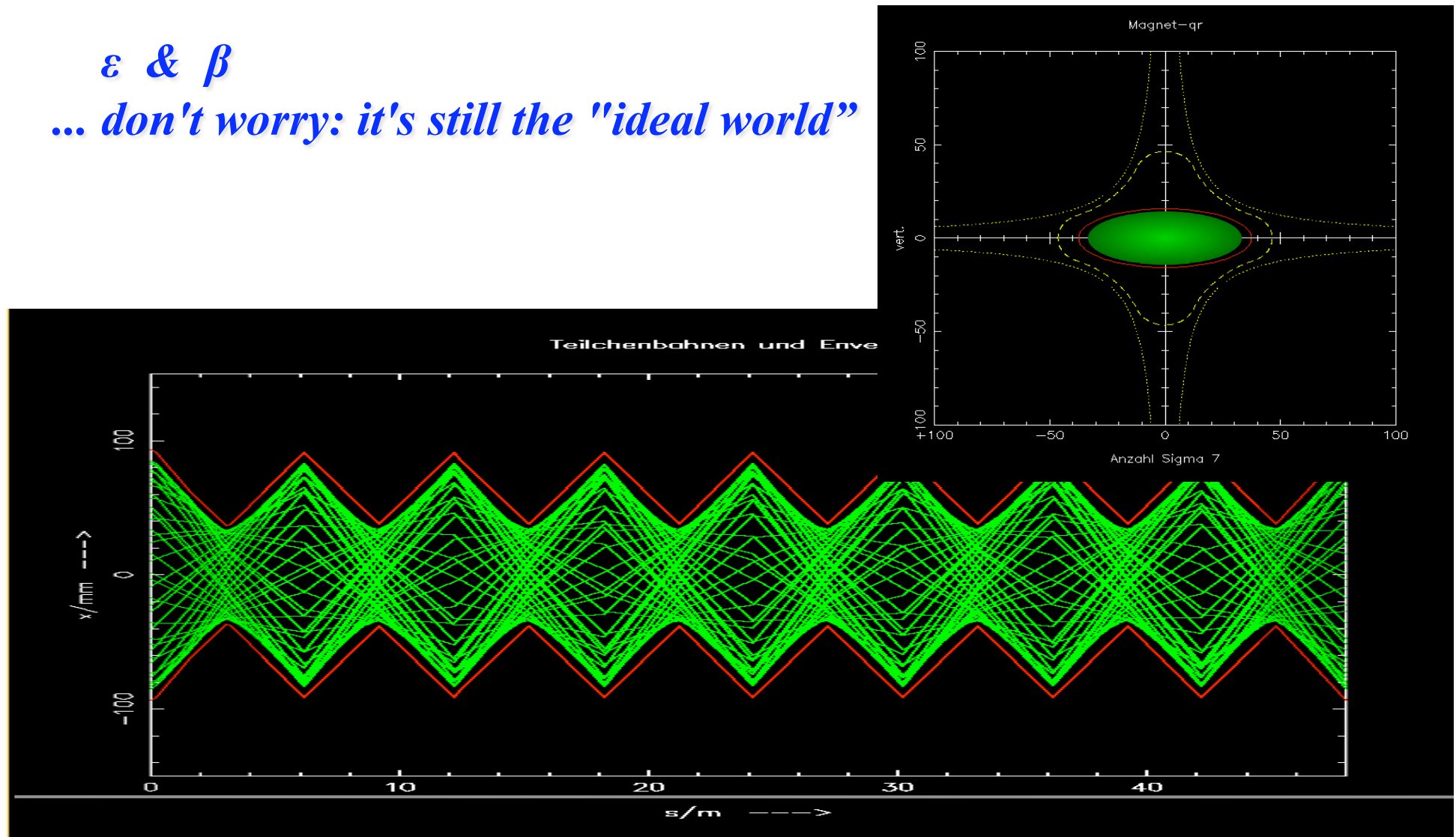


Introduction to Transverse Beam Optics

II.) Particle Trajectories, Beams & Bunch

ε & β

... don't worry: it's still the "ideal world"



4.) Solution of Trajectory Equations

$$x'' + K x = 0$$

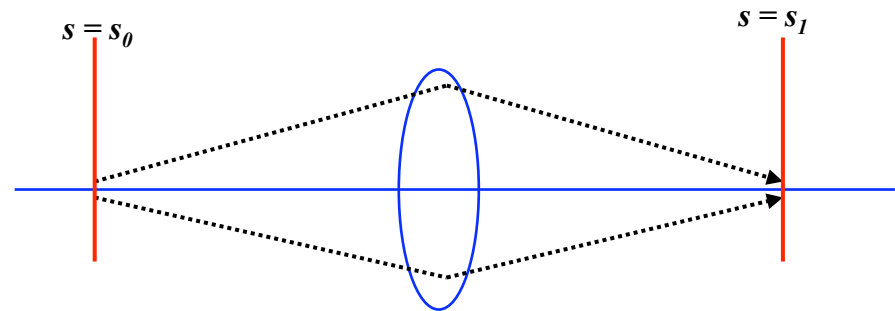
Hor. Focusing Quadrupole $K > 0$:

$$x(s) = x_0 \cdot \cos(\sqrt{|K|}s) + x'_0 \cdot \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}s)$$

$$x'(s) = -x_0 \cdot \sqrt{|K|} \cdot \sin(\sqrt{|K|}s) + x'_0 \cdot \cos(\sqrt{|K|}s)$$

For convenience expressed in matrix formalism:

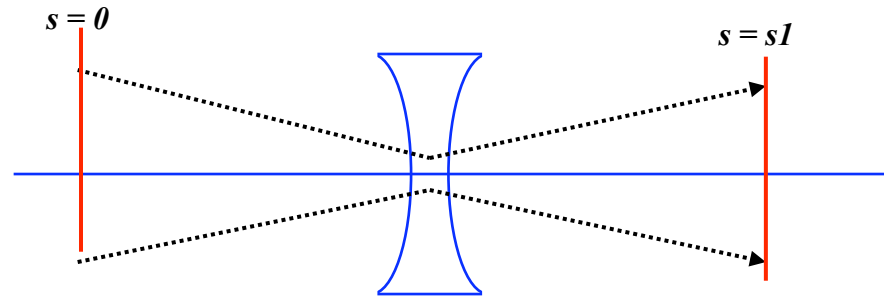
$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s1} = M_{foc} * \begin{pmatrix} x \\ x' \end{pmatrix}_{s0}$$



$$M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}s) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}s) \\ -\sqrt{|K|} \sin(\sqrt{|K|}s) & \cos(\sqrt{|K|}s) \end{pmatrix}_0$$

hor. defocusing quadrupole:

$$x'' - K x = 0$$



Remember from school:

$$f(s) = \cosh(s) \quad , \quad f'(s) = \sinh(s)$$

Ansatz: $x(s) = a_1 \cdot \cosh(\omega s) + a_2 \cdot \sinh(\omega s)$

$$M_{defoc} = \begin{pmatrix} \cosh \sqrt{|K|}l & \frac{1}{\sqrt{|K|}} \sinh \sqrt{|K|}l \\ \sqrt{|K|} \sinh \sqrt{|K|}l & \cosh \sqrt{|K|}l \end{pmatrix}$$

drift space:

$$K = 0$$

$$M_{drift} = \begin{pmatrix} 1 & l \\ 0 & 1 \end{pmatrix}$$

! with the assumptions made, the motion in the horizontal and vertical planes are independent „ ... the particle motion in x & y is uncoupled“

Combining the two planes:

Clear enough (hopefully ... ?) : a quadrupole magnet that is focussing o-in one plane acts as defocusing lens in the other plane ... et vice versa.

hor foc. quadrupole lens

$$M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}s) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}s) \\ -\sqrt{|K|} \sin(\sqrt{|K|}s) & \cos(\sqrt{|K|}s) \end{pmatrix}$$

matrix of the same magnet in the vert. plane:

$$M_{defoc} = \begin{pmatrix} \cosh \sqrt{|K|}l & \frac{1}{\sqrt{|K|}} \sinh \sqrt{|K|}l \\ \sqrt{|K|} \sinh \sqrt{|K|}l & \cosh \sqrt{|K|}l \end{pmatrix}$$

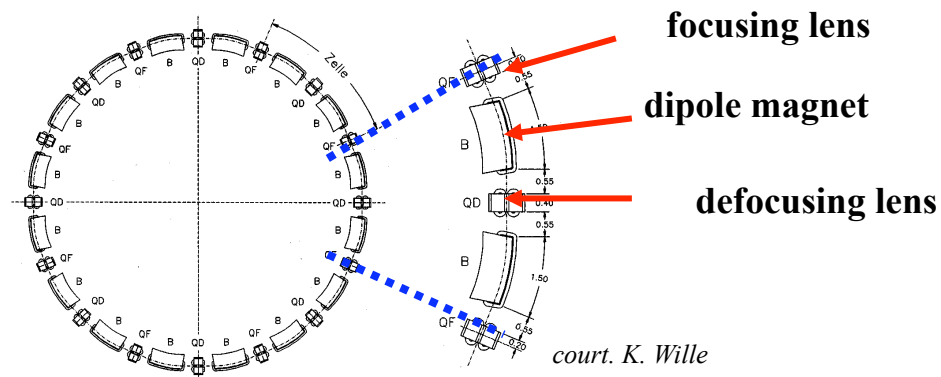
$$\begin{pmatrix} x \\ x' \\ y \\ y' \end{pmatrix}_f = \begin{pmatrix} \cos(\sqrt{|k|}s) & \frac{1}{\sqrt{|k|}} \sin(\sqrt{|k|}s) & 0 & 0 \\ -\sqrt{|k|} \sin(\sqrt{|k|}s) & \cos(\sqrt{|k|}s) & 0 & 0 \\ 0 & 0 & \cosh(\sqrt{|k|}s) & \frac{1}{\sqrt{|k|}} \sinh(\sqrt{|k|}s) \\ 0 & 0 & \sqrt{|k|} \sinh(\sqrt{|k|}s) & \cosh(\sqrt{|k|}s) \end{pmatrix} * \begin{pmatrix} x \\ x' \\ y \\ y' \end{pmatrix}_i$$

Transformation through a system of lattice elements

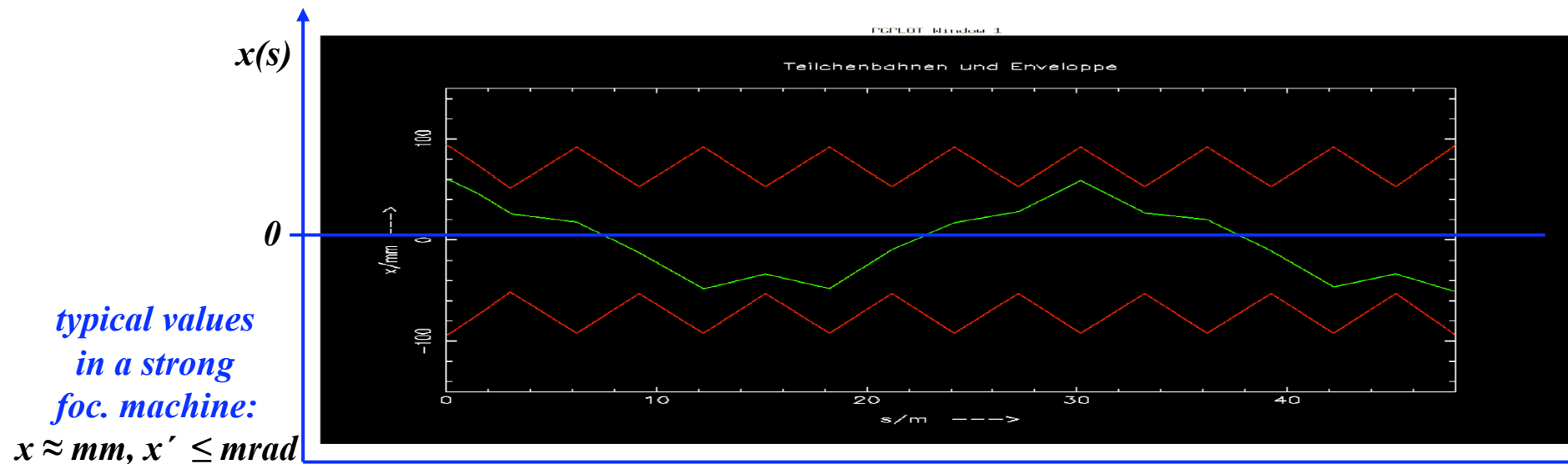
combine the single element solutions by multiplication of the matrices

$$M_{total} = M_{QF} * M_D * M_{QD} * M_{Bend} * M_D * \dots$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_2} = M(s_2, s_1) * \begin{pmatrix} x \\ x' \end{pmatrix}_{s_1}$$

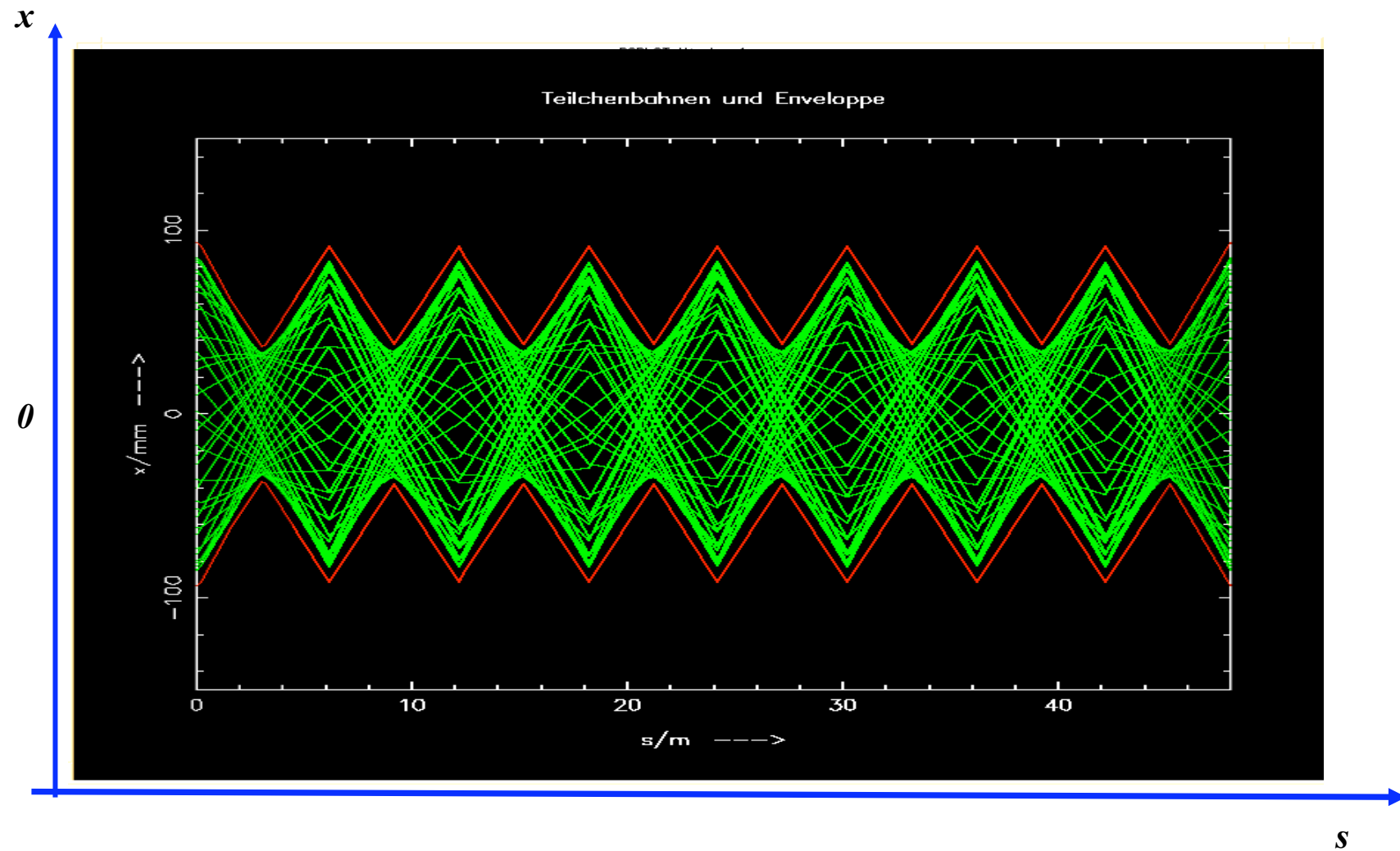


in each accelerator element the particle trajectory corresponds to the movement of a harmonic oscillator „



Question: what will happen, if the particle performs a second turn ?

... or a third one or ... 10^{10} turns



Astronomer Hill:

*differential equation for motions with periodic focusing properties
„Hill's equation“*

*Example: particle motion with
periodic coefficient*



equation of motion: $x''(s) - k(s)x(s) = 0$

*restoring force $\neq \text{const}$,
 $k(s)$ = depending on the position s
 $k(s+L) = k(s)$, periodic function*



*we expect a kind of quasi harmonic
oscillation: amplitude & phase will depend
on the position s in the ring.*

6.) The Beta Function

General solution of Hill's equation:

$$(i) \quad x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\psi(s) + \phi)$$

ε, Φ = integration **constants** determined by initial conditions

$\beta(s)$ **periodic function** given by **focusing properties** of the lattice $\leftrightarrow$ quadrupoles

$$\beta(s + L) = \beta(s)$$

Inserting (i) into the equation of motion ...

$$\psi(s) = \int_0^s \frac{ds}{\beta(s)}$$

$\Psi(s)$ = „**phase advance**“ of the oscillation between point „0“ and „s“ in the lattice.

For one complete revolution: number of oscillations per turn „**Tune**“

$$Q_y = \frac{1}{2\pi} \oint \frac{ds}{\beta(s)}$$

The Beta Function

Amplitude of a particle trajectory:

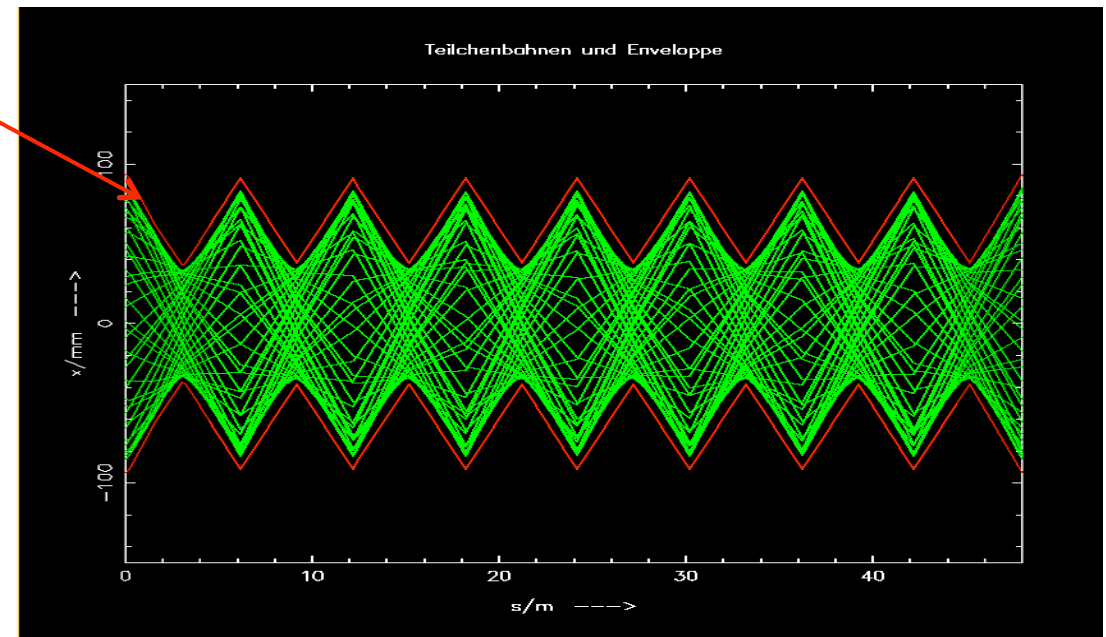
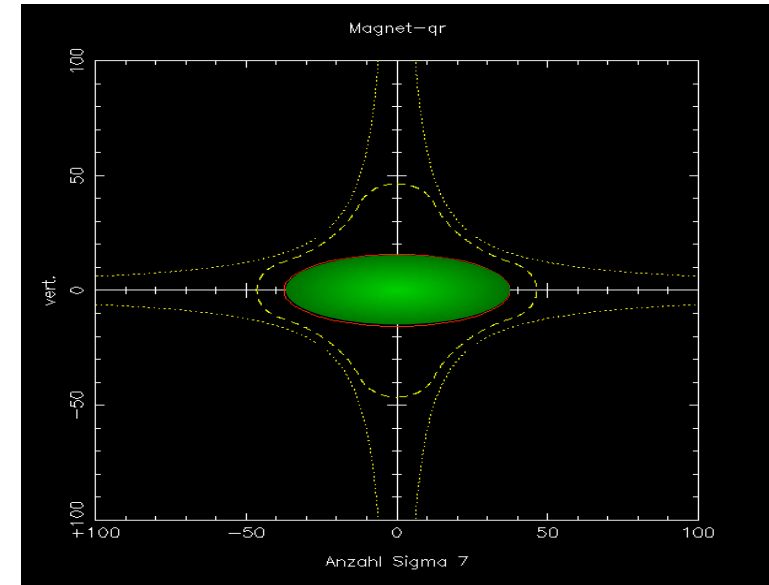
$$x(s) = \sqrt{\varepsilon} * \sqrt{\beta(s)} * \cos(\psi(s) + \varphi)$$

Maximum size of a particle amplitude

$$\hat{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)}$$

β determines the beam size
(... the envelope of all particle trajectories at a given position “s” in the storage ring.

It reflects the periodicity of the magnet structure.



7.) Beam Emittance and Phase Space Ellipse

general solution of Hill equation

$$\left\{ \begin{array}{l} (1) \quad \mathbf{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos(\psi(s) + \phi) \\ (2) \quad \mathbf{x}'(s) = -\frac{\sqrt{\varepsilon}}{\sqrt{\beta(s)}} \{ \alpha(s) \cos(\psi(s) + \phi) + \sin(\psi(s) + \phi) \} \end{array} \right.$$

from (1) we get

$$\cos(\psi(s) + \phi) = \frac{\mathbf{x}(s)}{\sqrt{\varepsilon} \sqrt{\beta(s)}}$$

$$\alpha(s) = \frac{-1}{2} \beta'(s)$$

$$\gamma(s) = \frac{1 + \alpha(s)^2}{\beta(s)}$$

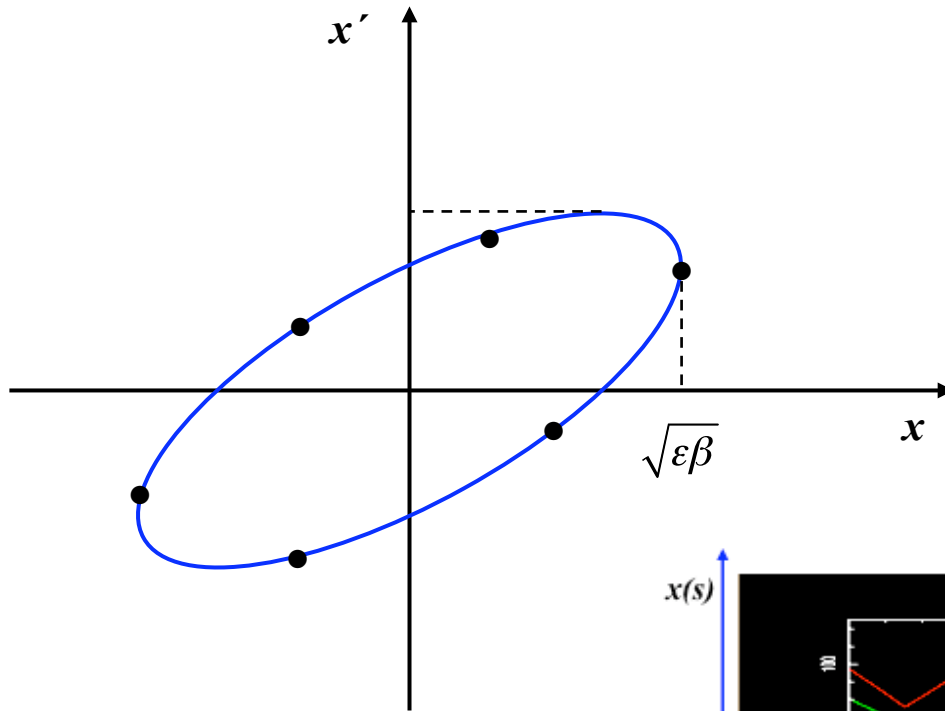
Insert into (2) and solve for ε

$$\varepsilon = \gamma(s) \mathbf{x}^2(s) + 2\alpha(s)\mathbf{x}(s)\mathbf{x}'(s) + \beta(s) \mathbf{x}'^2(s)$$

- * ε is a **constant** of the motion ... **it is independent of „s“**
- * **parametric representation of an ellipse in the $x \ x'$ space**
- * **shape and orientation of ellipse are given by α, β, γ**

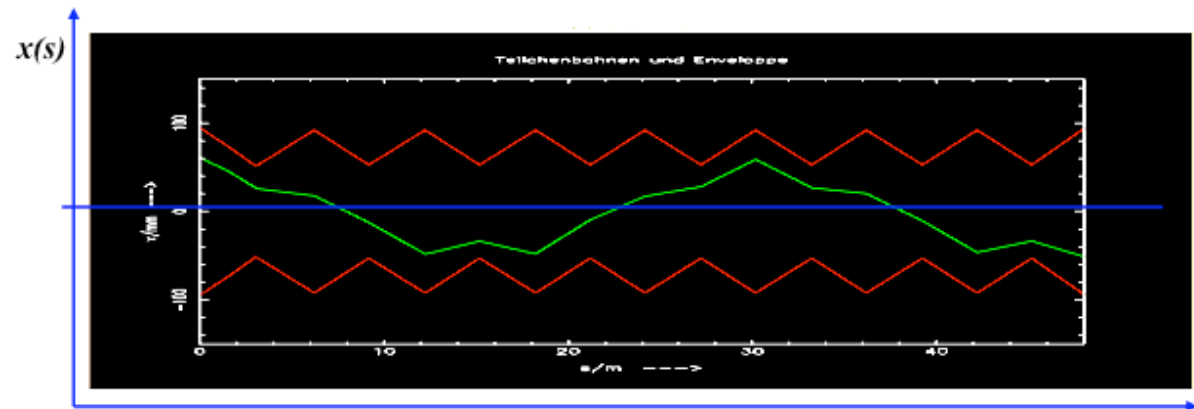
Beam Emittance and Phase Space Ellipse

$$\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$$



*Liouville: in reasonable storage rings
area in phase space is constant.*

$$A = \pi \cdot \varepsilon = \text{const}$$



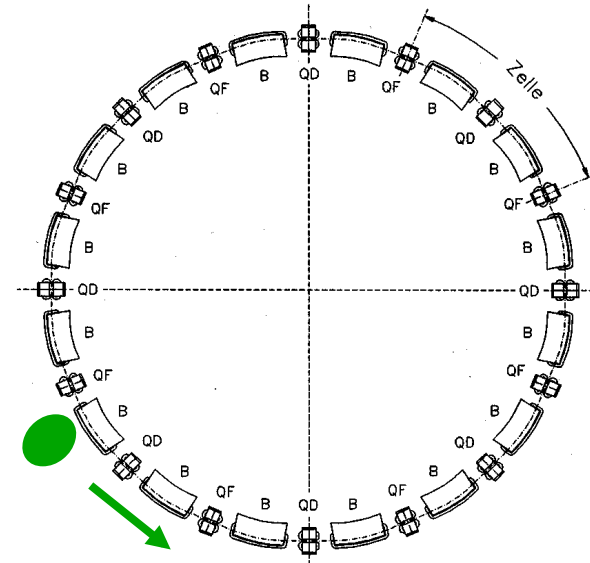
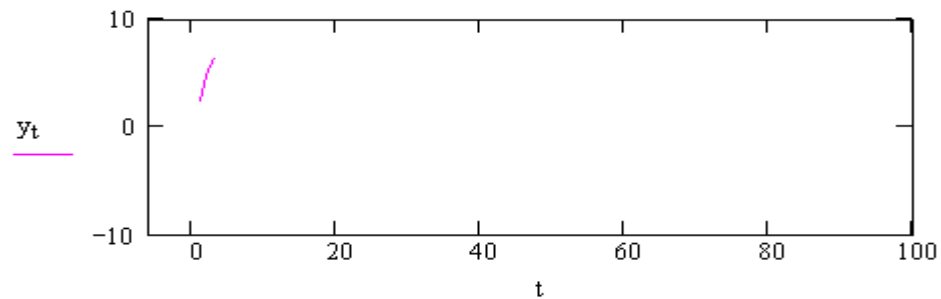
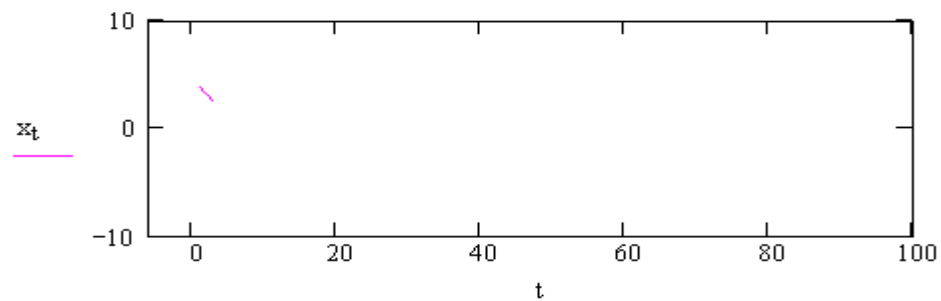
ε beam emittance = *woozilycity* of the particle ensemble, *intrinsic beam parameter*,
cannot be changed by the foc. properties.

Scientifiquely speaking: area covered in transverse x, x' phase space ... and it is constant !!!

Particle Tracking in a Storage Ring

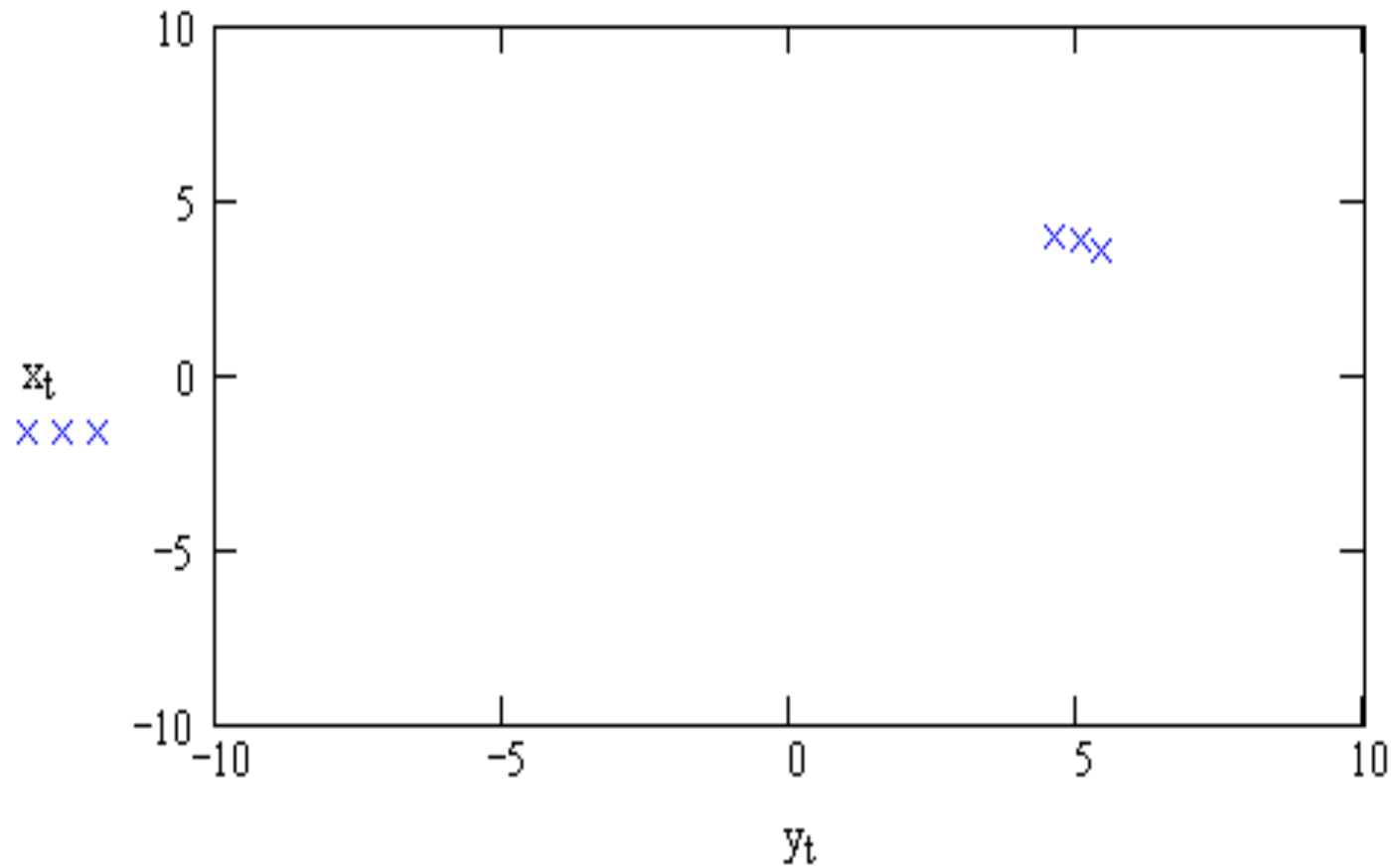
Calculate x, x' for each linear accelerator element according to matrix formalism

plot x, x' as a function of „s“



... and now the ellipse:

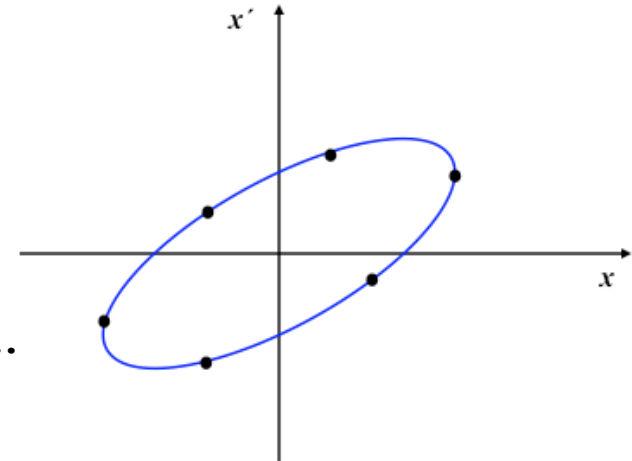
note for each turn x , x' at a given position „ s_1 “ and plot in the phase space diagram



Phase Space Ellipse

particle trajectory: $x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos \{ \psi(s) + \phi \}$

max. Amplitude: $\hat{x}(s) = \sqrt{\varepsilon \beta}$ $\longrightarrow$ x' at that position ...



... put $\hat{x}(s)$ into $\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$ and solve for x'

$$\varepsilon = \gamma \cdot \varepsilon \beta + 2\alpha \sqrt{\varepsilon \beta} \cdot x' + \beta x'^2$$

$\longrightarrow$ $x' = -\alpha \cdot \sqrt{\varepsilon / \beta}$

* A high β -function means a large beam size and a small beam divergence.
... et vice versa !!!

!

* In the middle of a quadrupole $\beta = \text{maximum}$,
 $\alpha = \text{zero}$ $\left. \vphantom{\begin{matrix} \beta = \text{maximum} \\ \alpha = \text{zero} \end{matrix}} \right\} x' = 0$

... and the ellipse is flat

Phase Space Ellipse

$$\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$$

$$\alpha(s) = -\frac{1}{2} \beta'(s)$$

$$\gamma(s) = \frac{1 + \alpha(s)^2}{\beta(s)}$$

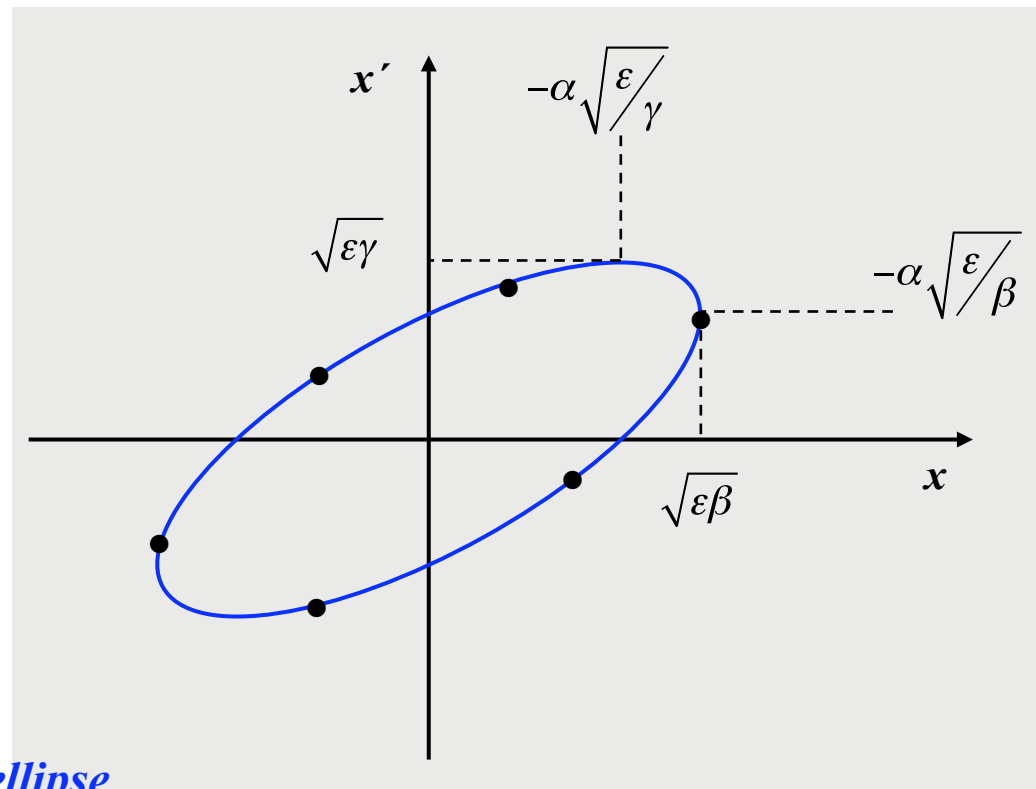
$$\longrightarrow \varepsilon = \frac{x^2}{\beta} + \frac{\alpha^2 x^2}{\beta} + 2\alpha \cdot x x' + \beta \cdot x'^2$$

$$\dots \text{ solve for } x' \quad x'_{1,2} = \frac{-\alpha \cdot x \pm \sqrt{\varepsilon\beta - x^2}}{\beta}$$

$$\dots \text{ and determine } \hat{x}' \text{ via: } \frac{dx'}{dx} = 0$$

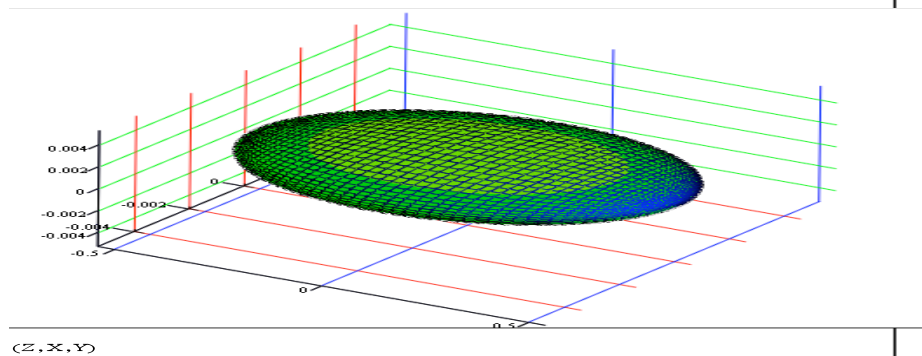
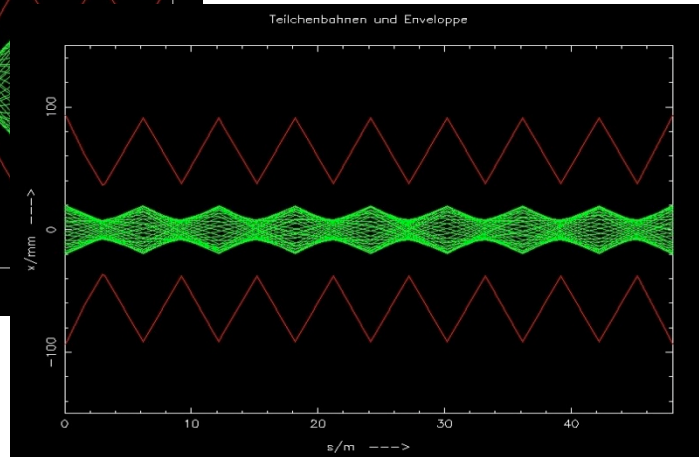
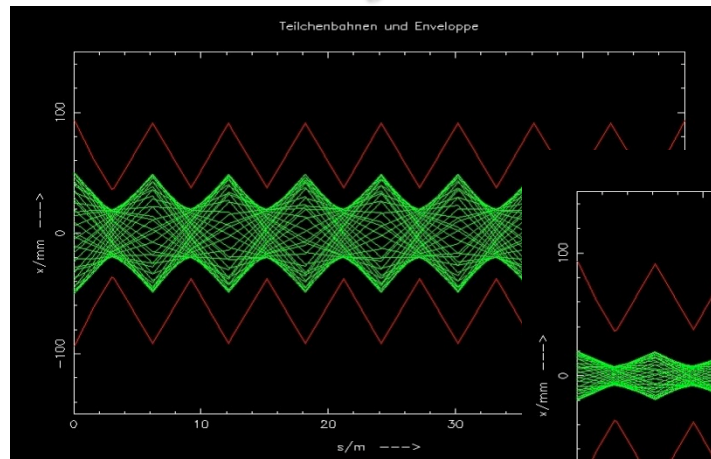
$$\longrightarrow \hat{x}' = \sqrt{\varepsilon\gamma}$$

$$\longrightarrow \hat{x} = \pm \alpha \sqrt{\varepsilon/\gamma}$$



*shape and orientation of the phase space ellipse
depend on the Twiss parameters β and γ*

Emittance of the Particle Ensemble:



0.04

-0.04

-0.1

0

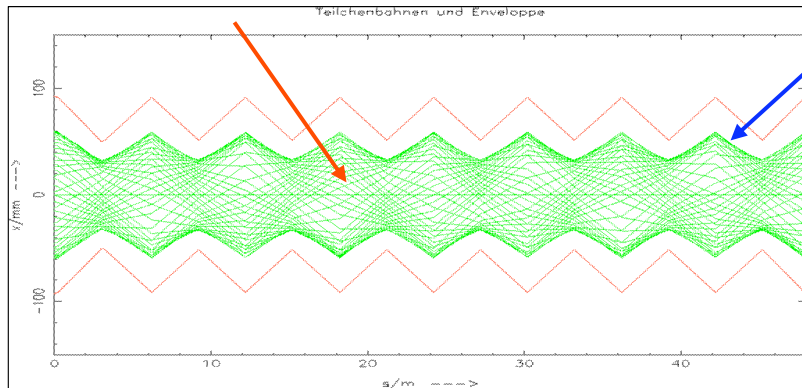
0.1

$x_{n,1}, x_{n,2}, x_{n,3}, x_{n,4}$
CAS Prague 2014

Emittance of the Particle Ensemble:

$$x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\Psi(s) + \phi)$$

$$\hat{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)}$$



single particle trajectories, $N \approx 10^{11}$ per bunch

**Gauß
Particle Distribution:**

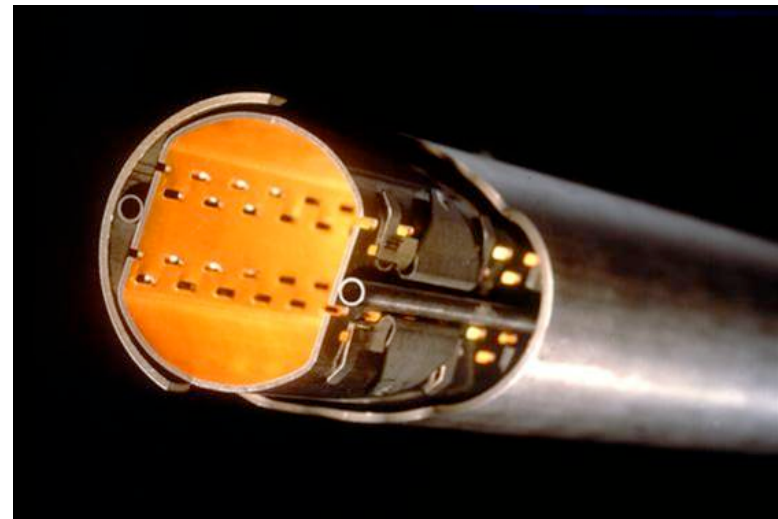
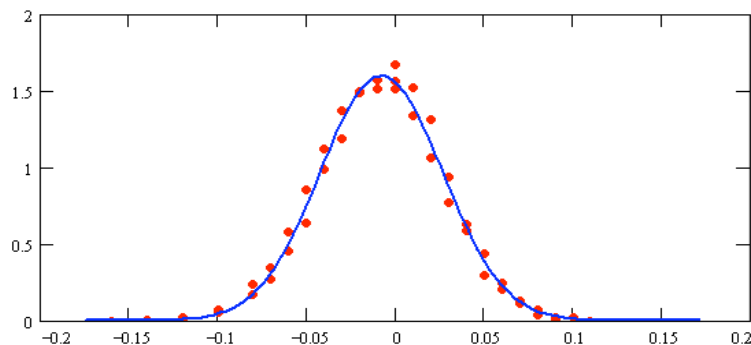
$$\rho(x) = \frac{N \cdot e}{\sqrt{2\pi} \sigma_x} \cdot e^{-\frac{1}{2} \frac{x^2}{\sigma_x^2}}$$

particle at distance 1σ from centre
 $\leftrightarrow 68.3 \%$ of all beam particles

LHC: $\beta = 180 m$

$$\varepsilon = 5 * 10^{-10} m rad$$

$$\sigma = \sqrt{\varepsilon * \beta} = \sqrt{5 * 10^{-10} m * 180 m} = 0.3 mm$$



aperture requirements: $r_0 = 12 * \sigma$

9.) Transfer Matrix M ... yes we had the topic already

*general solution
of Hill's equation*

$$\left\{ \begin{array}{l} x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos \{\psi(s) + \phi\} \\ x'(s) = \frac{-\sqrt{\varepsilon}}{\sqrt{\beta(s)}} [\alpha(s) \cos \{\psi(s) + \phi\} + \sin \{\psi(s) + \phi\}] \end{array} \right.$$

remember the trigonometrical gymnastics: $\sin(a + b) = \dots$ etc

$$x(s) = \sqrt{\varepsilon} \sqrt{\beta_s} (\cos \psi_s \cos \phi - \sin \psi_s \sin \phi)$$

$$x'(s) = \frac{-\sqrt{\varepsilon}}{\sqrt{\beta_s}} [\alpha_s \cos \psi_s \cos \phi - \alpha_s \sin \psi_s \sin \phi + \sin \psi_s \cos \phi + \cos \psi_s \sin \phi]$$

starting at point $s(0) = s_0$, where we put $\Psi(0) = 0$

$$\cos \phi = \frac{x_0}{\sqrt{\varepsilon \beta_0}},$$

$$\sin \phi = -\frac{1}{\sqrt{\varepsilon}} (x'_0 \sqrt{\beta_0} + \frac{\alpha_0 x_0}{\sqrt{\beta_0}})$$

inserting above ...

$$\underline{x(s) = \sqrt{\frac{\beta_s}{\beta_0}} \{ \cos \psi_s + \alpha_0 \sin \psi_s \} x_0 + \{ \sqrt{\beta_s \beta_0} \sin \psi_s \} x'_0}$$

$$\underline{x'(s) = \frac{1}{\sqrt{\beta_s \beta_0}} \{ (\alpha_0 - \alpha_s) \cos \psi_s - (1 + \alpha_0 \alpha_s) \sin \psi_s \} x_0 + \sqrt{\frac{\beta_0}{\beta_s}} \{ \cos \psi_s - \alpha_s \sin \psi_s \} x'_0}$$

which can be expressed ... for convenience ... *in matrix form* $\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \begin{pmatrix} x \\ x' \end{pmatrix}_0$

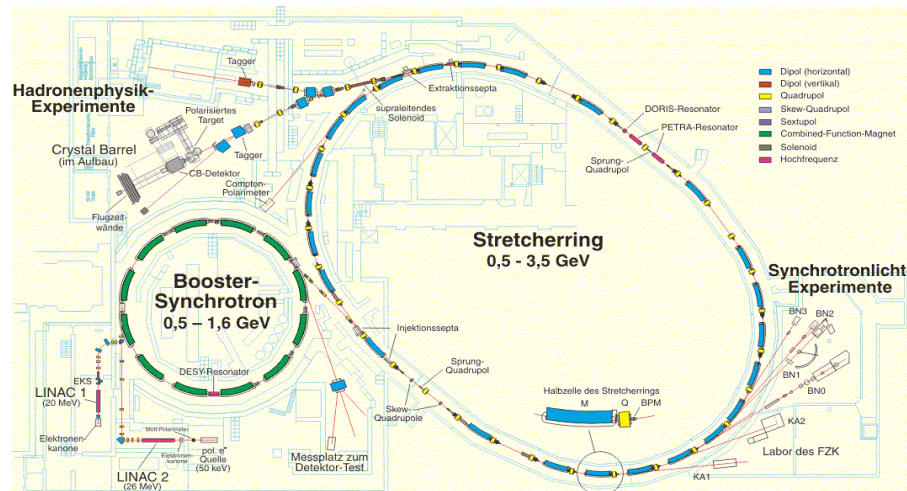
$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos \psi_s + \alpha_0 \sin \psi_s) & \sqrt{\beta_s \beta_0} \sin \psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos \psi_s - (1 + \alpha_0 \alpha_s) \sin \psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos \psi_s - \alpha_s \sin \psi_s) \end{pmatrix}$$

- * we can calculate *the single particle trajectories* between two locations in the ring, *if we know the α β γ at these positions.*
- * *and nothing but the α β γ at these positions.*

* ... !

10.) Periodic Lattices

$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos \psi_s + \alpha_0 \sin \psi_s) & \sqrt{\beta_s \beta_0} \sin \psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos \psi_s - (1 + \alpha_0 \alpha_s) \sin \psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos \psi_s - \alpha_s \sin \psi_s) \end{pmatrix}$$



ELSA Electron Storage Ring

„This *rather formidable looking matrix* simplifies considerably if we *consider one complete revolution ...*“

$$M(s) = \begin{pmatrix} \cos \psi_{turn} + \alpha_s \sin \psi_{turn} & \beta_s \sin \psi_{turn} \\ -\gamma_s \sin \psi_{turn} & \cos \psi_{turn} - \alpha_s \sin \psi_{turn} \end{pmatrix}$$

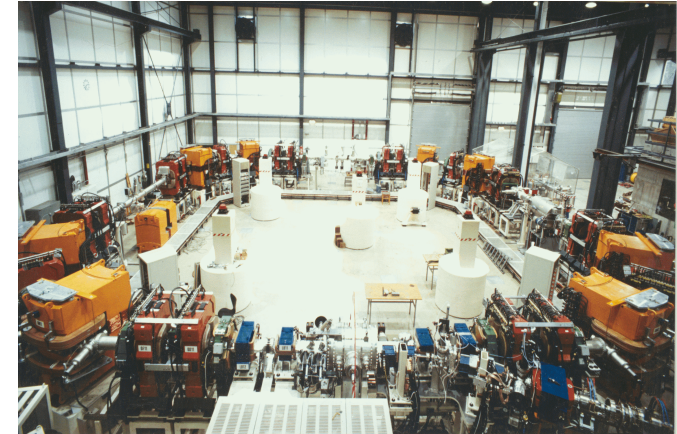
$$\psi_{turn} = \int_s^{s+L} \frac{ds}{\beta(s)} \quad \psi_{turn} = \text{phase advance per period}$$

$$Q = \frac{1}{2\pi} \oint \frac{ds}{\beta(s)}$$

Tune: Phase advance per turn in units of 2π

Stability Criterion:

Question: what will happen, if we do not make too many mistakes and your *particle performs one complete turn* ?



Matrix for 1 turn:

$$M = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix} = \underbrace{\cos\psi \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}}_{\mathbf{1}} + \sin\psi \underbrace{\begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}}_{\mathbf{J}}$$

Matrix for N turns:

$$M^N = (1 \cdot \cos\psi + J \cdot \sin\psi)^N = 1 \cdot \cos N\psi + J \cdot \sin N\psi$$

The motion for N turns remains bounded, if the elements of M^N remain bounded

$$\psi = \text{real} \quad \Leftrightarrow \quad |\cos\psi| \leq 1 \quad \Leftrightarrow \quad \text{Tr}(M) \leq 2$$

stability criterion proof for the disbelieving colleagues !!

Matrix for 1 turn:

$$M = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix} = \underbrace{\cos\psi \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}}_I + \sin\psi \underbrace{\begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}}_J$$

Matrix for 2 turns:

$$\begin{aligned} M^2 &= (I \cos\psi_1 + J \sin\psi_1)(I \cos\psi_2 + J \sin\psi_2) \\ &= I^2 \cos\psi_1 \cos\psi_2 + IJ \cos\psi_1 \sin\psi_2 + JI \sin\psi_1 \cos\psi_2 + J^2 \sin\psi_1 \sin\psi_2 \end{aligned}$$

now ...

$$I^2 = I$$

$$\left. \begin{aligned} IJ &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} * \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \\ JI &= \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} * \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \end{aligned} \right\} IJ = JI$$

$$J^2 = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} * \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} = \begin{pmatrix} \alpha^2 - \gamma\beta & \alpha\beta - \beta\alpha \\ -\gamma\alpha + \alpha\gamma & \alpha^2 - \gamma\beta \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} = -I$$

$$M^2 = I \cos(\psi_1 + \psi_2) + J \sin(\psi_1 + \psi_2)$$

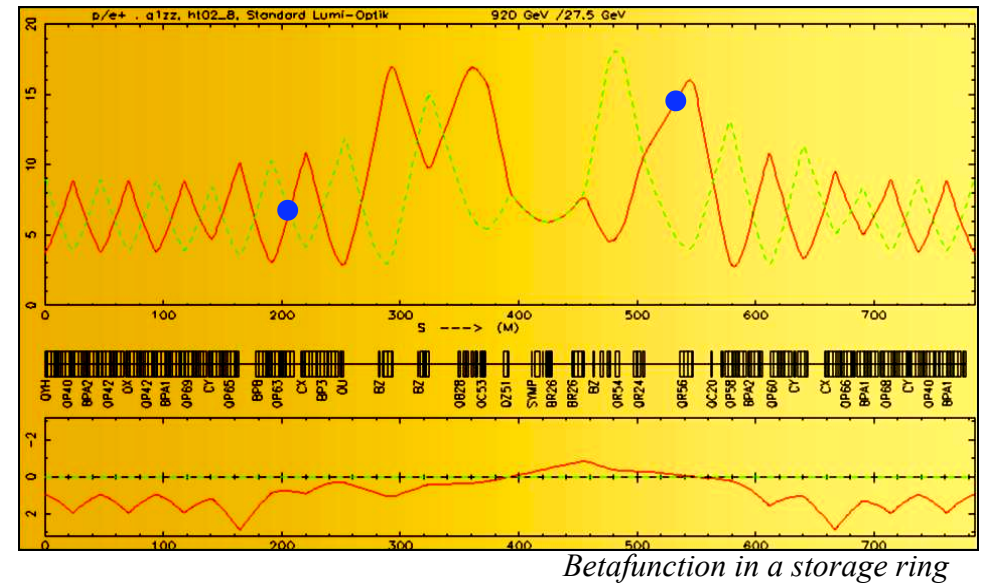
$$M^2 = I \cos(2\psi) + J \sin(2\psi)$$

11.) Transformation of α, β, γ

consider two positions in the storage ring: s_0, s

$$\begin{pmatrix} x \\ x' \end{pmatrix}_s = M * \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$

$$M = \begin{pmatrix} C & S \\ C' & S' \end{pmatrix}$$



since $\varepsilon = \text{const}$ (Liouville):

$$\varepsilon = \beta_s x'^2 + 2\alpha_s x x' + \gamma_s x^2$$

$$\varepsilon = \beta_0 x_0'^2 + 2\alpha_0 x_0 x_0' + \gamma_0 x_0^2$$

... remember $W = CS' - SC' = 1$

$$\left. \begin{aligned} \begin{pmatrix} x \\ x' \end{pmatrix}_0 &= M^{-1} * \begin{pmatrix} x \\ x' \end{pmatrix}_s \\ M^{-1} &= \begin{pmatrix} S' & -S \\ -C' & C \end{pmatrix} \end{aligned} \right\} \rightarrow \begin{aligned} x_0 &= S'x - Sx' \\ x_0' &= -C'x + Cx' \end{aligned} \quad \dots \text{inserting into } \varepsilon$$

$$\varepsilon = \beta_0 (Cx' - C'x)^2 + 2\alpha_0 (S'x - Sx')(Cx' - C'x) + \gamma_0 (S'x - Sx')^2$$

sort via x, x' and compare the coefficients to get

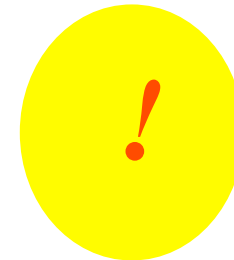
$$\beta(s) = C^2 \beta_0 - 2SC\alpha_0 + S^2 \gamma_0$$

$$\alpha(s) = -CC' \beta_0 + (SC' + S'C)\alpha_0 - SS' \gamma_0$$

$$\gamma(s) = C'^2 \beta_0 - 2S'C'\alpha_0 + S'^2 \gamma_0$$

in matrix notation:

$$\begin{pmatrix} \beta \\ \alpha \\ \gamma \end{pmatrix}_s = \begin{pmatrix} C^2 & -2SC & S^2 \\ -CC' & SC' + CS' & -SS' \\ C'^2 & -2S'C' & S'^2 \end{pmatrix} \cdot \begin{pmatrix} \beta_0 \\ \alpha_0 \\ \gamma_0 \end{pmatrix}$$



- 1.) *this expression is important*
- 2.) *given the twiss parameters α, β, γ at any point in the lattice we can transform them and calculate their values at any other point in the ring.*
- 3.) *the transfer matrix is given by the focusing properties of the lattice elements, the elements of M are just those that we used to calculate single particle trajectories.*
- 4.) *go back to point 1.)*

12.) Lattice Design:

„... how to build a storage ring“

$$B \rho = p / q$$

Circular Orbit: dipole magnets to define the geometry

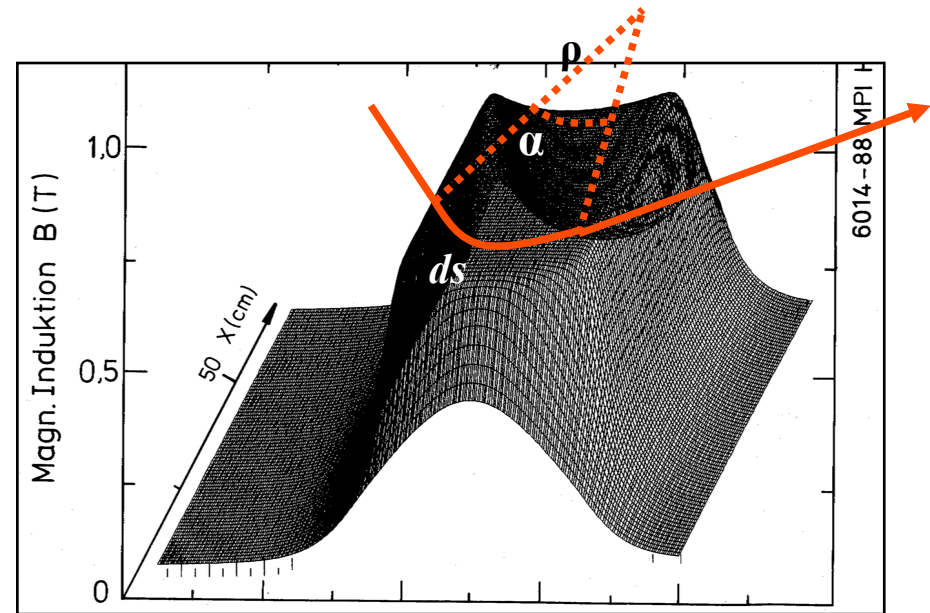
$$\alpha = \frac{ds}{\rho} \approx \frac{dl}{\rho} = \frac{Bdl}{B\rho}$$

The angle run out in one revolution must be 2π , so

$$\dots \text{for a full circle} \quad \alpha = \frac{\int Bdl}{B\rho} = 2\pi \quad \rightarrow \quad \int Bdl = 2\pi \frac{p}{q}$$

... defines the integrated dipole field around the machine.

Nota bene: $\frac{\Delta B}{B} \approx 10^{-4}$ is usually required !!



field map of a storage ring dipole magnet

Example LHC:



7000 GeV Proton storage ring
dipole magnets $N = 1232$
 $l = 15 \text{ m}$
 $q = +1 e$

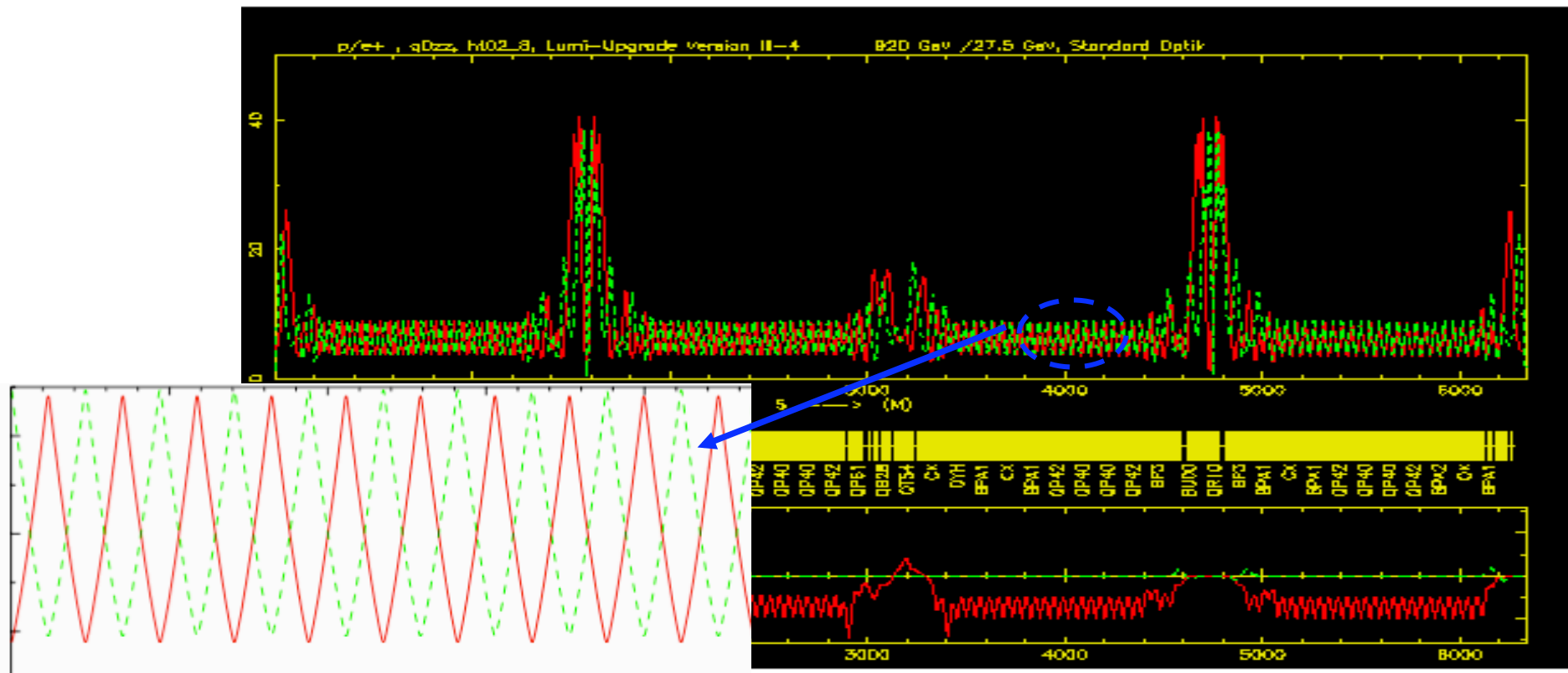
$$\int B \, dl \approx N \, l \, B = 2\pi \, p / e$$

$$B \approx \frac{2\pi \, 7000 \, 10^9 \, eV}{1232 \, 15 \, m \, 3 \, 10^8 \, \frac{m}{e}} = 8.3 \, \text{Tesla}$$

The FoDo-Lattice

A magnet structure consisting of focusing and defocusing quadrupole lenses in alternating order with **nothing** in between.

(**Nothing** = elements that can be neglected on first sight: drift, bending magnets, RF structures ... **and especially experiments...**)



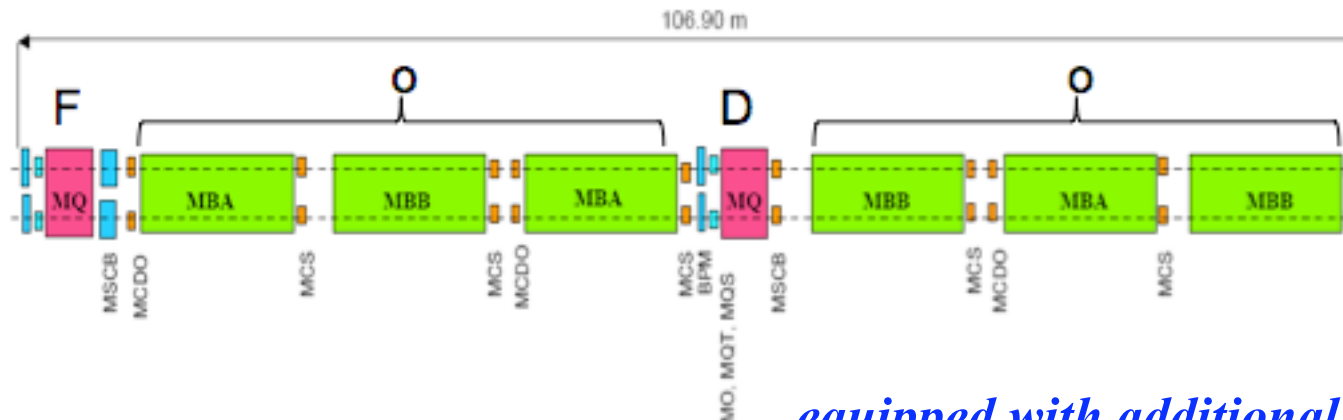
Starting point for the calculation: in the middle of a focusing quadrupole

Phase advance per cell $\mu = 45^\circ$,

→ calculate the twiss parameters for a periodic solution

LHC: Lattice Design

the ARC 90° FoDo in both planes

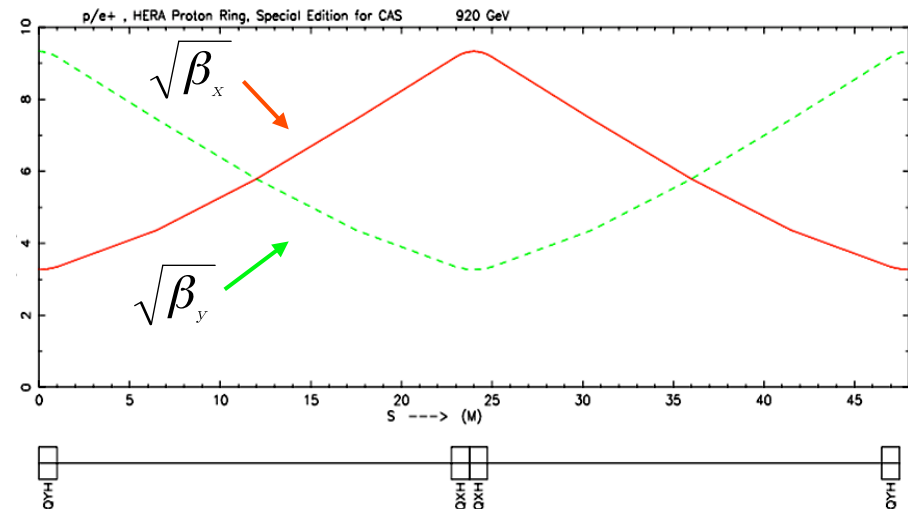
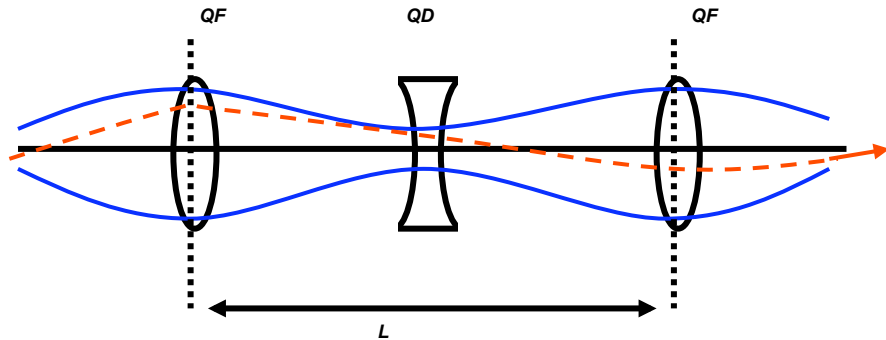


equipped with additional corrector coils



MB: main dipole
MQ: main quadrupole
MQT: Trim quadrupole
MQS: Skew trim quadrupole
MO: Lattice octupole (Landau damping)
MSCB: Skew sextupole
Orbit corrector dipoles
MCS: Spool piece sextupole
MCDO: Spool piece 8 / 10 pole
BPM: Beam position monitor + diagnostics

Periodic solution of a FoDo Cell



Output of the optics program:

<i>Nr</i>	<i>Type</i>	<i>Length</i> <i>m</i>	<i>Strength</i> <i>1/m2</i>	β_x <i>m</i>	α_x	ψ_x <i>1/2π</i>	β_y <i>m</i>	α_y	ψ_y <i>1/2π</i>
0	IP	0,000	0,000	11,611	0,000	0,000	5,295	0,000	0,000
1	QFH	0,250	-0,541	11,228	1,514	0,004	5,488	-0,781	0,007
2	QD	3,251	0,541	5,488	-0,781	0,070	11,228	1,514	0,066
3	QFH	6,002	-0,541	11,611	0,000	0,125	5,295	0,000	0,125
4	IP	6,002	0,000	11,611	0,000	0,125	5,295	0,000	0,125

$$Q_x = 0,125 \quad Q_y = 0,125$$

$$0,125 * 2\pi = 45^\circ$$

Can we understand, what the optics code is doing?

$$\text{matrices} \quad M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}l_q) & \frac{1}{\sqrt{|K|}}\sin(\sqrt{|K|}l_q) \\ -\sqrt{|K|}\sin(\sqrt{|K|}l_q) & \cos(\sqrt{|K|}l_q) \end{pmatrix} \quad M_{drift} = \begin{pmatrix} 1 & l_d \\ 0 & 1 \end{pmatrix}$$

strength and length of the FoDo elements

$$K = \pm 0.54102 \text{ m}^{-2}$$

$$l_q = 0.5 \text{ m}$$

$$l_d = 2.5 \text{ m}$$

The matrix for the **complete cell** is obtained by multiplication of the element matrices

$$M_{FoDo} = M_{qf h} * M_{ld} * M_{qd} * M_{ld} * M_{qf}$$

Putting the numbers in and **multiplying out** ...

$$M_{FoDo} = \begin{pmatrix} 0.707 & 8.206 \\ -0.061 & 0.707 \end{pmatrix}$$

The transfer matrix for one period gives us all the information that we need !

Phase advance per cell

$$M(s) = \begin{pmatrix} \cos \psi + \alpha \sin \psi & \beta \sin \psi \\ -\gamma \sin \psi & \cos \psi - \alpha \sin \psi \end{pmatrix} \rightarrow$$

$$\cos(\psi) = \frac{1}{2} \text{Trace}(M) = 0.707$$

$$\psi = \text{arc cos}\left(\frac{1}{2} \text{Trace}(M)\right) = \underline{\underline{45^\circ}}$$

hor β -function

hor α -function

$$\beta = \frac{M_{1,2}}{\sin \psi} = \underline{\underline{11.611 \text{ m}}}$$

$$\alpha = \frac{M_{1,1} - \cos \psi}{\sin \psi} = \underline{\underline{0}}$$

Resume':

transfer matrix in Twiss form

$$\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \begin{pmatrix} x \\ x' \end{pmatrix}_0$$

$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos \psi_s + \alpha_0 \sin \psi_s) & \sqrt{\beta_s \beta_0} \sin \psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos \psi_s - (1 + \alpha_0 \alpha_s) \sin \psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos \psi_s - \alpha_s \sin \psi_s) \end{pmatrix}$$

... and for the periodic case

$$M(s) = \begin{pmatrix} \cos \psi_{turn} + \alpha_s \sin \psi_{turn} & \beta_s \sin \psi_{turn} \\ -\gamma_s \sin \psi_{turn} & \cos \psi_{turn} - \alpha_s \sin \psi_{turn} \end{pmatrix}$$

beam emittance during acceleration

$$\varepsilon \propto \frac{1}{\beta \gamma}$$

dispersion

$$D(s) = \frac{x_i(s)}{\Delta p / p}$$