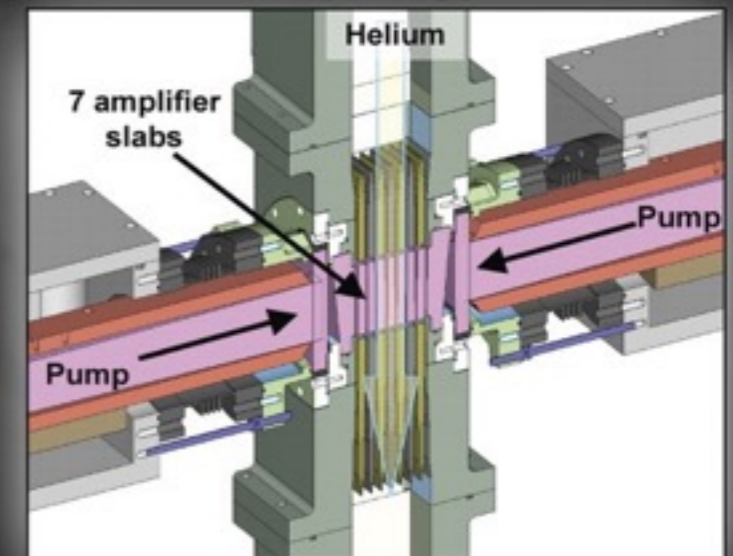
Water cooled,  
longitudinal pumped  
Yb: Fluorophosphate disk

**LALTI – Ohio**  
Dr. Linn Van Woerkom



Short pulse system  
with rep-rated targets

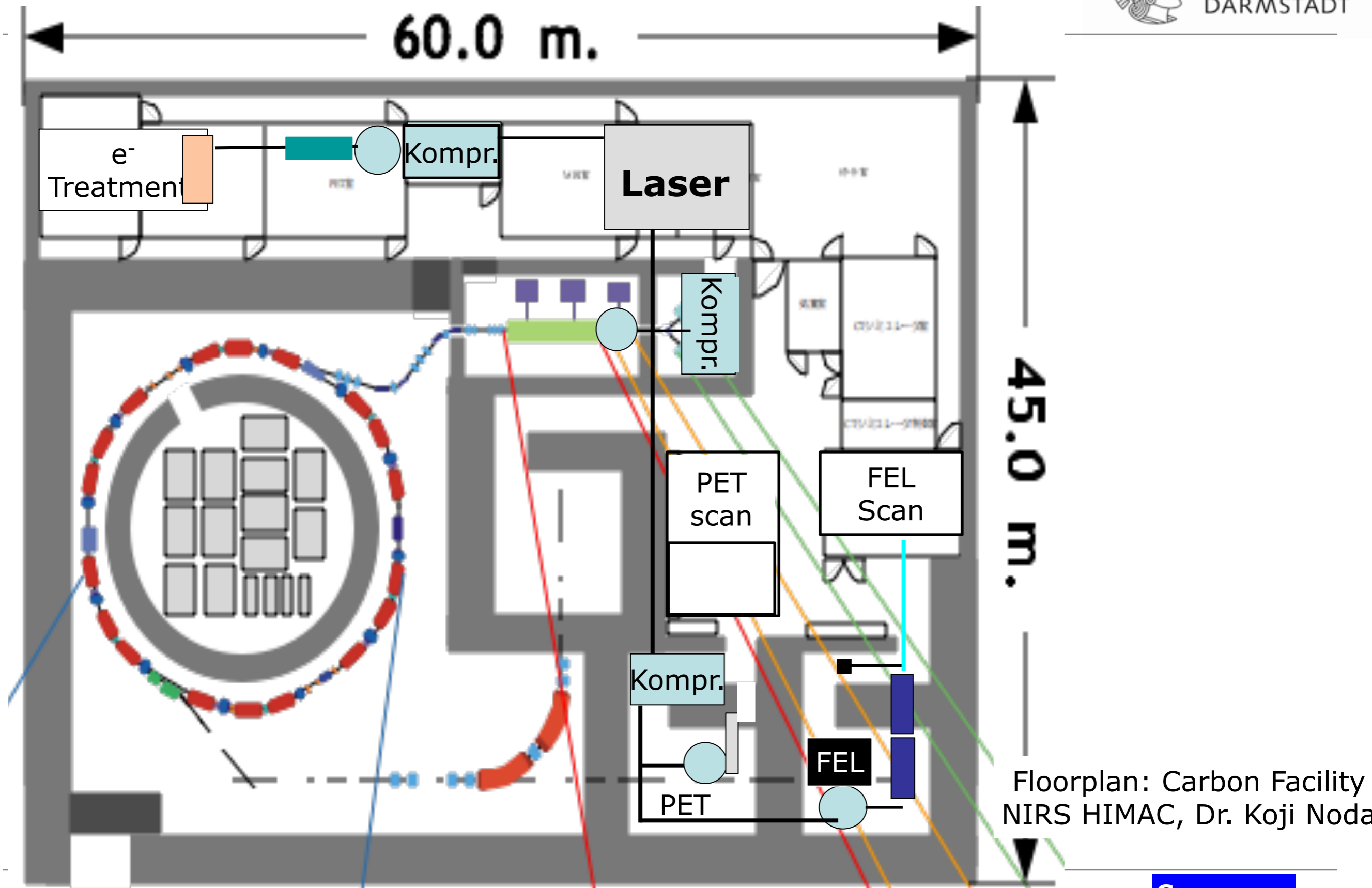
**Mercury – USA**  
Dr. Andy Bayramian



Gas cooled, face pumped  
Yb:S-FAP slab

- At an energy level of a few tens of Joules high-rep rate, DPSSL's are supposed to be available in a few years (Mercury, JAERI, POLARIS, LOA)  
you can buy 100 TW table top systems @ 10 Hz already
- „monochromatic“ ion beams have been observed for the first time
- Excellent beam emittance would allow for smaller apertures, higher gradients, smaller accelerators (PET Production, tumor therapy)
- the driving laser system could also be used for e-beam generation, PET-isotope production and coherent x-ray diagnostics at the same time

# Vision for a UHI Laser in Medicine



Floorplan: Carbon Facility  
NIRS HIMAC, Dr. Koji Noda

# Summary

- As of today laser acceleration has a theoretical potential to compete with conventional drivers for therapy
- Extremely high beam quality. How to preserve during transport?
- Laser injector into synchrotron
  - „good“ parameter match (based on PHELIX data)
  - 10 Hz Petawatt laser in reach
  - hard to compete with linac technology, space charge?!!
- „full energy laser“ scenario lacks data
  - small cones ( $\sim 2-3^\circ$ ), smaller production  $\Delta E/E (\sim 10^{-3} - 10^{-2})$
  - $> 100$  Hz laser systems, nm foils (feasibility?)
  - reproducibility, precision unknown
- Feasibility Experiment next step – GSI „best“ place to do it!

**Thank You !**

# High intense and energetic X-ray beams for phase contrast imaging in medical diagnostics: technical requirements and applications

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Paola Coan

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HABS



Munich-Centre for Advanced Photonics (MAP)  
Munich, LMU, Germany



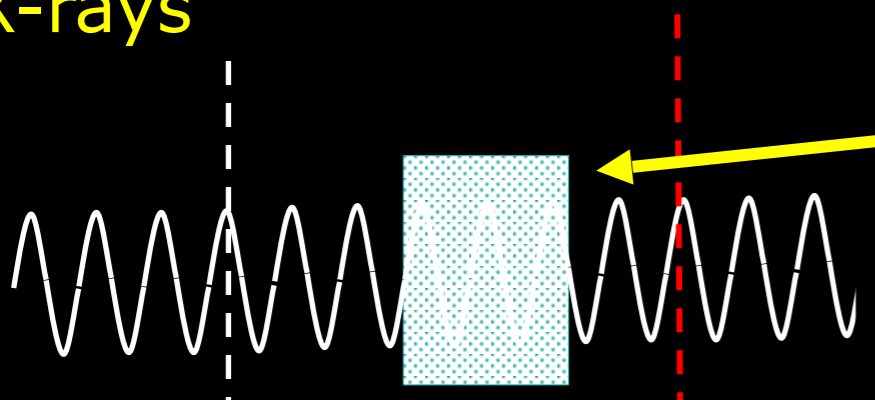
Biomedical Beamline  
European Synchrotron Radiation Facility  
Grenoble, France

**X-rays through the matter**

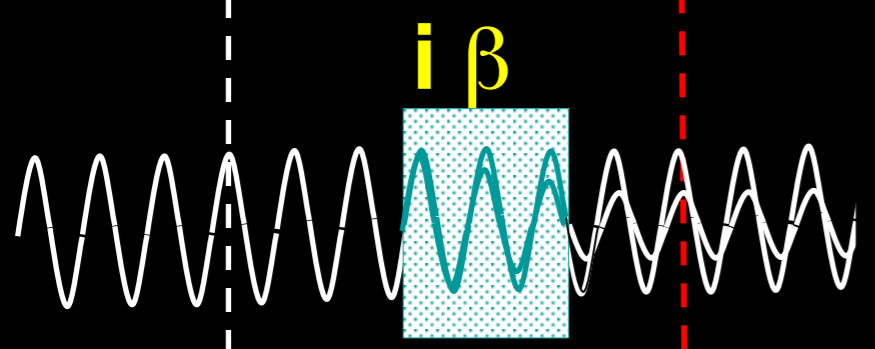
X-rays

$$n = 1 - \delta + i\beta$$

complex-valued index of refraction

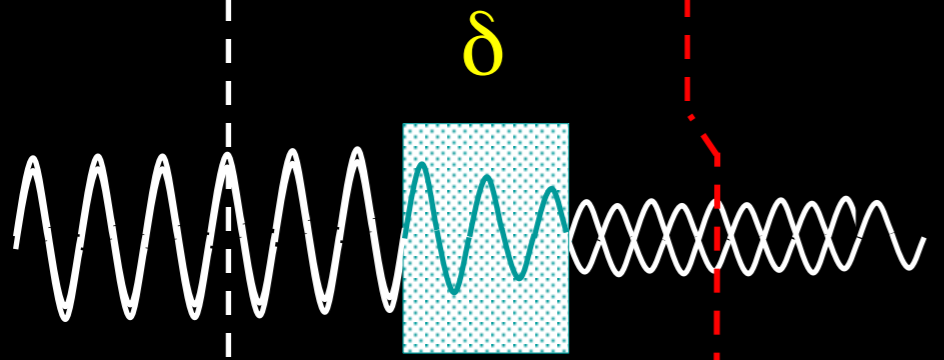


In the diagnostic X-ray energy range:



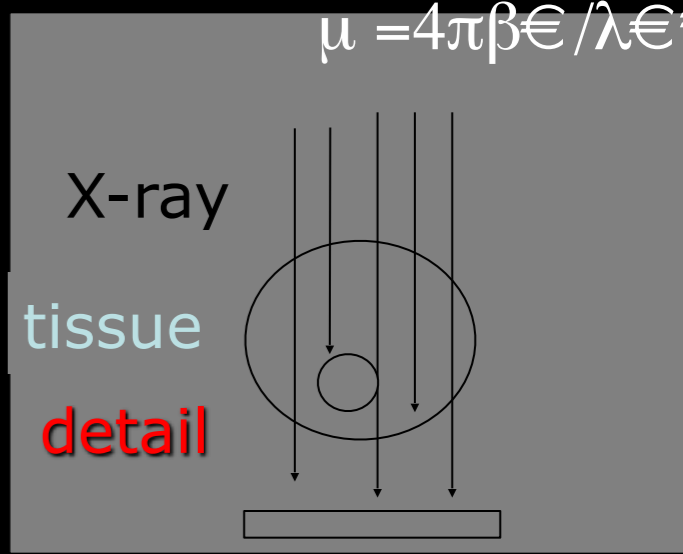
→ Amplitude modulation  
 - Photoelectric effect (+ Incoherent scattering)

**ABSORPTION-based RADIOLOGY**

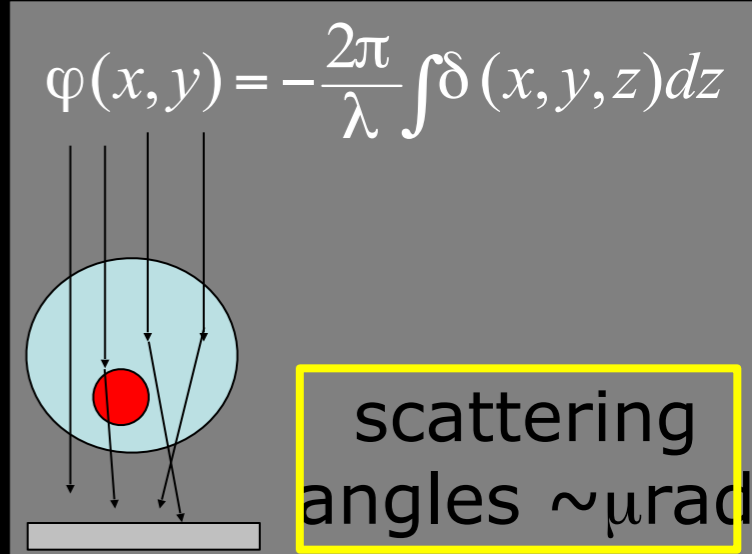


→ Phase shifts  
 - Elastic ("coherent") scattering

**PHASE CONTRAST IMAGING**



$$\mu = 4\pi\beta\epsilon / \lambda\epsilon_0$$



$$\varphi(x, y) = -\frac{2\pi}{\lambda} \int \delta(x, y, z) dz$$

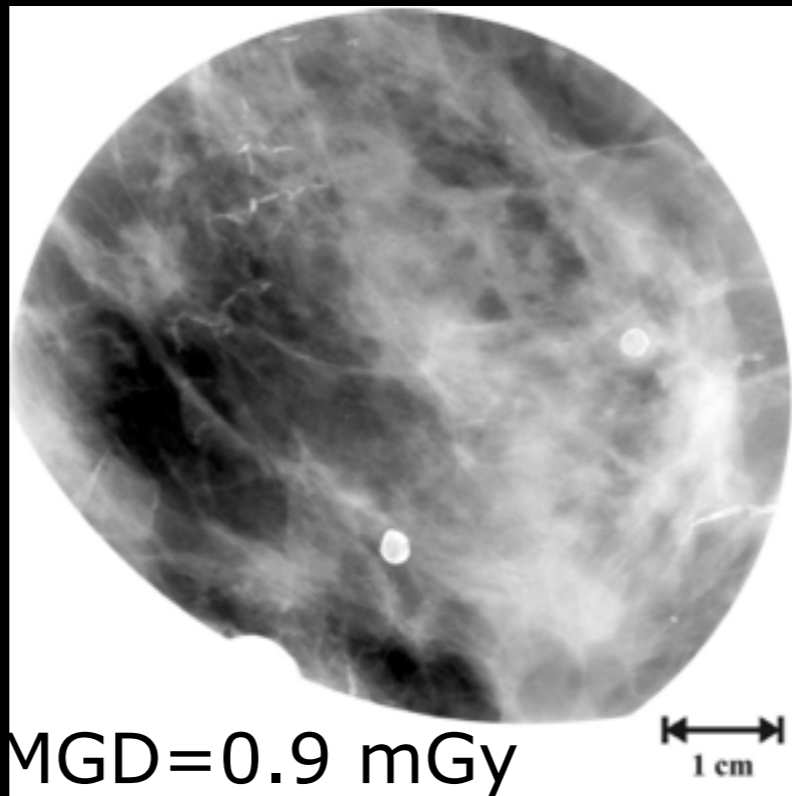
scattering angles  $\sim \mu\text{rad}$



Planar radiograph

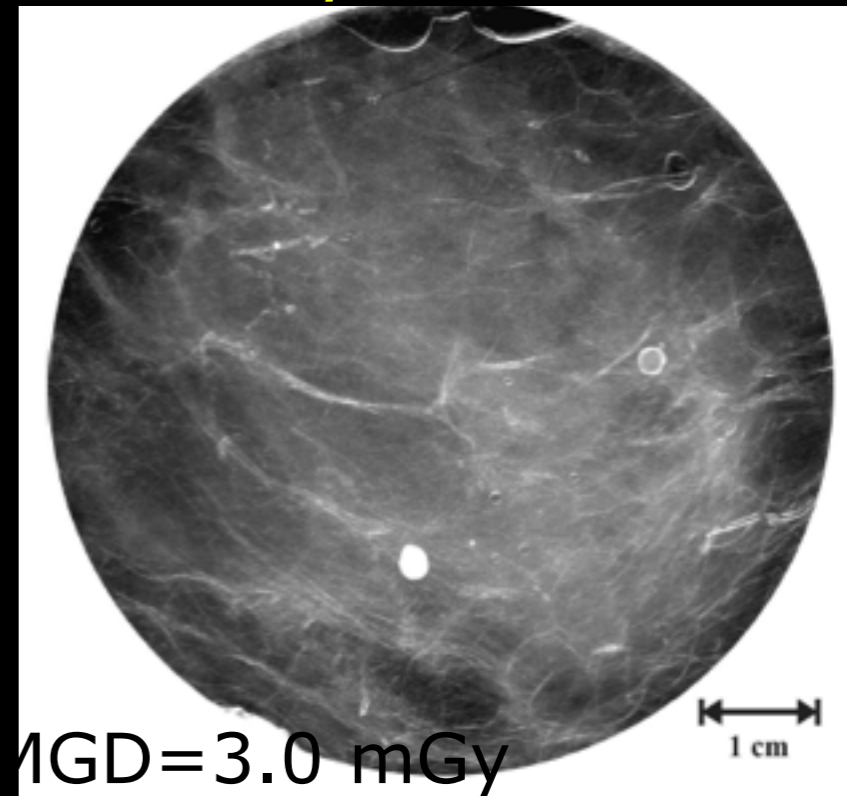
Multifocal lobular carcinoma

conventional



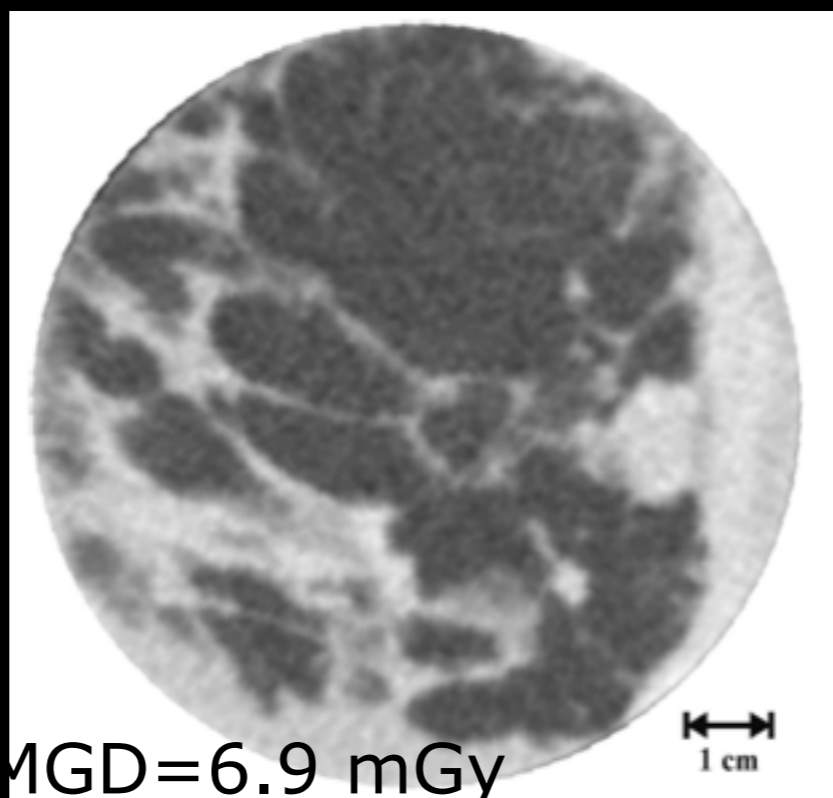
MGD=0.9 mGy  
26 kVp and 9 mAs

Analyzer-based

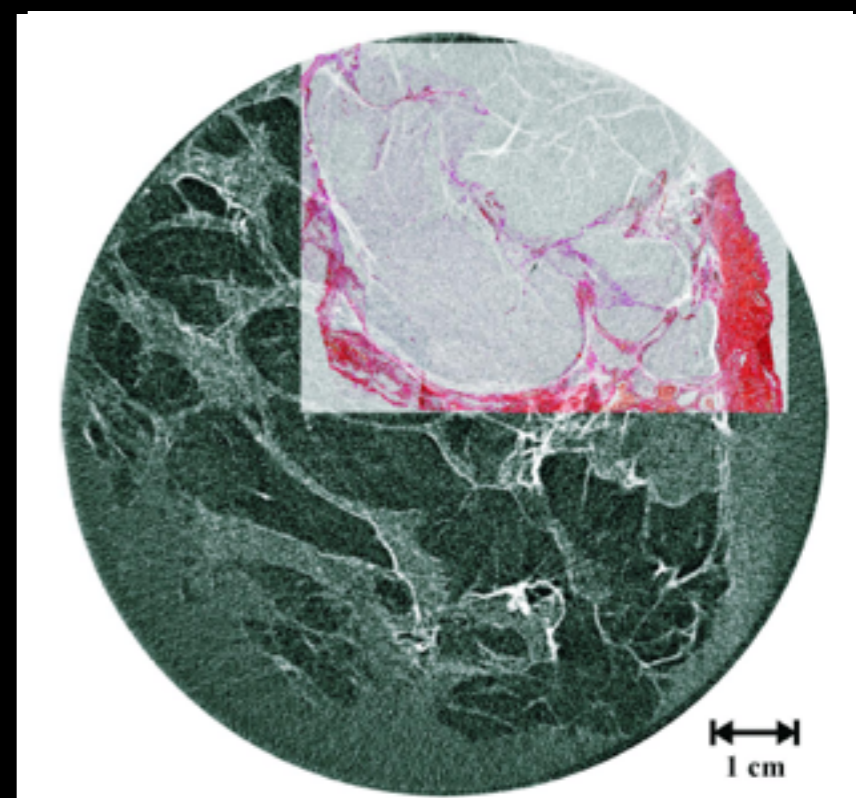


MGD=3.0 mGy  
30 keV

CT radiograph



MGD=6.9 mGy  
80 kVp and 50 mAs



30 keV