

Field Solvers



TECHNISCHE
UNIVERSITÄT
DARMSTADT

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Lecture 2 : Magnetoquasistatic Formulation and Discretisation

Overview

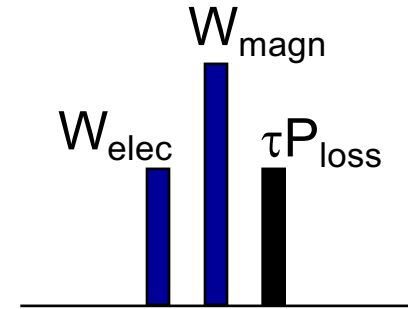


- **magnetoquasistatic formulation**
- discretisation in space
- boundary and symmetry conditions
- reduction to 2D models
- modelling of coils and permanent magnets

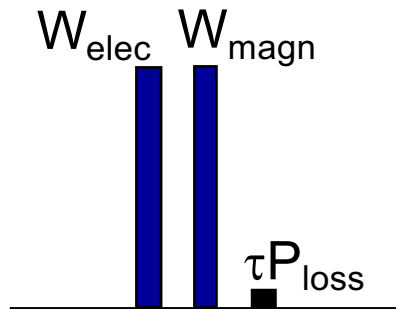
EM Field Simulation



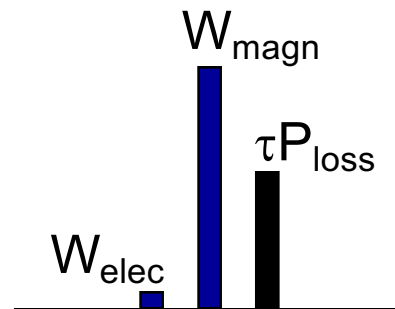
full Maxwell equations



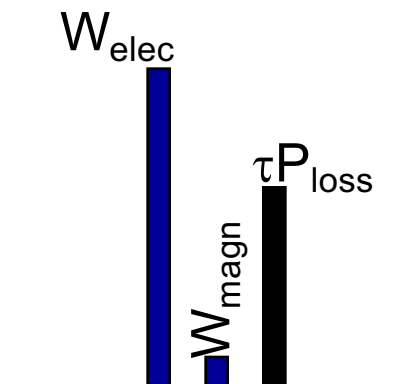
wave equation



magnetoquasistatics



electroquasistatics



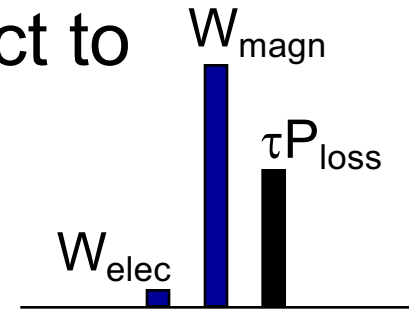
Magnetoquasistatics (1)



- neglect displacement currents with respect to conducting currents

- Ampère-Maxwell

$$\nabla \times \vec{H} = \vec{j} + \cancel{\frac{\partial \vec{D}}{\partial t}}$$



- magnetic vector potential \vec{A}
 - conservation of magnetic flux

$$\nabla \cdot \vec{B} = 0 \quad \rightarrow \quad \vec{B} = 0 + \nabla \times \vec{A}$$

- electric scalar potential φ (voltage)
 - Faraday-Lenz

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\nabla \times \frac{\partial \vec{A}}{\partial t} \quad \rightarrow \quad \vec{E} = -\frac{\partial \vec{A}}{\partial t} - \nabla \varphi$$

Magnetoquasistatics (2)



Ampère

$$\nabla \times \vec{H} = \vec{j}$$



$$\nabla \times (\nu \vec{B}) = \sigma \vec{E}$$



$$\nabla \times (\nu \nabla \times \vec{A}) + \sigma \frac{\partial \vec{A}}{\partial t} = -\sigma \underbrace{\nabla \varphi}_{\vec{j}_s}$$

$$\vec{B} = \mu \vec{H} = \frac{1}{\nu} \vec{H}$$

permeability

$$\vec{j} = \sigma \vec{E}$$

reluctivity

conductivity

source current density

parabolic partial differential equation

↔ elliptic PDEs (e.g. electrostatics,
magnetostatics)

↔ hyperbolic PDEs (e.g. wave equation)

Physical Meaning (1)



flux

$$\phi = \int_S \vec{B} \cdot d\vec{S}$$



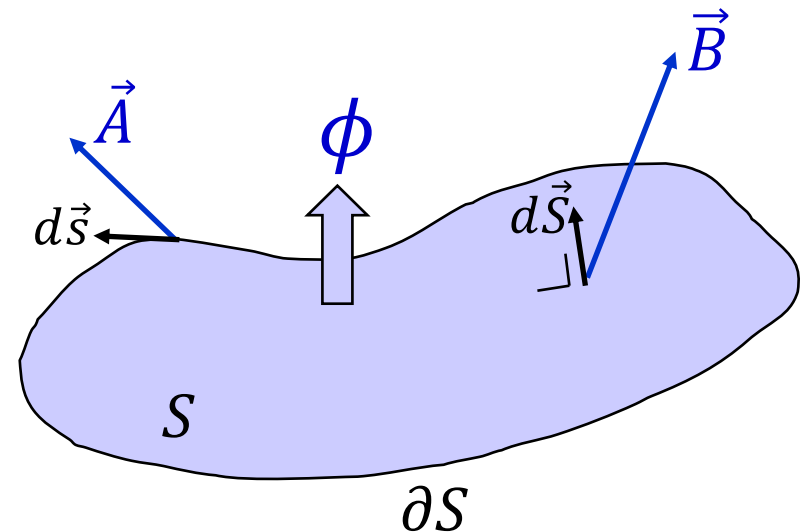
definition magnetic vector potential

$$\phi = \int_S \nabla \times \vec{A} \cdot d\vec{S}$$



Stokes

$$\phi = \oint_{\partial S} \vec{A} \cdot d\vec{s}$$



induced voltage

$$u_{\text{ind}} = - \frac{d}{dt} \oint_{\partial S} \vec{A} \cdot d\vec{s}$$

Physical Meaning (2)

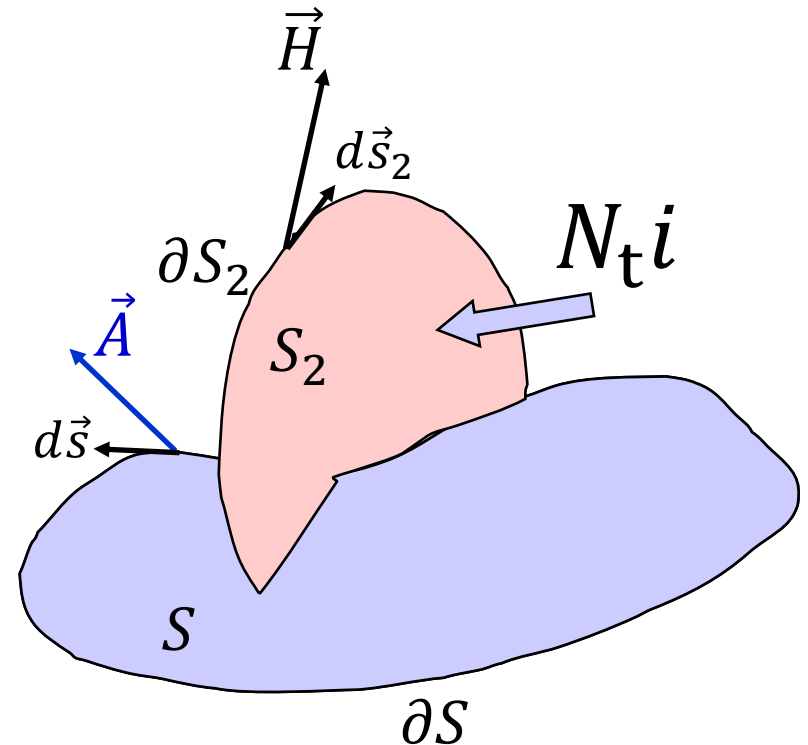
flux $\phi = \oint_{\partial S} \vec{A} \cdot d\vec{s}$

induced voltage

$$u_{\text{ind}} = - \frac{d}{dt} \oint_{\partial S} \vec{A} \cdot d\vec{s}$$

Ampère

$$N_{\text{t}}i = \oint_{\partial S_2} \vec{H} \cdot d\vec{s}_2$$



make (a lot of) drawings (also if they look stupid)

Overview



- magnetoquasistatic formulation
- **discretisation in space**
- boundary and symmetry conditions
- reduction to 2D models
- modelling of coils and permanent magnets

Spatial Discretisation (1)



- weighted residual approach

$$\nabla \times (v \nabla \times \vec{A}) + \sigma \frac{\partial \vec{A}}{\partial t} = \vec{J}_s \quad \text{in } \Omega$$

$$\int_V \left(\nabla \times (v \nabla \times \vec{A}) + \sigma \frac{\partial \vec{A}}{\partial t} \right) \cdot \vec{w}_i \, dV = \int_V \vec{J}_s \cdot \vec{w}_i \, dV \quad \forall \vec{w}_i(x, y, z)$$

$\vec{w}_i(x, y, z)$ vectorial „weighting functions“
vectorial „test functions“

- scalar product : $(\vec{\alpha}, \vec{\beta}) = \int_V \vec{\alpha} \cdot \vec{\beta} \, dV$

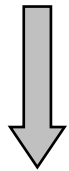
Spatial Discretisation (2)



→ weak formulation

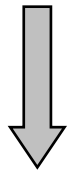
do not study derivations by heart,
know how to get from one to the next line!

$$\int_V \left(\nabla \times (v \nabla \times \vec{A}) + \sigma \frac{\partial \vec{A}}{\partial t} \right) \cdot \vec{w}_i \, dV = \int_V \vec{J}_s \cdot \vec{w}_i \, dV \quad \forall \vec{w}_i(x, y, z)$$



$$(\nabla \times \vec{v}) \cdot \vec{w} = \nabla \cdot (\vec{v} \times \vec{w}) + \vec{v} \cdot \nabla \times \vec{w}$$

$$\int_V \left(\nabla \cdot (v \nabla \times \vec{A} \times \vec{w}_i) + v \nabla \times \vec{A} \cdot \nabla \times \vec{w}_i + \sigma \frac{\partial \vec{A}}{\partial t} \cdot \vec{w}_i \right) dV = \int_V \vec{J}_s \cdot \vec{w}_i \, dV$$



Gauss

$$\oint_{\partial V} \underbrace{v \nabla \times \vec{A} \times \vec{w}_i}_{\vec{H}} \cdot d\vec{S} + \int_V \left(\underbrace{v \nabla \times \vec{A} \cdot \nabla \times \vec{w}_i}_{\text{only first derivative required}} + \sigma \frac{\partial \vec{A}}{\partial t} \cdot \vec{w}_i \right) dV = \int_V \vec{J}_s \cdot \vec{w}_i \, dV$$

only first derivative required → „weak“ formulation

Spatial Discretisation (3)

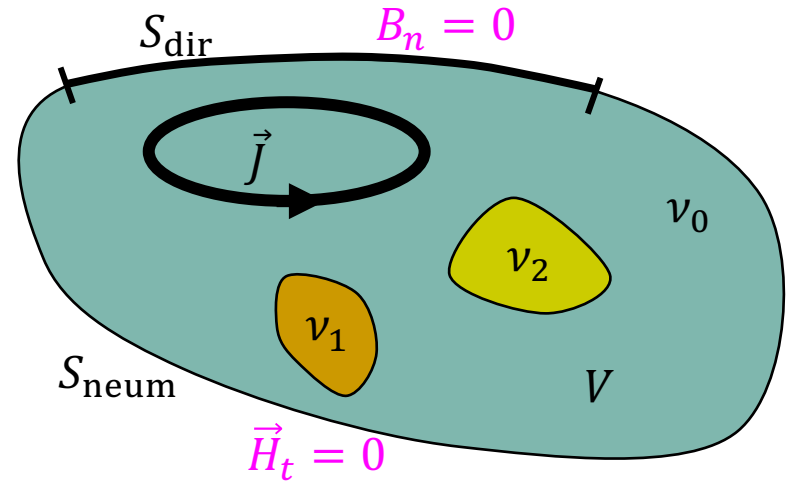


Dirichlet BC at S_{dir}

$$\vec{A} \times \vec{n} = \vec{A}_{\text{dir}} \times \vec{n} \iff \vec{B} \cdot \vec{n} = B_n$$

homogeneous Neumann BC at S_{neum}

$$\vec{H}_t = (\nu \nabla \times \vec{A}) \times \vec{n} = 0$$



$$\underbrace{\int_{S_{\text{neum}}} \overbrace{(\nu \nabla \times \vec{A}) \times \vec{w}_i}^{\vec{H}} \cdot d\vec{S}}_{= 0} + \underbrace{\int_{S_{\text{dir}}} (\nu \nabla \times \vec{A}) \times \vec{w}_i \cdot d\vec{S}}_{= 0} \quad \forall \vec{w}_i(x, y, z)$$

„natural“
„essential“

boundary condition
boundary condition

$\forall \vec{w}_i : \vec{w}_i \times \vec{n} = 0 \text{ at } S_{\text{dir}}$

Spatial Discretisation (4)

- discretisation

$$\vec{A}(x, y, z) = \sum_j u_j \vec{v}_j(x, y, z) \quad \vec{v}_j(x, y, z) \times \vec{n} = 0 \quad \text{at } S_{\text{dir}}$$

$\left\{ \begin{array}{l} \vec{v}_j(x, y, z) \\ u_j \end{array} \right.$ „shape/form functions“, „trial functions“
unknowns, degrees of freedom

- Ritz-Galerkin method

$$\vec{v}_j(x, y, z) = \vec{w}_j(x, y, z)$$

- Petrov-Galerkin method

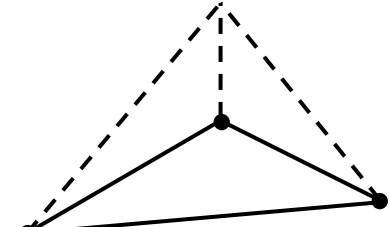
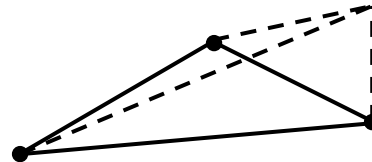
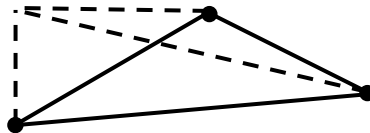
$$\vec{v}_j(x, y, z) \neq \vec{w}_j(x, y, z)$$

Spatial Discretisation (5)

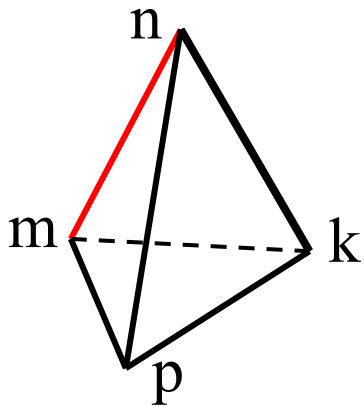


▪ nodal shape functions

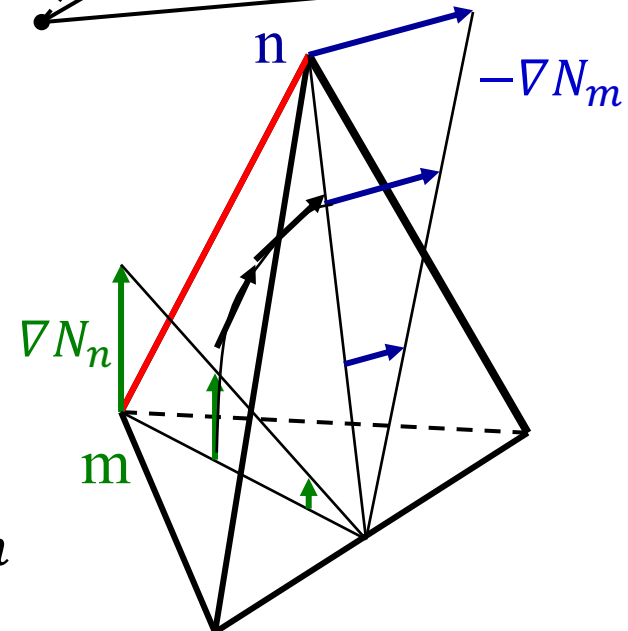
$$\varphi(x, y) = u_1 N_1(x, y) + u_2 N_2(x, y) + u_3 N_3(x, y)$$



▪ edge shape functions



$$\vec{v}(\vec{x}) = N_m \nabla N_n - N_n \nabla N_m$$



Spatial Discretisation (6)



▪ discretization

$$\int_V \left(\nu \nabla \times \vec{A} \cdot \nabla \times \vec{w}_i + \sigma \frac{\partial \vec{A}}{\partial t} \cdot \vec{w}_i \right) dV = \int_V \vec{J}_s \cdot \vec{w}_i dV \quad \forall \vec{w}_i(x, y, z)$$

$$\vec{A}(x, y, z) = \sum u_j \vec{w}_j(x, y, z)$$

`curlcurl_ll.m`

`edgemass_ll.m`

$$\sum_j \left(u_j \underbrace{\int_V \nu \nabla \times \vec{w}_j \cdot \nabla \times \vec{w}_i dV}_{= k_{ij}} + \frac{du_j}{dt} \underbrace{\int_V \sigma \vec{w}_j \cdot \vec{w}_i dV}_{= m_{ij}} \right) = \underbrace{\int_V \vec{J}_s \cdot \vec{w}_i dV}_{= f_i}$$

$$[k_{ij}][u_j] + [m_{ij}] \left[\frac{du_j}{dt} \right] = [f_i]$$

K and M symmetric,
semi-positive-definite

Spatial Discretisation (7)



$$\sum_j \left(u_j \int_V v \nabla \times \vec{w}_j \cdot \nabla \times \vec{w}_i \, dV + \frac{du_j}{dt} \int_V \sigma \vec{w}_j \cdot \vec{w}_i \, dV \right)$$

$$\nabla \times \vec{w}_j = \sum_q c_{jq} \vec{z}_q$$

$$\sum_j \left(u_j \sum_p \sum_q c_{ip} c_{jq} \int_V v \vec{z}_q \cdot \vec{z}_p \, dV + \frac{du_j}{dt} \int_V \sigma \vec{w}_j \cdot \vec{w}_i \, dV \right) =$$

$\hat{\mathbf{a}}_j$

$\mathbf{M}_{v,p,q}^{\text{FE}}$

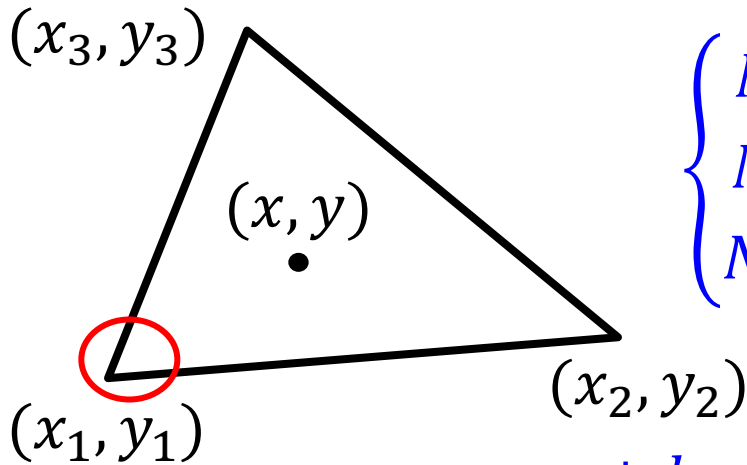
$\mathbf{M}_{\sigma,i,j}^{\text{FE}}$

$\hat{\mathbf{j}}_{s,i}$

The diagram shows a triangular element with nodes labeled n (top), m (left), and p (bottom). A red triangle is formed by these nodes. Several black arrows radiate from the nodes, representing vector fields. A blue arrow labeled 'n' points upwards from node n. A green arrow labeled 'm' points to the left from node m. A yellow arrow labeled 'p' points downwards from node p. A dashed line connects nodes m and p.

$$\tilde{\mathbf{C}} \mathbf{M}_v^{\text{FE}} \mathbf{C} \hat{\mathbf{a}} + \mathbf{M}_\sigma^{\text{FE}} \frac{d\hat{\mathbf{a}}}{dt} = \hat{\mathbf{j}}_s$$

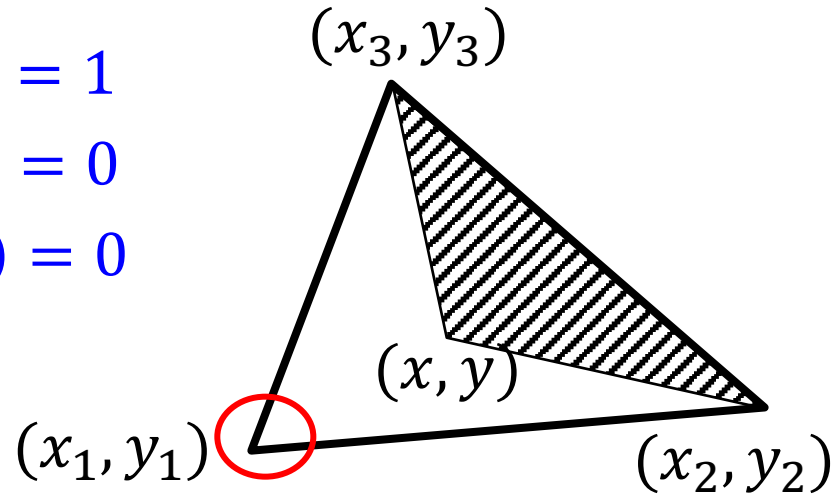
2D Nodal Shape Functions



$$\begin{cases} N_i(x_i, y_i) = 1 \\ N_i(x_j, y_j) = 0 \\ N_i(x_k, y_k) = 0 \end{cases}$$

$$N_i(x, y) = \frac{a_i + b_i x + c_i y}{2A_{ijk}}$$

$$\begin{cases} a_i = x_j y_k - x_k y_j \\ b_i = y_j - y_k \\ c_i = x_k - x_j \\ A_{ijk} = \text{element area} \end{cases}$$



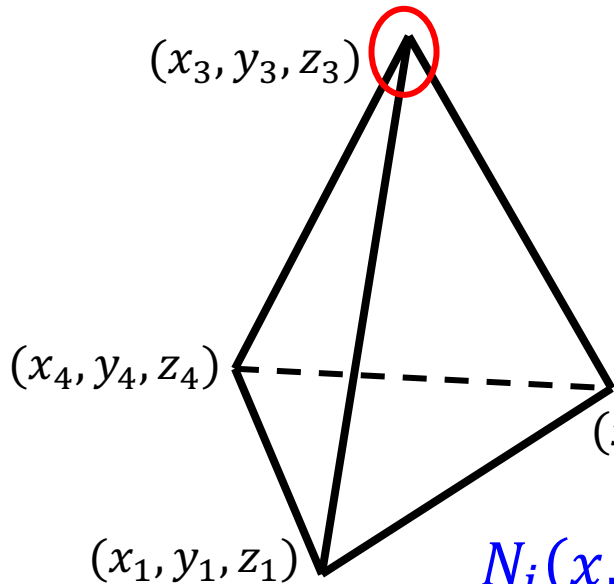
surface coordinates

$$N_i(x, y) = \frac{A_{pjk}}{A_{ijk}}$$

partition of unity

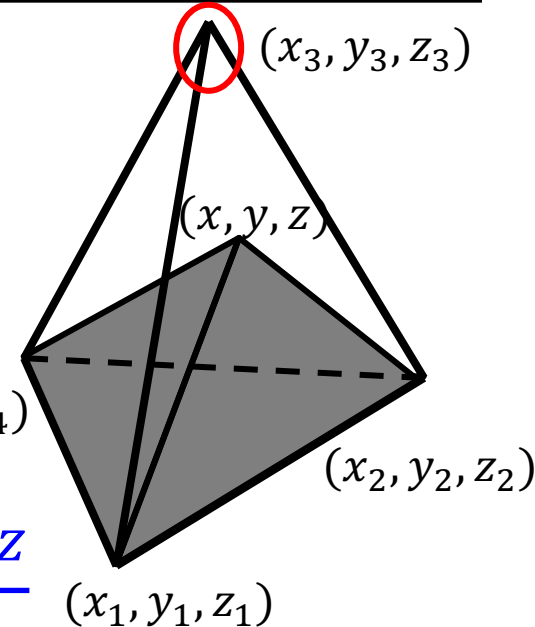
$$\sum_{i=1}^3 N_i(x, y) = 1, \quad \forall (x, y)$$

3D Nodal Shape Functions



$$\begin{cases} N_i(x_i, y_i, z_i) = 1 \\ N_i(x_j, y_j, z_j) = 0 \\ N_i(x_k, y_k, z_k) = 0 \\ N_i(x_\ell, y_\ell, z_\ell) = 0 \end{cases}$$

$$N_i(x, y, z) = \frac{a_i + b_i x + c_i y + d_i z}{6V_{ijkl}}$$



$$\begin{cases} a_i = \\ b_i = \\ c_i = \\ d_i = \\ V_{ijkl} = \text{element volume} \end{cases}$$

volume coordinates $N_i(x, y, z) = \frac{V_{pjkl}}{V_{ijkl}}$

partition of unity $\sum_{i=1}^4 N_i(x, y, z) = 1, \forall (x, y, z)$

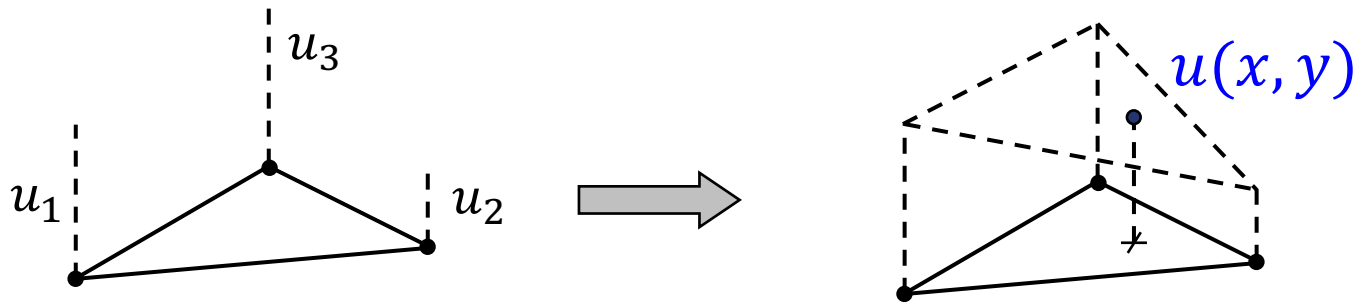
Dofs or Nodal Values ?



• **solution** $\varphi(x, y, z) = \varphi_{\text{dir}}(x, y, z) + \sum_j u_j N_j(x, y, z)$

–interpolation :

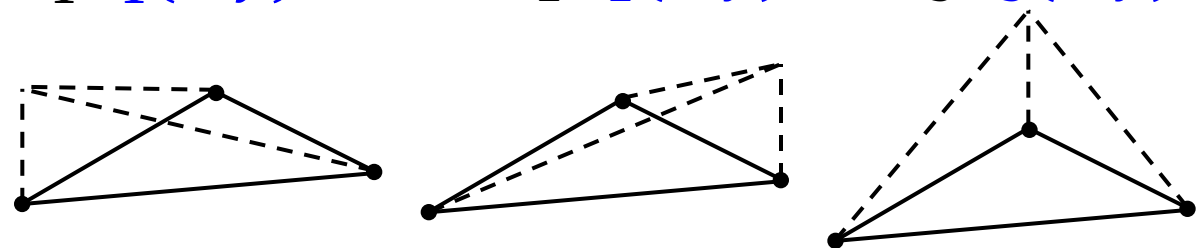
nodal values & interpolating functions



–series development :

degrees of freedom & shape functions

$$u_1, u_2, u_3 \quad \longrightarrow \quad u_1 N_1(x, y) \quad + u_2 N_2(x, y) \quad + u_3 N_3(x, y)$$



Vectorial Shape Functions



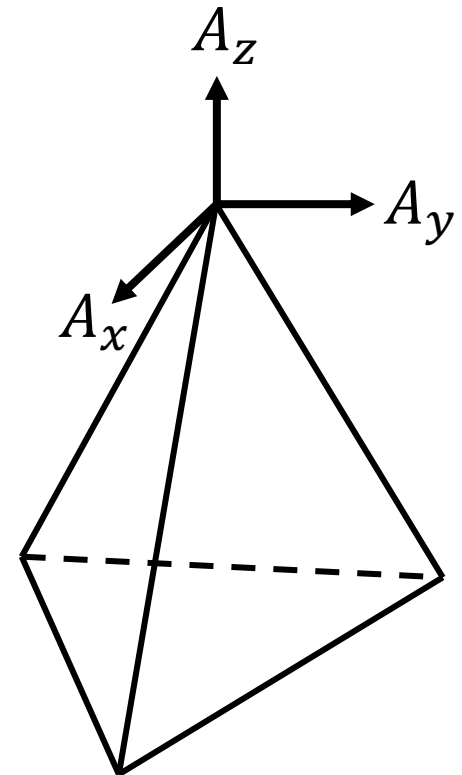
- **vector field** $\vec{A} = A_x \vec{e}_x + A_y \vec{e}_y + A_z \vec{e}_z$

– interpolate each component separately by *scalar* shape functions

$$\vec{A} = \left(\sum_j u_{x,j} N_j \right) \vec{e}_x + \left(\sum_j u_{y,j} N_j \right) \vec{e}_y + \left(\sum_j u_{z,j} N_j \right) \vec{e}_z$$

– BUT

- too much continuity (both normal and tangential components)
- spurious modes



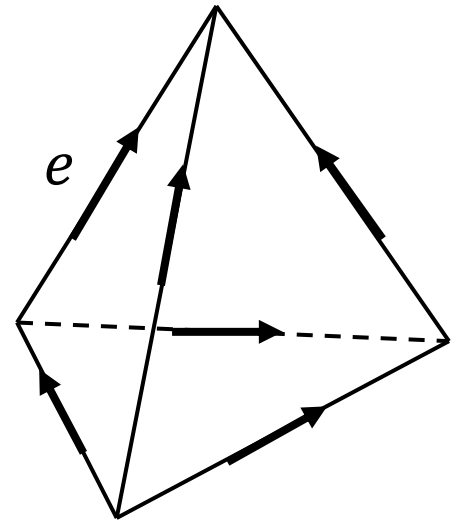
Edge Elements



- shape functions defined by

$$\begin{cases} \int_e \vec{w}_e \cdot d\vec{s} = 1 \\ \int_{e'} \vec{w}_e \cdot d\vec{s} = 0, e \neq e' \end{cases}$$

–linear combination $\vec{E}(\vec{x}) = \sum_{j=1}^E u_j \vec{w}_{e_j}(\vec{x})$



- features tangential continuity (1-form)
- physical meaning of the degrees of freedom

$$\int_e \vec{E} \cdot d\vec{s} = \sum_{j=1}^E u_j \int_e \vec{w}_{e_j} \cdot d\vec{s} = u_e \int_e \vec{w}_{e_j} \cdot d\vec{s} = u_e$$

- voltage drop along the edge

Face(t) Elements

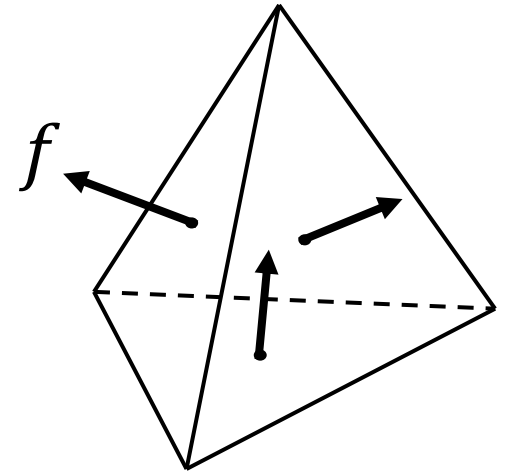


- shape functions defined by

$$\begin{cases} \int_f \vec{w}_f \cdot d\vec{A} = 1 \\ \int_{f'} \vec{w}_f \cdot d\vec{A} = 0, f' \neq f \end{cases}$$

–linear combination

$$\vec{B}(\vec{x}) = \sum_{j=1}^F u_j \vec{w}_{f_j}(\vec{x})$$



- features normal continuity (2-form)

–physical meaning of the degrees of freedom

$$\int_f \vec{B} \cdot d\vec{A} = \sum_{j=1}^F u_j \int_f \vec{w}_{f_j} \cdot d\vec{A} = u_f \int_f \vec{w}_{f_j} \cdot d\vec{A} = u_f$$

- flux through the facet

Volume Elements



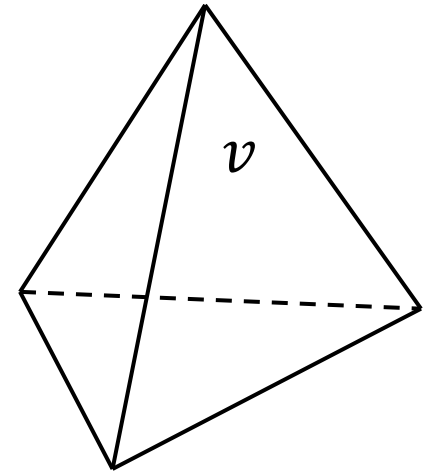
- shape functions defined by

$$\int_v w_v dV = 1$$

$$\int_{v'} w_v dV = 0, v' \neq v$$

–linear combination

$$q(\vec{x}) = \sum_{j=1}^V u_j w_{v_j}(\vec{x})$$



- commonly discontinuous (3-form)
- physical meaning of the degrees of freedom

$$\int_v \rho dV = \sum_{j=1}^V u_j \int_v w_{v_j} dV = u_v \int_v w_{v_j} dV = u_v$$

- charge within the element

Canonical Construction



- **construction**

- nodal shape functions

$$N_m, N_n, N_p, N_q$$

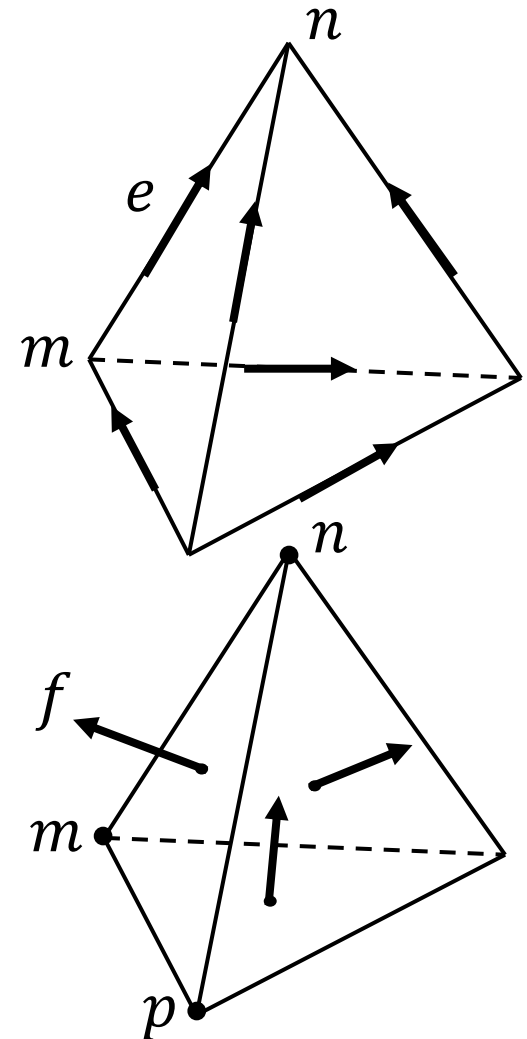
- edge elements

$$\vec{w}_e(\vec{x}) = N_m \nabla N_n - N_n \nabla N_m$$

- facet elements

$$\begin{aligned} \vec{w}_f(\vec{x}) = & 2(N_m \nabla N_n \times \nabla N_p \\ & + N_n \nabla N_p \times \nabla N_m \\ & + N_p \nabla N_m \times \nabla N_n) \end{aligned}$$

- volume elements $w_v(\vec{x}) = \frac{1}{V}$



Whitney Complex

- Whitney elements (for a given mesh)

- W^0 : nodal elements
- W^1 : edge elements
- W^2 : facet elements
- W^3 : volume elements

- Whitney complex

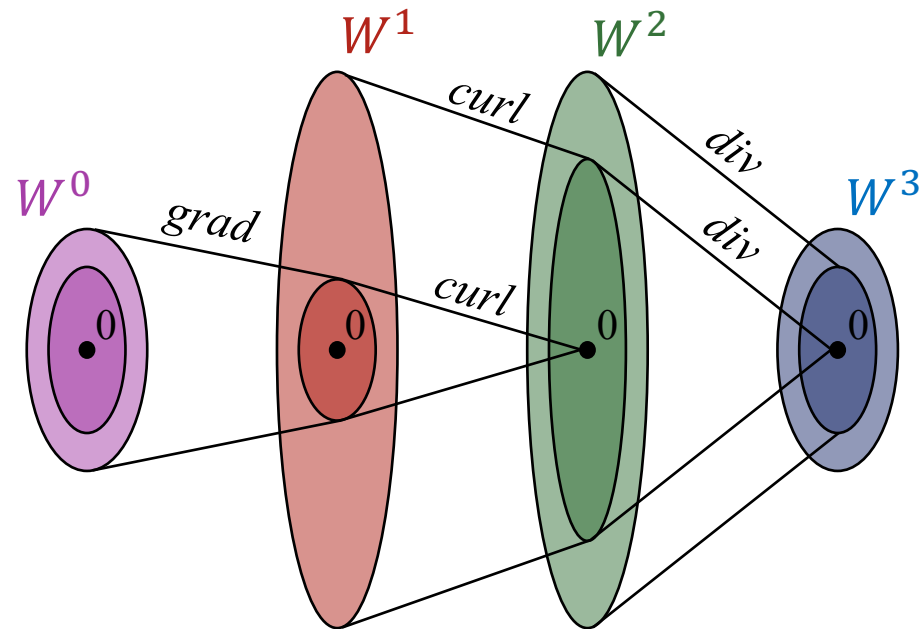
$$\text{grad } W^0 \subset W^1$$

$$\text{curl } W^1 \subset W^2$$

$$\text{div } W^2 \subset W^3$$

$$\text{curl grad } W^0 = 0$$

$$\text{div curl } W^1 = 0$$



2D Higher-Order Elements

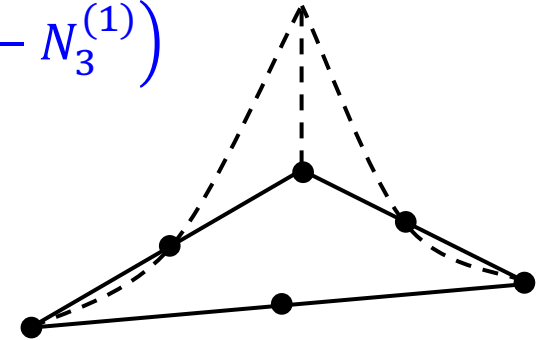
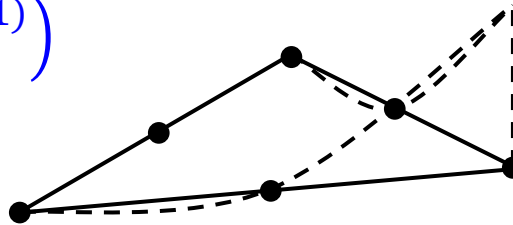
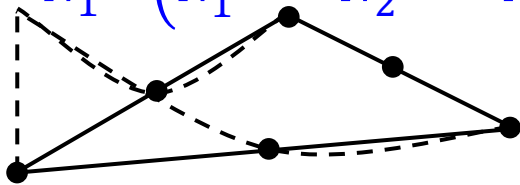


- nodal elements

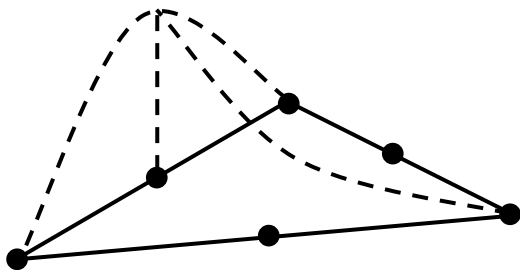
$$N_3^{(2)} = N_3^{(1)} \left(-N_1^{(1)} - N_2^{(1)} + N_3^{(1)} \right)$$

$$N_2^{(2)} = N_2^{(1)} \left(-N_1^{(1)} + N_2^{(1)} - N_3^{(1)} \right)$$

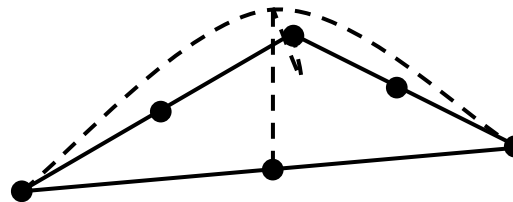
$$N_1^{(2)} = N_1^{(1)} \left(N_1^{(1)} - N_2^{(1)} - N_3^{(1)} \right)$$



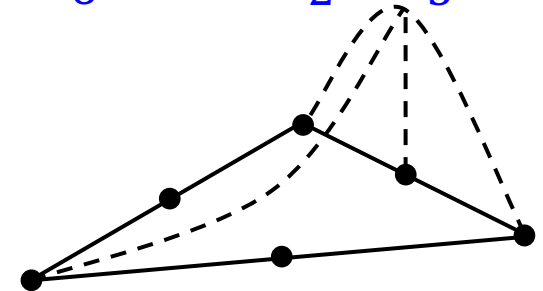
$$N_4^{(2)} = 4N_3^{(1)} N_1^{(1)}$$



$$N_5^{(2)} = 4N_1^{(1)} N_2^{(1)}$$



$$N_6^{(2)} = 4N_2^{(1)} N_3^{(1)}$$



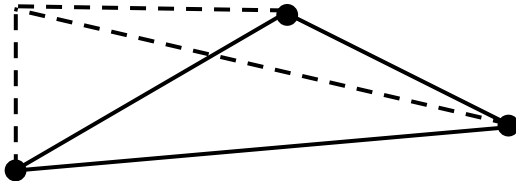
degrees of freedom == nodal values

Hierarchical Finite Elements

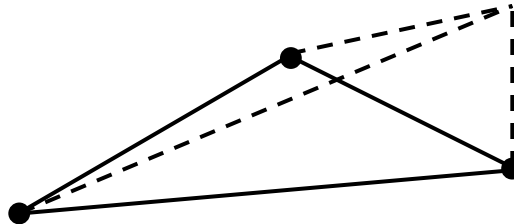


• nodal elements

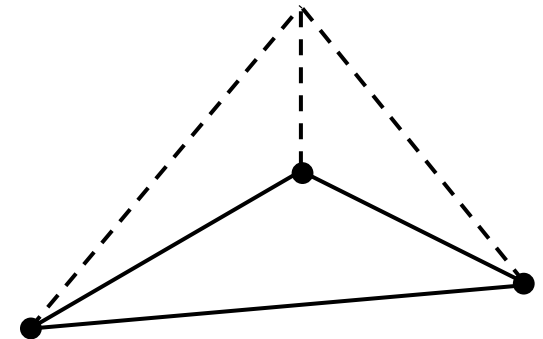
$$N_1^{(2)} = N_1^{(1)}$$



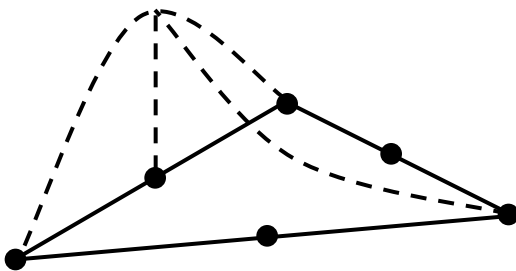
$$N_2^{(2)} = N_2^{(1)}$$



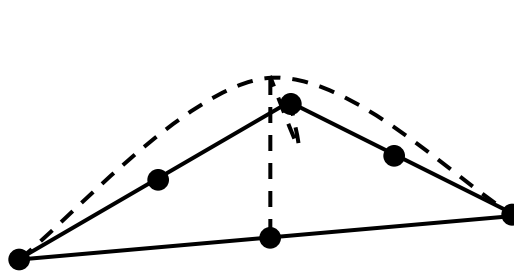
$$N_3^{(2)} = N_3^{(1)}$$



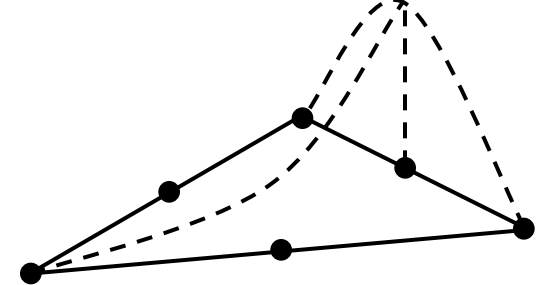
$$N_4^{(2)} = 4N_3^{(1)} N_1^{(1)}$$



$$N_5^{(2)} = 4N_1^{(1)} N_2^{(1)}$$



$$N_6^{(2)} = 4N_2^{(1)} N_3^{(1)}$$



degrees of freedom \neq nodal values

Hierarchical Finite Elements

- first-order system of equations

$$\left[K^{(11)} \right] \left[u^{(1)} \right] = \left[f^{(1)} \right]$$

- second-order system of equations

$$\begin{bmatrix} K^{(11)} & K^{(12)} \\ K^{(21)} & K^{(22)} \end{bmatrix} \begin{bmatrix} u^{(1)} \\ u^{(2)} \end{bmatrix} = \begin{bmatrix} f^{(1)} \\ f^{(2)} \end{bmatrix}$$



first-order system embedded in second-order system

Overview



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- discretisation in space
- **boundary and symmetry conditions**
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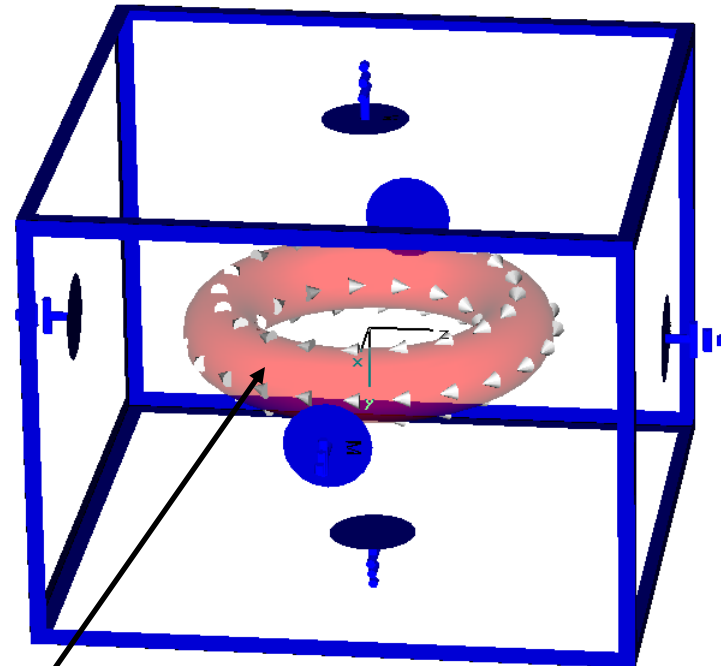
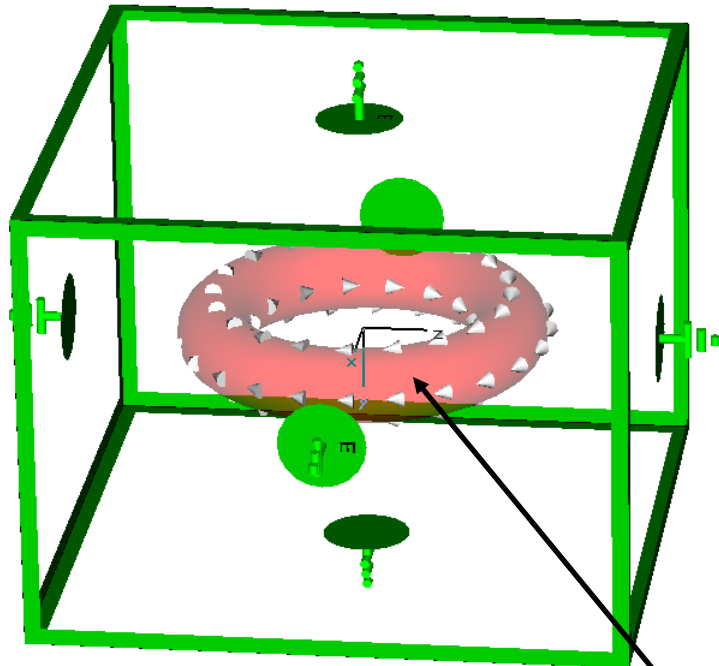
Boundary Conditions (1)

	electric BC „flux wall“ „current gate“ 	magnetic BC „flux gate“ „current wall“ 
definition	$\vec{E}_t = 0$	$\vec{H}_t = 0$
electric current	$\vec{J}_n \neq 0$	$\vec{J}_n = 0$
magnetic flux	$\vec{B}_n = 0$	$\vec{B}_n \neq 0$
magnetic vector potential formulation	Dirichlet BC	Neumann BC
magnetic scalar potential formulation	Neumann BC	Dirichlet BC

Boundary Conditions (2)

electric boundary conditions

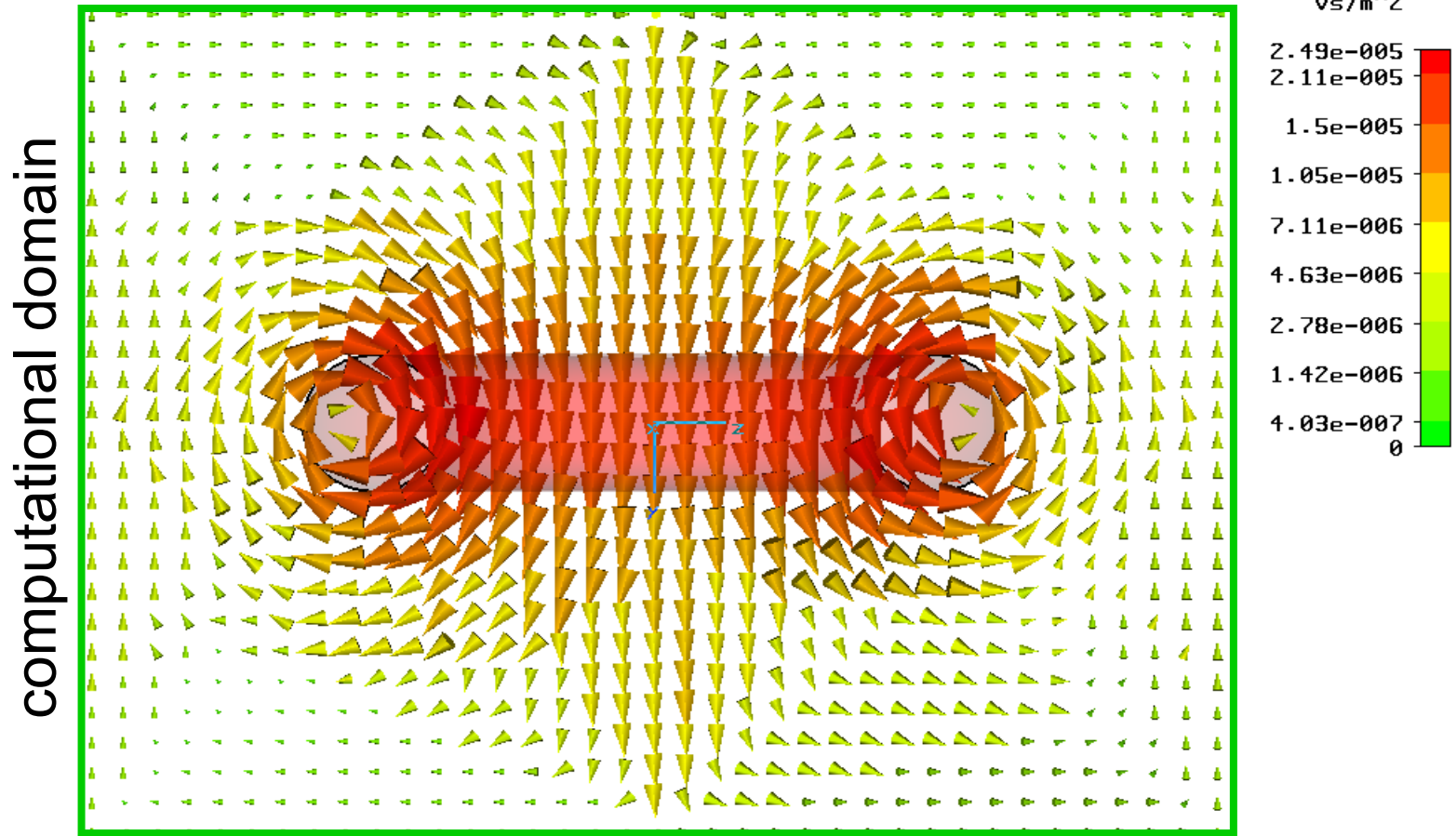
magnetic boundary conditions



coil

Boundary Conditions (3)

electric boundary conditions



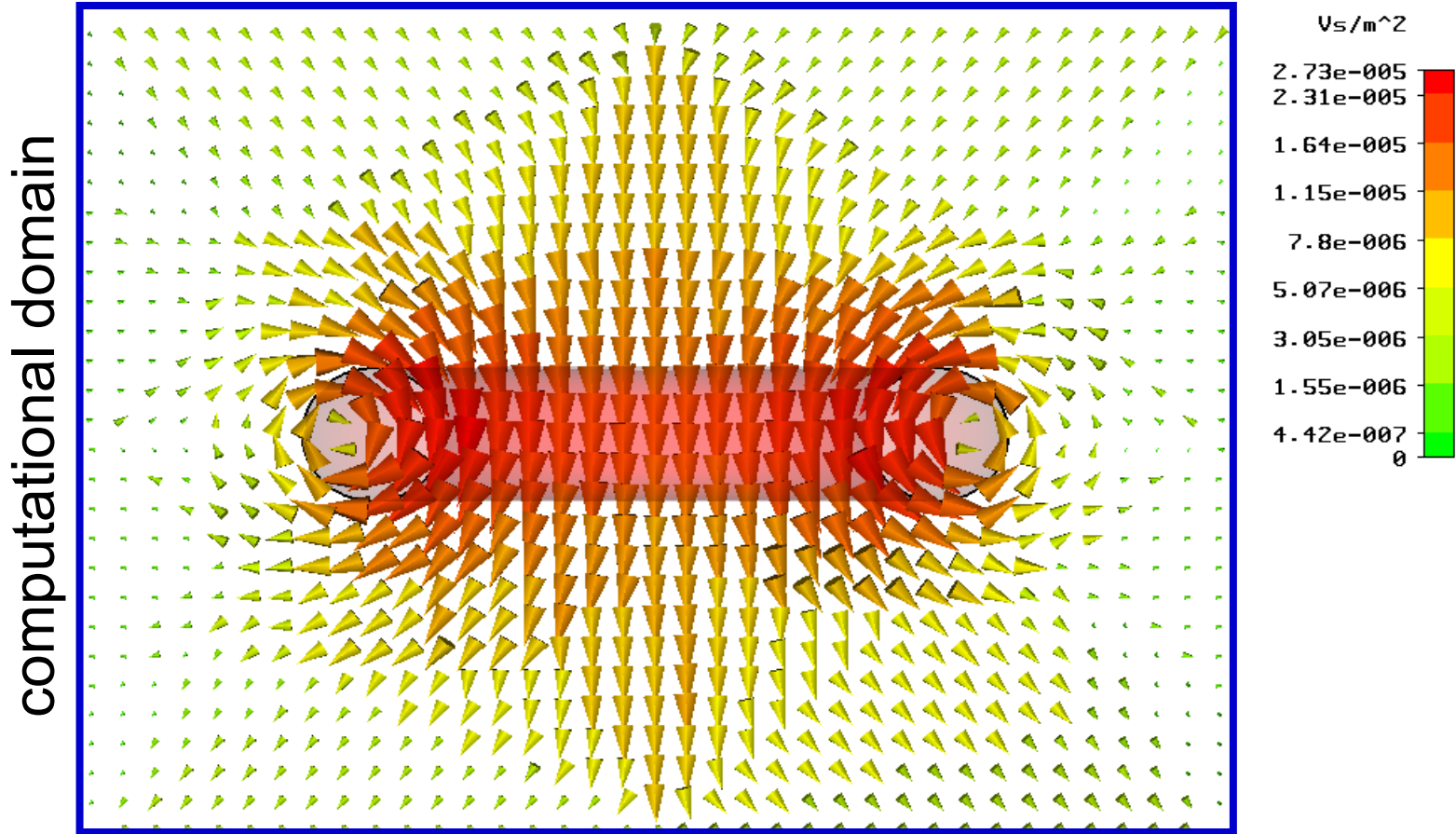
Type = B-Field

Plane at x = 0

Maximum-2d = 2.49377e-005 Vs/m² at 2.36848e-015 / 0.222222 / -2.66667

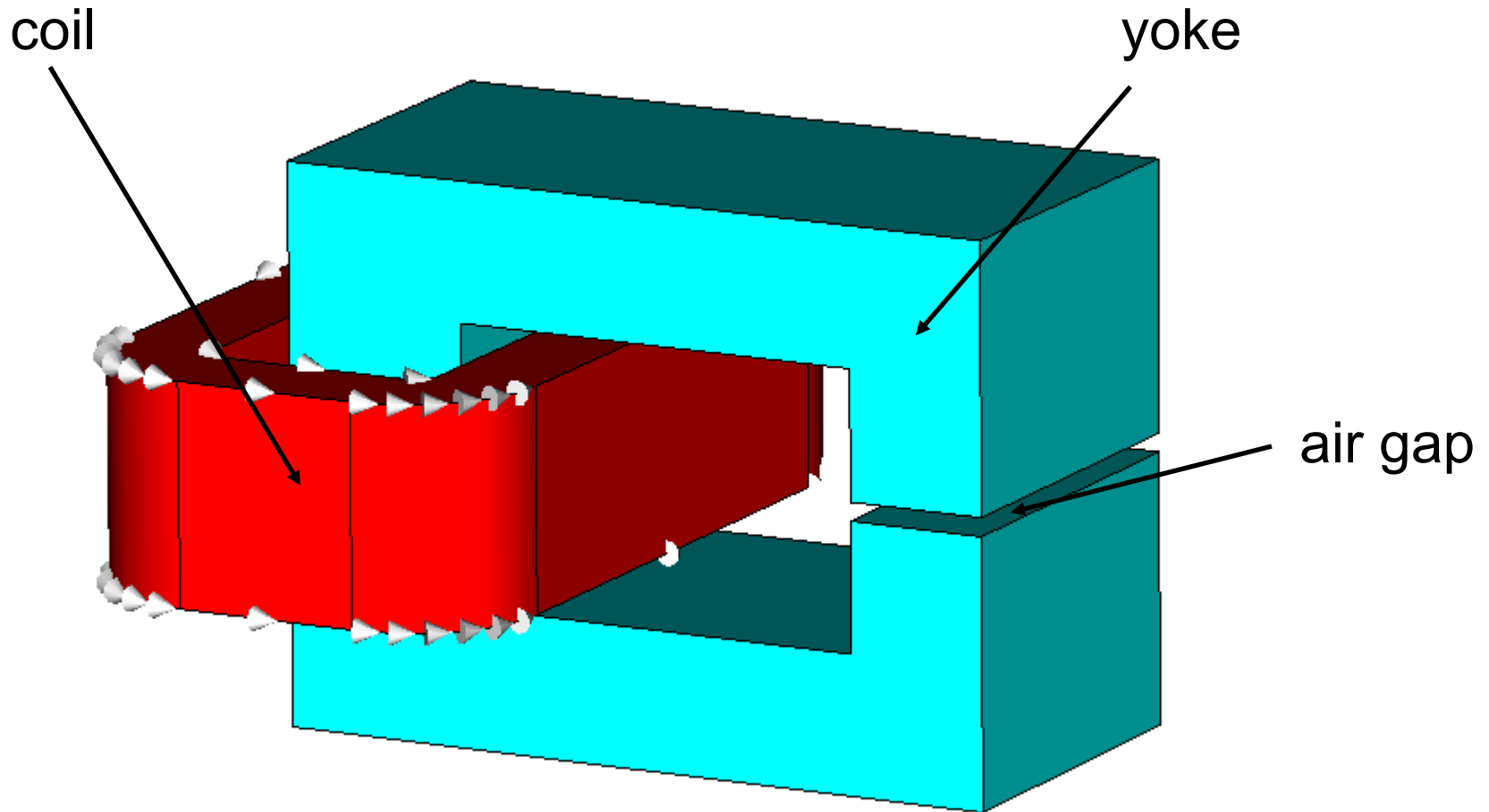
Boundary Conditions (4)

magnetic boundary conditions





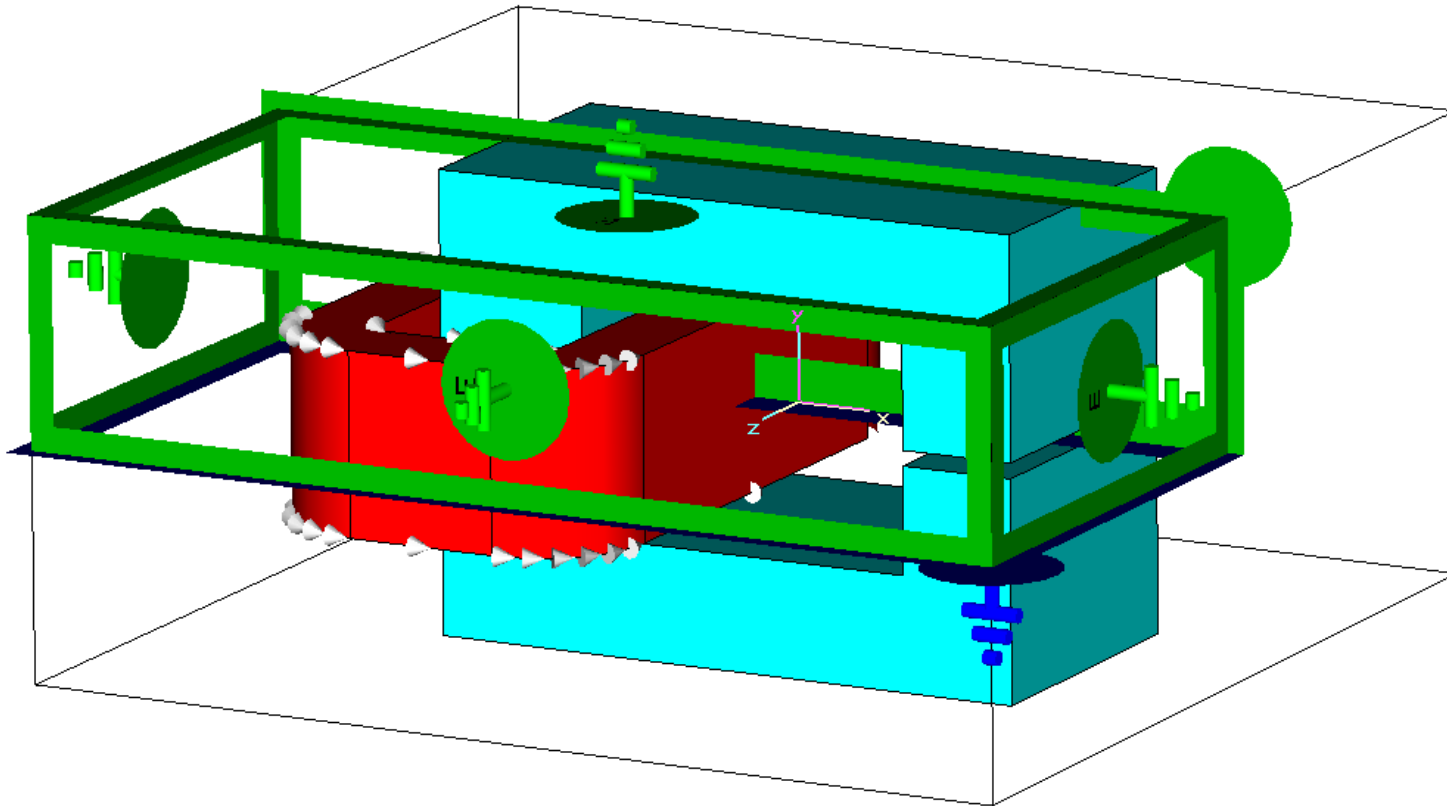
Type = B-Field
Plane at x = 0
Maximum-2d = 2.73322e-005 Vs/m² at 2.36848e-015 / 0.222222 / -2.66667

Symmetries (1)

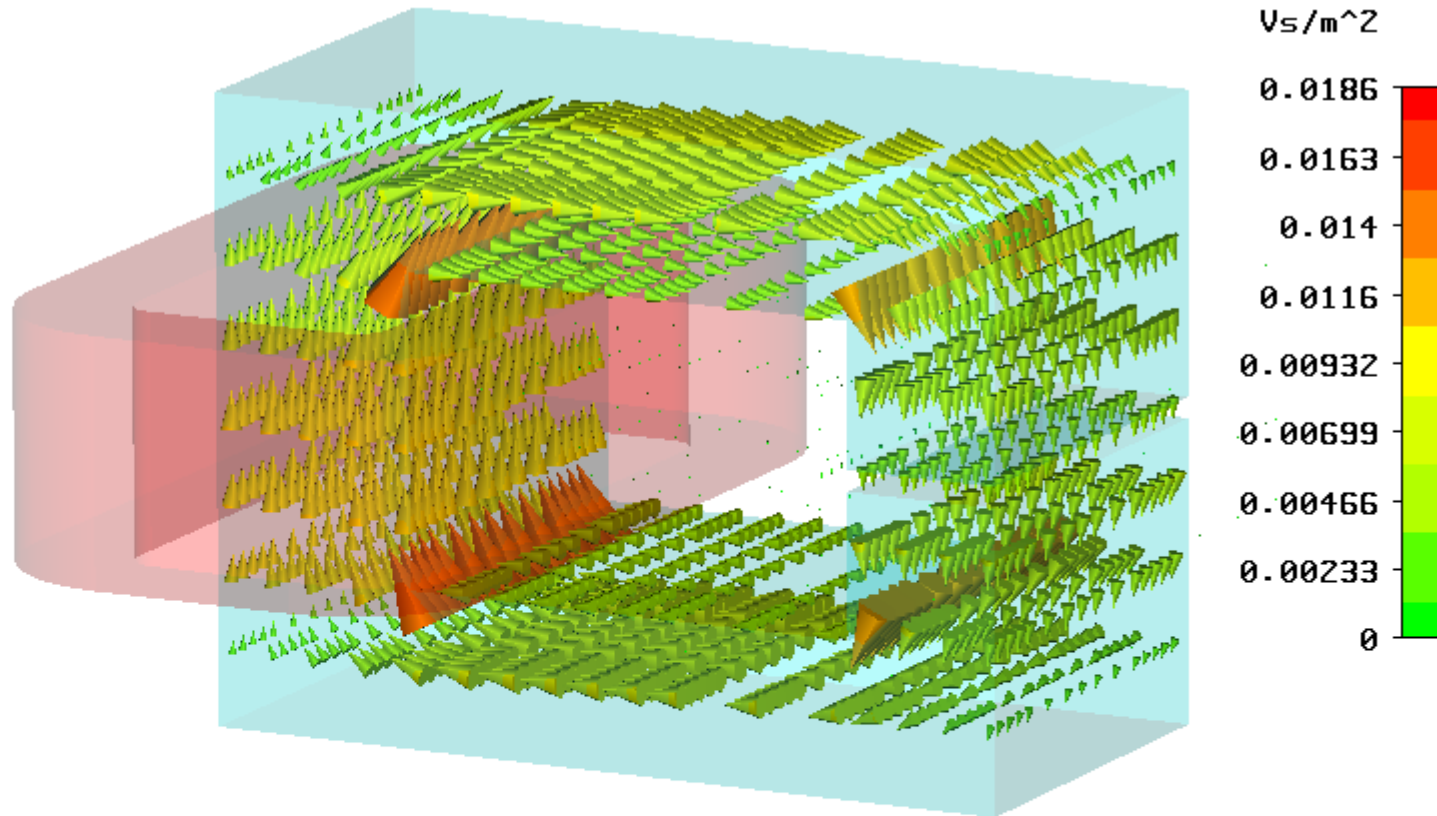


Symmetries (2)

-  electric boundary conditions
-  magnetic boundary conditions



Symmetries (3)



Inserting Boundary Conditions



$$\begin{bmatrix} \mathbf{K}_{bb} & \mathbf{K}_{bc} \\ \mathbf{K}_{cb} & \mathbf{K}_{cc} \end{bmatrix} \begin{bmatrix} \mathbf{u}_b \\ \mathbf{u}_c \end{bmatrix} = \begin{bmatrix} \mathbf{f}_b \\ \mathbf{f}_c \end{bmatrix}$$

potentials living at Dirichlet boundaries

insert essential
boundary conditions



$$\begin{bmatrix} \mathbf{K}_{bb} & \mathbf{K}_{bc} & 0 \\ \mathbf{K}_{cb} & \mathbf{K}_{cc} & \mathbf{B}_{cq} \\ 0 & \mathbf{B}_{qc} & 0 \end{bmatrix} \begin{bmatrix} \mathbf{u}_b \\ \mathbf{u}_c \\ \mathbf{y}_q \end{bmatrix} = \begin{bmatrix} \mathbf{f}_b \\ \mathbf{f}_c \\ 0 \end{bmatrix}$$

eliminate known
potentials



$$\mathbf{K}_{bb} \mathbf{u}_b = \mathbf{f}_b - \mathbf{K}_{bc} \mathbf{u}_c$$

$$\mathbf{K} \quad \mathbf{f} \quad \xrightarrow{\text{bdrycond_shrink.m}} \quad \mathbf{K}_{bb} \quad \mathbf{f}_b - \mathbf{K}_{bc} \mathbf{u}_c$$

backslash

$$\mathbf{u} \quad \leftarrow \quad \text{bdrycond_inflate.m} \quad \leftarrow \quad \mathbf{u}_b$$

Overview



- magnetoquasistatic formulation
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Reduction to 2D Models (1)



▪ 2D cartesian models

$$\vec{J} = (0, 0, J_z(x, y))$$

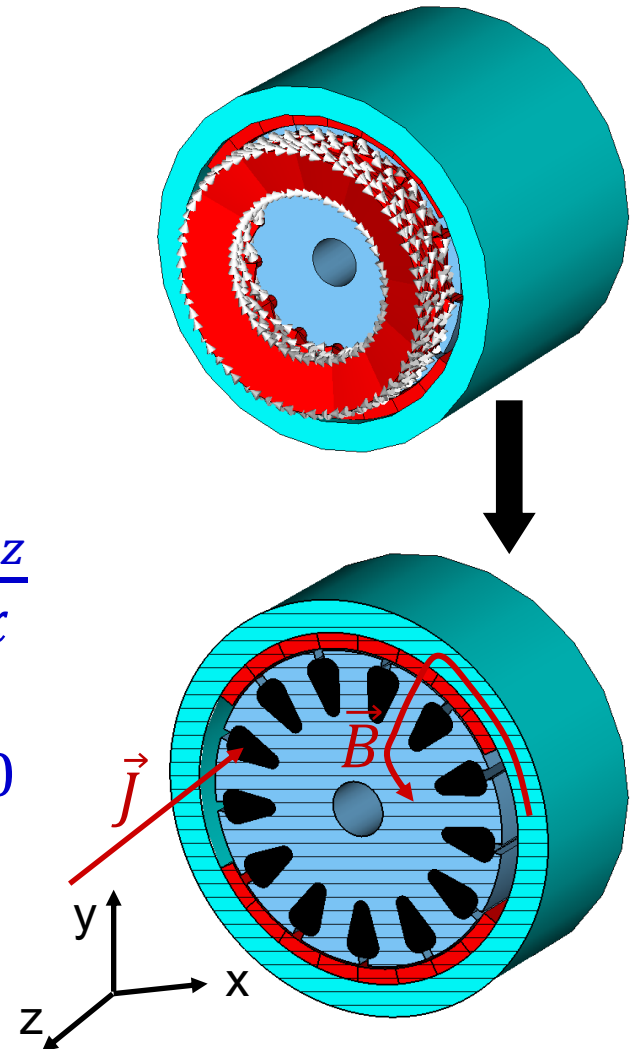
$$\vec{B} = (B_x(x, y), B_y(x, y), 0)$$

$$\vec{A} = (0, 0, A_z(x, y))$$

$$\vec{B} = \nabla \times \vec{A} \iff B_x = \frac{\partial A_z}{\partial y}; \quad B_y = -\frac{\partial A_z}{\partial x}$$

$$\nabla \cdot \vec{B} = \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} = \frac{\partial^2 A_z}{\partial x \partial y} - \frac{\partial^2 A_z}{\partial x \partial y} = 0$$

$$\frac{\partial A_z}{\partial z}, \frac{\partial B_x}{\partial z}, \frac{\partial B_y}{\partial z}, \frac{\partial J_z}{\partial z} = 0 \quad \text{but} \quad \frac{\partial \varphi}{\partial z} \neq 0$$



Reduction to 2D Models (2)



2D cartesian models

$$B_x = \frac{\partial A_z}{\partial y}; \quad B_y = -\frac{\partial A_z}{\partial x}$$

$$H_x = \nu_x \frac{\partial A_z}{\partial y}; \quad H_y = -\nu_y \frac{\partial A_z}{\partial x}$$

$$J_z = \frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} = J_{s,z} - \sigma \frac{\partial A_z}{\partial t}$$

anisotropic material

$$\begin{bmatrix} H_x \\ H_y \end{bmatrix} = \begin{bmatrix} \nu_x & 0 \\ 0 & \nu_y \end{bmatrix} \begin{bmatrix} B_x \\ B_y \end{bmatrix}$$

Ampère
+ Faraday-Lenz

$$-\frac{\partial}{\partial x} \left(\nu_y \frac{\partial A_z}{\partial x} \right) - \frac{\partial}{\partial y} \left(\nu_x \frac{\partial A_z}{\partial y} \right) + \sigma \frac{\partial A_z}{\partial t} = J_{s,z}$$

$$\longleftrightarrow \quad \nabla \times (\nu \nabla \times \vec{A}) + \sigma \frac{\partial \vec{A}}{\partial t} = \vec{J}_s$$

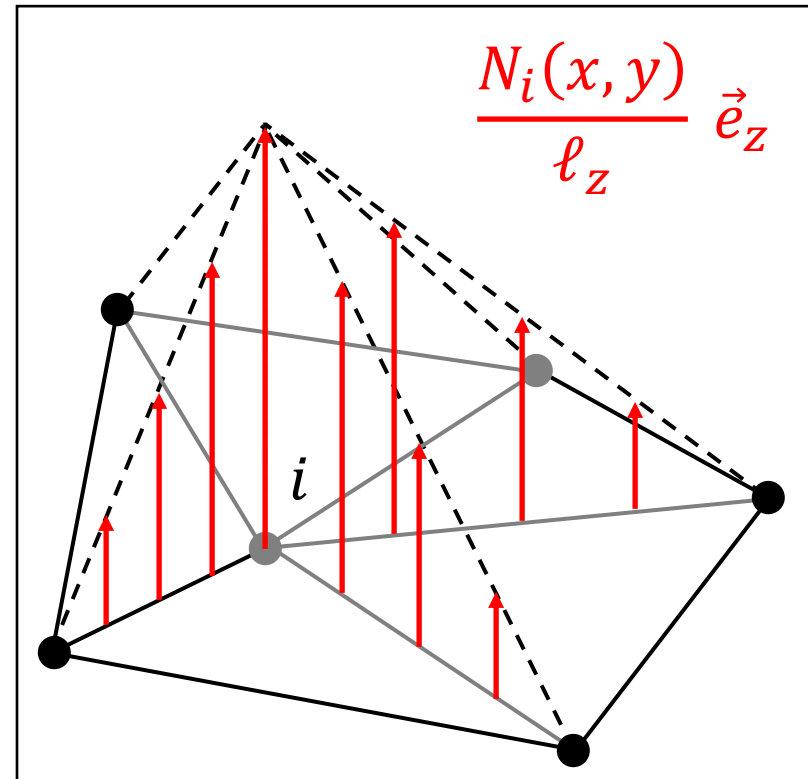
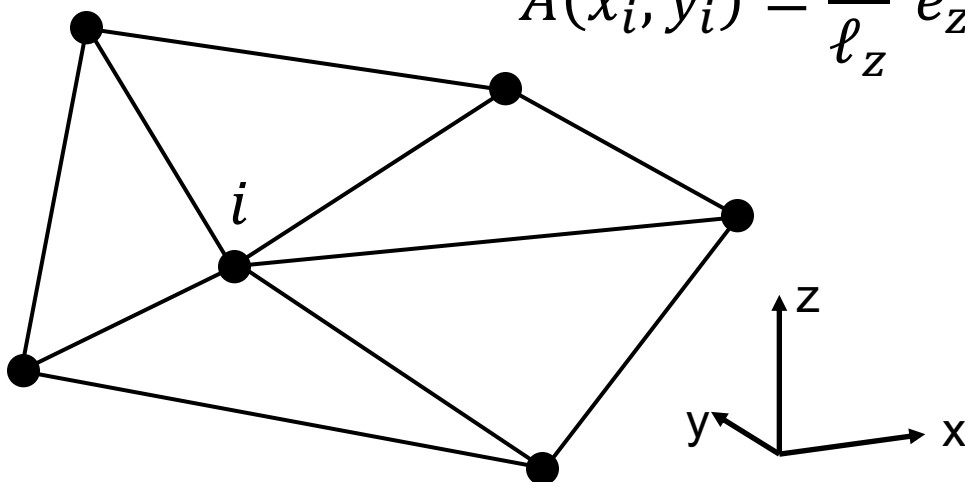
2D Discretisation (1)



$$\sum_j \left(u_j \int_V v \nabla \times \vec{w}_j \cdot \nabla \times \vec{w}_i \, dV + \frac{du_j}{dt} \int_V \sigma \vec{w}_j \cdot \vec{w}_i \, dV \right) = \int_V \vec{J}_s \cdot \vec{w}_i \, dV$$

$$\vec{A} = \sum_j u_j \vec{w}_j = \sum_j u_j \frac{N_j(x, y)}{\ell_z} \vec{e}_z$$

$$\vec{A}(x_i, y_i) = \frac{u_i}{\ell_z} \vec{e}_z$$



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Coil Model (1)



assumptions

- homogeneous current distribution
- no eddy currents

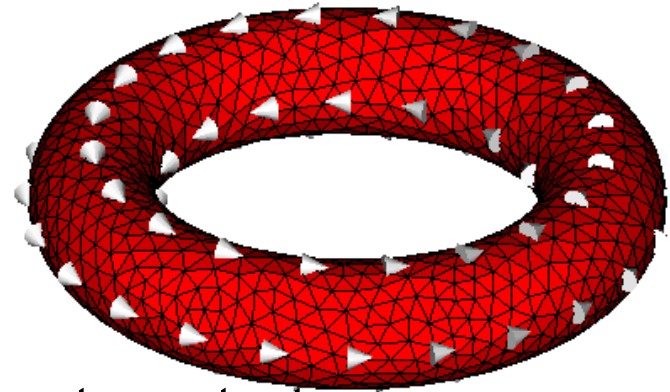
notice (*model*)

- there will be an induced voltage !!
- current density not ct when cross-section not constant

winding function $\vec{\chi}_{\text{coil},q}(x, y, z)$ [1/m²]

- computed geometrically
- by field solution (lecture V11)

$$\vec{J}_S(x, y, z, t) = \sum_{q=1}^{n_{\text{coil}}} \underbrace{\vec{\chi}_{\text{coil},q}(x, y, z)}_{\vec{J}_{\text{coil},q}(x, y, z, t)} i_q(t)$$



current_Pstr.m

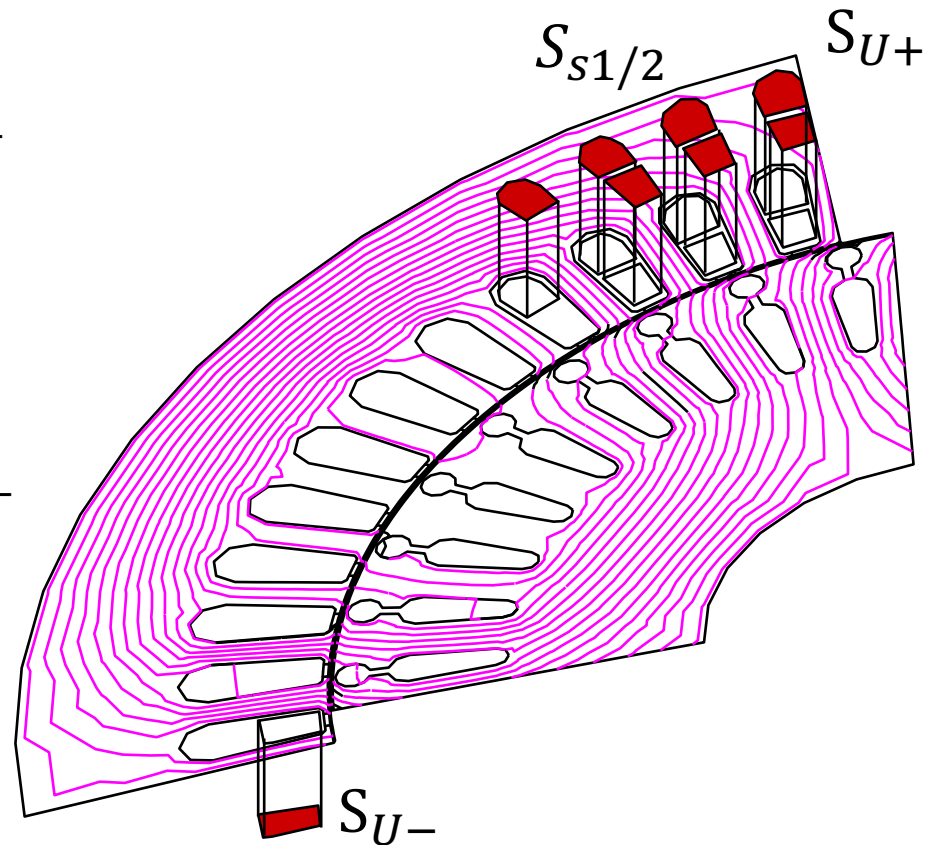


Coil Model (2)



▪ in 2D: $\vec{J}(x, y, t) = (0, 0, J_z(x, y, t))$

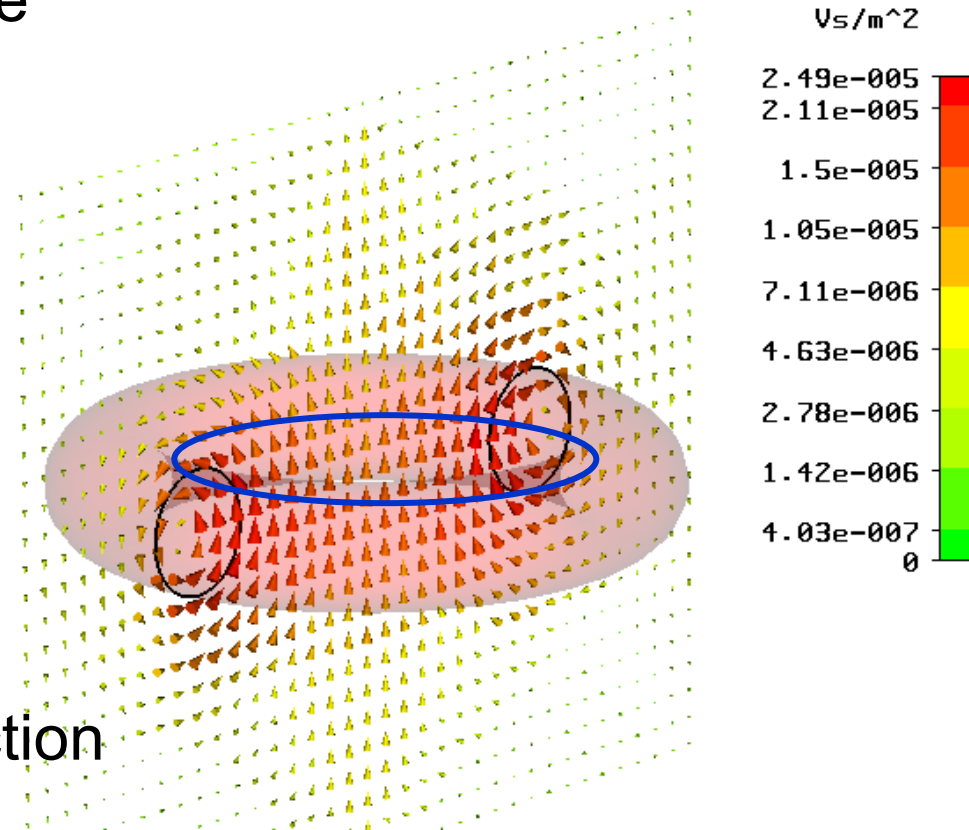
$$\left\{ \begin{array}{l} \vec{\chi}_{\text{coil},U}(x, y) = +\frac{N_t}{S_{s1/2}} \vec{e}_z \quad \text{in } S_{U+} \\ \vec{\chi}_{\text{coil},U}(x, y) = -\frac{N_t}{S_{s1/2}} \vec{e}_z \quad \text{in } S_{U-} \\ \vec{\chi}_{\text{coil},U}(x, y) = 0 \quad \text{in } S_{2D} \setminus S_{U+} \setminus S_{U-} \end{array} \right.$$



Coil Model (3)

- induced voltage \sim flux linkage
- which flux is linked?
- for a single path
$$\phi = \oint_{\partial S} \vec{A} \cdot d\vec{s}$$
- for a coil
- integrating along the coil
- average at the coil cross-section

$$\psi_{\text{coil},q}(t) = \int_V \vec{A}(x, y, z, t) \cdot \vec{\chi}_{\text{coil},q}(x, y, z) dV$$

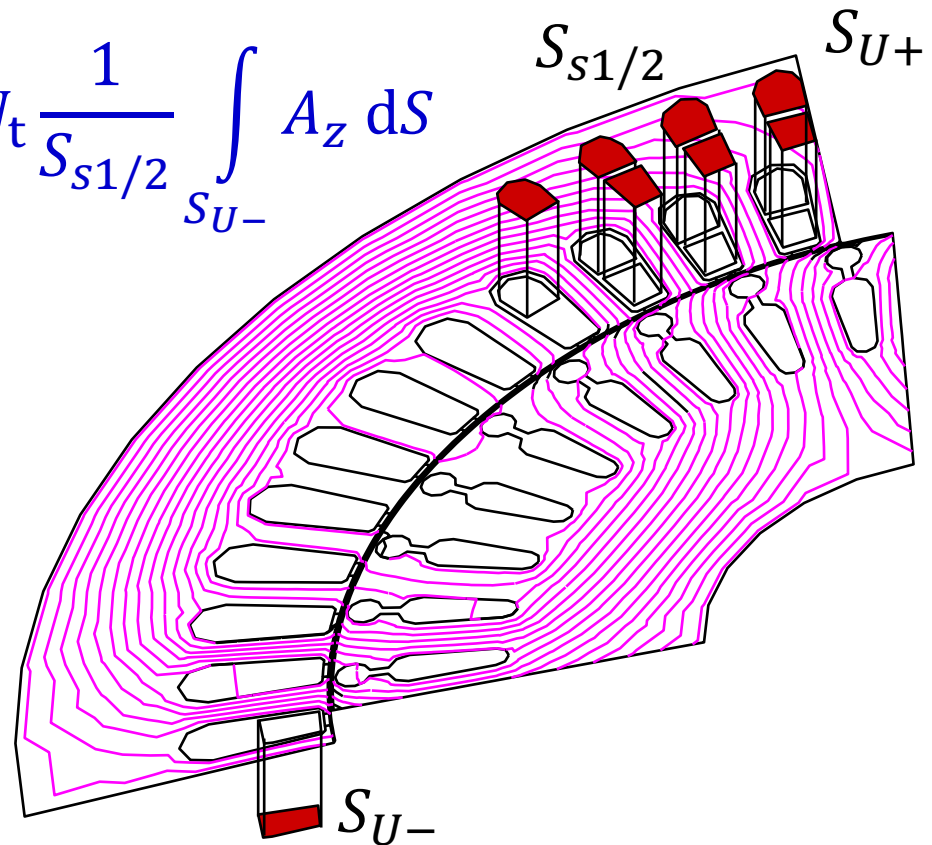


Coil Model (4)



▪ in 2D:
$$\psi_{\text{coil},q} = \int_V \vec{A} \cdot \vec{\chi}_{\text{coil},q} dV$$

$$\psi_{\text{coil},U} = N_t \frac{1}{S_{S1/2}} \int_{S_{U+}} A_z dS - N_t \frac{1}{S_{S1/2}} \int_{S_{U-}} A_z dS$$



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Permanent Magnet (1)



scalar and linear permanent-magnet model

$$B = B_r + \mu H$$

$$H = H_m + \nu B$$

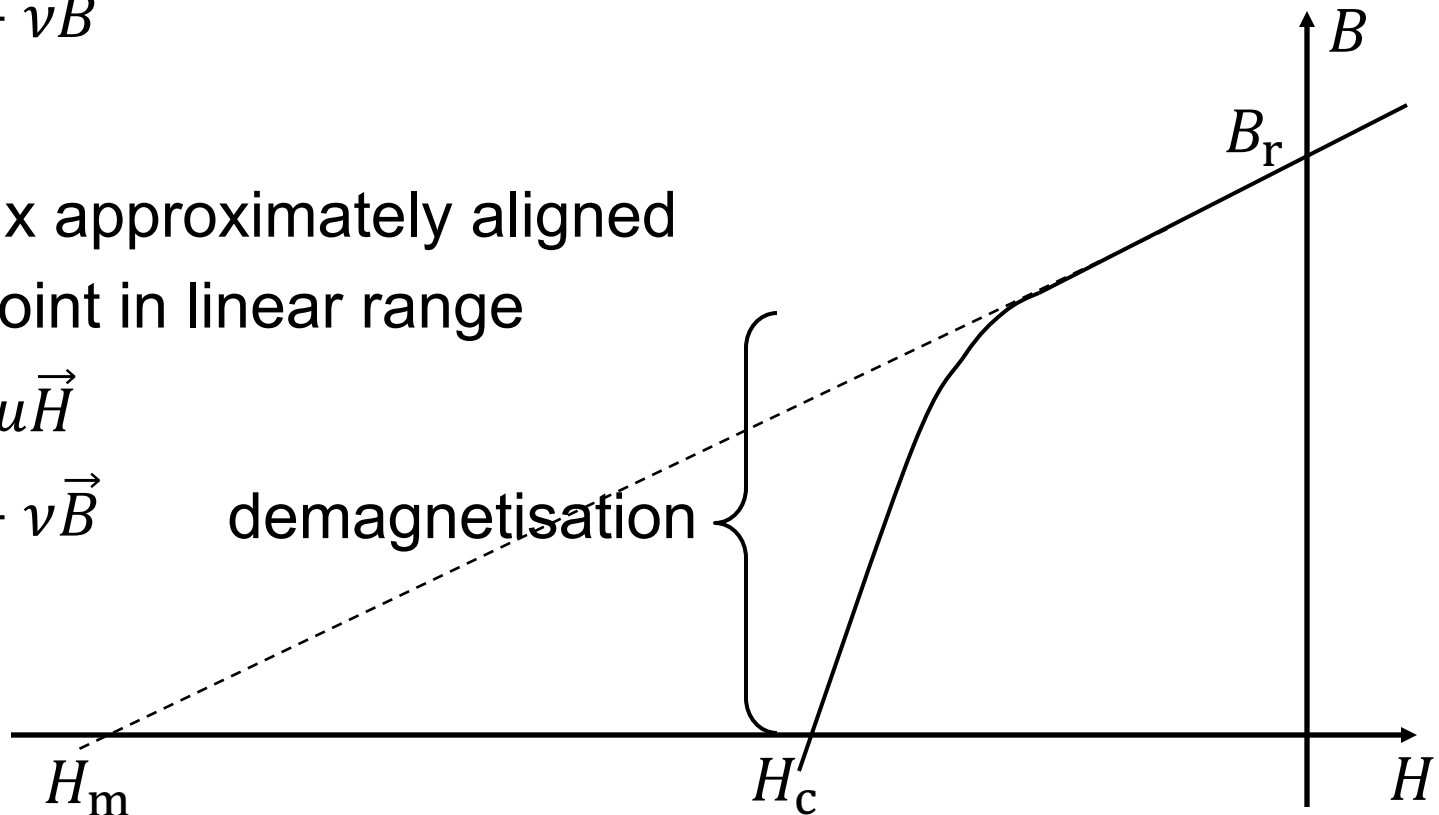
assumptions

- resulting flux approximately aligned
- operation point in linear range

$$\vec{B} = \vec{B}_r + \mu \vec{H}$$

$$\vec{H} = \vec{H}_m + \nu \vec{B}$$

demagnetisation



Permanent Magnet (2)



magnetoquasistatic formulation:

$$\vec{H} = \vec{H}_m + \nu \vec{B}$$

$$\vec{H} = \vec{H}_m + \nu \nabla \times \vec{A}$$

$$\nabla \times (\nu \nabla \times \vec{A}) = \vec{J} - \underbrace{\nabla \times \vec{H}_m}_{\vec{J}_m}$$

\vec{J}_m magnetisation current

in 2D:

$$-\frac{\partial}{\partial x} \left(\nu_y \frac{\partial A_z}{\partial x} \right) - \frac{\partial}{\partial y} \left(\nu_x \frac{\partial A_z}{\partial y} \right) = J_z - \frac{\partial H_{m,y}}{\partial x} + \frac{\partial H_{m,x}}{\partial y}$$

Permanent Magnet (3)



discretisation:

$$\text{RHS} = \int_V \vec{J}_m \cdot \vec{w}_i \, dV$$

$$\text{RHS} = - \int_V \nabla \times \vec{H}_m \cdot \vec{w}_i \, dV$$

$$\text{RHS} = \int_V \nabla \cdot (\vec{w}_i \times \vec{H}_m) \, dV - \int_V \vec{H}_m \cdot \nabla \times \vec{w}_i \, dV$$

$$\text{RHS} = \underbrace{\oint_{\partial V} \vec{w}_i \times \vec{H}_m \cdot d\vec{S}}_{=0 \text{ when no PMs at the model boundary}} - \int_V \vec{H}_m \cdot \nabla \times \vec{w}_i \, dV$$

=0 when no PMs at the model boundary

Permanent Magnet (4)



discretisation (in 2D):

$$\text{RHS} = - \int_V \vec{H}_m \cdot \nabla \times \vec{w}_i \, dV$$



$$\vec{w}_i = \frac{N_i(x, y)}{\ell_z} \vec{e}_z$$

$$\text{RHS} = \int_V \frac{1}{\ell_z} \left(\frac{\partial N_i}{\partial y} H_{mx} - \frac{\partial N_i}{\partial x} H_{my} \right) dV$$

$$\text{RHS} = \int_{\emptyset V} \left(\frac{\partial N_i}{\partial y} H_{mx} - \frac{\partial N_i}{\partial x} H_{my} \right) dS$$

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