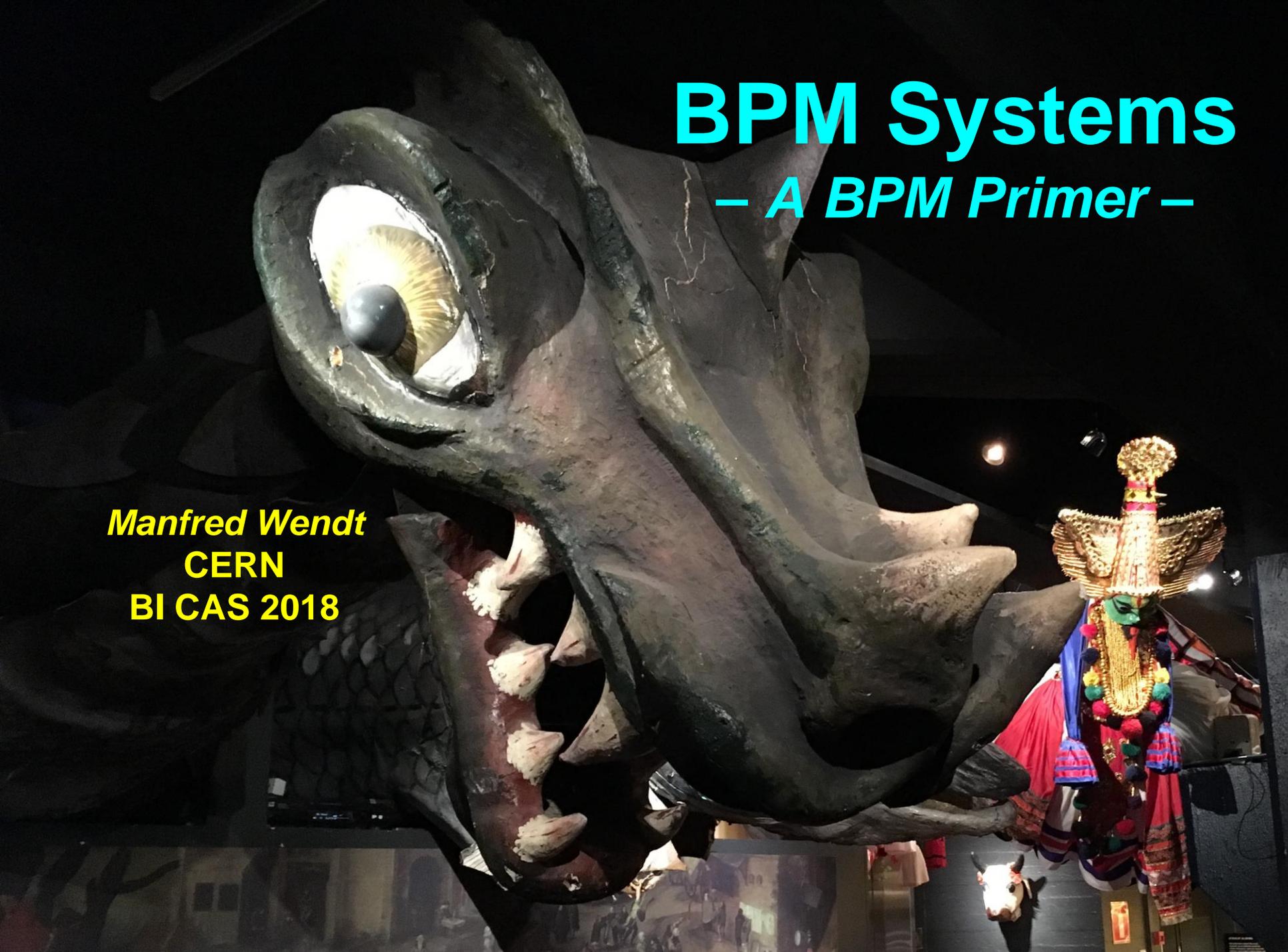


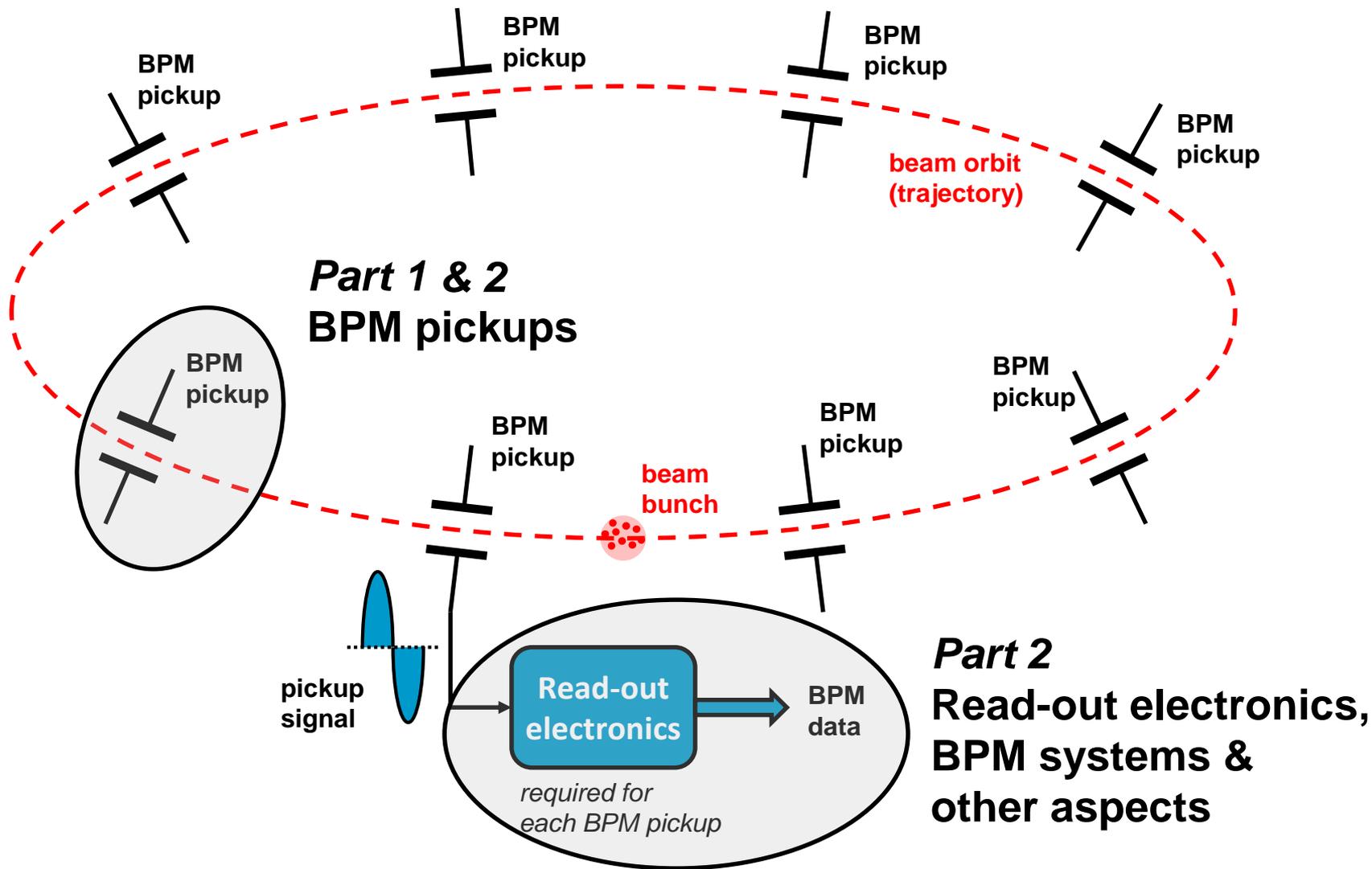
BPM Systems

– *A BPM Primer* –

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CERN
BI CAS 2018



A Beam Position Monitoring System



Learning Objectives

- **Introduction into BPM systems**
 - From an engineering perspective
 - **How to design and build a BPM system!?**
 - Focus on the principle of operation
 - **Often in graphical format**
 - Practical formulas as necessary, but no lengthy derivations
 - Stick to the fundamentals
 - **Expert and exotic stuff: please study the related papers**
- **Very personal selection of BPM topics**
 - Focus on popular BPM components and subsystems
 - **BPM pickups, such as buttons, stripline, cavity BPM, etc.**
 - **BPM electronics, e.g. RF signal conditioning, digital signal processing**
 - Will not cover BPM applications and beam measurements
 - **Covered at the CAS introduction and advanced courses**
- **Give the design basics for the key elements of a BPM system**
 - So lets start...

BPM Systems

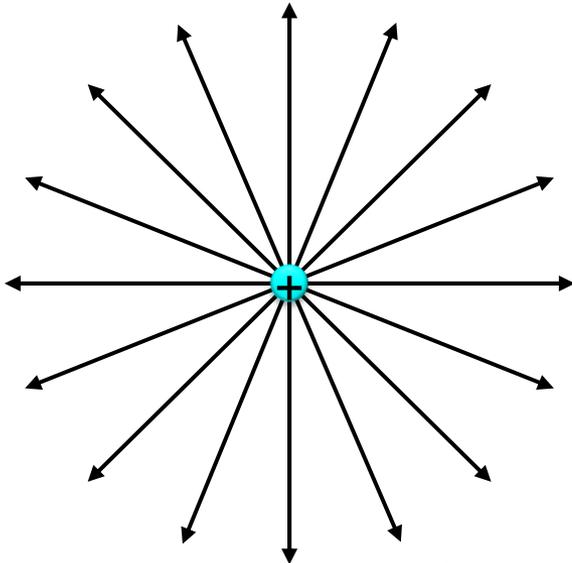
- are based on **Beam Position Monitors (BPM)**, which are beam detectors located along the accelerator
 - **BPM: Beam Position Monitor**
 - **Beam pickup with signal processing (read-out) electronics**
 - Often colleagues just refer to the beam pickup as BPM
 - **BPMs are typically located near each quadrupole magnet**
 - **Use 4 or more BPMs per betatron oscillation period**
- **deliver beam orbit (trajectory) information**
 - **Non-invasive monitoring based on the EM-field of the passing beam**
 - **Synchronized BPMs deliver beam timing information**
 - **Beam orbit measurement**
 - turn-by-turn, batch-by-batch, bunch-by-bunch, or averaged over many turns
 - **Beam phase or time-of-flight (TOF) information in linacs**
- **are a powerful beam diagnostics tool**
 - **Machine commissioning, characterization of the beam optics, measurement of beam parameters, trouble-shooting,...**

BPM Systems Part 1

- **Principle of operation**
 - Wall currents and the electrostatic BPM pickup
- **Bunch Beam Signals**
- **The image charge (current) model for BPMs**
 - Position characteristic in a circular beam pipe
 - Numerical analysis and correction of non-linearities
- **BPM pickups**
 - Split-plane BPM
 - Button BPM
 - Stripline BPM
 - Cavity BPM
- **Summary of part 1**

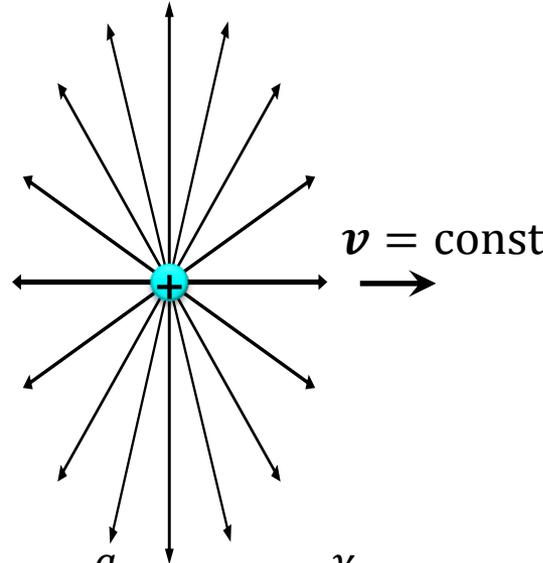
E-field of a moving point charge

Point charge
at rest



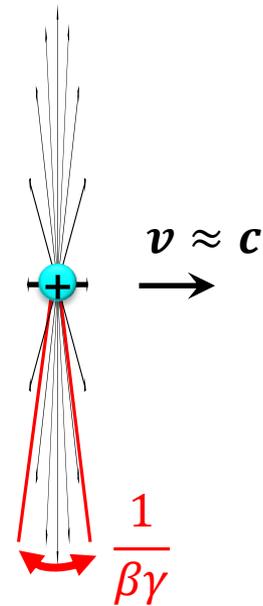
$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon} \rightarrow \vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon} \frac{q}{r^2} \hat{r}$$

Moving
point charge



$$\vec{E}(\vec{r}, t) = \frac{q}{4\pi\epsilon_0} \frac{\gamma}{r^3 \left(1 + \frac{v_r^2 \gamma^2}{c^2}\right)^{3/2}} \vec{r}$$

Relativistic moving
point charge

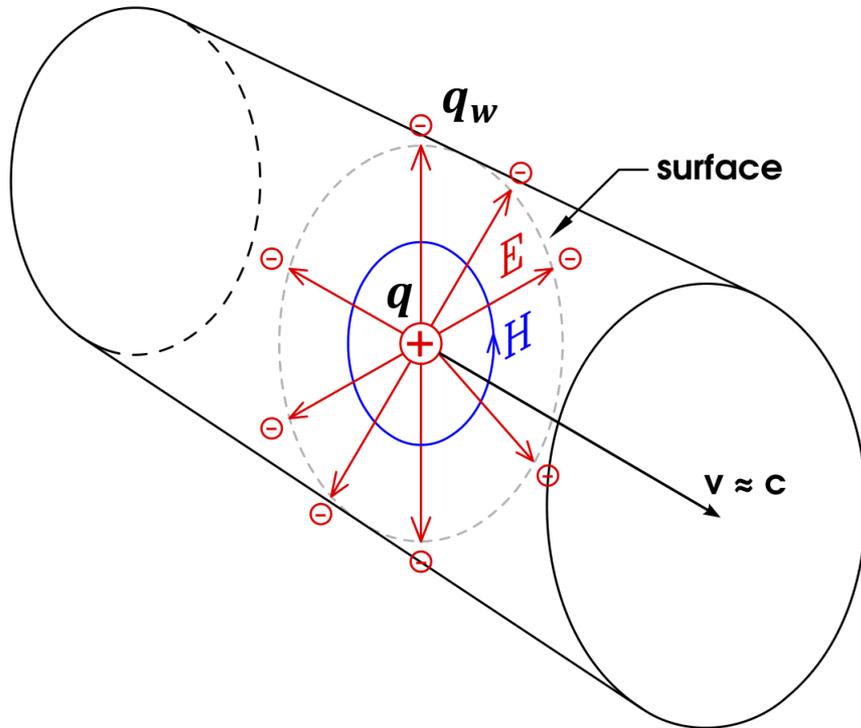


- A relativistic moving charge has a “pancake”-like EM-field

$$\beta = \frac{v}{c} \approx 1 \quad E'_z = \gamma E_z$$

- Only Transverse Electric Magnetic (TEM) field components
 - almost no longitudinal field components

Wall current



- **Single proton in a perfect conducting cylindrical beam pipe of radius r**
 - Travelling with: $\beta = \frac{v}{c} \approx 1$
 - Image charges q_w along the azimuth of the beam pipe wall:

$$q = -q_w$$

- **Wall current density:**

$$j_w(t) = -\frac{i_b(t)}{2\pi r}$$

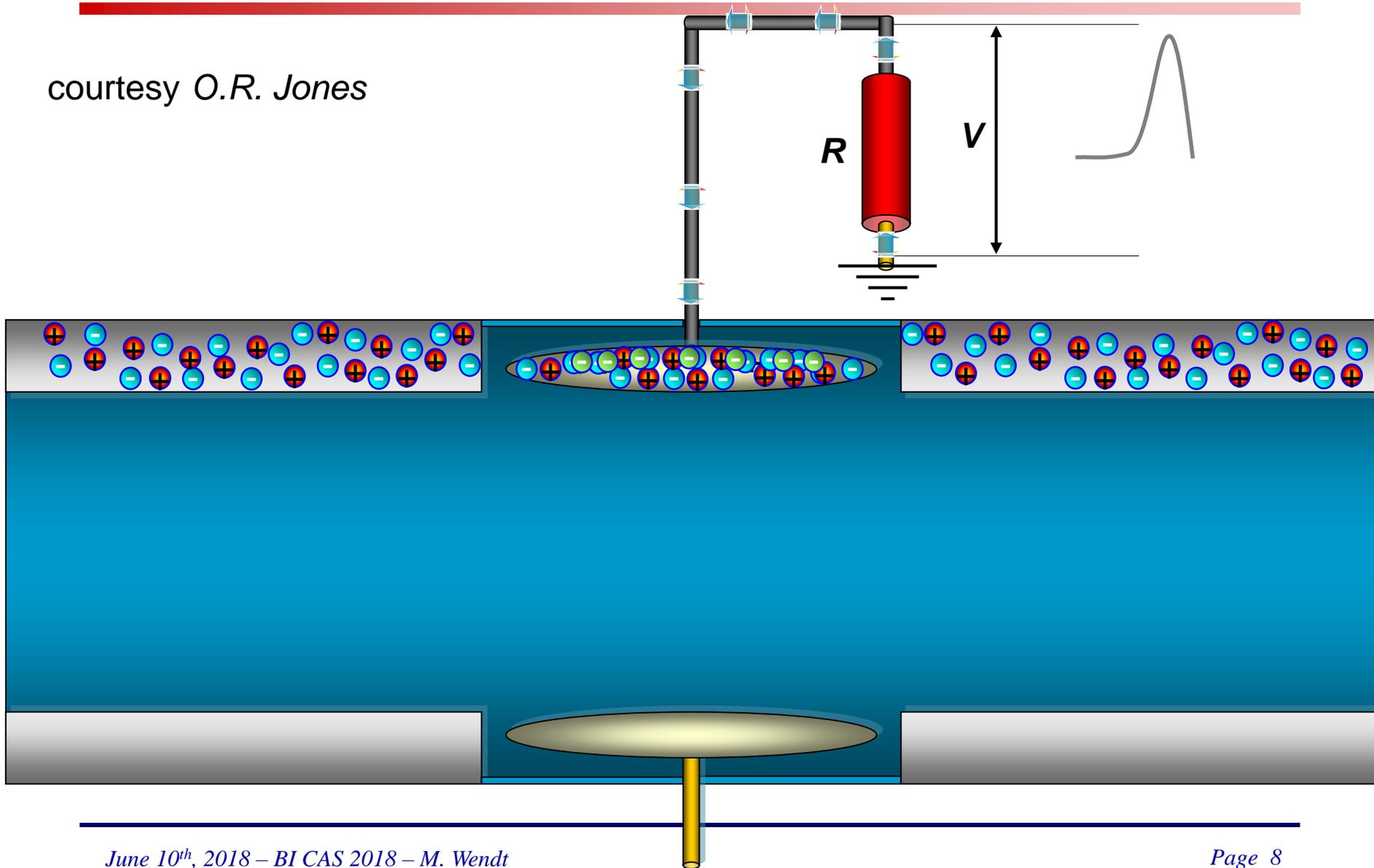
- **Beam current:**

$$i_b(t) = \frac{dq}{dt}$$

$$q = Ne$$

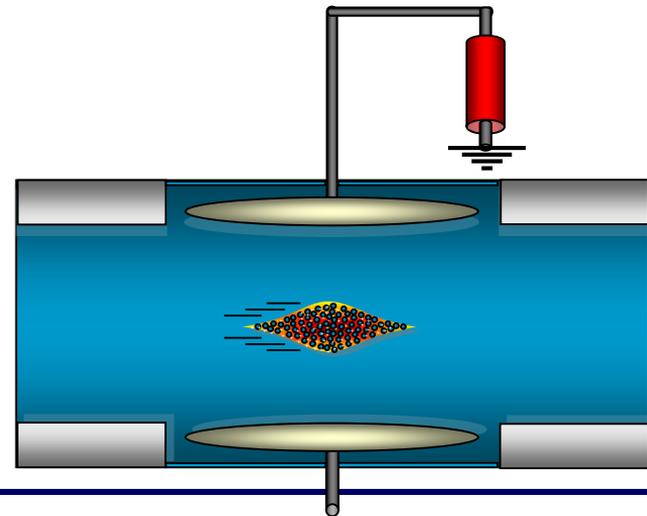
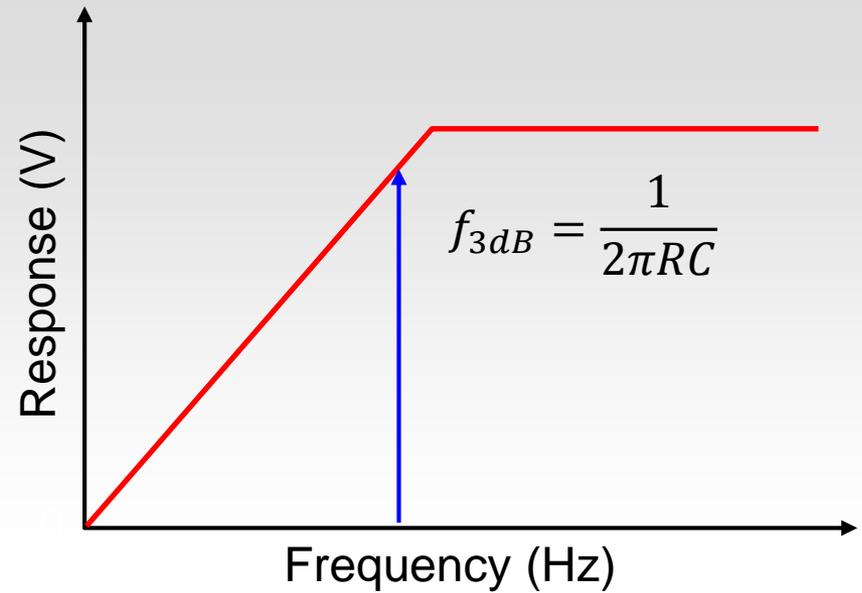
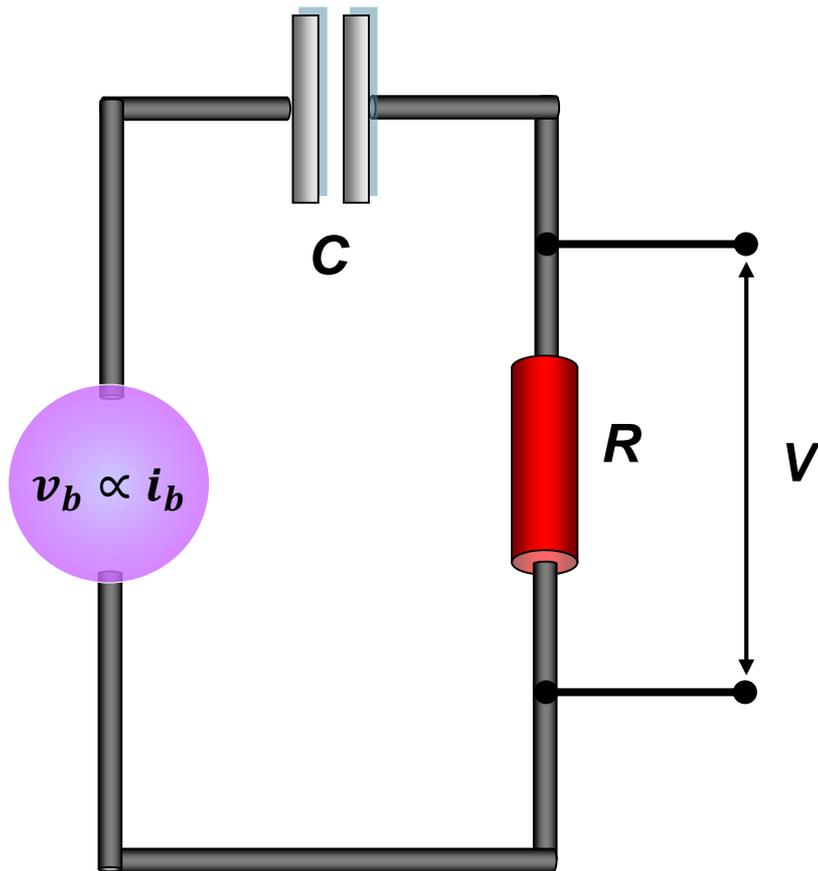
Electrostatic Beam Monitor

courtesy O.R. Jones



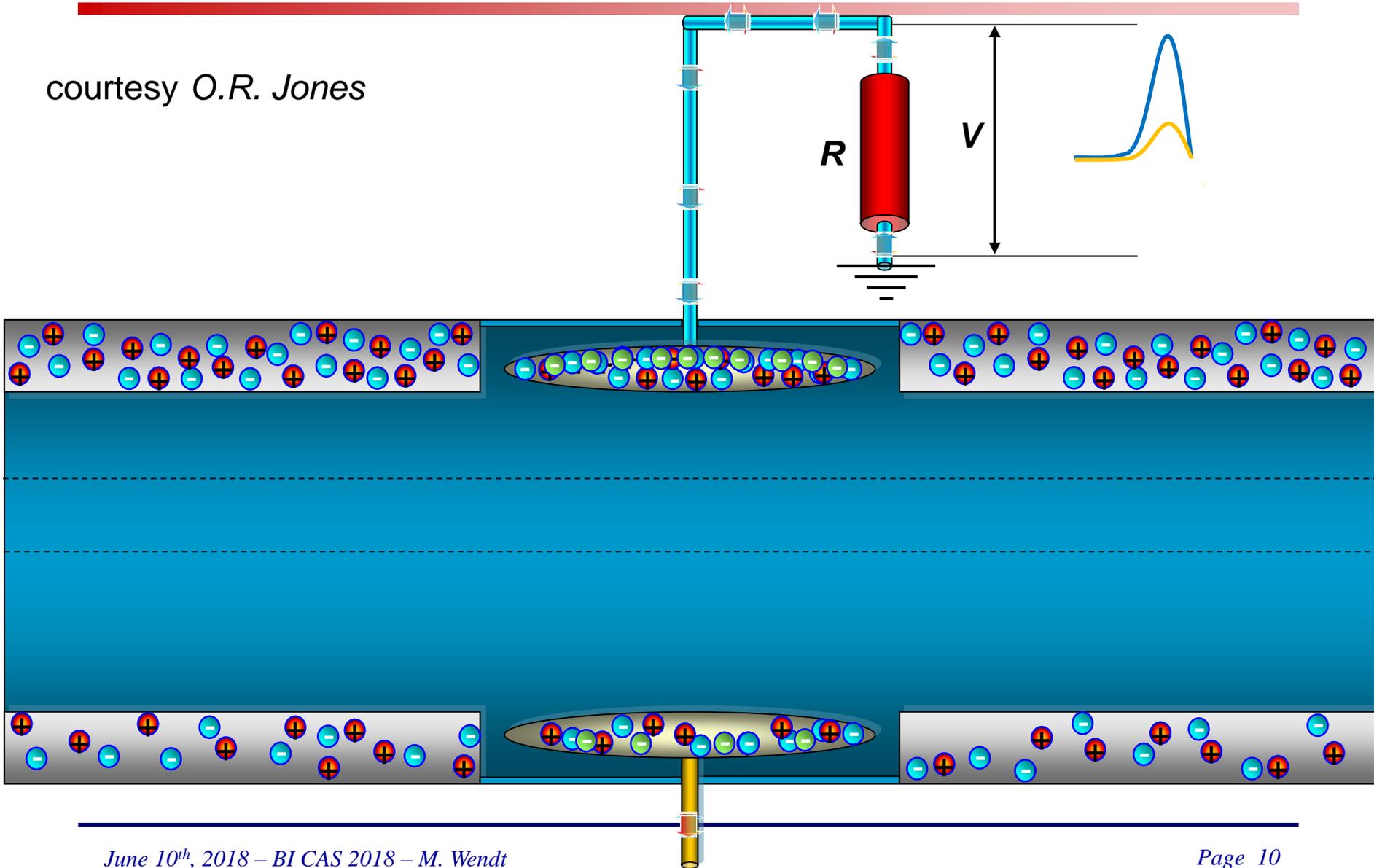
Beam response – High-pass

courtesy O.R. Jones



Electrostatic BPM

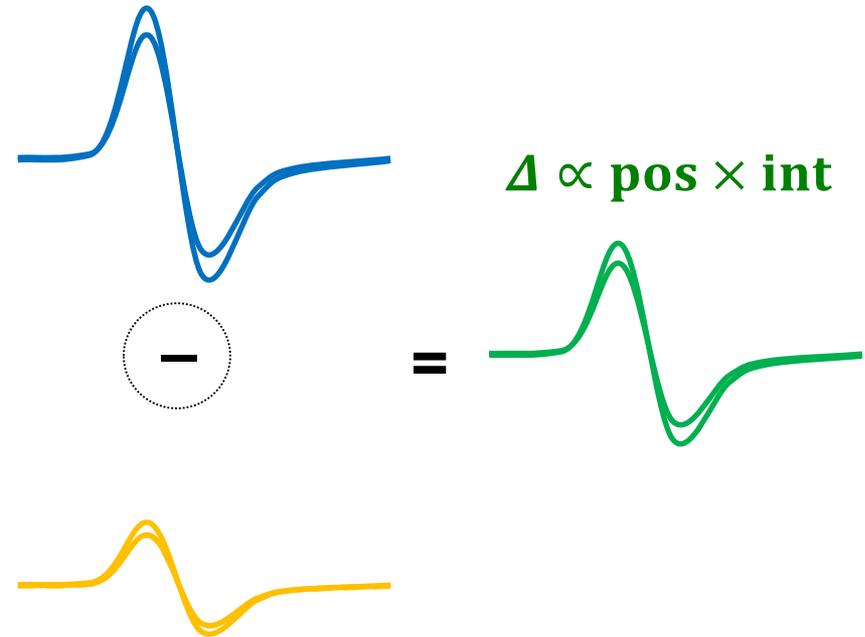
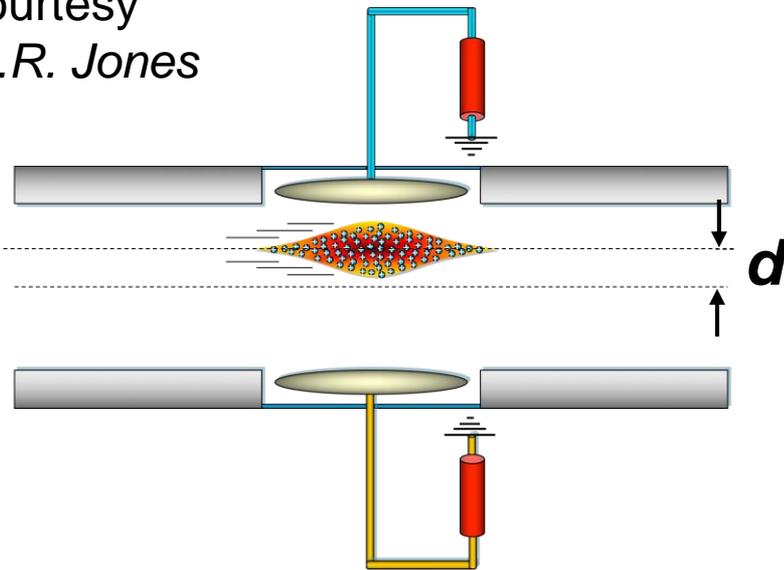
courtesy O.R. Jones



Beam Position Monitor Principle

- The BPM principle is based on symmetry
 - The beam displacement d is detected by a pair of symmetrical arrange electrodes

courtesy
O.R. Jones



- BUT: Hoops, wait a minute, not so fast...
What happens if the beam / bunch intensity changes?!
 - The Δ -signal still contains beam intensity information!
 - Need to “normalize” the Δ -signal

EM Pickup for Beam Position Measurements

- The BPM pickup detects the beam positions by means of:

- **identifying asymmetries** of the signal amplitudes from two symmetrically arranged electrodes A & B :

$$\text{norm. beam pos.} \propto \frac{A - B}{A + B}$$

- **broadband pickups, e.g. buttons, striplines**

- Detecting dipole-like eigenmodes of a beam excited, passive resonator:

- **narrowband pickups, e.g. cavity BPM**

- BPM electrode transfer impedance:**

$$V_{elec}(\mathbf{x}, \mathbf{y}, \omega) = s(\mathbf{x}, \mathbf{y}, \omega) Z(\omega) I_b(\omega)$$

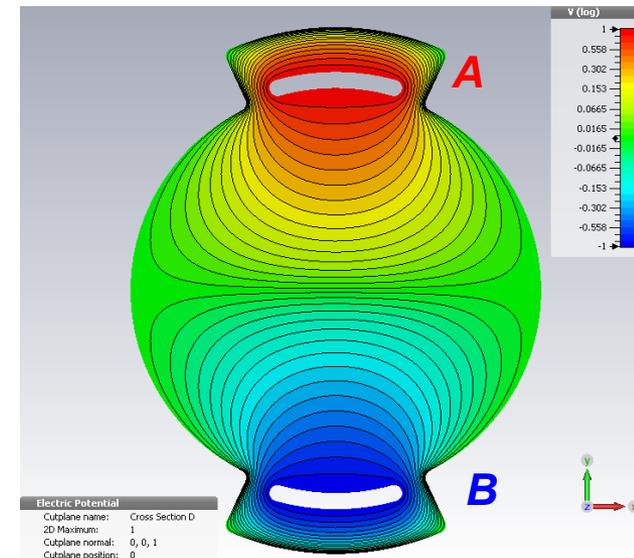
beam position

applies only for resonant pickups

frequency depending transfer impedance of the BPM electrode

beam intensity

- The beam displacement or sensitivity function $s(\mathbf{x}, \mathbf{y})$ is frequency independent for broadband pickups ($@ v \approx c$)

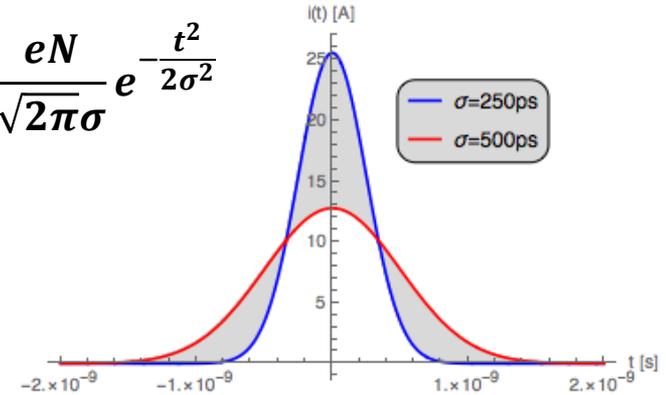


Beam Current Signals

- **Single bunch with Gaussian distribution of the particles (long.)**

- Number of particles: N
- Bunch length (time): σ

$$i_b(t) = \frac{eN}{\sqrt{2\pi}\sigma} e^{-\frac{t^2}{2\sigma^2}}$$



- **Bunches spaced by an equidistant time T**

- **Fourier series expansion:** $i_b(t) = \langle I_{DC} \rangle + 2\langle I_{DC} \rangle \sum_{m=1}^{\infty} A_m \cos(m\omega t)$

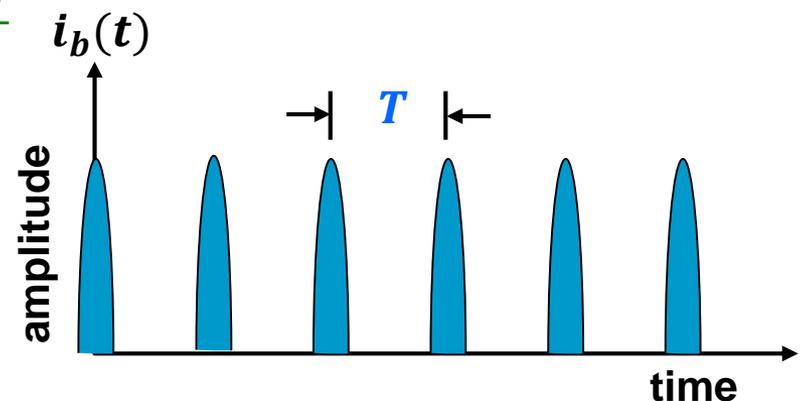
- Average beam current: $\langle I_{DC} \rangle = \frac{eN}{T}$

- Frequency harmonics spaced by: $\omega = 2\pi f$

- Time distance: $T = \frac{2\pi}{\omega}$

➤ **Note:** Understand the integration time of the BPM electronics

- Harmonic amplitude factor: $A_m = \dots$



Beam Signals: Harmonic Amplitude Factors

Bunch shape	Harmonic amplitude factor A_m	Comments
δ -function (point charge)	1	For all harmonics
Gaussian	$\exp\left[\frac{(m\omega\sigma)^2}{2}\right]$	$\sigma = \text{RMS bunch length}$
parabolic	$3\left(\frac{\sin\alpha}{\alpha^3} - \frac{\cos\alpha}{\alpha^2}\right)$	$\alpha = m\pi W/T$
$(\cos)^2$	$\frac{\sin(\alpha - 2)\frac{\pi}{2}}{(\alpha - 2)\pi} + \frac{\sin\frac{\alpha\pi}{2}}{\frac{\alpha\pi}{2}} + \frac{\cos(\alpha + 2)\frac{\pi}{2}}{(\alpha + 2)\pi}$	$\alpha = 2mW/T$
triangular	$\frac{2(1 - \cos\alpha)}{\alpha^2}$	$\alpha = m\pi W/T$
square	$\frac{\sin\alpha}{\alpha}$	$\alpha = m\pi W/T$

Normaization: $A_m \rightarrow 1$ for $\omega \rightarrow 0$ T : bunch period W : bunch length at base

Beam Signal Frequency Response

- Use *Fourier* transformation
 - instead of *Fourier* series expansion with infinite sums
- Examples: *Gaussian* and **raised cosine (cos²)** pulse
 - Time domain

$$i_{Gauss}(t) = \frac{eN}{\sqrt{2\pi\sigma}} e^{-\frac{t^2}{2\sigma^2}}$$

$$I_{Gauss}(f) = eN e^{-2(\pi f\sigma)^2}$$

$$i_{RaisedCos}(t) = \begin{cases} \frac{eN}{2w} \left(1 + \cos \frac{\pi t}{w} \right), & -w < t < w \\ 0, & \text{elsewhere} \end{cases}$$

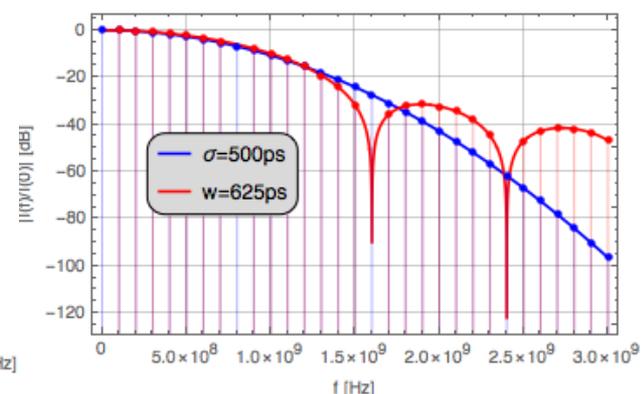
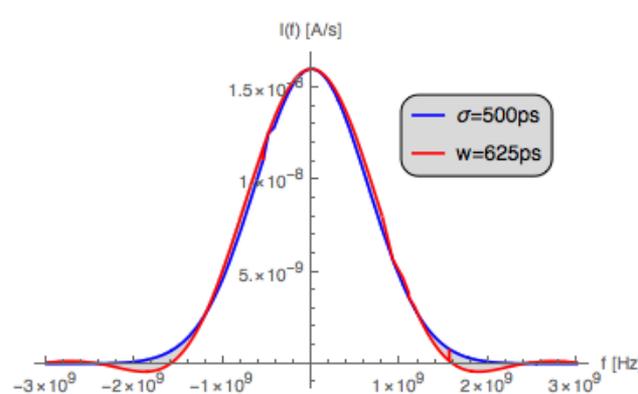
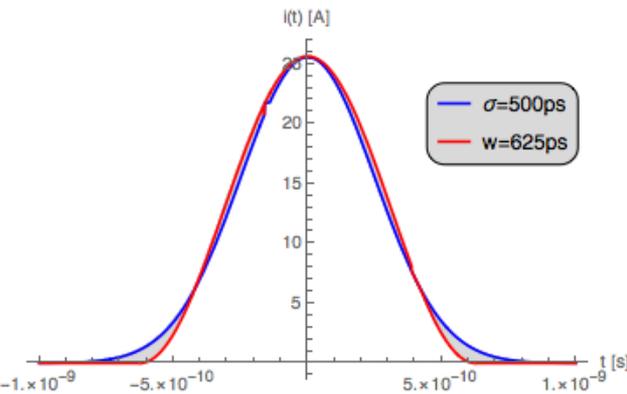
$$I_{RaisedCos}(f) = \frac{eN \sin(2\pi f w)}{2\pi f w (1 - 4f^2 w^2)}$$

Time domain ($N=10^{11}$)

Frequency domain

Frequency domain

$$\log_{10}[\text{magn.}(f)] / \log_{10}[\text{magn.}(0)]$$



More on Image Charges and Image Currents

- **Relativistic beams $v \approx c$:**
 - **Electrostatic problem of a line charge in a conductive circular cylinder**
- **Solution based on the image charge method:**

- **Image current density**

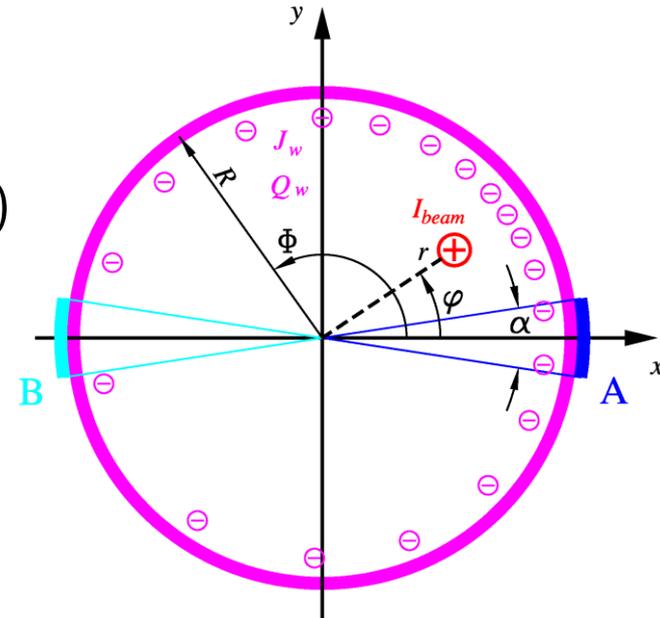
$$J_w(R, \Phi_w) = -\frac{I_{beam}}{2\pi R} \left[1 + 2 \sum_{n=1}^{\infty} \left(\frac{r}{R}\right)^n \cos n(\Phi_w - \varphi) \right] = -\frac{I_{beam}}{2\pi R} \frac{R^2 - r^2}{R^2 + r^2 - 2r \cos(\Phi_w - \varphi)}$$

- **Image current integrated on BPM electrode A**

$$I_A = R \int_{-\alpha/2}^{+\alpha/2} J_w(R, \Phi_w) d\Phi_w = -\frac{I_{beam}}{2\pi} s_A(r/R, \varphi, \alpha)$$

$$s_A(r/R, \varphi, \alpha) = \alpha + 4 \sum_{n=1}^{\infty} \frac{1}{n} \left(\frac{r}{R}\right)^n \cos(n\varphi) \sin\left(\frac{n\alpha}{2}\right)$$

- **Similar solution for electrode B**
or for vertically arranged electrodes



Normalized Beam Position Characteristic

- **Pair of symmetric horizontal electrodes:**

$$A = I_A \text{ (right electrode)} \quad B = I_B \text{ (left electrode)}$$

- **Horizontal and vertical beam position:**

$$x = r \cos \varphi \quad y = r \sin \varphi$$

- **Normalized beam position characteristic Δ/Σ (horizontal):**

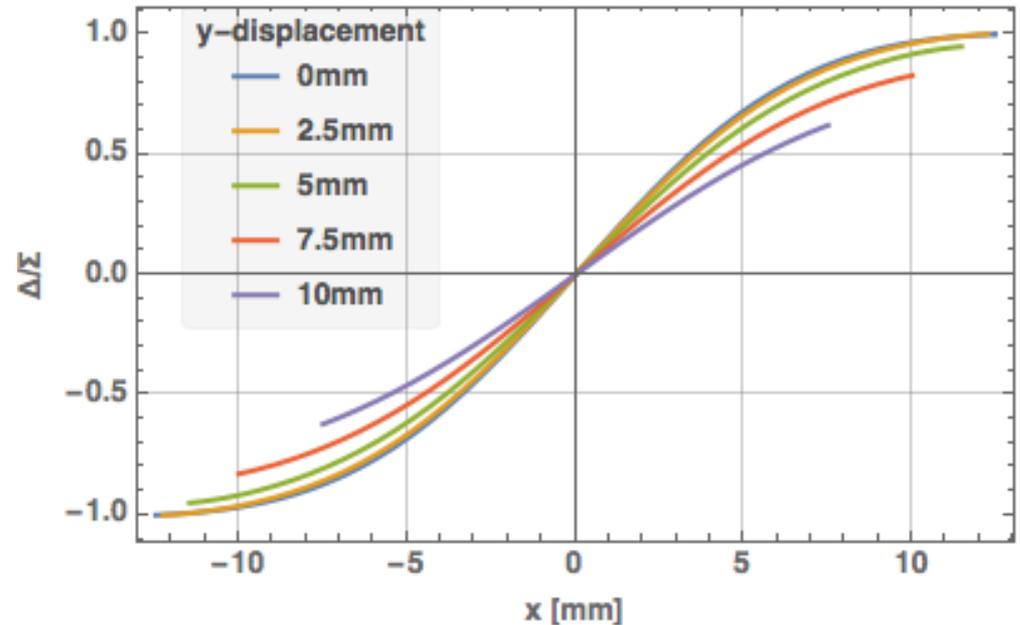
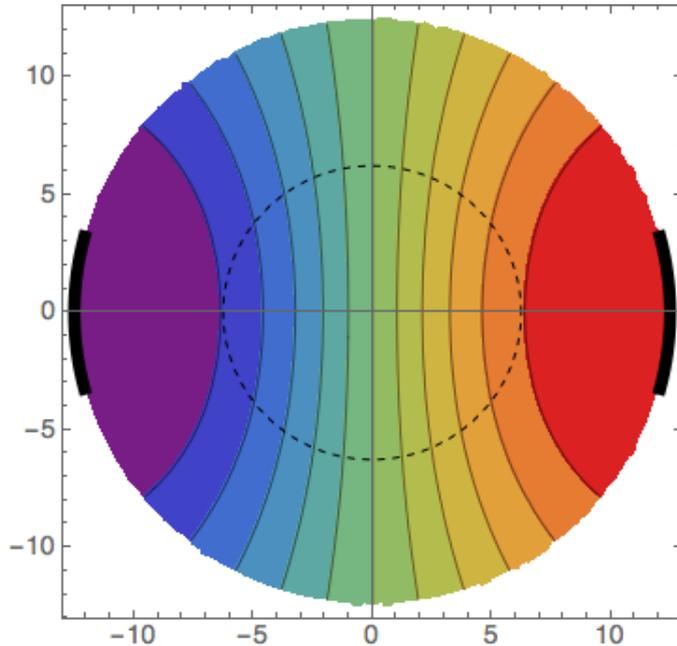
- Approximation and **closed form solution**

$$\begin{aligned} \text{hor. position} = \frac{\Delta}{\Sigma} &= \frac{A - B}{A + B} = \frac{4 \sin \frac{\alpha}{2} x}{\alpha R} + \text{higher-order terms} \\ &= \frac{f(x, y, R, \alpha) - f(-x, y, R, \alpha)}{f(x, y, R, \alpha) + f(-x, y, R, \alpha)} \end{aligned}$$

$$f(x, y, R, \alpha) = \pi \int_{-\alpha/2}^{+\alpha/2} J_w(R, \Phi_w) d\Phi_w = \tan^{-1} \frac{[(R+x)^2 + y^2] \tan\left(\frac{\alpha}{4}\right) - 2Ry}{x^2 + y^2 - R^2} + \tan^{-1} \frac{[(R+x)^2 + y^2] \tan\left(\frac{\alpha}{4}\right) + 2Ry}{x^2 + y^2 - R^2}$$

Example

- Position characteristic for $R = 12.5$ mm, $\alpha = 30^\circ$:



- Approximations**

- **Point-like electrodes:**

$$\text{hor. position} = \frac{A - B}{A + B} = \frac{1}{1 + \frac{x^2 + y^2}{R^2}} \frac{2}{R} x, \quad \alpha \rightarrow 0$$

- **Small beam positions:**

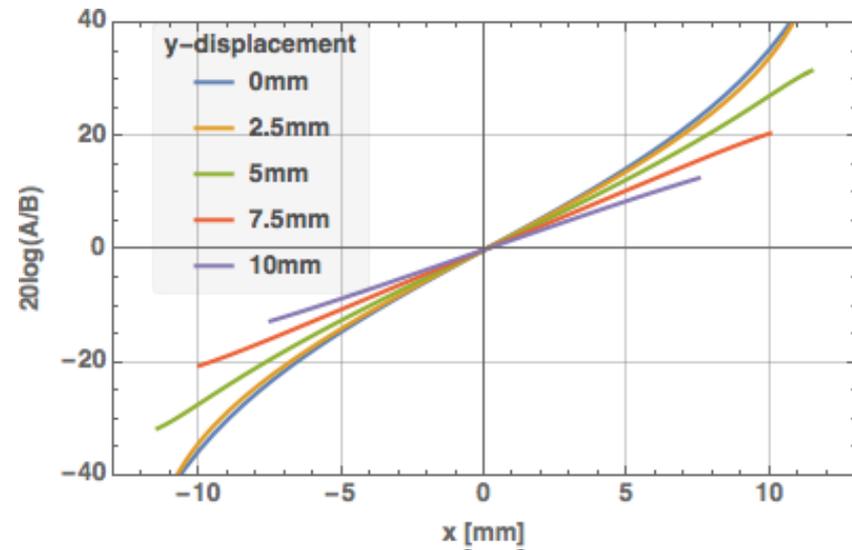
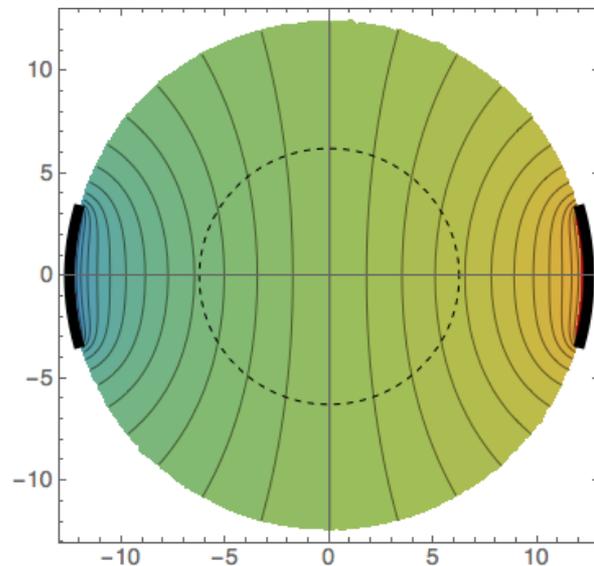
$$\text{hor. position} = \frac{A - B}{A + B} \approx \frac{2}{R} x, \quad (x^2 + y^2) \ll R^2$$

Logarithmic Ratio Normalization

- Results in a more linear position characteristic

$$\begin{aligned} \text{hor. position} = 20\log_{10}\left(\frac{A}{B}\right) &= 20\log_{10}\left(\frac{\sin\frac{\alpha}{2}x}{\alpha R} + \text{higher-order terms}\right) \\ &= 20\log_{10}\left[\frac{f(x, y, R, \alpha)}{f(-x, y, R, \alpha)}\right] \end{aligned}$$

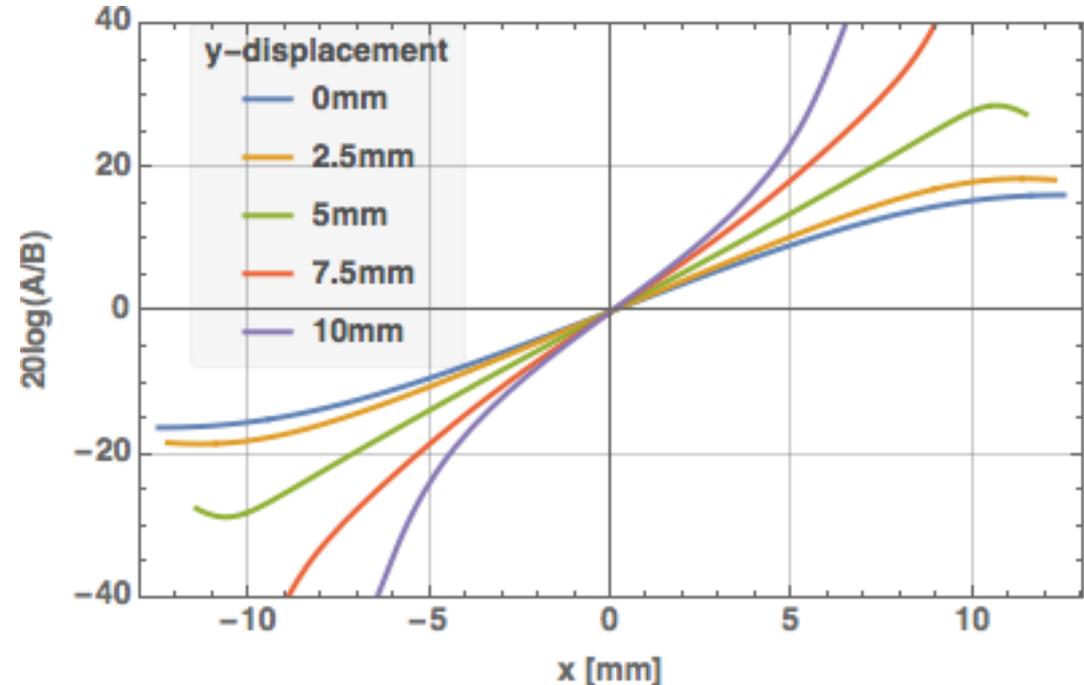
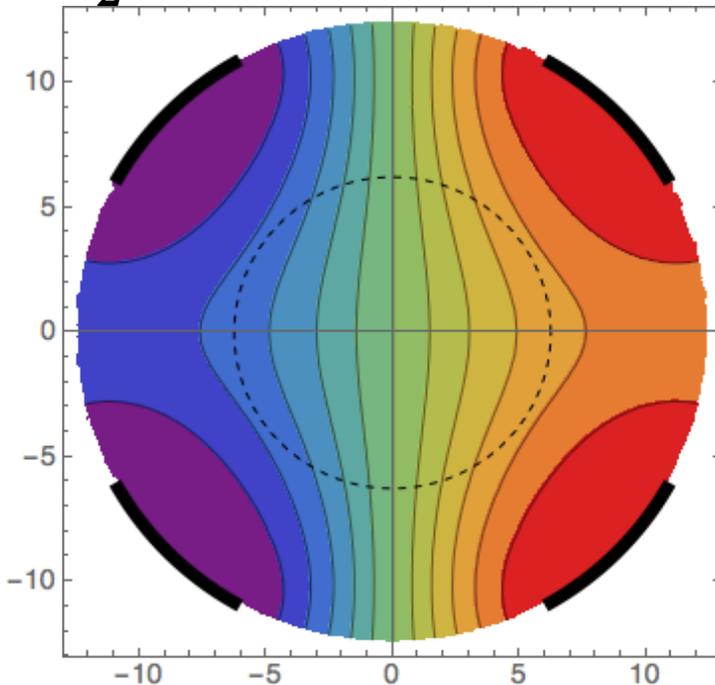
- The example for $R = 12.5$ mm, $\alpha = 30^\circ$ gives a sensitivity of 2.75 dB/mm near the origin



BPM with 45° Rotated Electrodes

- Popular in synchrotron light and damping rings, and in the LHC!
 - Prevents that sync light, or collision debris, hits the electrodes
 - More non-linear position behavior
 - Reduced position sensitivity (this example: 1.94 dB/mm)
 - Same mathematical procedure, but with coordinate rotation of 45°

$$\Delta/\Sigma = \text{const.}$$



Numerical Analysis (1)

- For most button or stripline type BPMs: 2D E-static analysis

- Relativistic beam $1/\gamma^2 \ll (\sigma_\ell/R)^2$ long bunches $\sigma_\ell \gg R$

- Solve the 2D E-static potential problem

- **BUT: without varying ρ**
for many positions $\vec{r} \in (x, y)$

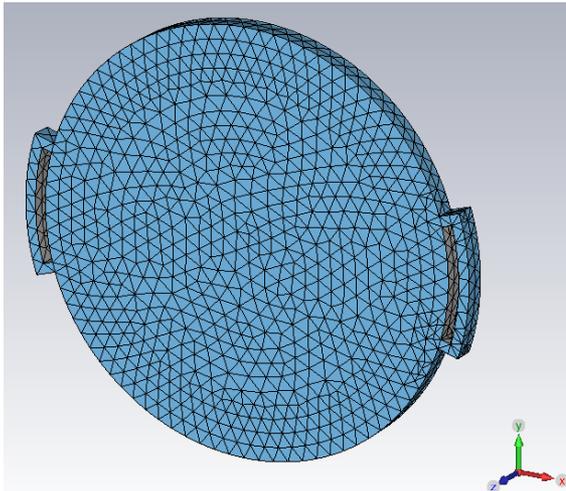
line charge density

$$\nabla_{\perp}^2 \Phi_{elec}(r) = \frac{\rho}{\phi} \delta(\vec{r} - \vec{r}_0)$$

potential

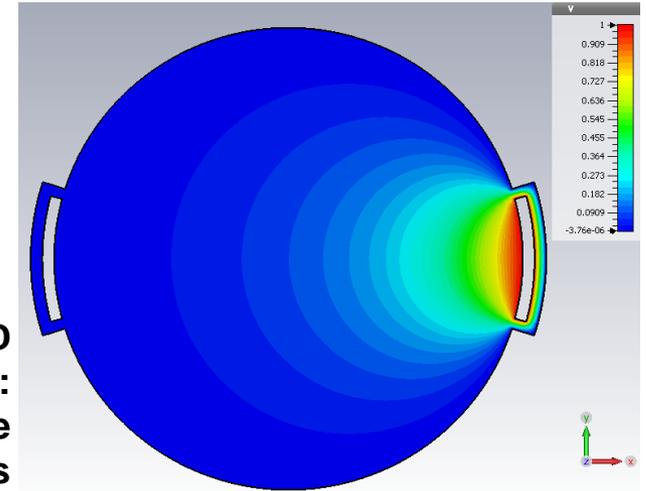
- Apply Green's reciprocity theorem

- and solve the Laplace equation in 2D: $\nabla_{\perp}^2 \Phi_{elec}(r) = 0 \rightarrow \Phi_{elec}(x, y)$



2D “slice”,
prepared with
tetrahedral mesh
for the numerical
analysis

Result of the quasi-2D
numerical analysis:
equipotentials of the
right electrodes



Numerical Analysis (2)

- For a horizontal BPM

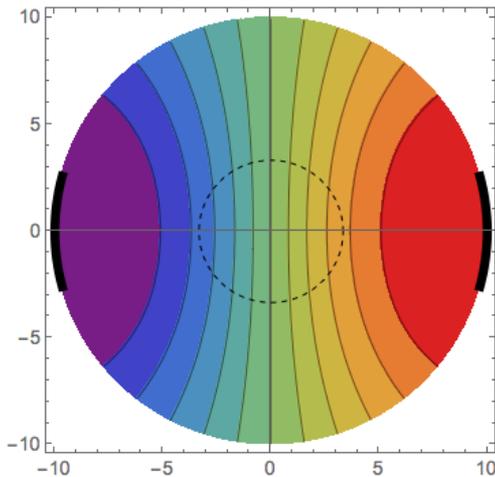
- Mirror the potential field $\Phi_{elec}(x, y) = \Phi_A(x, y) = \Phi_B(-x, y)$

- And then combine the fields $\Phi_{hor}(x, y) = \frac{\Phi_A - \Phi_B}{\Phi_A + \Phi_B} = \frac{\Delta}{\Sigma}$

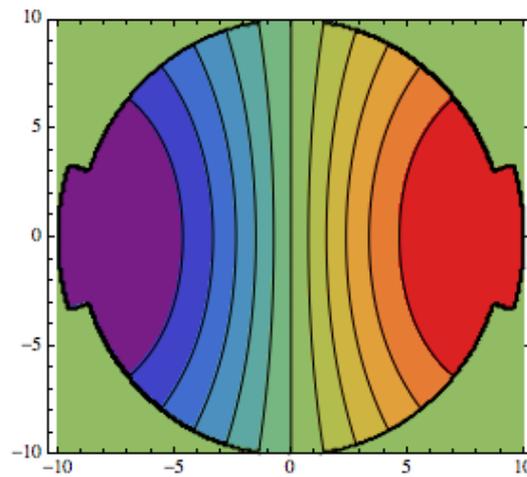
- Can be performed for any symmetric arrangement of the electrodes

- Comparison:

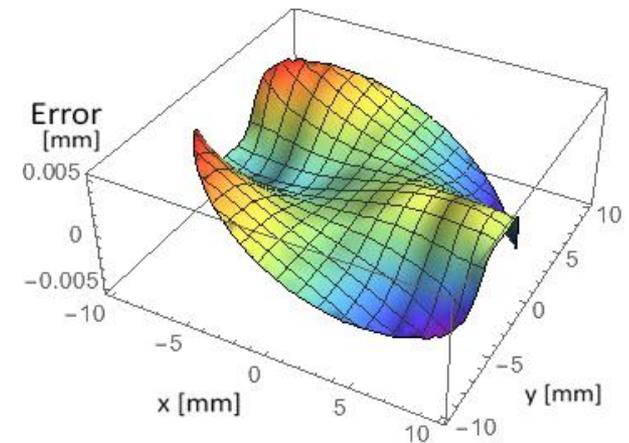
analytical analysis



numerical analysis



analytical - numerical



1D Non-linear Correction

- The 2D electrostatic analysis enables
 - optimization of the characteristic impedance of a stripline
 - Coverage factor, centered beam sensitivity $s(x = 0, y = 0)$
 - Correction of the non-linear position behavior in 1D or 2D
 - By look-up tables or a polynomial fit function

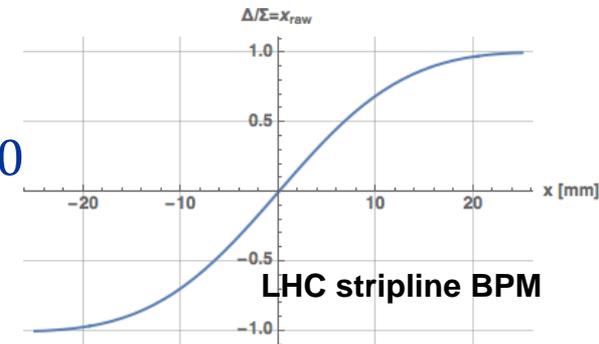
- Find a 1D polynomial fit function for horizontal beam displacements $x \neq 0, y = 0$

- Relationship between raw (measured) and true beam position:
$$x_{raw} = f(x)$$
- Fit the inverted function by a polynomial of power p

- Larger p , better fit
 - A too high p may lead to an unstable fit

- In most cases a x-y symmetry is given, and the same correction polynomial can be applied to the vertical axes

- In practice, the quadratic correction area has to be limited, e.g. $\mathbb{R} = 40\%$

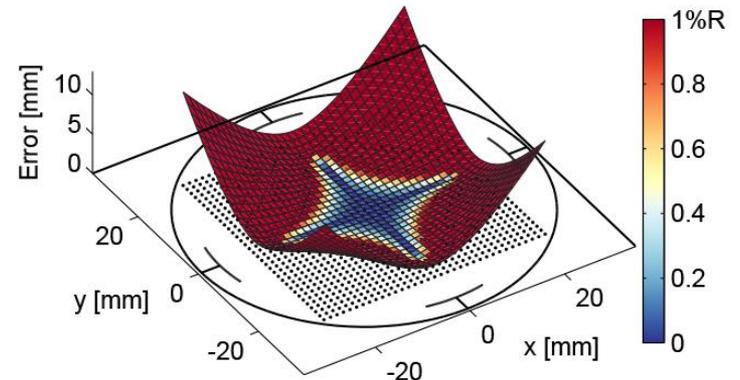


$$x_{bpm}^{1D} = \sum_{i=0}^p c_i x_{raw}^i = U_p(x_{raw}) \approx x$$

2D Non-linear Correction

- **1D non-linear correction example:**

- Remaining errors for the LHC stripline BPM applying correction polynomials $U_5(x), U_5(y)$ at an area $\mathbb{R} = 68\%$



- **Find a 2D polynomial fit function for $x, y \neq 0$**

- Raw and true beam position are given by:

with: $f = g$ and $y = f(y_{raw}, x_{raw})$

➤ **Notice the swap!**



- Fit a 2D surface polynomial for $f(x, y)$ of power p and q for x and y

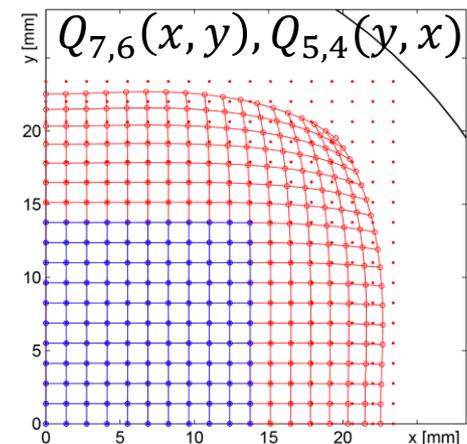
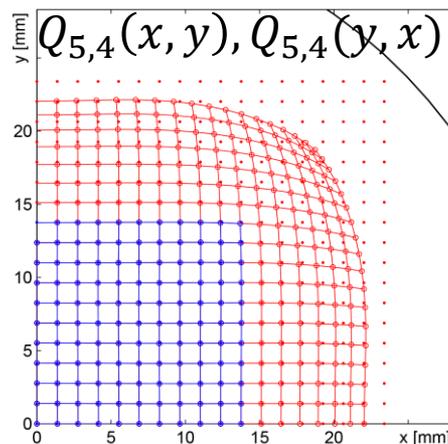
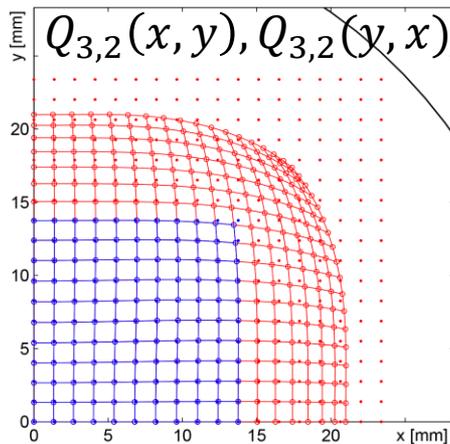
$$\begin{cases} x_{bpm}^{2D} = \sum_{i,j=0}^{p,q} (c_{ij} x_{raw}^i y_{raw}^j) = Q_{p,q}(x_{raw}, y_{raw}) \approx x \\ y_{bpm}^{2D} = \sum_{i,j=0}^{p,q} (c_{ij} y_{raw}^i x_{raw}^j) = Q_{p,q}(y_{raw}, x_{raw}) \approx y \end{cases}$$

Examples for 2D corrections

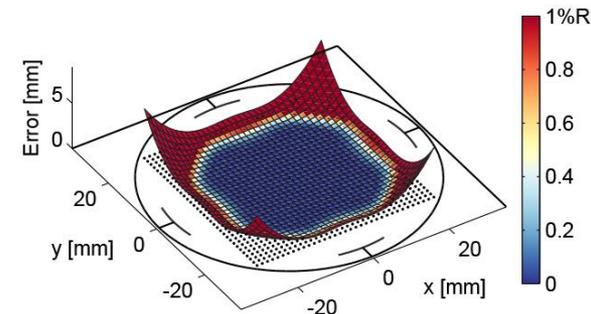
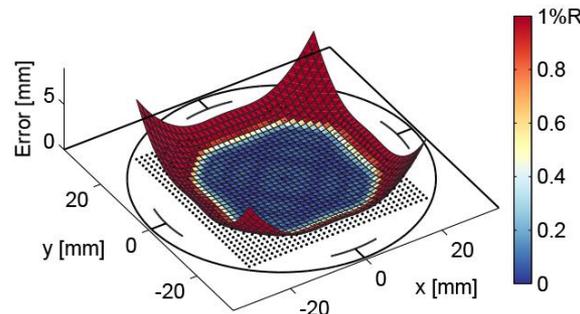
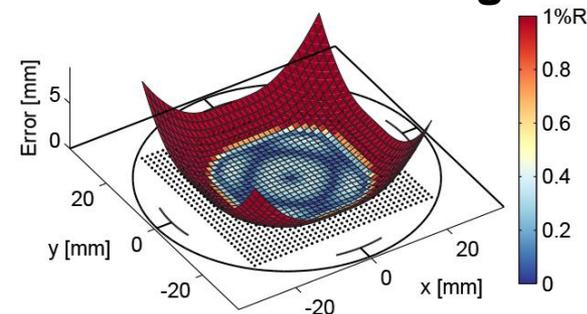
- **LHC stripline BPM**

- **Pin-cushion maps ($\mathbb{R} = 40\%$)**

- For symmetry reasons some cross-terms are 0, or very small (negligible)



- **Remaining errors:**

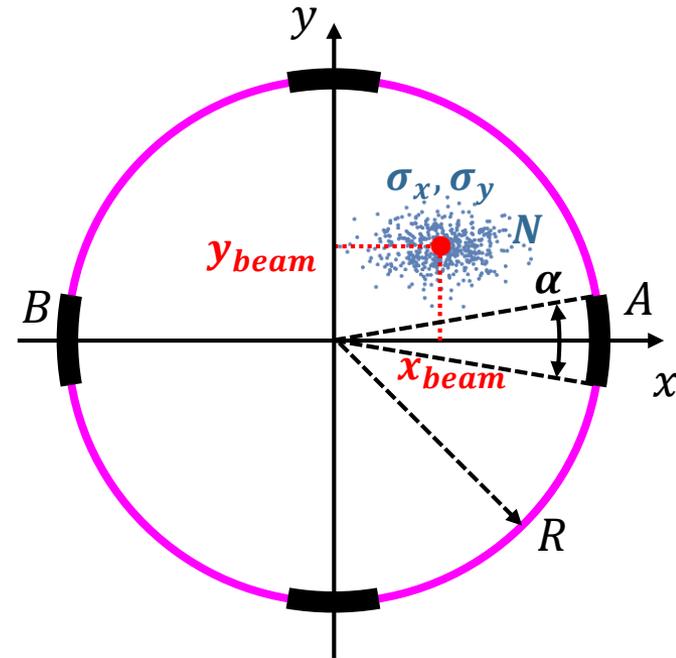


Higher Order Moments

- Particle \leftrightarrow BPM EM-fields are linear and time-invariant
 - Superposition principles apply
 - Many particles \rightarrow BPM detects (**approximately**) the center-of-charge
 - Point charge $q = eN$

- But: Non-linear effects allow the detection of higher moments:
 - Electrode A signal to N particles:

$$I_A = -\frac{I_{beam}}{2\pi} \left[\alpha + 4 \sum_{i=1}^N \sum_{n=1}^{\infty} \frac{1}{n} \left(\frac{r_i}{R}\right)^n \cos(n\varphi_i) \sin\left(\frac{n\alpha}{2}\right) \right]$$



- After some math gym follows:

monopole moment \propto intensity (common mode)

$$I_A = -\frac{I_{beam}}{\pi} \left[\underbrace{\alpha}_{\text{monopole moment}} + \underbrace{\frac{2}{R} \sin\left(\frac{\alpha}{2}\right) x_{beam}}_{\text{dipole moment} \propto \text{position}/R} + \underbrace{\frac{1}{R^2} \sin(\alpha) (\sigma_x^2 - \sigma_y^2 + x_{beam}^2 - y_{beam}^2)}_{\text{quadrupolar moment} \propto (\Delta\text{size} + \Delta\text{pos})/R^2} + \dots \right]$$

Effects of the Beam Size

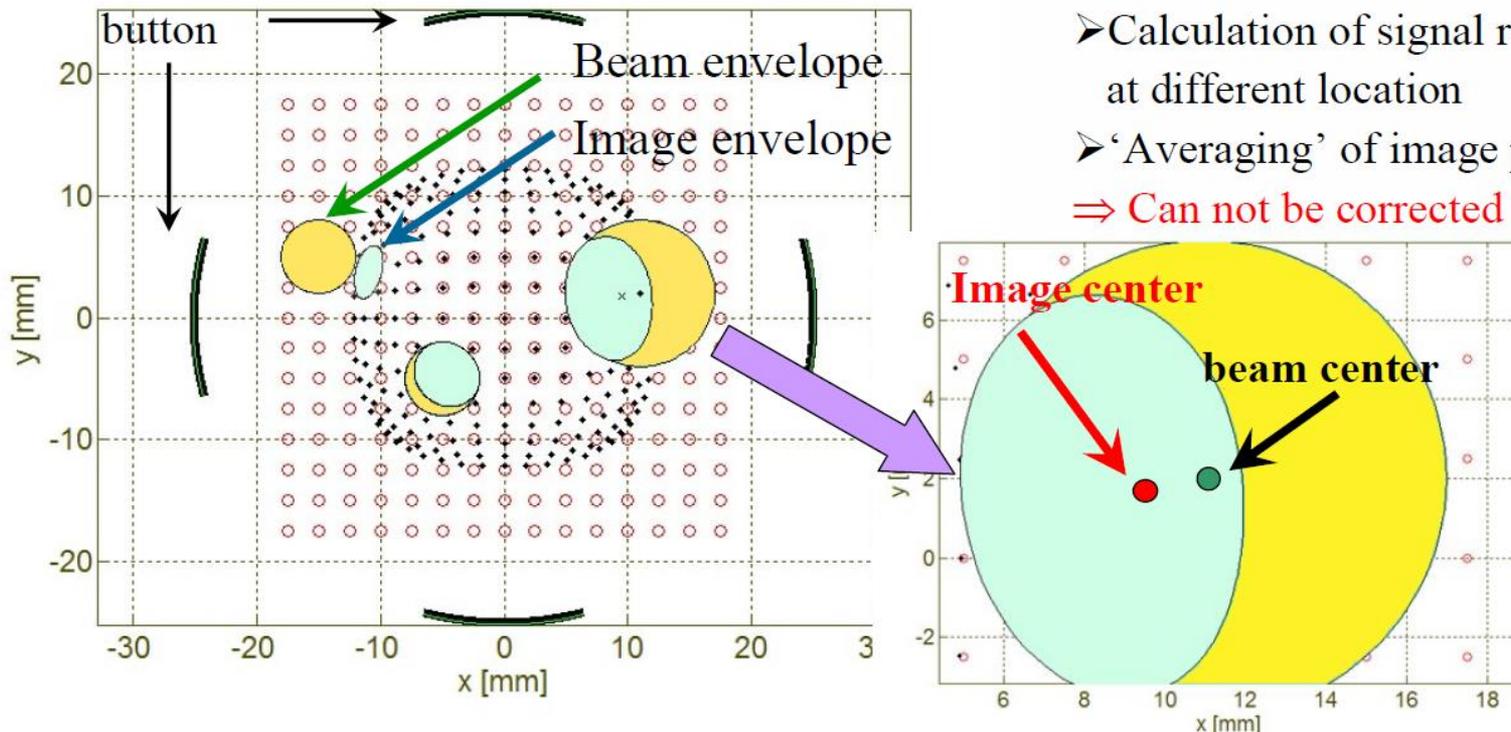
Ideal 2-dim model:

Due to the non-linearity, the beam size enters in the position reading.

courtesy *M. Bozzolan*

Finite beam size:

- Calculation of signal response at different location
 - 'Averaging' of image position
- ⇒ Can not be corrected !



Remark: For most LINACs: Linearity is less important, because beam has to be centered
→ correction as feed-forward for next macro-pulse.

BPM with Linear Position Response

- **Split-plane tube electrode**
 - radius: R
 - length: $\ell(\Phi_w) = \ell(1 + \cos \Phi_w)$
 - **A beam of**
 - charge: q_b , position: (r, φ)
- induces an image charge on the electrode:**

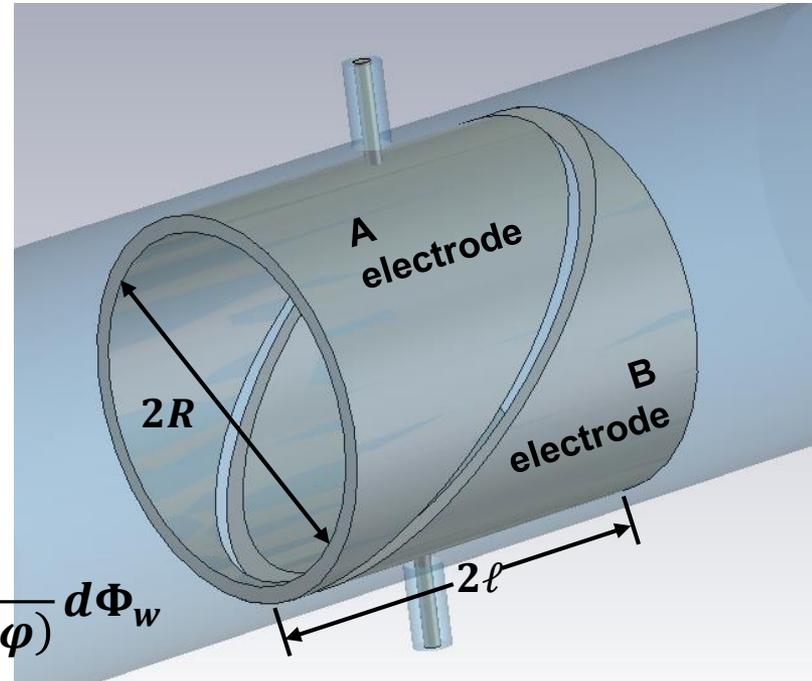
$$Q_{elec} = -q_b \ell \int_0^{2\pi} \frac{(1 + \cos \Phi_w)(R^2 + r^2)}{R^2 + r^2 - 2Rr \cos(\Phi_w - \varphi)} d\Phi_w$$

- with a linear position response

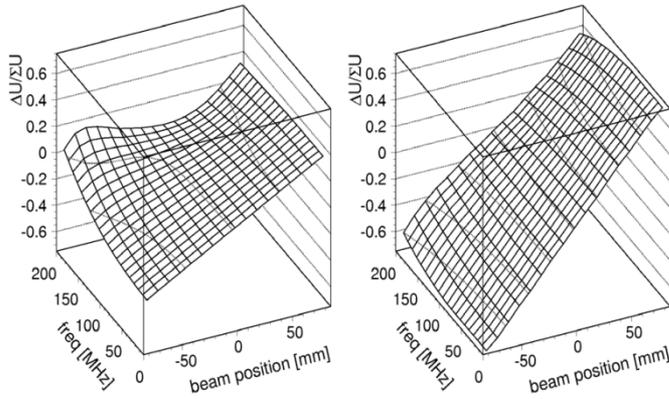
$$Q_{elec} = -q_b \ell \left(1 + \frac{r \cos \varphi}{R} \right) = -q_b \ell \left(1 + \frac{x}{R} \right)$$

⇒

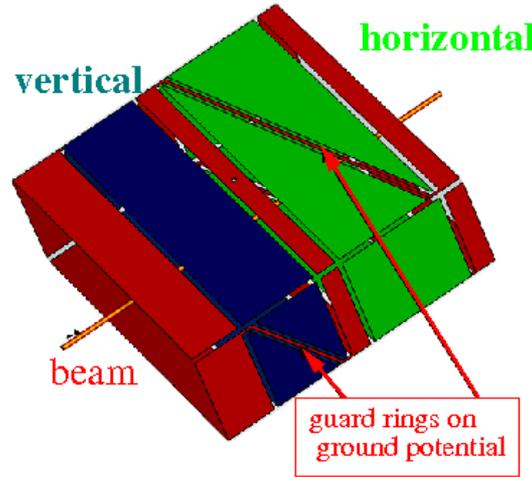
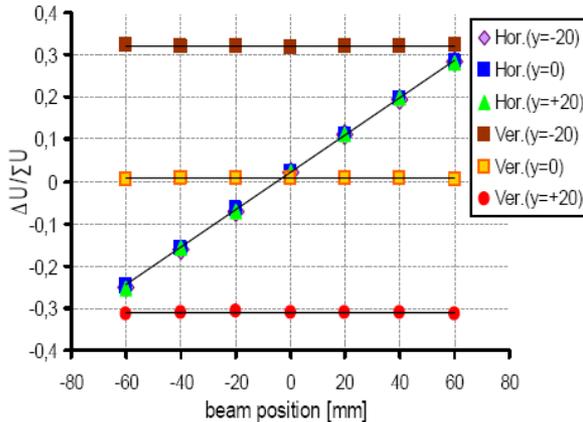
$$\text{hor. position} = \frac{\Delta}{\Sigma} = \frac{A - B}{A + B} = \frac{x}{R}$$



“Shoe-box” BPM



without ground guards with ground guards



- Split-plane (“shoe-box”) BPM with ground-guards

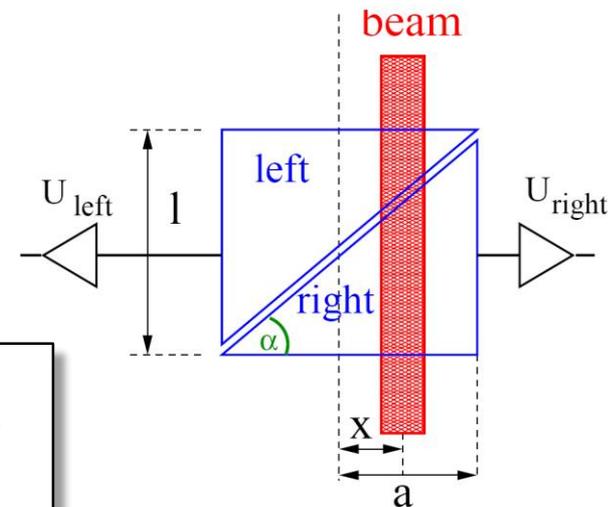
- Improved sensitivity and reduced cross-talk
- Improved linearity

courtesy P. Kowina

$$V_{left} = (a + x) \tan \alpha$$

$$V_{right} = (a - x) \tan \alpha$$

$$x = a \frac{V_{right} - V_{left}}{V_{right} + V_{left}} = \frac{\Delta}{\Sigma}$$



Spit-Plane BPM

+ Pros

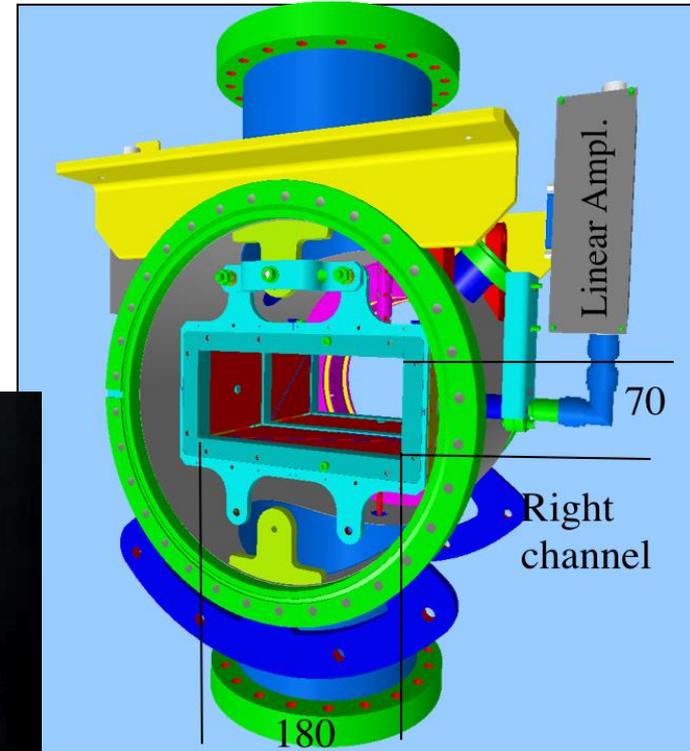
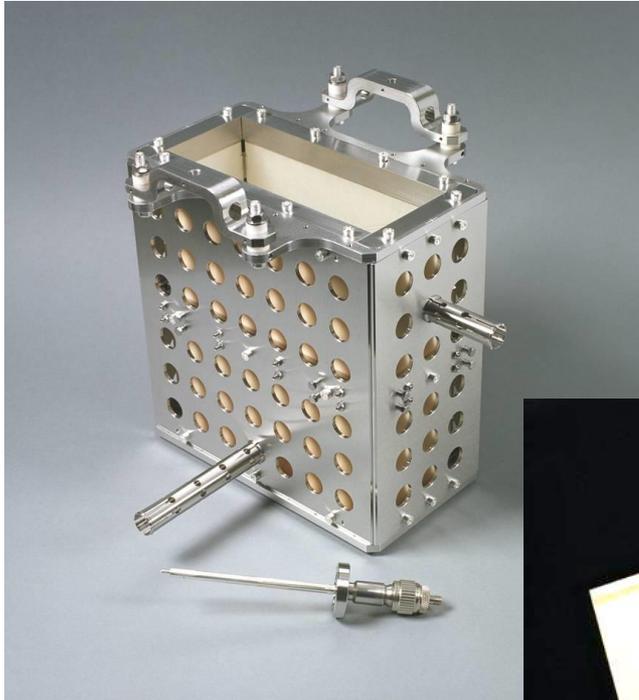
- **Linear position response**
 - **No correction for non-linearities required, simple analog read-out electronics without post-processing**
- **high transfer impedance at low frequencies**
 - **Good match for long bunches, low- β beam**
 - **High signal levels, allows operation at low frequencies with high load impedance**

- Cons

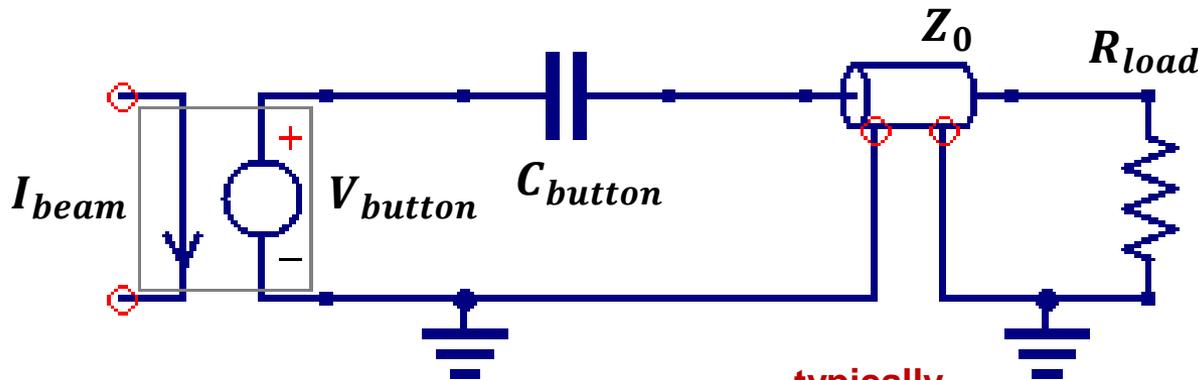
- **Complicated mechanics, requires sufficient real-estate**
- **High capacitive source impedance, reflective**
- **Eigenmodes at low frequencies, high beam coupling impedance**
 - **Limited to low frequency operation, typically <300 MHz**

“Shoe-box” BPM Example

Technical realization at the HIT synchrotron of 46 m length for 7 MeV/u \rightarrow 440 MeV/u
BPM clearance: 180×70 mm², standard beam pipe diameter: 200 mm.

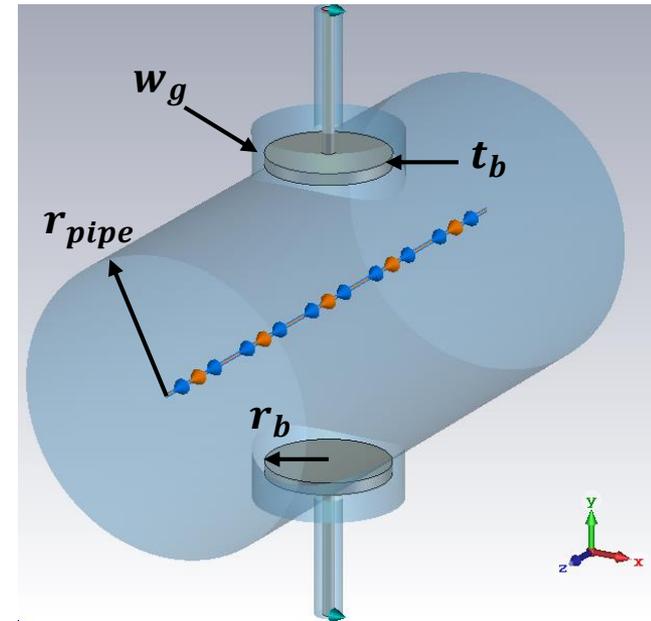


Button BPM Equivalent Circuit



$$Z_{button}(\omega) = \frac{V_{button}(\omega)}{s(x, y)I_{beam}(\omega)} = \phi R_{load} \frac{\omega_1}{\omega_2} \frac{j\omega/\omega_1}{1 + j\omega/\omega_1}$$

typically
50 Ω



$$\omega_1 = \frac{1}{R_{load}C_{button}} \quad \omega_2 = \frac{c}{2r_{button}}$$

time constant 1/electrode transit time

$$\phi = \frac{r_{button}}{4r_{pipe}} = s(x=0, y=0) \quad \text{coverage factor}$$

button capacitance, very simplified
typically 2...15 pF

$$C_{button} \approx \frac{2\pi\epsilon_0}{\ln\left(\frac{r_{button} + w_{gap}}{r_{button}}\right)} t_{button} + \dots$$

Increase	Effect
d_{button} (larger area & C_{button})	Higher signal level Lower resolution
t_{button} (larger volume)	Lower wake impedance Lower resolution, higher weight
w_{gap} (more trapped modes)	Higher resolution Higher wake impedance

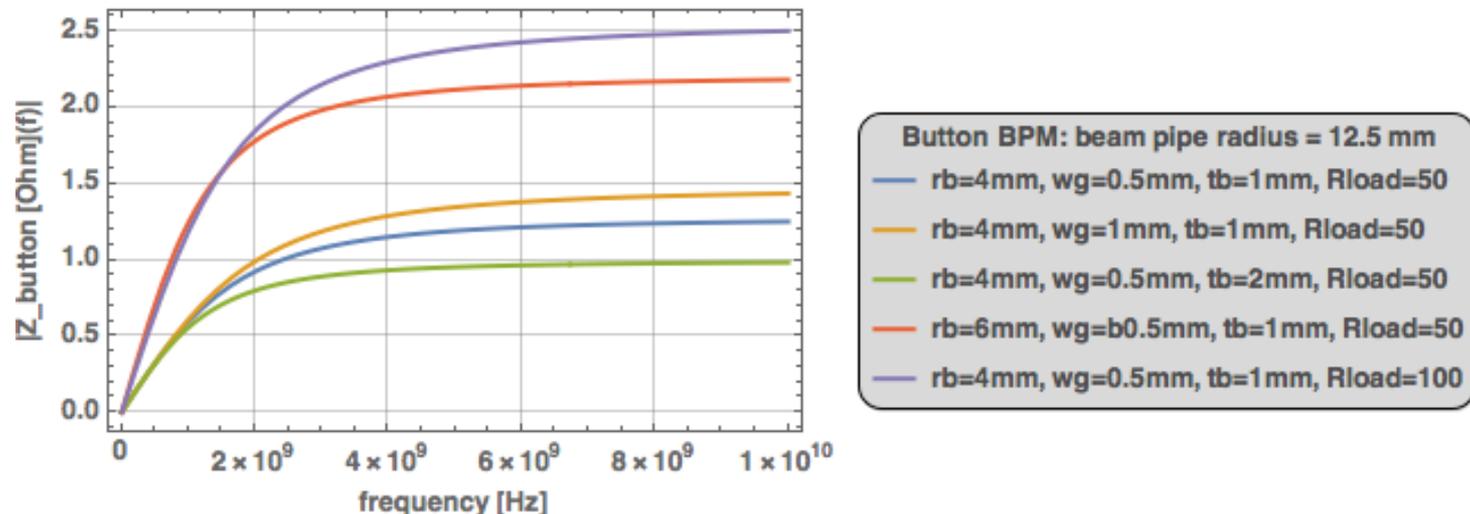
Button BPM

+ Pros

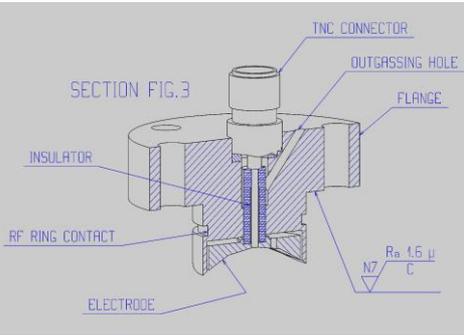
- Robust simple construction, cost effective, minimum real-estate
 - RF UHV feedthrough and button is a single element

- Cons

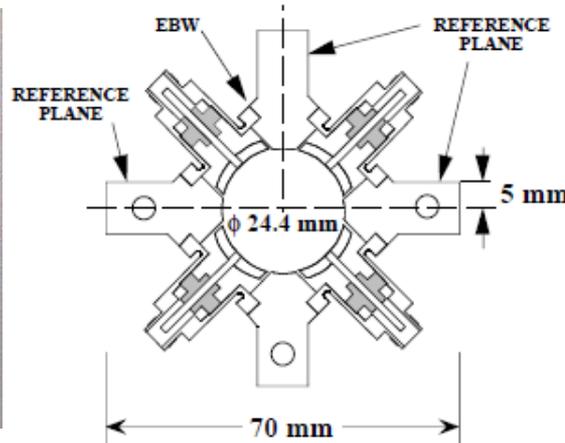
- High-pass characteristic with high cut-off frequency
 - Typically ≥ 500 MHz,
 - bad match to operate with long bunches, or at low frequencies
- Capacitive source impedance, reflective



Examples of Button BPMs



LHC Button BPM (CERN)



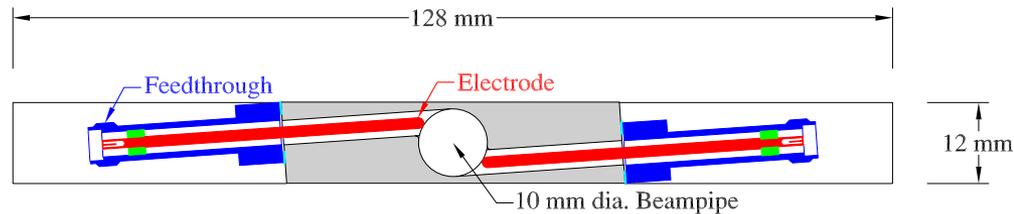
ATF Button BPM (KEK)



NSLS-II Button BPM (BNL)

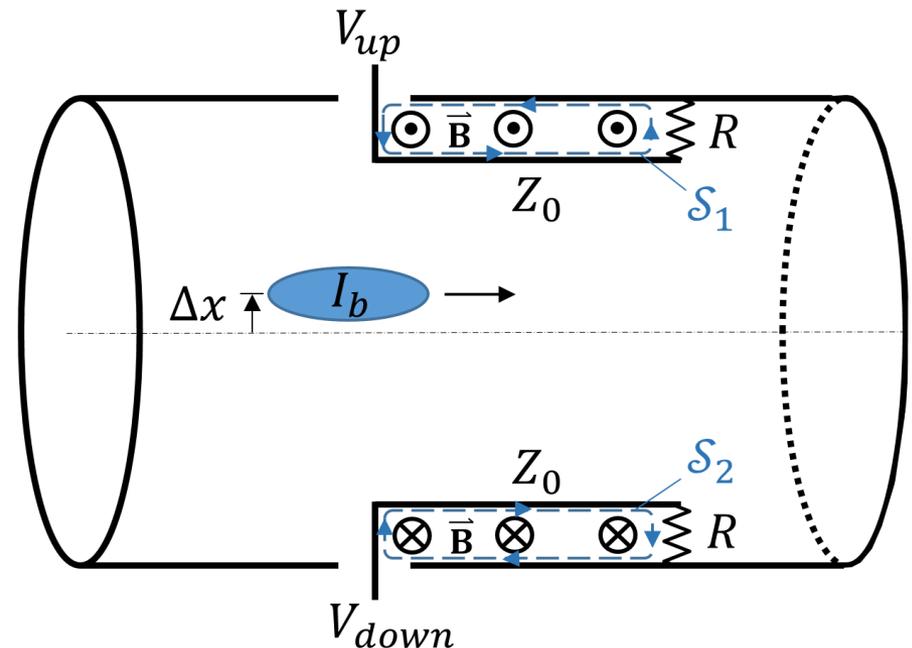


**TTFII Button BPM (DESY)
located inside an undulator magnet**

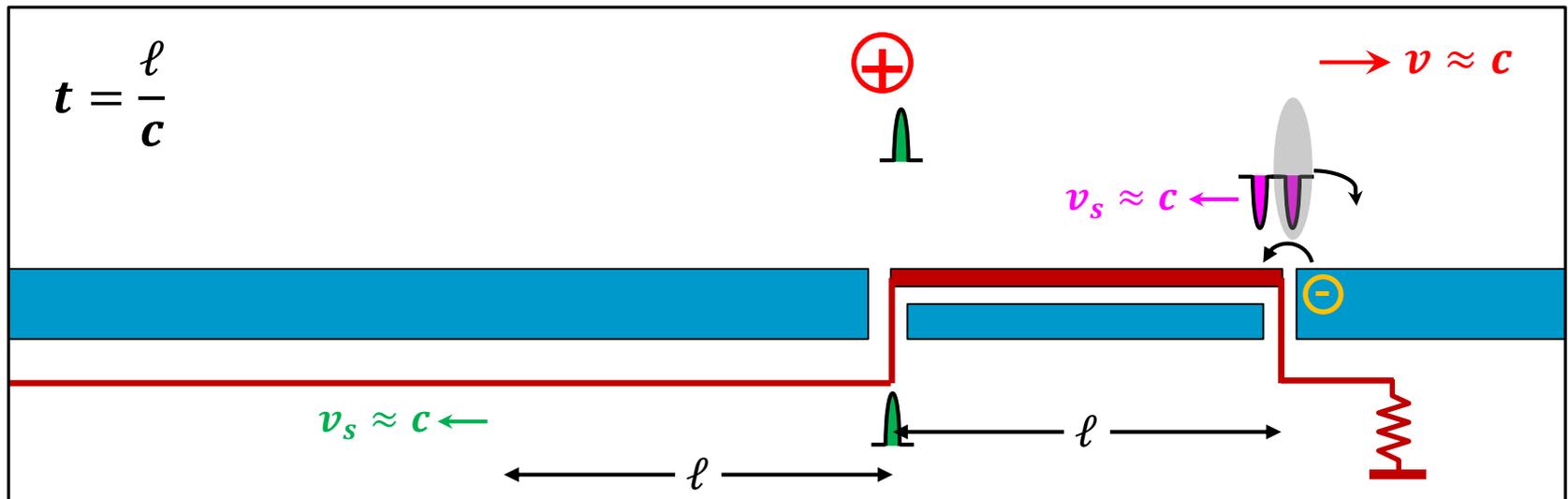
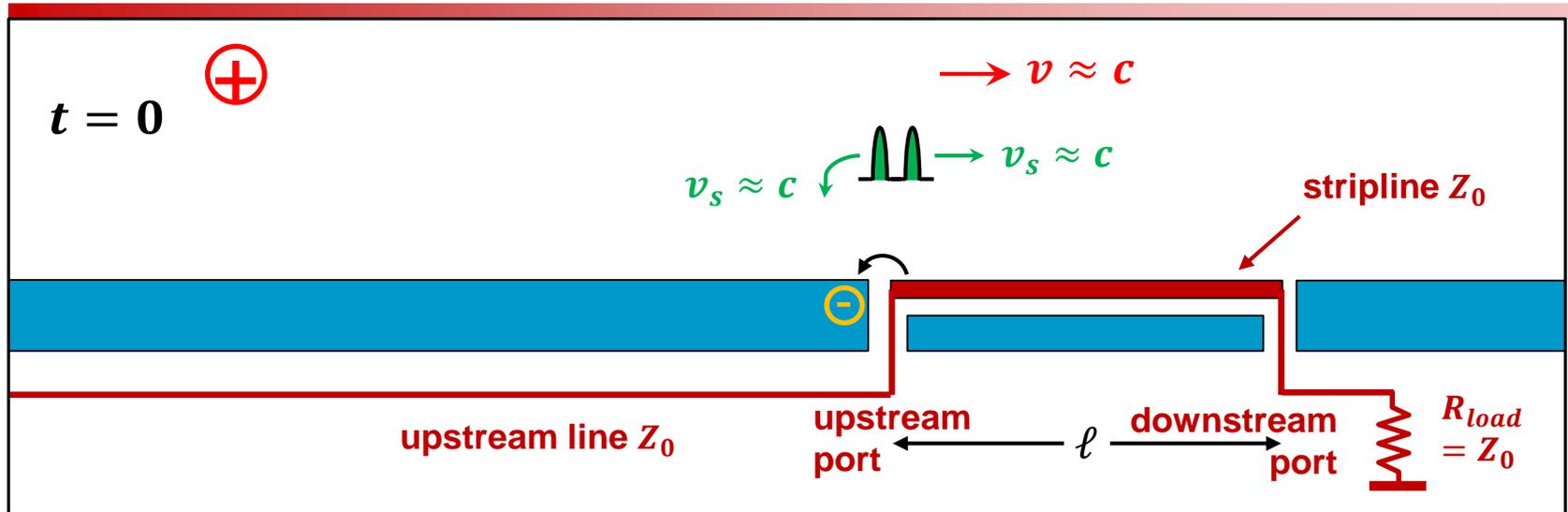


Stripline BPM

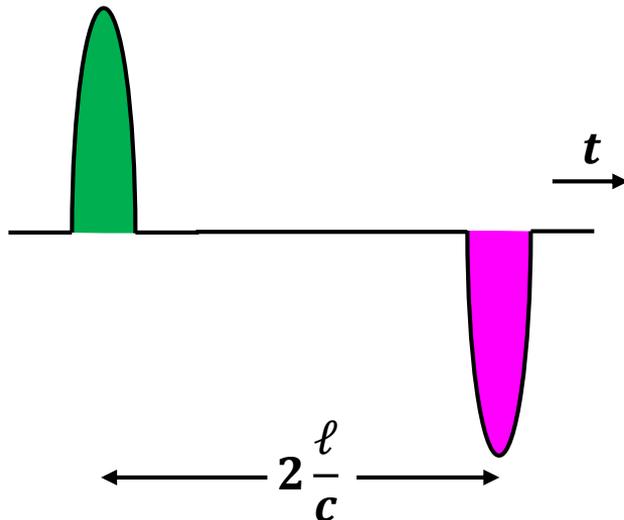
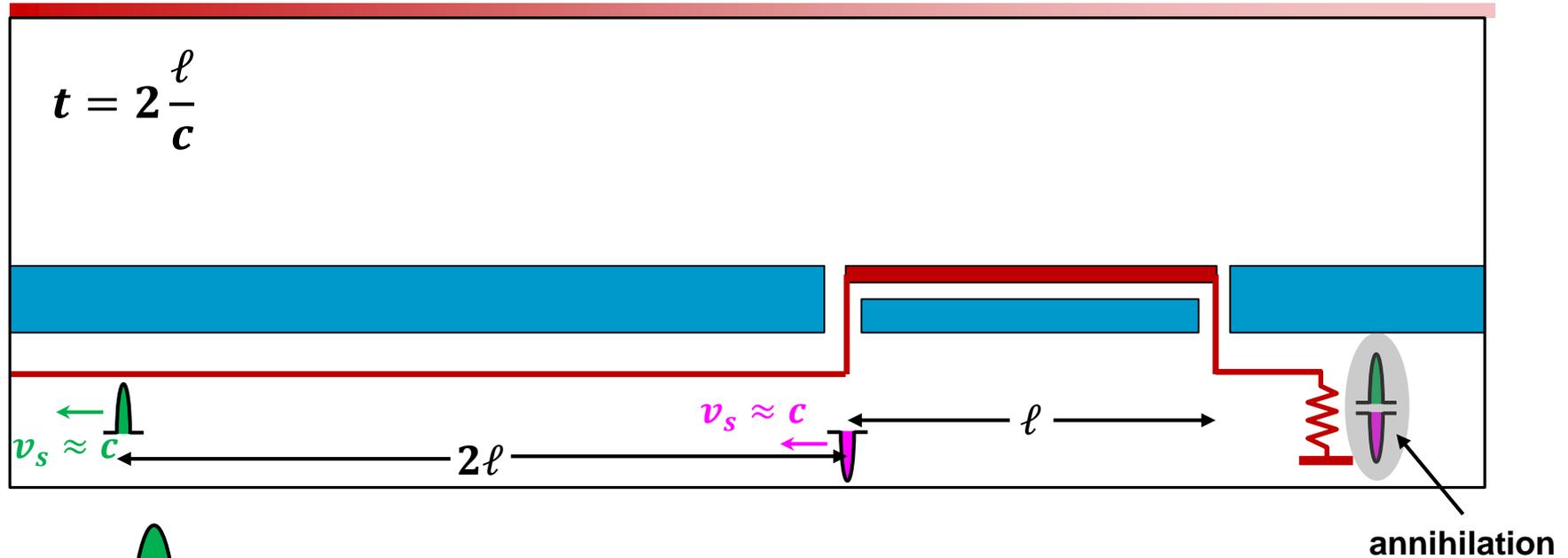
- Strip transmission-line of characteristic impedance $Z_0 = 50 \Omega$
- No dielectric materials ($\epsilon_r = 1$, vacuum)
 - TEM fields of the signal, except near the transition at the ports
- Coupling to the TEM field of the beam
 - Electric and magnetic field coupling
- Only the upstream port delivers a beam signal!
 - Directional coupler
- Both ports are terminated
 - Sometimes the downstream port is shorted to ground
 - No directivity



Stripline BPM Principle (1)

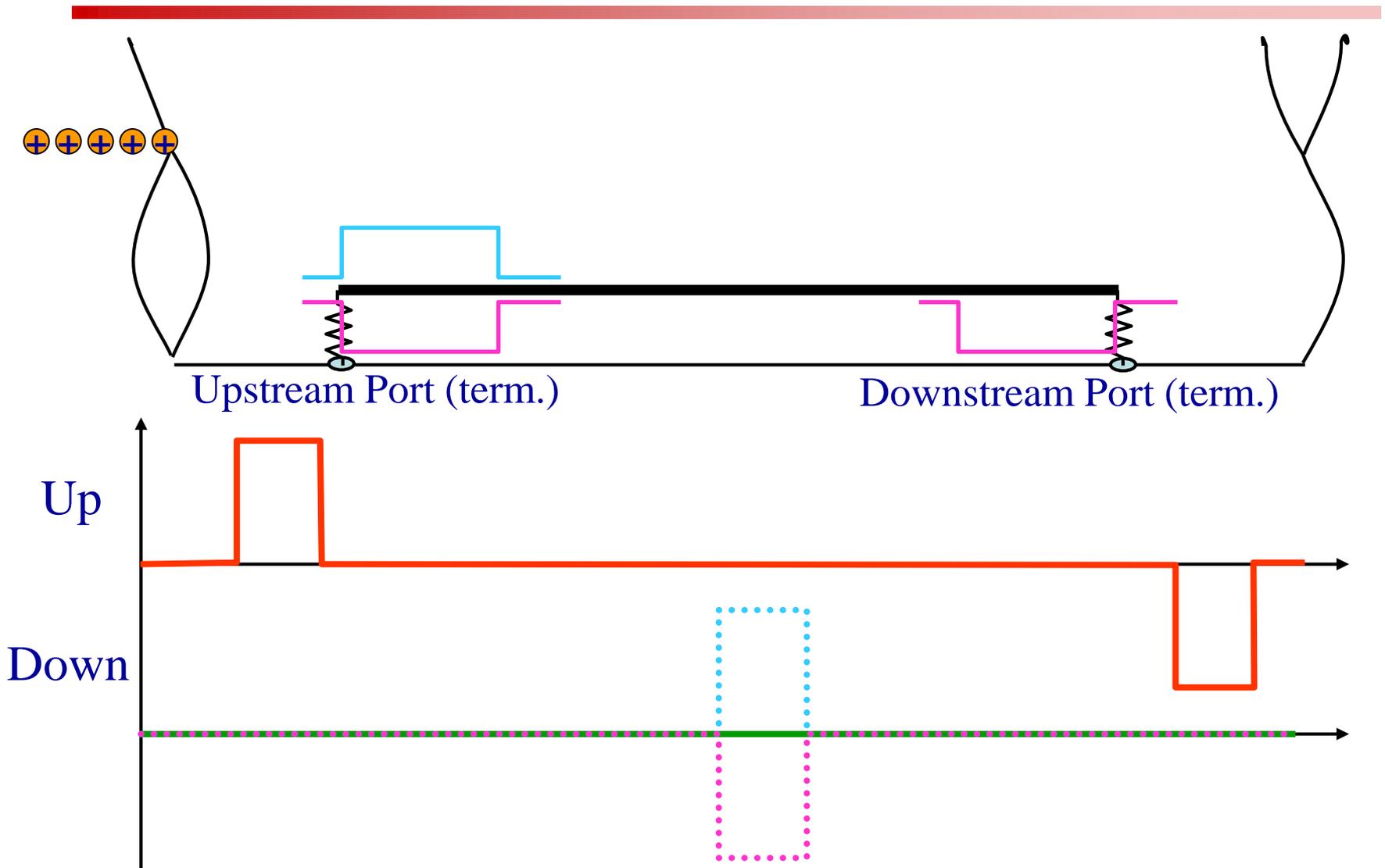


Stripline BPM Principle (2)

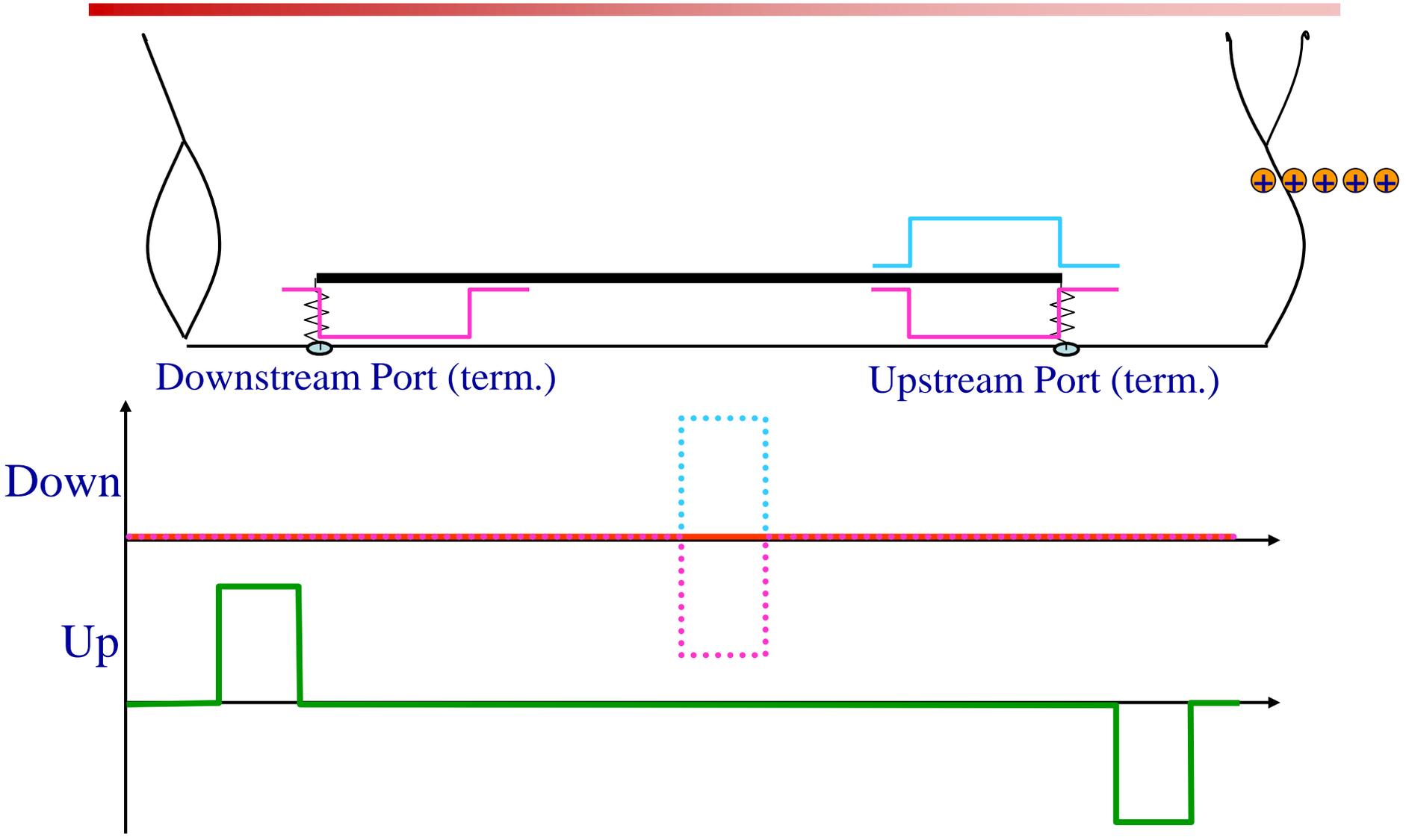


- **Stripline BPM response:**
 - δ -doublet pulse,
space by $t = 2 \ell / c$

Operational Principle: Beam from Left



Operational Principle: Beam from Right



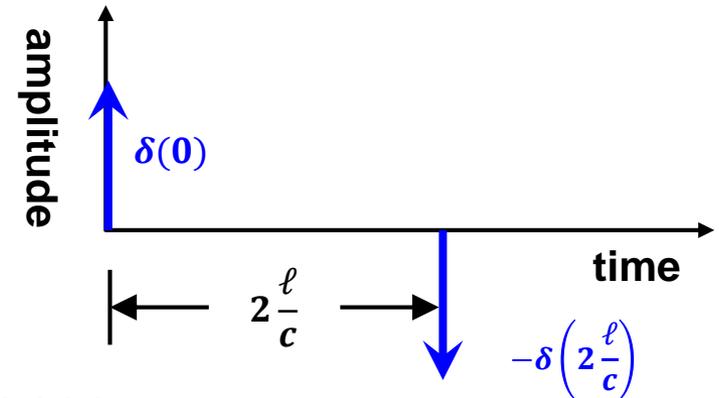
Transfer Impedance

- Time domain impulse response: δ -doublet pulse

characteristic impedance
typically 50Ω

$$z(t) = \phi \frac{Z_0}{2} \left[\delta(t) - \delta\left(t - 2\frac{\ell}{c}\right) \right]$$

coverage
factor



- Frequency domain transfer impedance

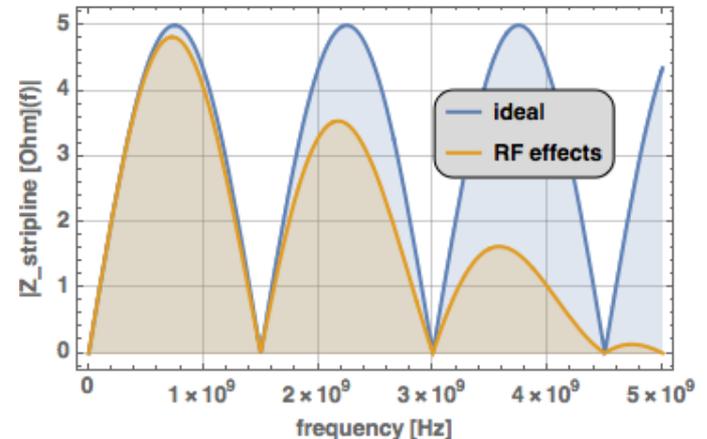
$$Z(\omega) = j\phi Z_0 e^{j\frac{\pi}{2}} e^{-j\omega\frac{\ell}{c}} \sin\left(\omega\frac{\ell}{c}\right)$$

- Lobes at:** $f_c = (2n - 1) \frac{c}{4\ell}$
- 3dB lobe bandwidth:**

$$f_{lo} = \frac{f_c}{2} \quad f_{BW} = f_{hi} - f_{lo} = f_c$$

$$f_{hi} = 3 f_{lo}$$

$Z_0 = 50 \Omega, \ell = 100 \text{ mm}, \phi = 0.1$



Stripline BPM

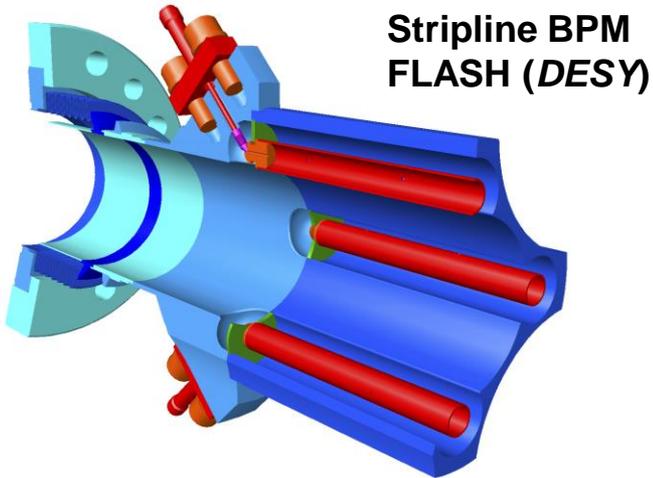
+ Pros

- **Well defined, and high transfer impedance function**
 - **For striplines with $l \gg w$, TEM operation**
 - **Can be “tuned” to match a dominant beam frequency**
 - Optimizing the length of the stripline
- **Matched characteristic impedance with 50 Ω source**
 - **Only for terminated ports**
 - Shorted downstream port: 0 Ω source, no directivity
 - **Reduced reflection effects between generator (stripline BPM) and read-out electronics!**
- **Directivity**
 - **Separate beam position measurement of counter-rotating beams**
 - e.g. near the IP in a ring collider

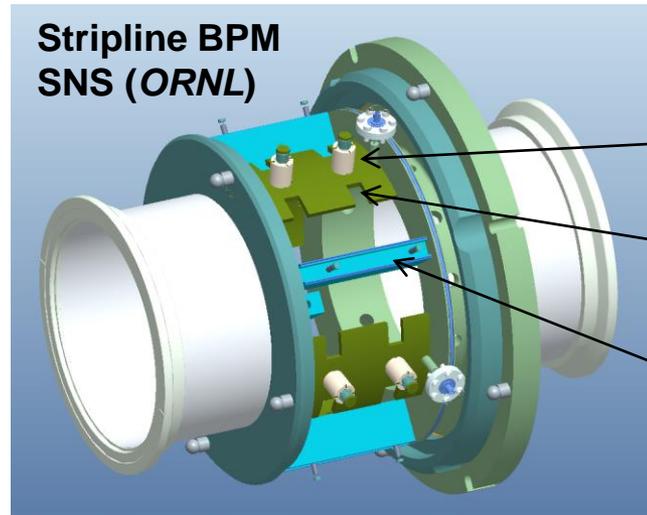
- Cons

- **Complicated, fragile mechanics, high costs, requires sufficient real-estate**
 - **More RF feedthroughs required, impedance control of the stripline, strip-electrode may require a support structure**

Examples of Stripline BPMs



**Stripline BPM
FLASH (DESY)**

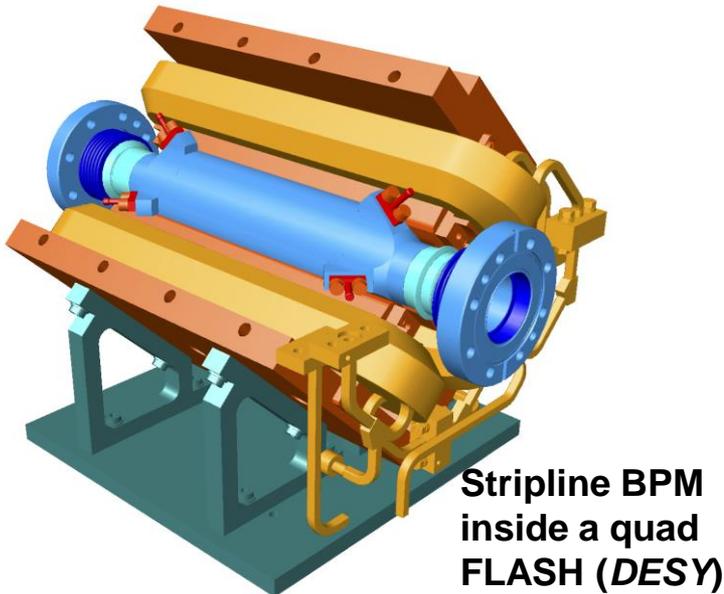


**Stripline BPM
SNS (ORNL)**

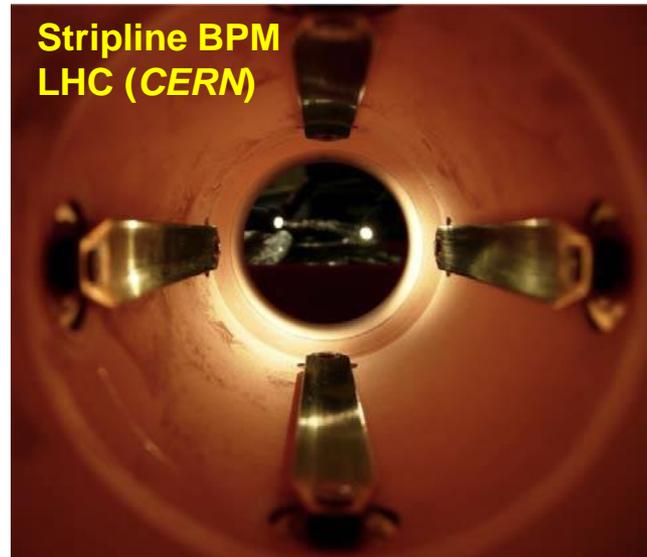
Ceramic posts
hold the electrode

Impedance Match
at the post

Inner-shielding bar
reduces electrode
to electrode
coupling



**Stripline BPM
inside a quad
FLASH (DESY)**



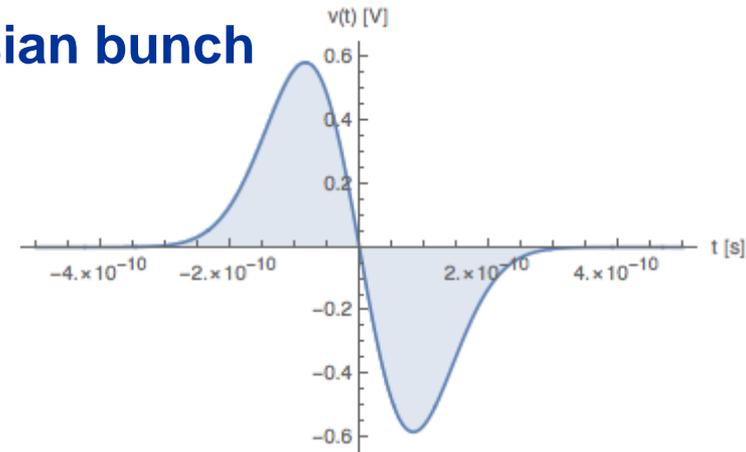
**Stripline BPM
LHC (CERN)**

Bunch Signals from broadband BPMs

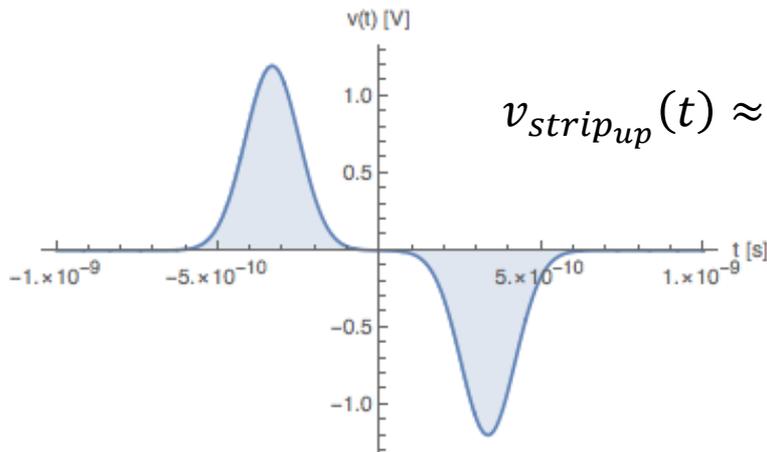
- Button BPM output signal to a Gaussian bunch

$$v_{button}(t) \approx \frac{A_{button} R_{load}}{\pi r_{pipe} \beta c} \frac{di_{beam}(t)}{dt}$$

$$= \frac{r_{button}^2 R_{load}}{2 r_{pipe} \beta c} \frac{eN}{\sqrt{2\pi} \sigma^3} t e^{-\frac{t^2}{2\sigma^2}}$$



- Stripline BPM output signal to a Gaussian bunch
 - at the upstream port



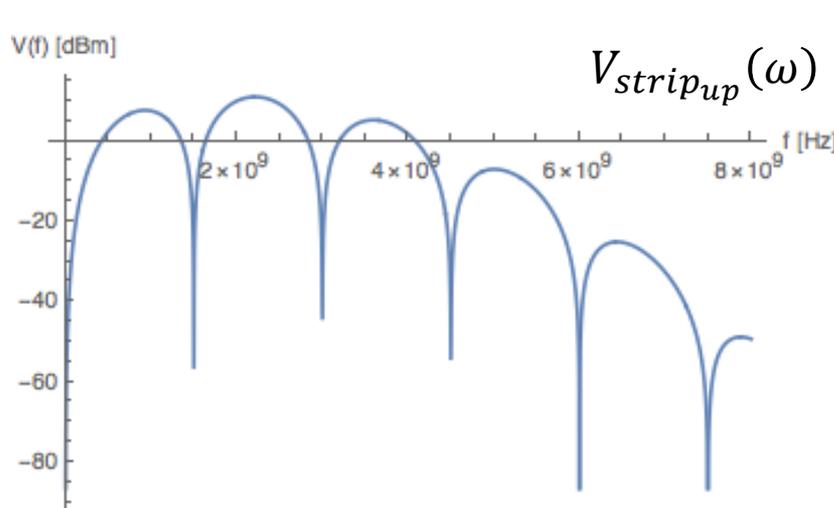
$$v_{strip_{up}}(t) \approx \frac{\phi Z_0}{2} \frac{eN}{\sqrt{2\pi} \sigma} \left\{ \exp\left[-\frac{(t + \tau)^2}{2\sigma^2}\right] - \exp\left[-\frac{(t - \tau)^2}{2\sigma^2}\right] \right\}$$

$$\text{with: } \tau = \frac{\ell}{2c} \left(\frac{1}{\beta} + \frac{1}{\beta_{strip}} \right)$$

$$Z_0 = R_{load} = 50 \Omega$$

Limits of the Analytical Analysis

- **Stripline BPM response in the frequency domain**



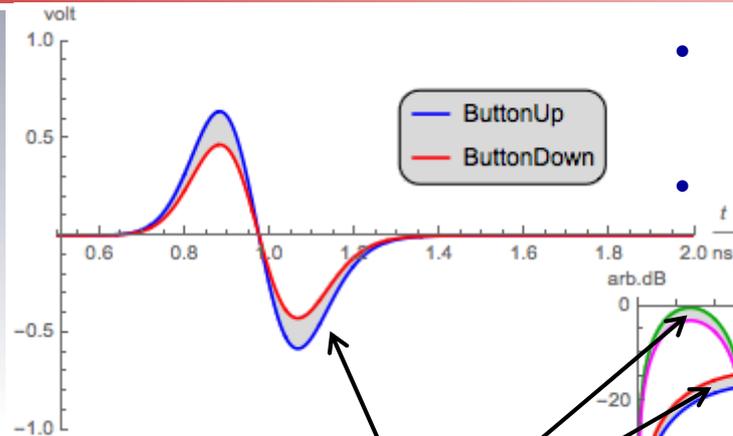
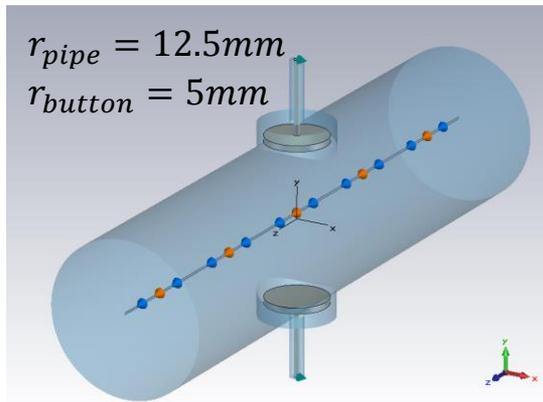
$$V_{strip_{up}}(\omega) \approx \sqrt{2}\phi Z_0 \frac{eN}{T} A(\omega) \sin \left[\frac{\omega \ell}{2c} \left(\frac{1}{\beta} + \frac{1}{\beta_{strip}} \right) \right]$$

$$\text{with: } A(\omega) = \exp \left(-\frac{\omega^2 \sigma^2}{2} \right)$$

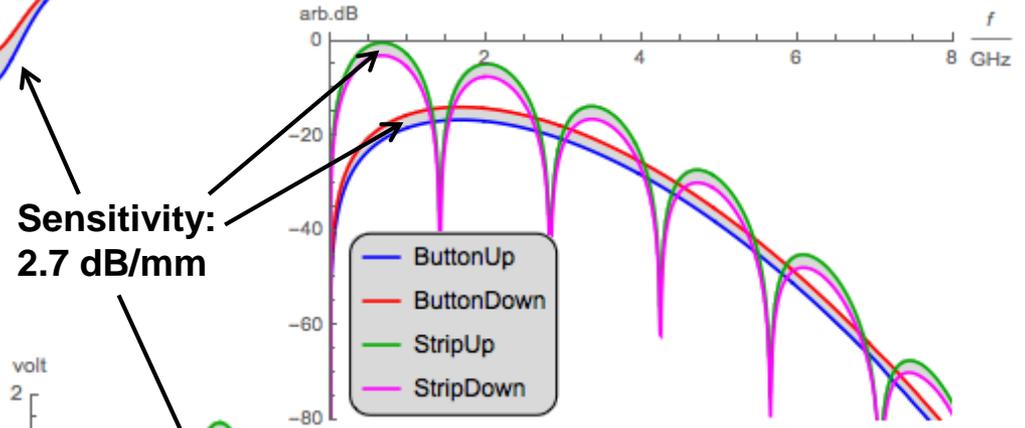
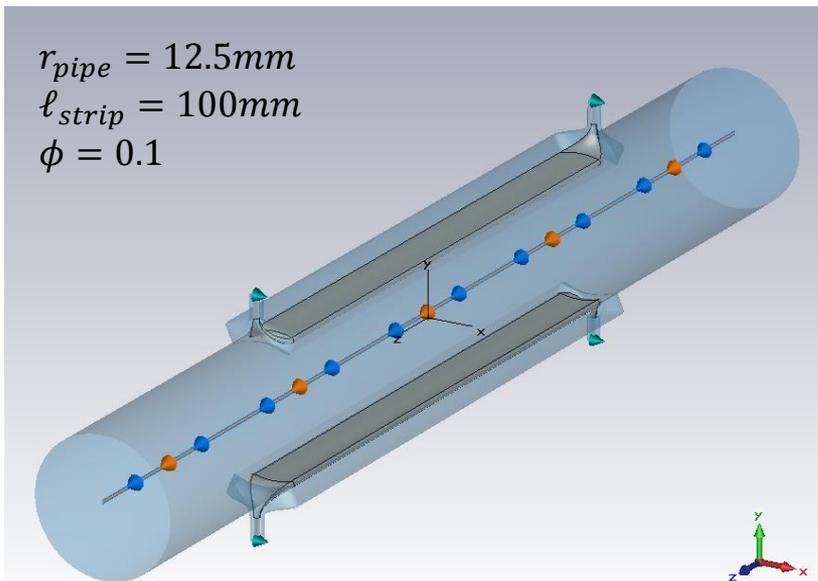
$$T = \frac{1}{f} = \frac{2\pi}{\omega}$$

- **No closed form expression which includes important details**
 - high frequency effects a the feedthrough transition
 - attenuation and dispersion of long coaxial cables
 - effects of low- β beams
- **Numerical analysis of the EM problem**
 - With a beam field as stimulus signal

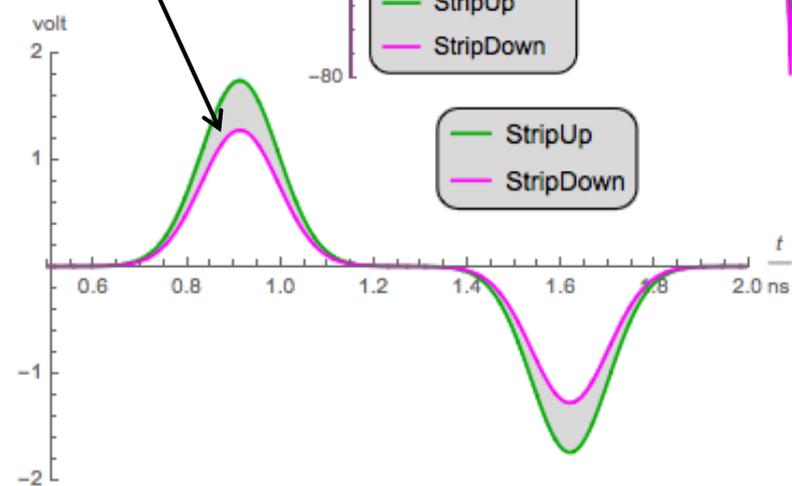
Numerically generated BPM Signals



- **Gaussian bunch:**
 - $N = 6.24e8, \sigma_s = 25\text{mm}, \beta = 1$
 - vertical beam offset=1mm
- **BPM sensitivity:**
 - e.g.: 2.7 dB / mm

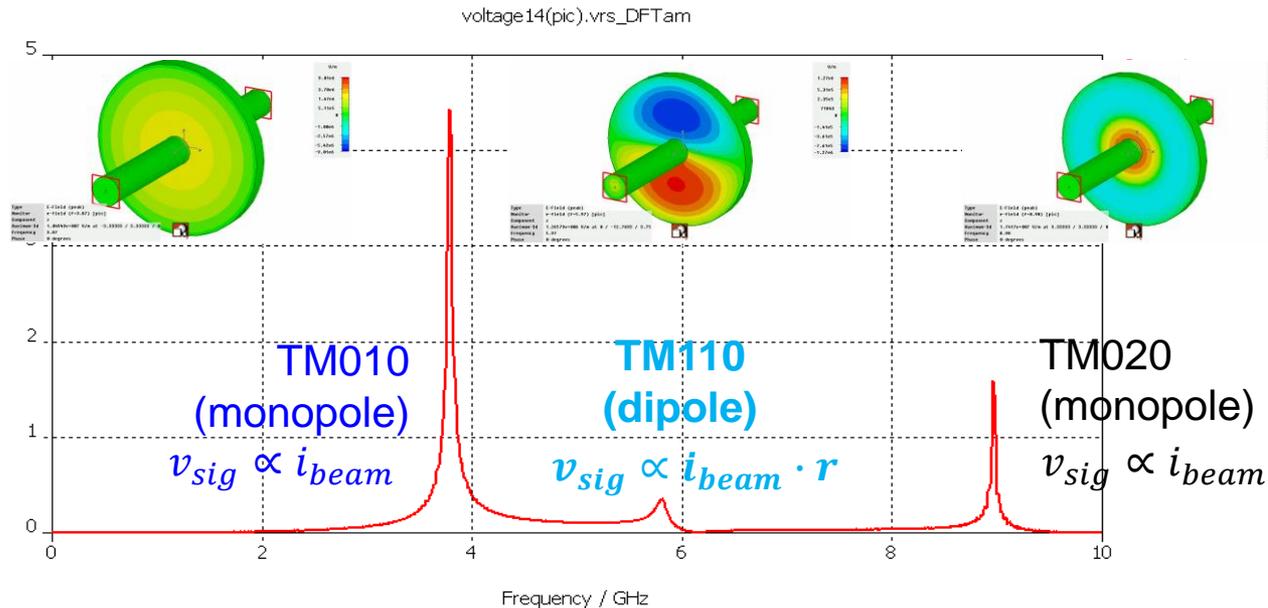


Sensitivity:
2.7 dB/mm



Resonant Cavity BPM

- Based on a beam-excited, passive resonator
 - Often a cylindrical “pillbox” cavity is used
 - Operating on the TM₁₁₀ dipole-eigenmode offers a higher resolution potential than comparable broadband BPMs (button, stripline).
 - No common-mode Σ signal, only a difference Δ signal
 - High transfer impedance, typically in the k Ω /mm range



courtesy
D. Lipka

Towards a Cavity BPM...

- **Eigenmodes in a brick-style resonator**

- 1st step towards a cavity BPM

- **Unfortunately you need to go through the math of the modal expansion of the vector potential Ψ ...** 😊

Laplace equation:

$$k_0^2 = \omega^2 \epsilon_0 \mu_0$$

k_0 : free space wave number

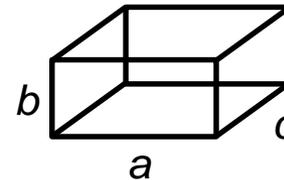
$$\Delta \Psi + k_0^2 \epsilon_r \mu_r \Psi = 0$$

$$k_0 = 2\pi / \lambda_0$$

λ_0 : free space wave length

Product ansatz (*Cartesian* coordinates):

$$\Psi = X(x)Y(y)Z(z)$$



standing waves

General solution (field components):

$$\Psi = \left\{ \begin{array}{l} A \cos(k_x x) + B \sin(k_x x) \\ \hat{A} e^{-jk_x x} + \hat{B} e^{-jk_x x} \end{array} \right\} \left\{ \begin{array}{l} C \cos(k_y y) + D \sin(k_y y) \\ \hat{C} e^{-jk_y y} + \hat{D} e^{-jk_y y} \end{array} \right\} \left\{ \begin{array}{l} E \cos(k_z z) + F \sin(k_z z) \\ \hat{E} e^{-jk_z z} + \hat{F} e^{-jk_z z} \end{array} \right\}$$

separation condition:

$$k_x^2 + k_y^2 + k_z^2 = k_0^2 \epsilon_r \mu_r$$

$$k_x = \frac{m\pi}{a} \quad k_y = \frac{n\pi}{b} \quad k_z = \frac{p\pi}{c}$$

travelling waves

Eigen frequencies:

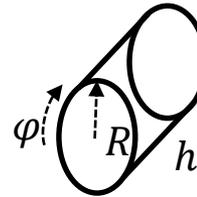
$$f_{mnp} = \frac{c_0}{2\pi \epsilon_r \mu_r} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 + \left(\frac{p\pi}{c}\right)^2}$$

Cylindrical “Pillbox” Cavity Resonator

- Same procedure, but now with cylindrical functions ☹️☹️

Product ansatz (cylindrical coordinates):

$$\Psi = R(\rho)F(\varphi)Z(z)$$



standing waves

General solution (field components):

$$\Psi = \left\{ \begin{array}{l} A J_m(k_r \rho) + B N_m(k_r \rho) \\ \hat{A} H_m^{(2)}(k_r \rho) + \hat{B} H_m^{(2)}(k_r \rho) \end{array} \right\} \left\{ \begin{array}{l} C \cos(m\varphi) + D \sin(m\varphi) \\ \hat{C} e^{-jm\varphi} + \hat{D} e^{-jm\varphi} \end{array} \right\} \left\{ \begin{array}{l} E \cos(k_z z) + F \sin(k_z z) \\ \hat{E} e^{-jk_z z} + \hat{F} e^{-jk_z z} \end{array} \right\}$$

travelling waves

$J_m, N_m, H_m^{(1,2)}$: cylindrical functions (Bessel, Hankel, Neumann)

^(1,2)see Abramowitz and Stegun

separation condition:

$$k_r^2 + k_z^2 = k_0^2 \epsilon_r \mu_r$$

Eigen frequencies:

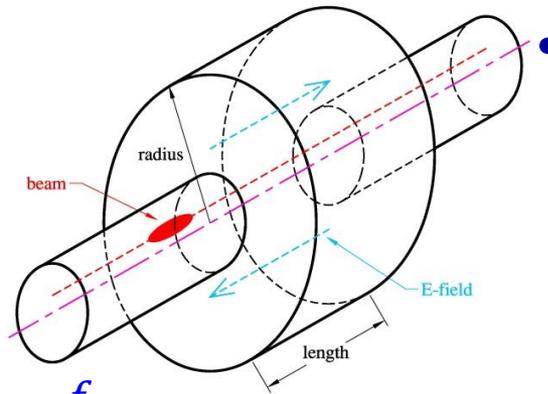
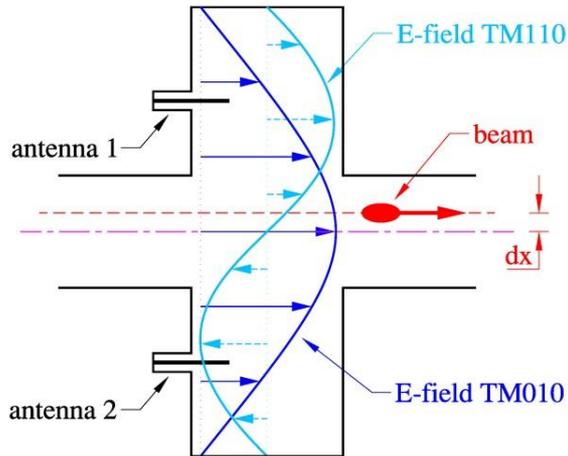
$$f_{TMmnp} = \frac{c_0}{2\pi \epsilon_r \mu_r} \sqrt{\left(\frac{j_{mn}}{R}\right)^2 + \left(\frac{p\pi}{h}\right)^2}$$

$$f_{TEmp} = \frac{c_0}{2\pi \epsilon_r \mu_r} \sqrt{\left(\frac{j'_{mn}}{R}\right)^2 + \left(\frac{p\pi}{h}\right)^2}$$

j_{mn} being the n^{th} root of $J_m(x)$

j'_{mn} being the n^{th} root of $J'_m(x)$

Cavity BPM



• Beam couples to:

$$E_z = C J_1 \left(\frac{j_{111} r}{R} \right) e^{i\omega t} \cos \varphi$$

dipole (TM_{110}) and monopole (TM_{010}) & other modes

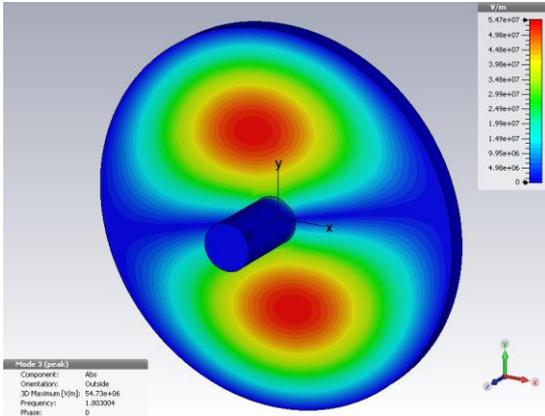
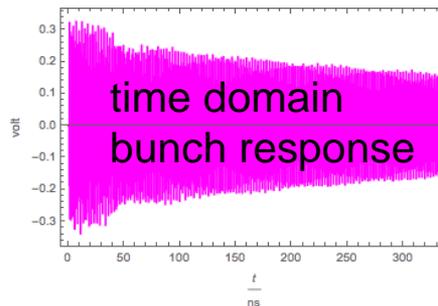
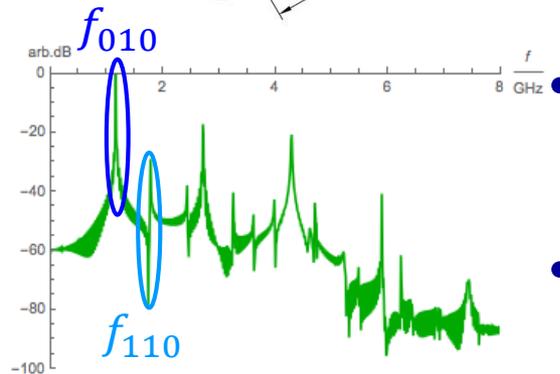
• Common mode (TM_{010}) frequency discrimination

• Decaying RF signal response

– Position signal: TM_{110}

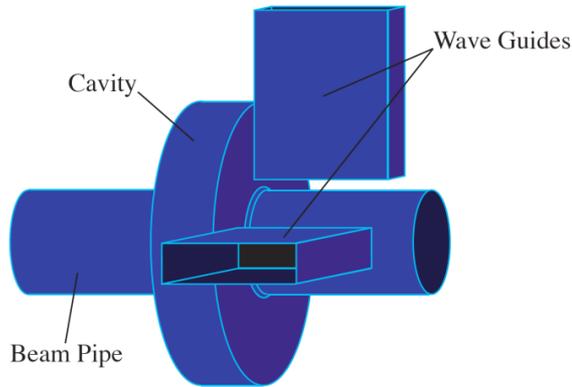
➤ Requires normalization and a phase reference

– Intensity signal: TM_{010}

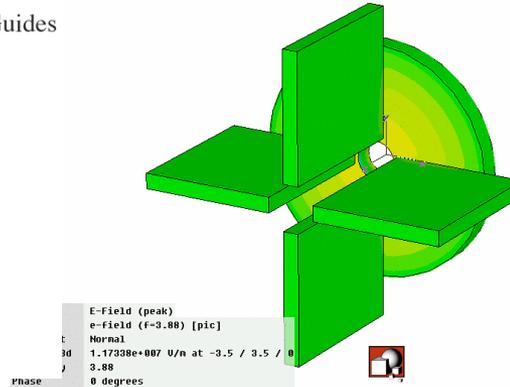


$$\begin{aligned} r_{pipe} &= 12.5\text{mm} \\ r_{cav} &= 100\text{mm} \\ \ell_{cav} &= 10\text{mm} \end{aligned}$$

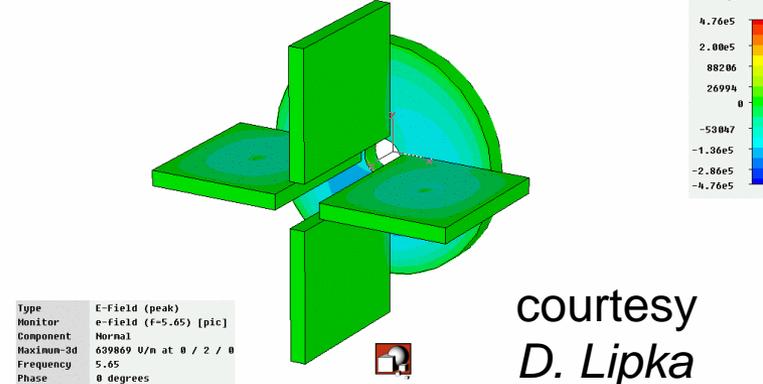
Common-mode free Cavity BPMs



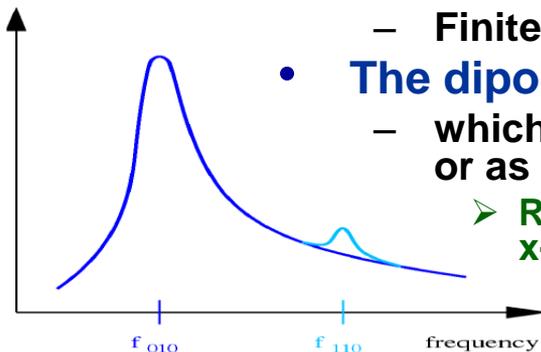
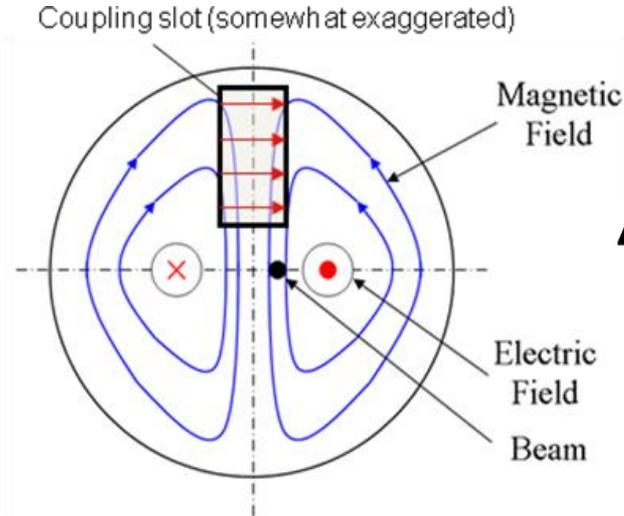
monopole mode



dipole mode



courtesy
D. Lipka



- Add slot-coupled waveguide TE_{01} -mode high-pass filter

$$f_{010} < f_{10} = \frac{1}{2a\sqrt{\epsilon\mu}} < f_{110}$$

between cavity and coaxial output port.

- Finite Q of TM_{010} still leaks into TM_{110} !
- The dipole mode has two polarizations
 - which will orientate along imperfections, or as wanted along the coupling slots
 - Requires tight tolerances to minimize x-y cross talk

Cavity BPM

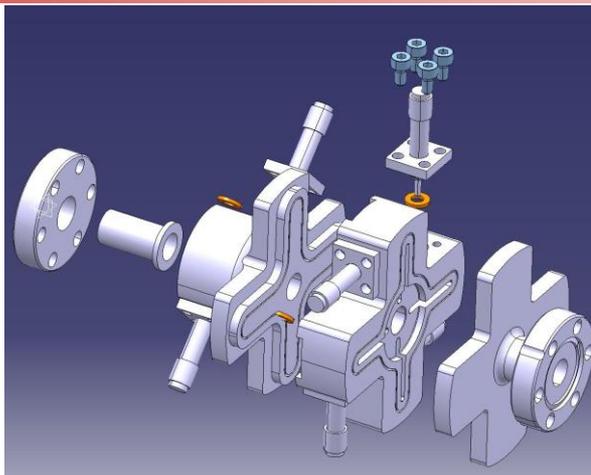
+ Pros

- No or minimum common mode signal contribution in the Δ -signal
 - Frequency discrimination of dipole (TM₁₁₀) and monopole (TM₀₁₀) modes
- High resolution potential
 - High shunt (transfer) impedance of the TM₁₁₀ mode
 - Even for lower Q tuning of the TM₁₁₀ mode
 - Sub- μm signal pass resolution potential

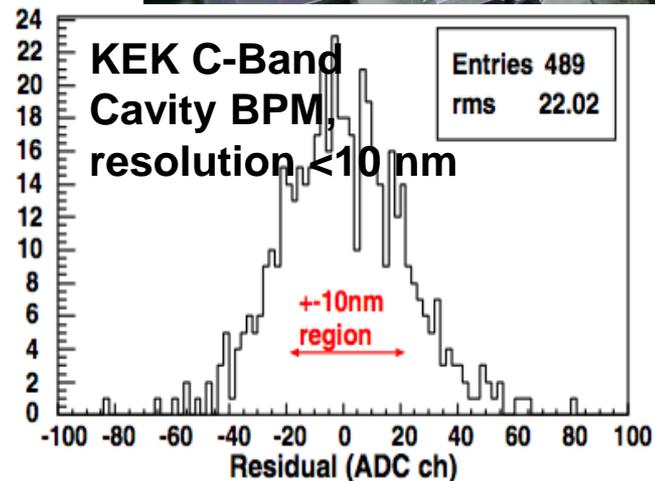
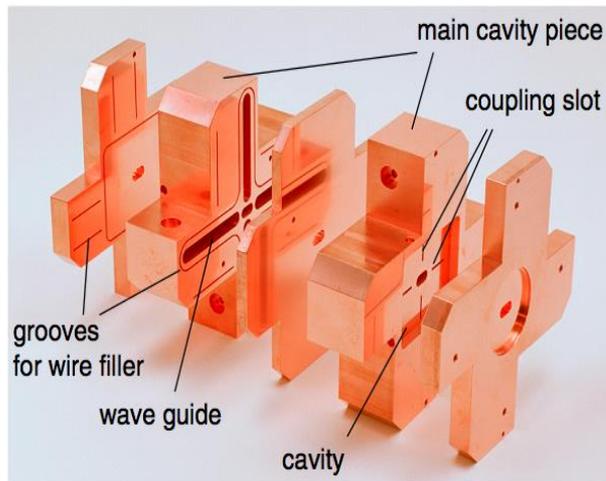
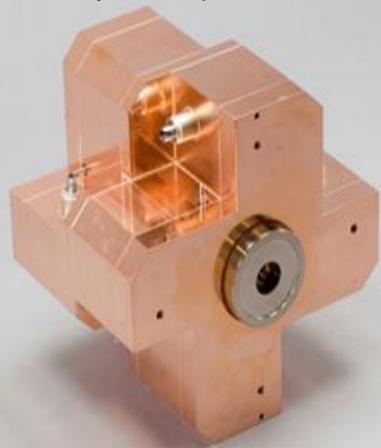
- Cons

- High beam coupling impedance
 - No free lunch: high impedance may cause beam break-up and/or instabilities
 - No or very limited use in ring accelerators
- Requires a reference monopole mode (TM₀₁₀) resonator
 - Beam phase and intensity
- Limited position range
 - ~half aperture
- Requires advanced RF read-out electronics
- High-Q resonator may not be suitable for single bunch position measurements

Examples of Cavity BPMs



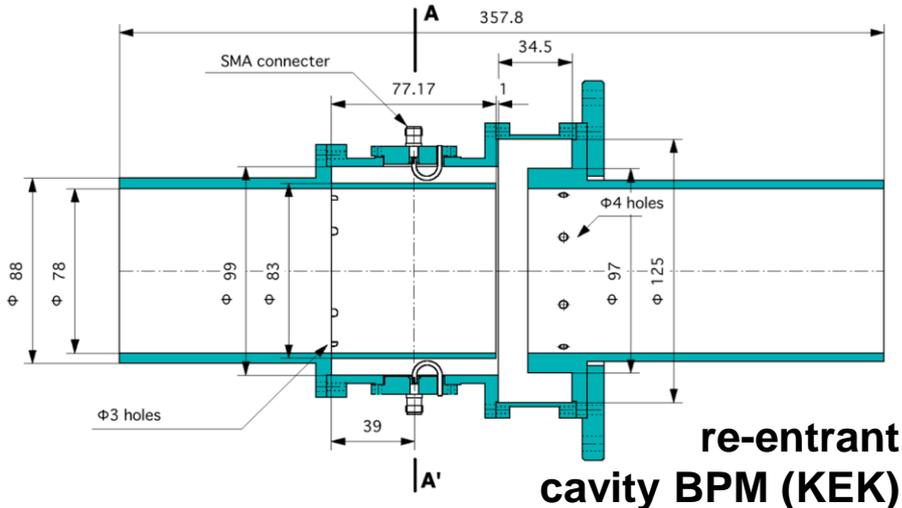
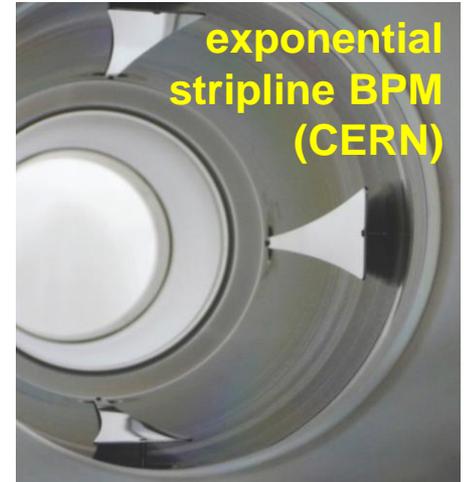
Courtesy of D. Lipka & Y. Honda



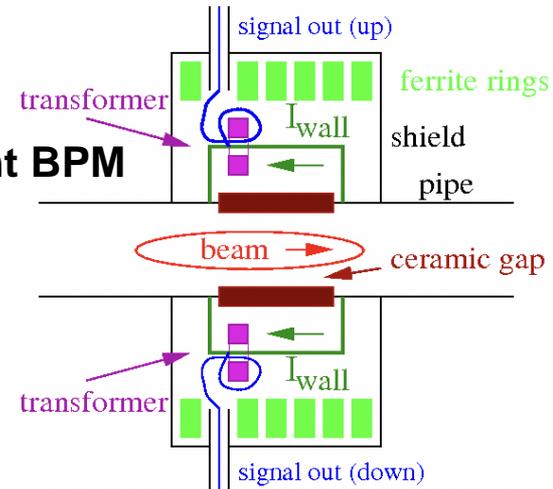
Other Types of BPM Pickups

- Less popular, but sometimes better suited for a specific application
 - Stripline BPM with shorted downstream ports
 - Exponentially tapered stripline BPM
 - Re-entrant cavity BPM
 - Resonant stripline of button BPM
 - Inductive BPM, ...

➤ In common: based on symmetry



inductive wall-current BPM (CERN)



Summary Part 1

- **What have we learned?!**
 - It is hard to stay up on Sunday morning and to follow a BPM session...
- **What else have we learned,... about BPMs?!**
 - **BPMs are based on symmetry**
 - **Detect the asymmetric beam position by a pair of symmetrically arranged electrodes or by a symmetric resonant, beam driven cavity**
 - **BPM pickups are non-invasive**
 - **Detect the beam's center-of-charge by electromagnetic coupling**
 - **Broadband BPM electrode (stripline, button, etc.) signals**
 - **contain information of the beam position and the beam intensity**
 - Need to “normalize” to extract a position signal
 - **May have a non-linear position characteristic**
 - **Have a broad spectral response**
 - **Resonant cavity BPMs**
 - **Have a frequency selective beam response**
 - **Have a high resolution potential**
 - No or only little “common mode” signal contribution in the dipole mode