



Instabilities Part IV: Electron cloud – build up and effects on beam dynamics

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Outline



We will look into the description and the impact of **electron cloud**. We will discuss the conditions for an electron cloud to build in the vacuum chamber of an accelerator and mitigation/suppression techniques. We will also show some examples linked to **electron cloud effects** such as beam induced instability and incoherent effects.

Part 4: Electron cloud – Build up and effects on beam dynamics

- Electron cloud build up
 - Electron production and multiplication
 - Observation in accelerator rings
- Scrubbing and other techniques of mitigation/suppression
- E-cloud induced instabilities and incoherent effects







- We have learned about the concept of particles, macroparticles and particle distributions as well as some peculiarities of multiparticle dynamics in accelerators.
- We have learned about the basic **concept of wake fields** and how these can be characterized as a **collective effect** in that they depend on the particle distribution.
- We have learned the **impact of these** in the longitudinal and transverse planes.
- We are ready to look into a new, but popular ©, source of collective effects, i.e. the **electron cloud**

Part 4: Electron cloud –

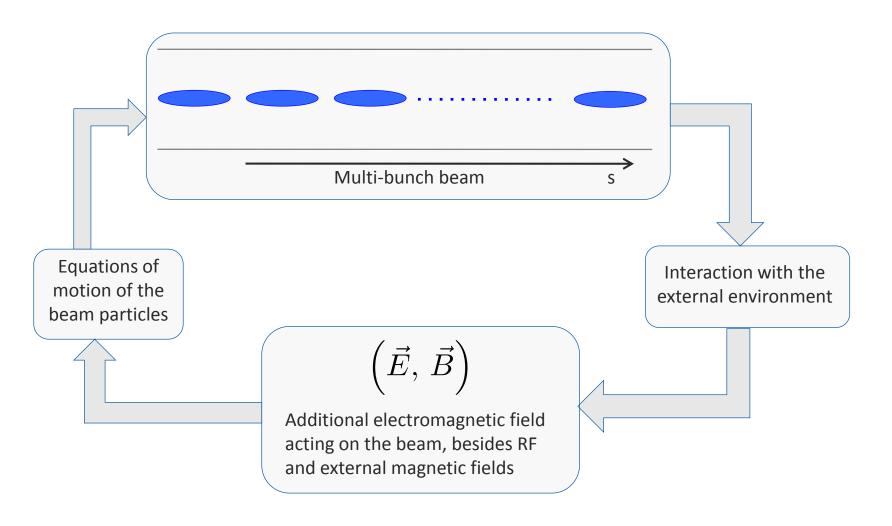
Build up and effects on beam dynamics

- Electron cloud build up
 - Electron production and multiplication
 - Observation in accelerator rings
- Techniques of mitigation/suppression
- E-cloud induced instabilities and incoherent effects



Reminder: The instability loop







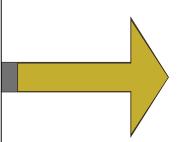
Reminder



Interaction of the beam with the external environment

Pure EM interaction

- Maxwell's equations
 - The beam as the source term
 - Boundary conditions given by the chamber in which the beam is propagating
- Generation of wake functions/impedances



 $\left(ec{E},\,ec{B}
ight)$

Additional electromagnetic field acting on the beam, besides RF and external magnetic fields



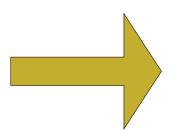
Different type of interaction possible



Interaction of the beam with the external environment

The electron cloud

- Electron production and accumulation
- Poisson's equation with
 - The electron cloud as the source term
 - Boundary conditions given by the chamber in which the electron cloud builds up



 $\left(ec{E},\,ec{B}
ight)$

Additional electromagnetic field acting on the beam, besides RF and external magnetic fields





Generation of charged particles inside the vacuum chamber (primary, or seed, electrons)

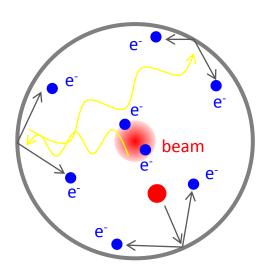
Residual gas

Desorption

ionization

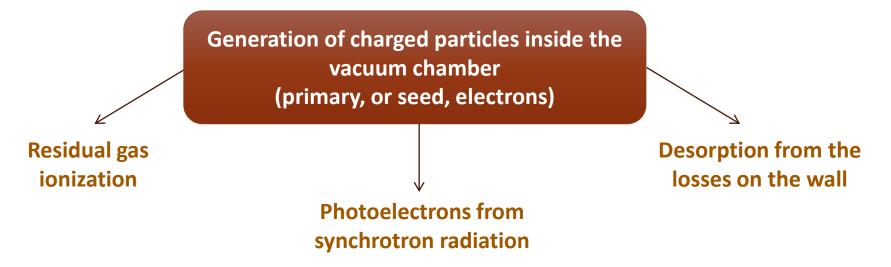
Desorption from the losses on the wall

Photoelectrons from synchrotron radiation









- Gas ionization and wall desorption produce both electrons and ions (the former one with the same rate, the second one with different rates depending on the desorption yields), photoemission is only a source of electrons
- The dominant mechanism depends upon e.g.
 - o Beam type and parameters (e.g. lepton vs hadrons, beam energy)
 - o Vacuum level
 - o Design (material, shape), roughness, cleanness of the inner surface of chamber

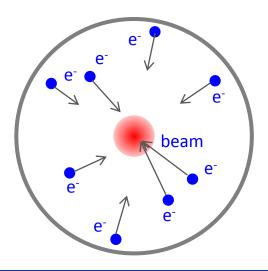




Generation of charged particles inside the vacuum chamber (primary, or seed, electrons)



Acceleration of primary electrons in the beam field



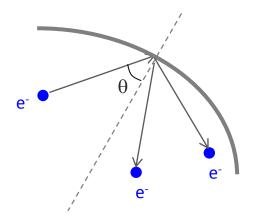




Generation of charged particles inside the vacuum chamber (primary, or seed, electrons)



- Acceleration of primary electrons in the beam field
- Secondary electron production when hitting the wall



Secondary Electron Yield (SEY) $\delta_{\text{max}} = 1.3$, normal incidence δ_{max} =1.3, θ =90° 2.5 2 SEY 1.5 0.5 200 300 0 100 400 500 600 700 E_p (eV)

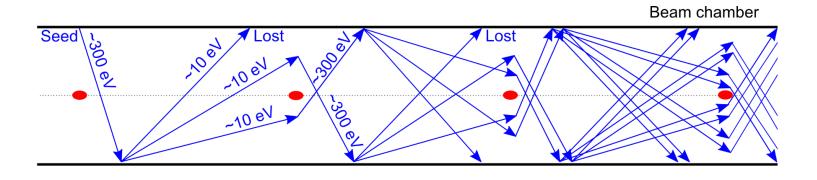




Generation of charged particles inside the vacuum chamber (primary, or seed, electrons)



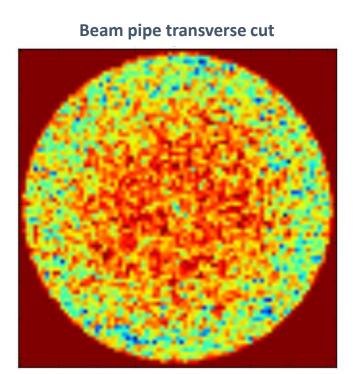
- Acceleration of primary electrons in the beam field
- Secondary electron production when hitting the wall
 - Avalanche electron multiplication if SEY > 1

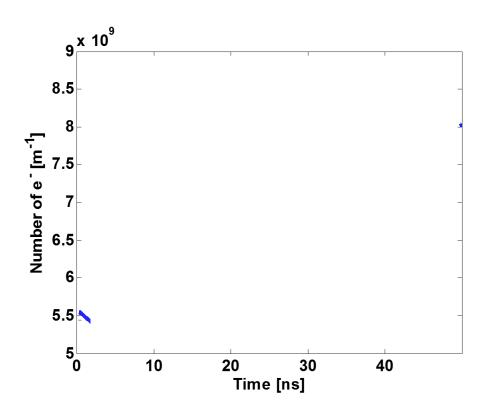


Bunch spacing (e.g. 25 ns)

Time





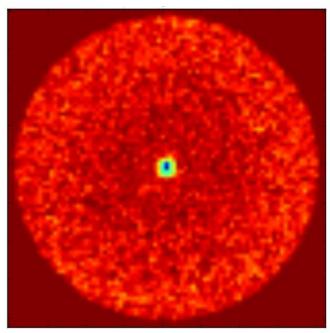


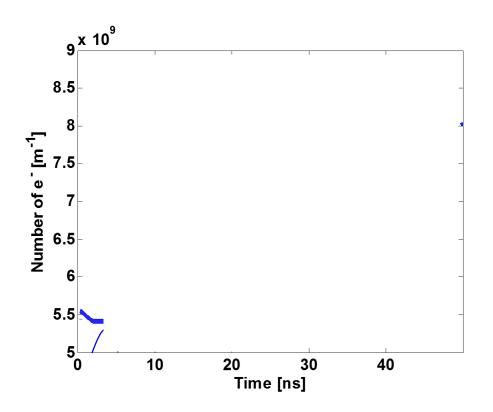
Assume an initial distribution of electrons (from any of the mechanisms discussed before)









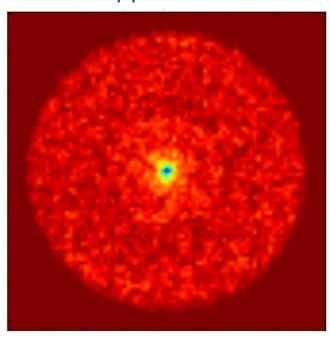


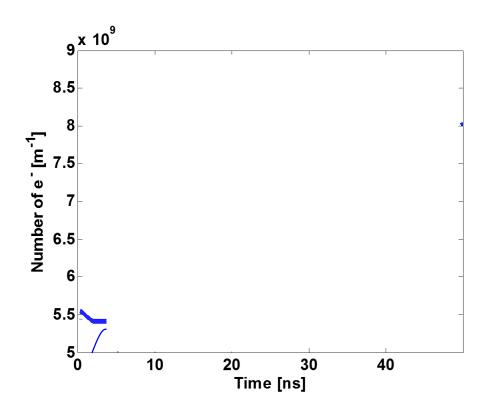
"Pinch" of electrons when bunch is passing







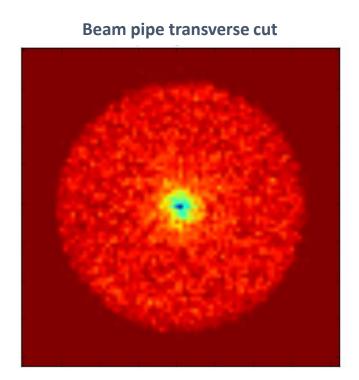


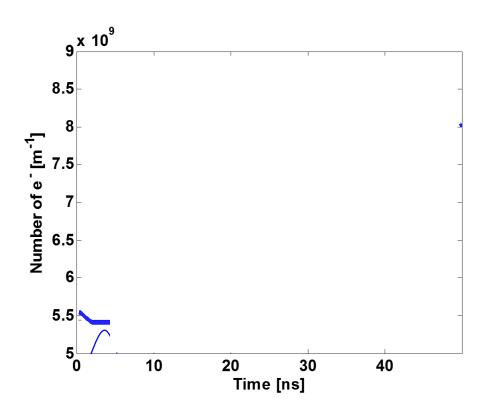


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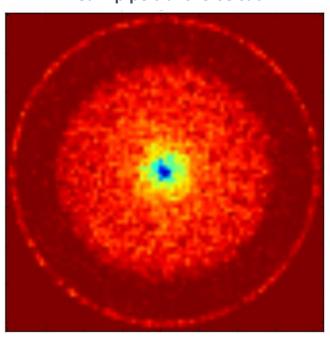


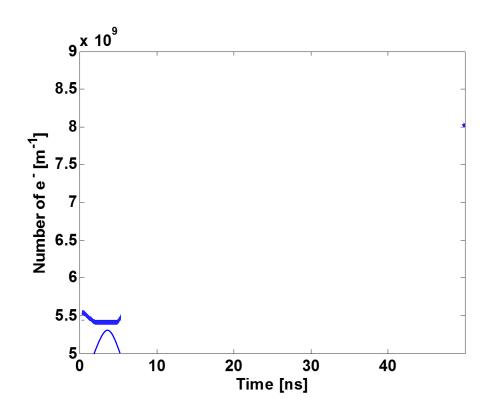
"Pinch" of electrons when bunch is passing









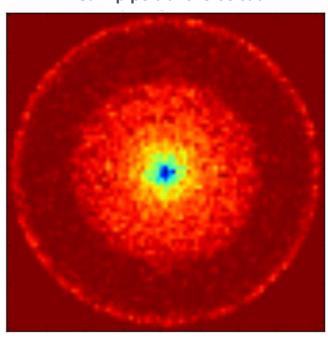


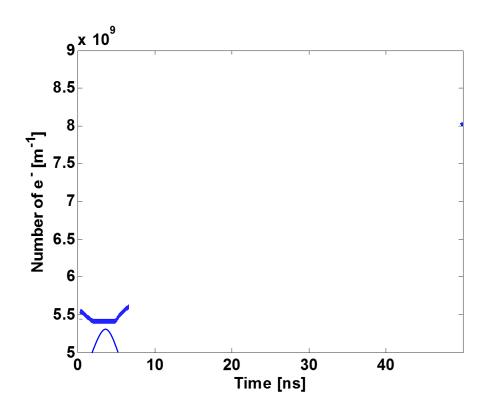
Few high energy (>100 eV) electrons reach the chamber wall already on the falling edge of the bunch and start producing secondaries







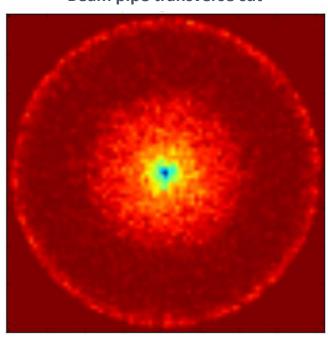


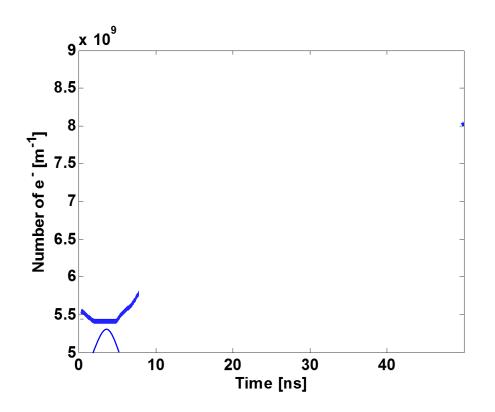








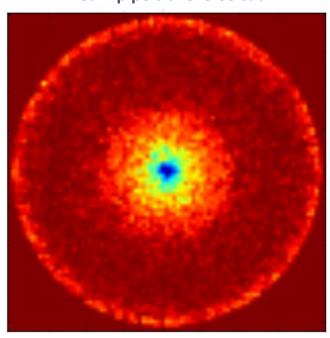


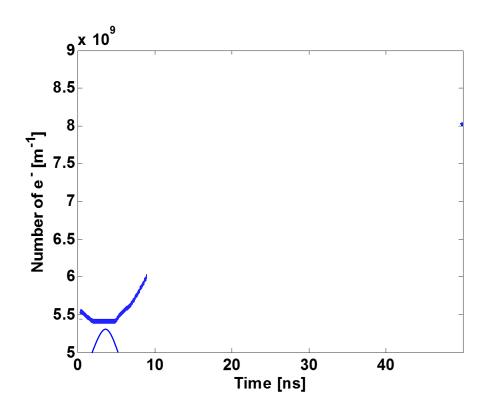








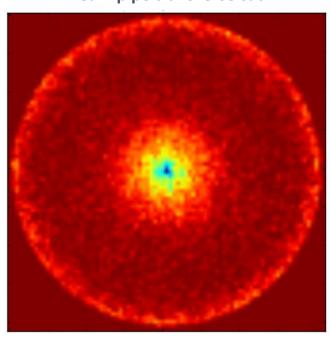


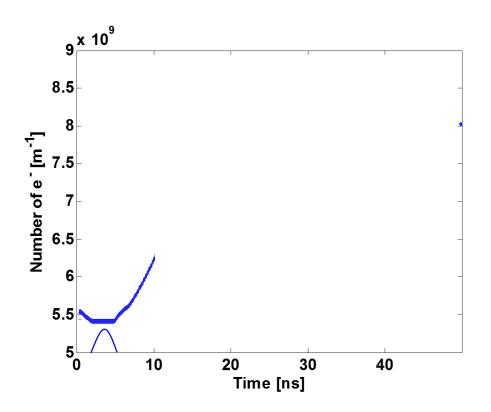








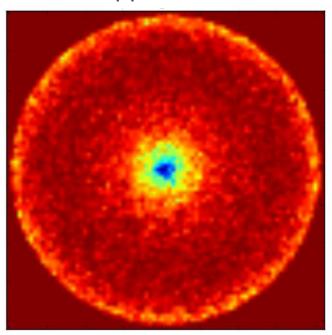


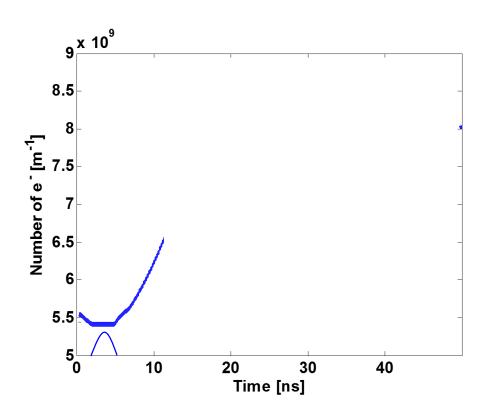










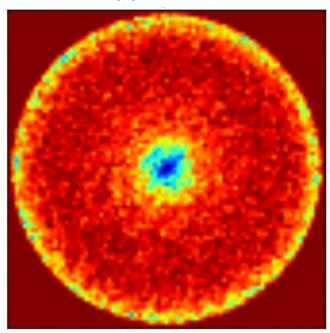


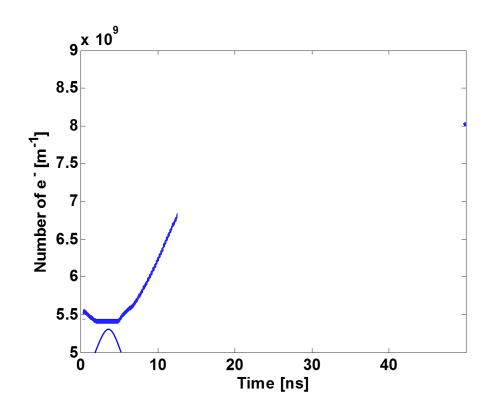
As they are produced, the emitted secondaries form a halo near the chamber wall because they have low energy (up to 10 eV)









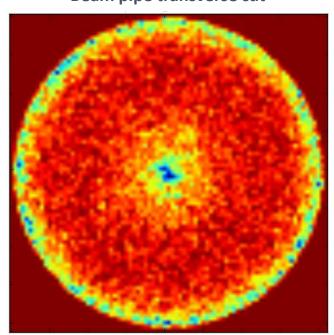


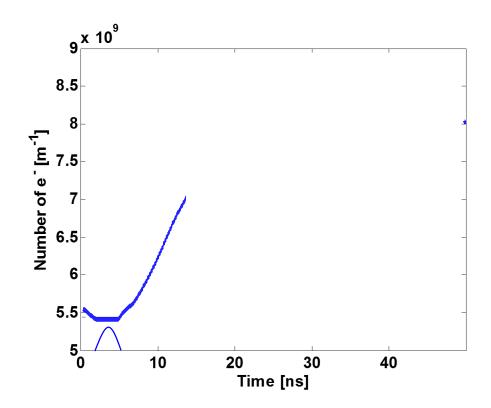
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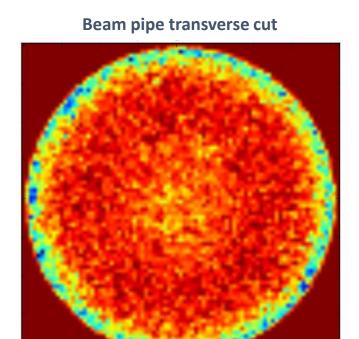


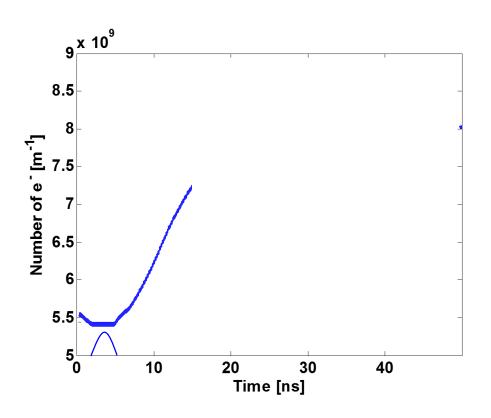


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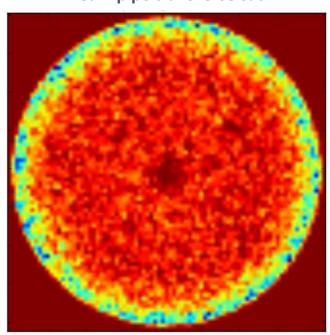


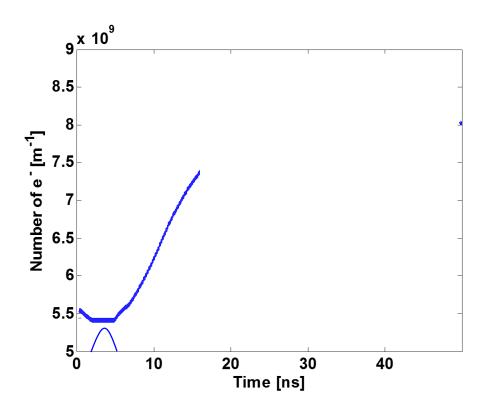
While the halo gets more and more populated, the center is gradually depleted









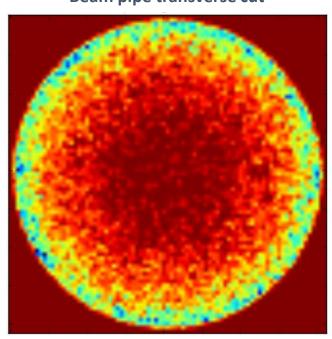


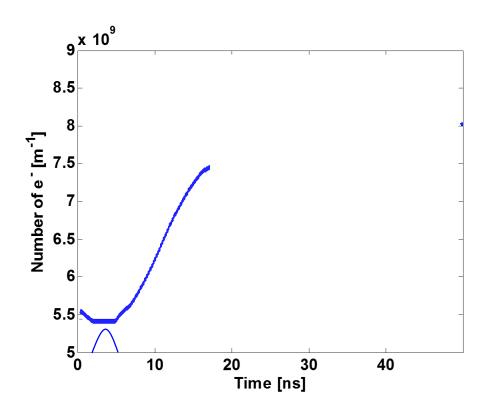
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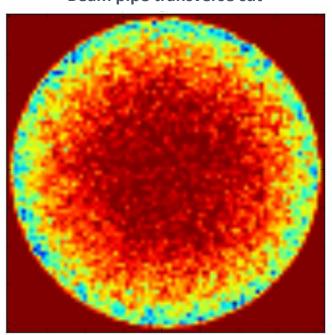


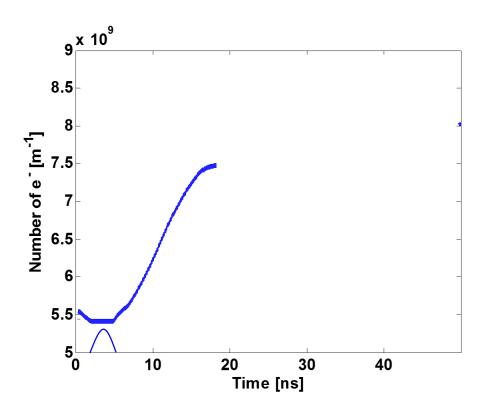
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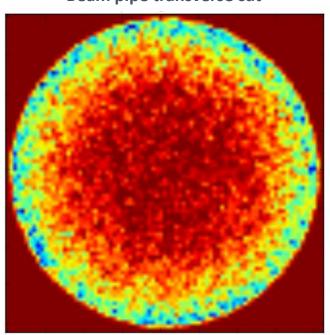


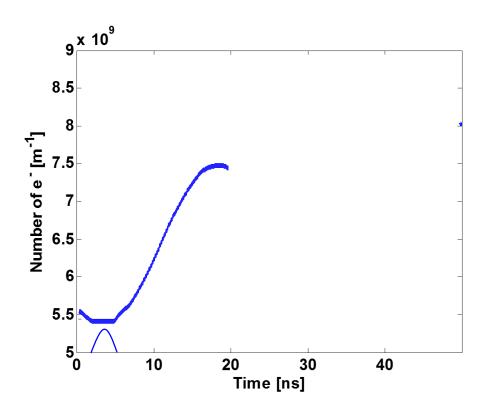
The center is strongly depleted No more secondaries are produced because there are no longer high energy electrons reaching the walls







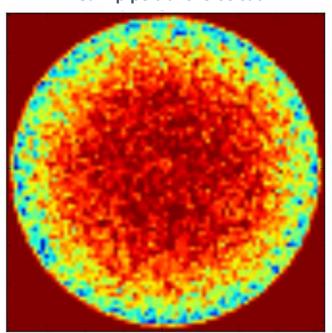


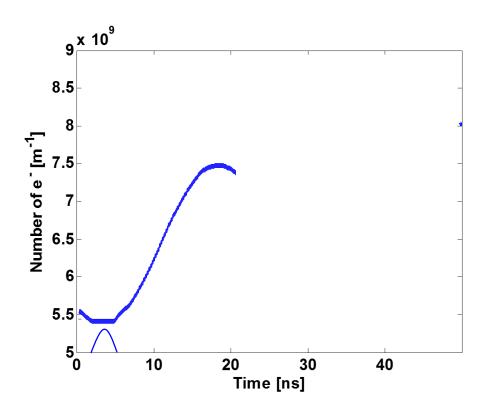








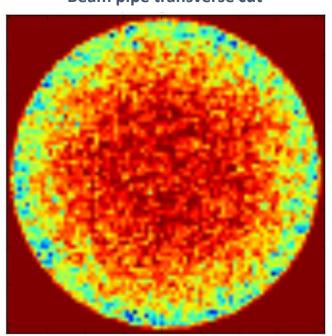


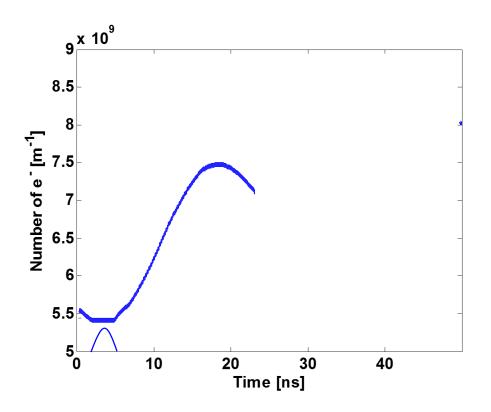








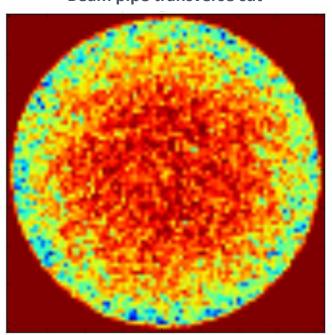


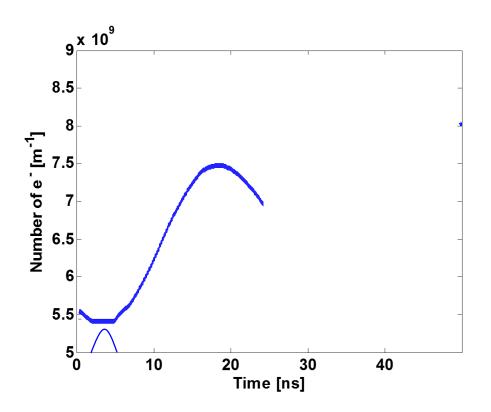






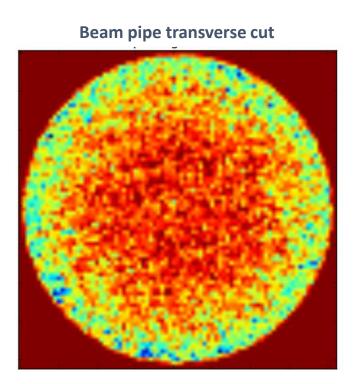


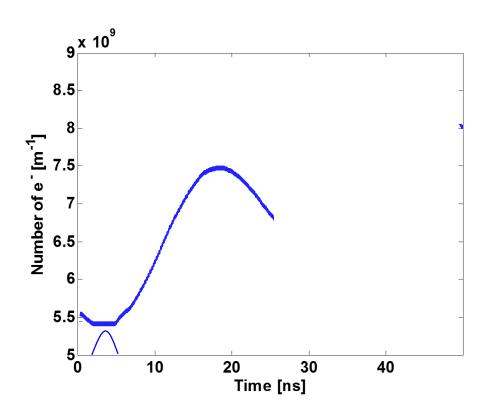








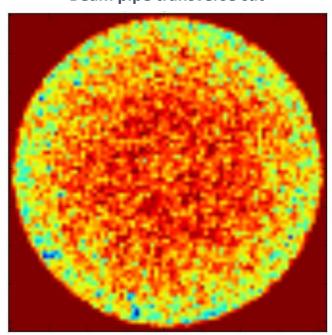


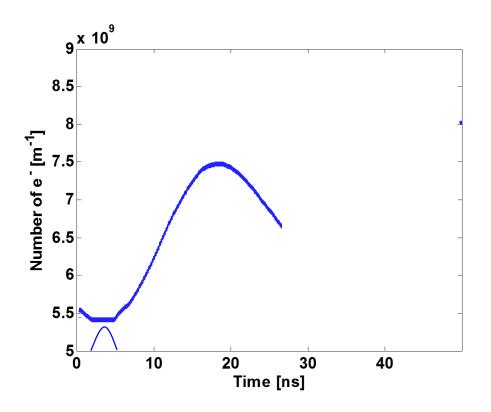






Beam pipe transverse cut



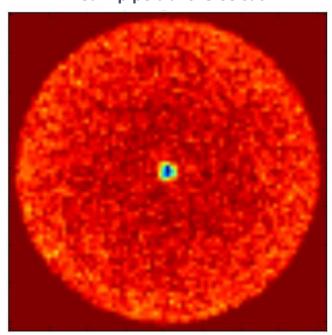


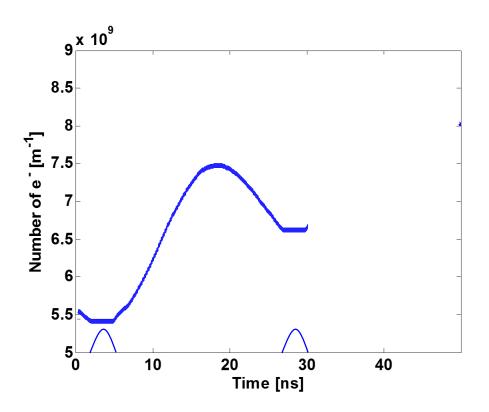
No more secondaries are produced because there are no longer high energy electrons reaching the walls Some low energy electrons are absorbed at the walls while the center gets repopulated





Beam pipe transverse cut



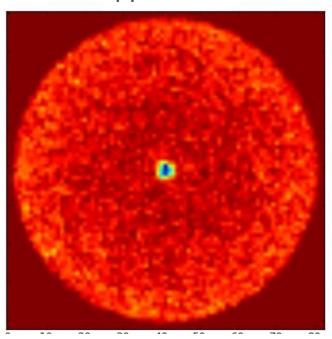


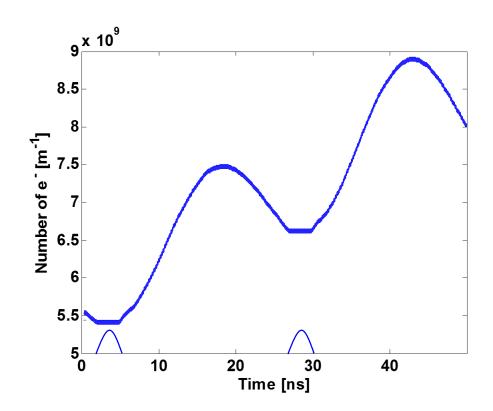
But then the next bunch comes, there is a new pinch and the whole process starts all over





Beam pipe transverse cut



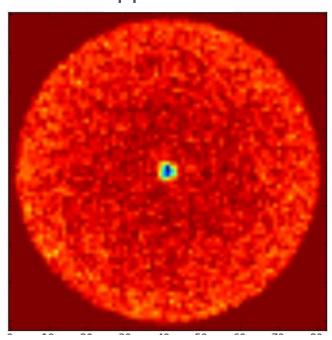


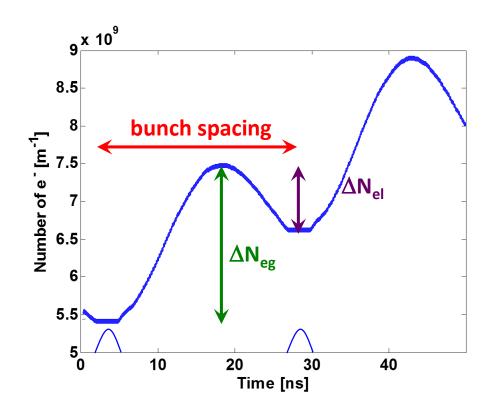
And it all repeats until the next bunch comes





Beam pipe transverse cut

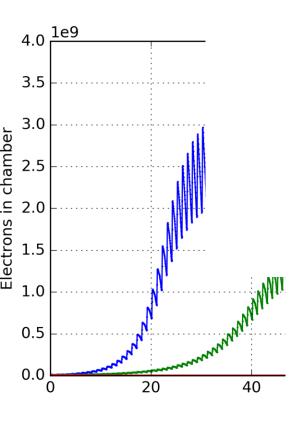




- Electrons generated (ΔN_{eg}) depend on bunch charge, chamber radius and surface SEY
- Electrons lost (ΔN_{el}) depend on chamber radius and probability of reflection at low energy
- Balance between the two depends on bunch spacing







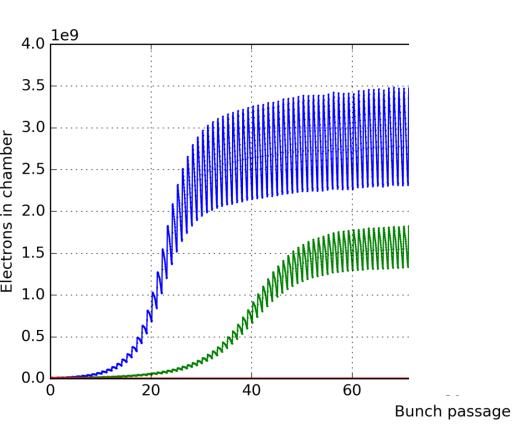
 $\delta_{
m max}$ (SEY) - 1.7 - 1.5

Bunch passage

 Bunch after bunch, the e-cloud grows exponentially (if SEY above a certain threshold value)





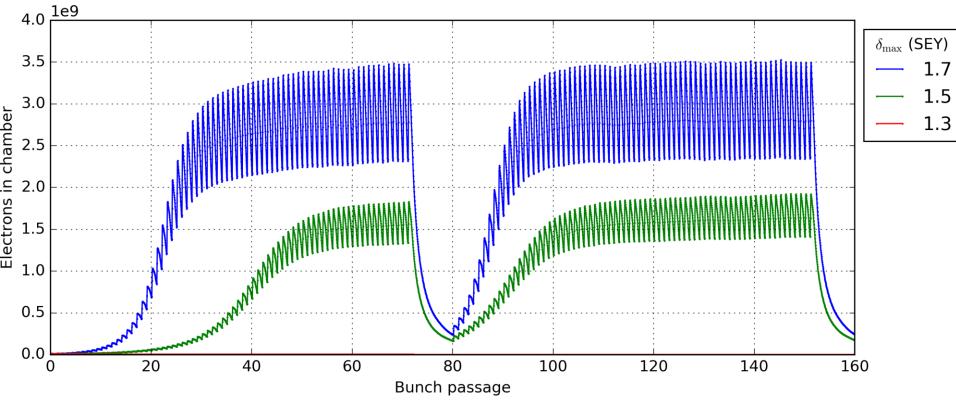


 $\delta_{
m max}$ (SEY) - 1.7 - 1.5

- Bunch after bunch, the e-cloud grows exponentially (if SEY above a certain threshold value)
- The exponential rise stops when the space charge of the electrons becomes significant → At this point electron generation and loss compensate each other







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- The exponential rise stops when the space charge of the electrons becomes significant → At this point electron generation and loss compensate each other
- The electron cloud decays in the gaps between trains







- We have learned that **electrons are generated** in the vacuum chamber of an accelerator when the beam passes.
- We have learned that
 - The number of electrons can grow because of secondary electron emission at the chamber walls
 - The process at some point saturates because of the electron cloud space charge
 - A significant electron density builds up in the machine while bunches are passing → electron cloud
- Once the machine operates with electron cloud, what do we observe?

Part 4: Electron cloud –

Build up and effects on beam dynamics

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- Scrubbing and other techniques of mitigation/suppression
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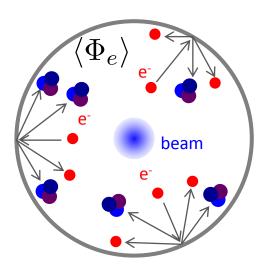


Electron cloud effects



The presence of an e-cloud inside an accelerator ring is revealed by several **typical signatures**

- ✓ Fast pressure rise, outgassing
- ✓ Additional heat load
- ✓ Baseline shift of the pick-up electrode signal
- ✓ Synchronous phase shift along the bunch train due to energy loss



$$\Delta P \propto \int \eta_e(E) \langle \Phi_e(E) \rangle dE$$

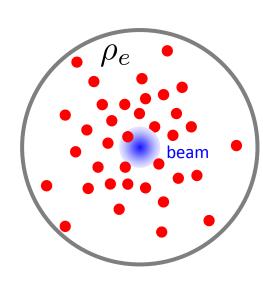
$$\Delta W = \int \langle \Phi_e(E) \rangle E dE$$

Electron cloud effects



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- ✓ Synchronous phase shift along the bunch train due to energy loss
- ✓ Tune shift along the bunch train
- ✓ Coherent instability
 - Single bunch effect affecting the last bunches of a train
 - Coupled bunch effect
- ✓ Poor beam lifetime and emittance growth



Machine

observables



Electron cloud effects



The presence of an e-cloud inside an accelerator ring is revealed by several **typical signatures**

- ✓ Fast pressure rise, outgassing
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- ✓ Baseline shift of the pick-up electrode signal
- ✓ Synchronous phase shift along the bunch train due to energy loss?
- ✓ Tune shift along the bunch train
- ✓ Coherent instability
 - Single bunch effect affecting the last bunches of a train
 - Coupled bunch effect
- Poor beam lifetime and emittance growth
- Active monitoring: signal on dedicated electron detectors (e.g. strip monitors) and retarding field analysers

Machine observables

Beam observables



Electron cloud effects: pressure rise



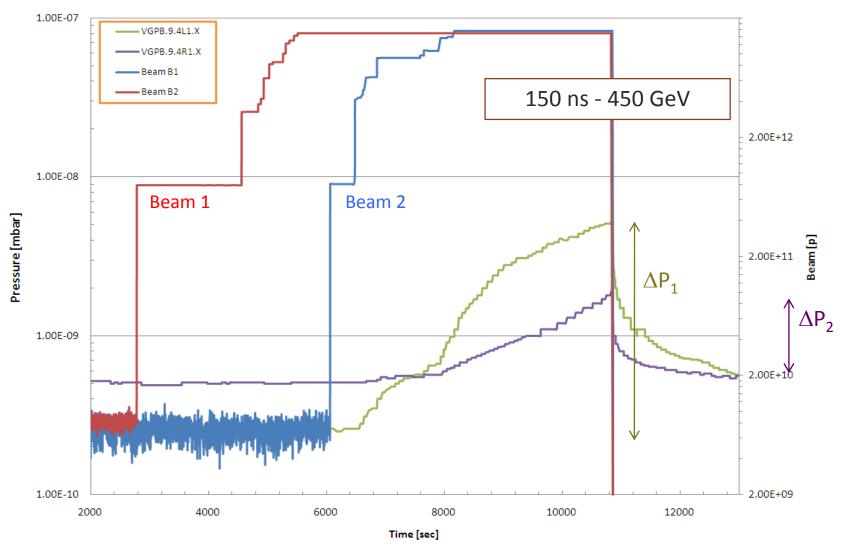
Early LHC operation

- Routine operation with 150 ns beams started in Summer 2010
- Electron cloud made its first appearance as a **pressure rise** in the common chamber in presence of both beams, i.e. for effectively lower bunch spacings



Electron cloud effects: pressure rise



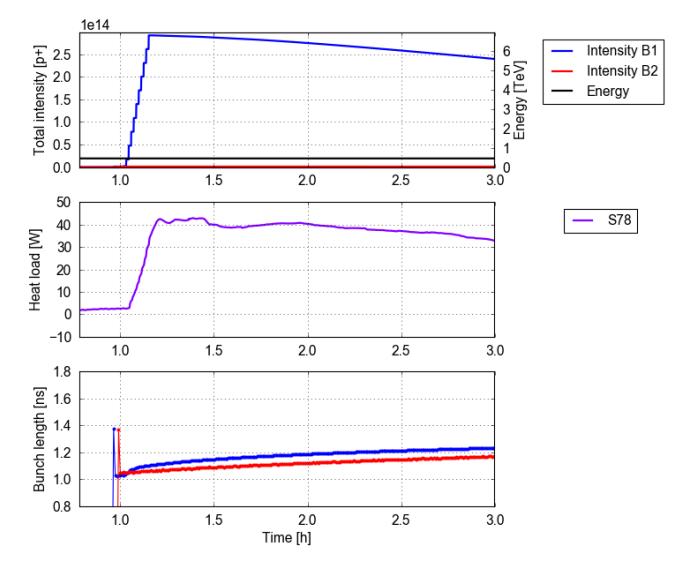




Electron cloud effects: heat load



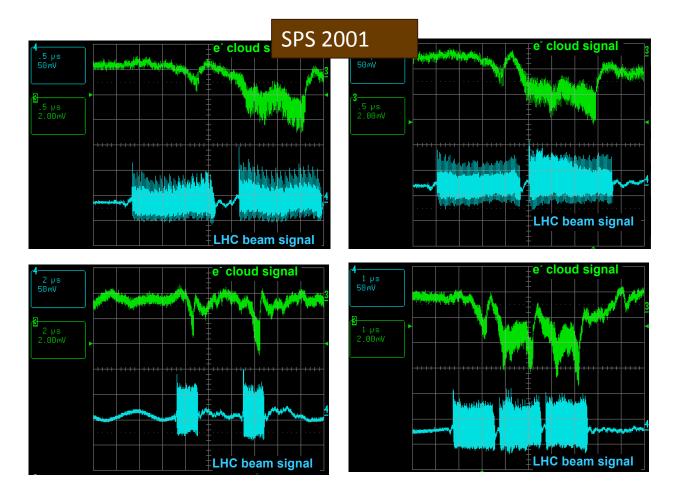
⇒ Heat load on the LHC beam screen of the cold arcs





Electron cloud effects: pick up signal





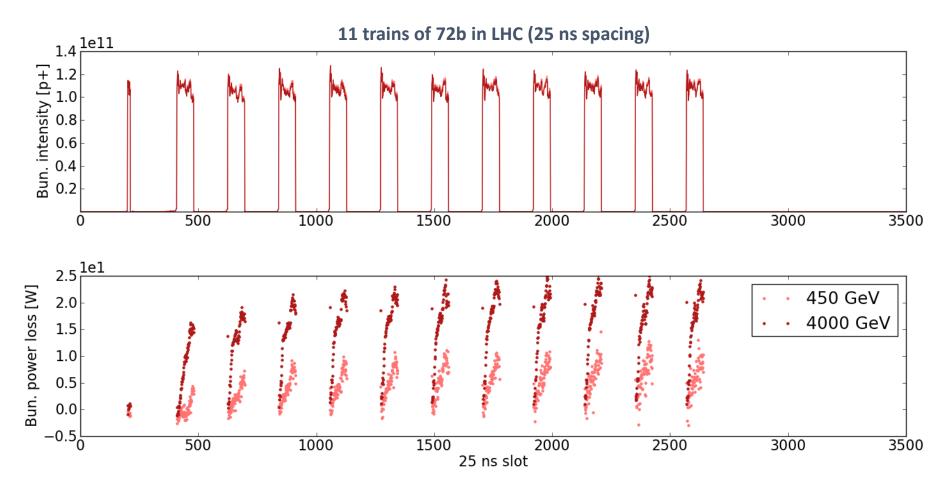
- The electron cloud signal first appeared in the SPS on the signal from a pick up as a shift of the baseline (depending on the charge collected by the electrodes)
- Correlation with train structure, length, gap were immediately apparent.



Electron cloud effects: stable phase shift



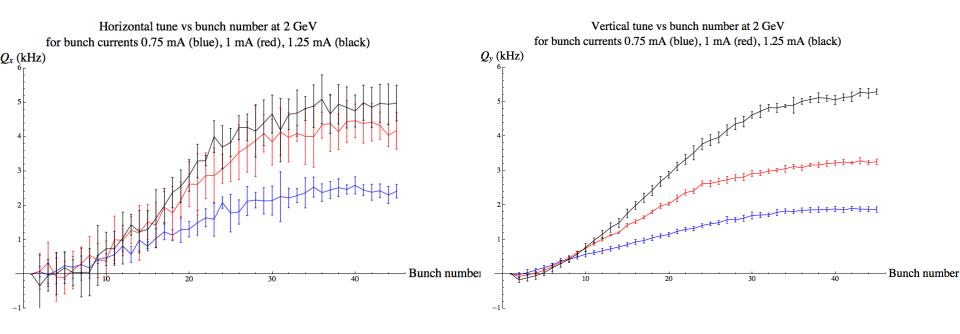
- ⇒ Bunch-by-bunch phase shift reveals the shape of the e-cloud build up
- ⇒ Larger electron cloud at 4 TeV is due to photoelectrons





Electron cloud effects: tune shift





- Horizontal and vertical tune shifts along a 46 bunch train in Cesr-TA (Cornell facility used for electron cloud studies) taken during a positron run
- Higher currents lead to stronger electron cloud.







- We have learned that **electron clouds** can build up in the vacuum chamber of an accelerator operating in a certain range of beam parameters.
- Electron clouds are associated to many detrimental effects, like pressure rise, additional heat load, tune and stable phase shift, beam degradation through instability and emittance growth
- How can we avoid or cure it?

Part 4: Electron cloud –

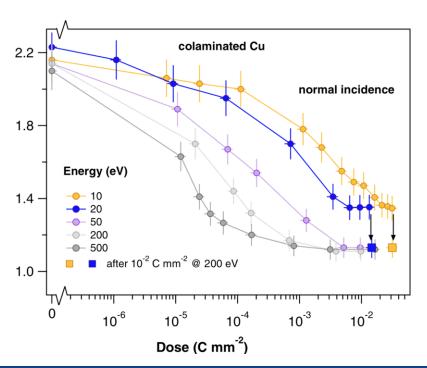
Build up and effects on beam dynamics

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- Fortunately, the SEY of a surface is not a fixed property but it becomes lower under electron bombardment (scrubbing)
- Laboratory measurements show that
 - SEY decreases quickly at the beginning of the process, then slows down
 - Electrons with different energies have different 'scrubbing efficiency'
 - The 'final' value of SEY depends on material, e⁻ energy, temperature, vacuum composition, more?

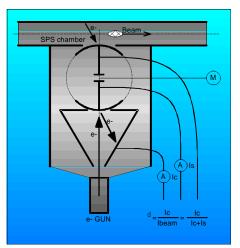




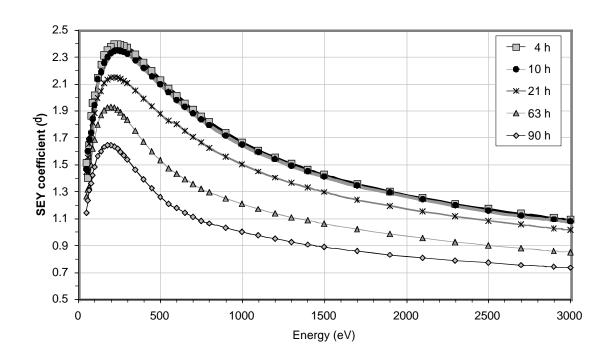


Beam-induced scrubbing

 Has been measured directly at the SPS with a Stainless Steel rotatable sample exposed to the beam or to SEY measurement device (2004)



Schematic view of the in-situ SEY detector installed in the SPS

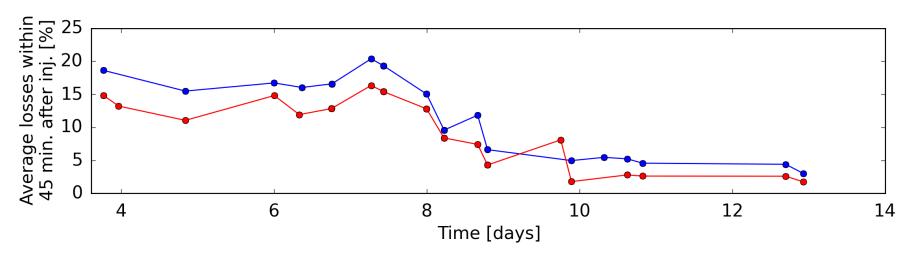






Beam-induced scrubbing

- Has been measured directly at the SPS with a Stainless Steel rotatable sample exposed to the beam or to SEY measurement device (2004)
- Is revealed by improving accelerator conditions over time, e.g. decrease of pressure rise, heat load, stable phase shift, general improvement of beam quality (lower losses, less emittance growth)



Example: Reduction of losses in LHC over 9 days of scrubbing (no clear reduction visible in first phase due to increasing length of the injected trains)





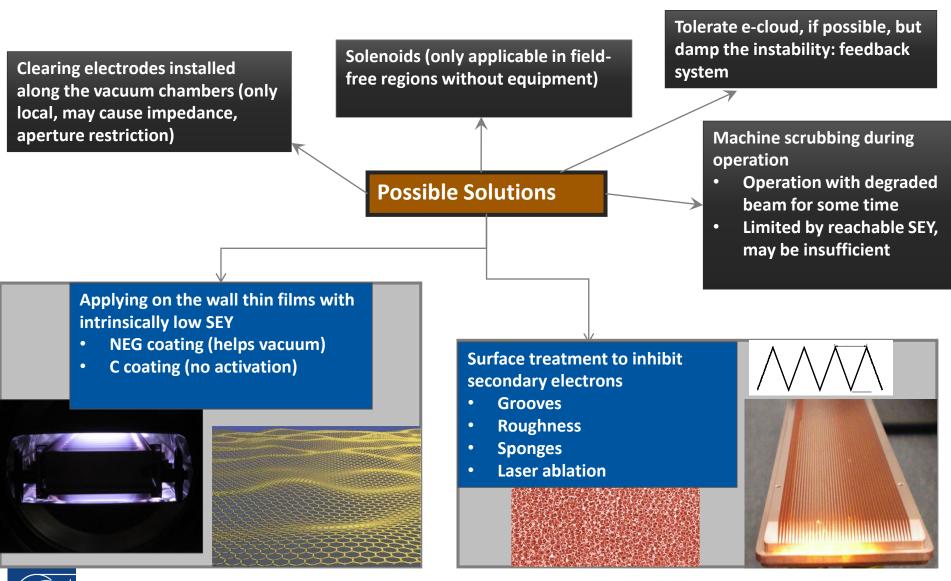
Beam-induced scrubbing

- o Has been measured directly at the SPS with a Stainless Steel rotatable sample exposed to the beam or to SEY measurement device (2004)
- Is revealed by improving accelerator conditions over time, e.g. decrease of pressure rise, heat load, stable phase shift, general improvement of beam quality (lower losses, less emittance growth)
- ⇒ Many accelerators rely nowadays on beam induced scrubbing to reach their desired performance!



Mitigation/suppression techniques









- We have learned that **electron clouds** can build up in the vacuum chamber of an accelerator operating in a certain range of beam parameters.
- They are the origin of many detrimental effects, like pressure rise, additional heat load, beam degradation through instability and emittance growth.
- They can be self-healing through beam induced scrubbing or they can be avoided by design (surface coating/treatment, solenoids, clearing electrodes).
- What is the mechanism through which an electron cloud degrades the beam?

Part 4: Electron cloud –

Build up and effects on beam dynamics

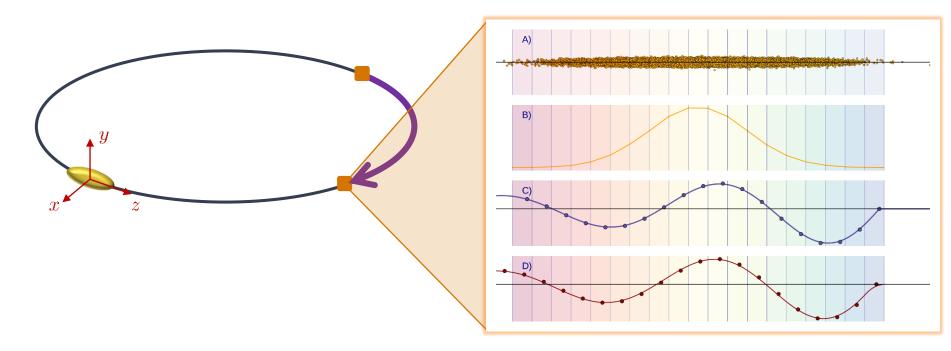
- Electron cloud build up
 - Electron production and multiplication
 - Observation in accelerator rings
- Scrubbing and other techniques of mitigation/suppression
- E-cloud induced instabilities and incoherent effects



Accelerator beam system - wakefields



 Our first 'real' collective interaction from impedances



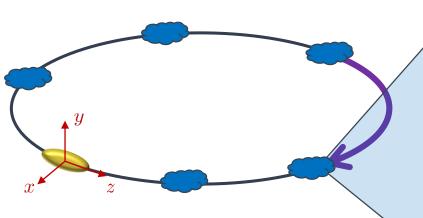


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Accelerator beam system – electron clouds



 Two stream collective interaction – much more complicated



For this we need to solve

$$\Delta \phi(x,y)_{p^{+}} = -\frac{\rho_{p^{+}}(x,y)}{\varepsilon_{0}}$$

$$\Delta \phi(x,y)_{e^{-}} = -\frac{\rho_{e^{-}}(x,y)}{\varepsilon_0}$$

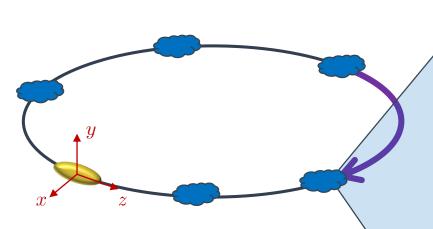
and apply the corresponding kicks to the cloud and the beam



Electron clouds in a drift section



 Two stream collective interaction – much more complicated

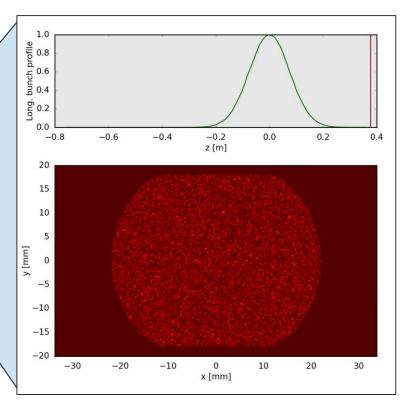


For this we need to solve

$$\Delta \phi(x,y)_{p^{+}} = -\frac{\rho_{p^{+}}(x,y)}{\varepsilon_{0}}$$

$$\Delta \phi(x,y)_{e^-} = -\frac{\rho_{e^-}(x,y)}{\varepsilon_0}$$

and apply the corresponding kicks to the cloud and the beam



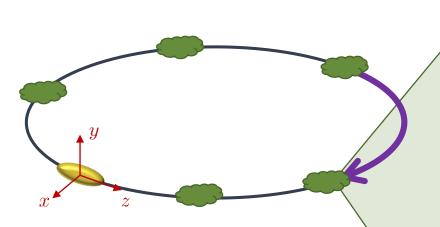
 Beam passage leads to a pinch of the cloud which in turn acts back on the beam – differently each turn



Electron clouds in a bending magnet



 Two stream collective interaction – much more complicated

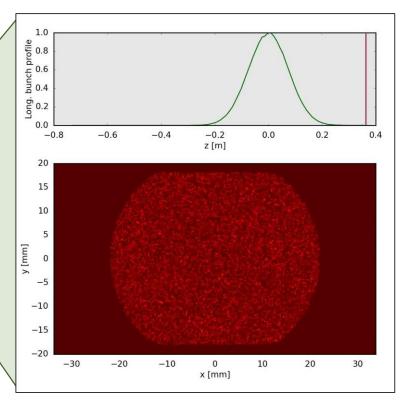


For this we need to solve

$$\Delta \phi(x,y)_{p^{+}} = -\frac{\rho_{p^{+}}(x,y)}{\varepsilon_{0}}$$

$$\Delta \phi(x,y)_{e^-} = -\frac{\rho_{e^-}(x,y)}{\varepsilon_0}$$

and apply the corresponding kicks to the cloud and the beam



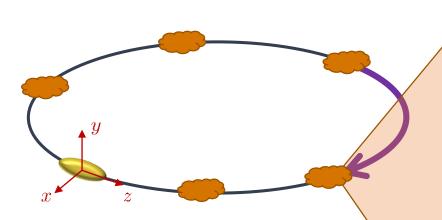
 Beam passage leads to a pinch of the cloud which in turn acts back on the beam – differently each turn



Electron clouds in a quadrupole magnet



 Two stream collective interaction – much more complicated



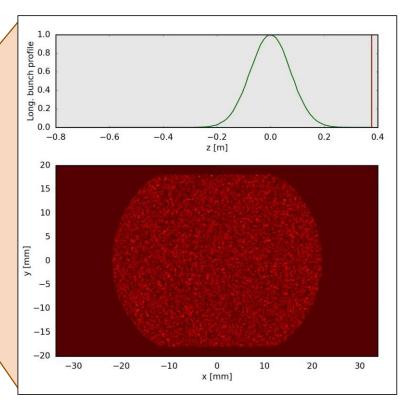
For this we need to solve

$$\Delta \phi(x,y)_{p^{+}} = -\frac{\rho_{p^{+}}(x,y)}{\varepsilon_{0}}$$

$$\Delta \phi(x,y) = -\frac{\rho_{e^{-}}(x,y)}{\varepsilon_{0}}$$

$$\Delta \phi(x,y)_{e^{-}} = -\frac{\rho_{e^{-}}(x,y)}{\varepsilon_0}$$

and apply the corresponding kicks to the cloud and the beam



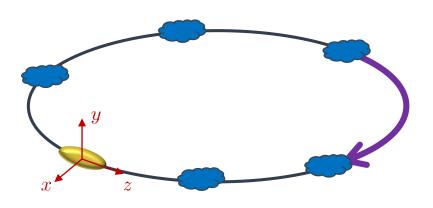
 Beam passage leads to a pinch of the cloud which in turn acts back on the beam – differently each turn



Accelerator beam system – e-cloud



$$\mathcal{M} = \begin{pmatrix} \sqrt{\beta_1} & 0 \\ -\frac{\alpha_1}{\sqrt{\beta_1}} & \frac{1}{\sqrt{\beta_1}} \end{pmatrix} \begin{pmatrix} \cos(\Delta\mu_i) & \sin(\Delta\mu_i) \\ -\sin(\Delta\mu_i) & \cos(\Delta\mu_i) \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{\beta_0}} & 0 \\ \frac{\alpha_0}{\sqrt{\beta_0}} & \sqrt{\beta_0} \end{pmatrix}$$



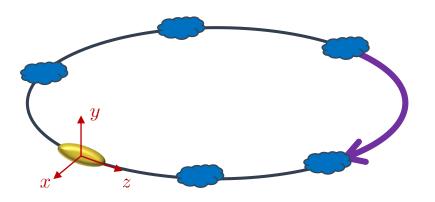
- Basic loop of tracking with electron clouds:
 - Transport beam along segment to interaction point

Accelerator beam system – e-cloud



$$\mathcal{M} = \begin{pmatrix} \sqrt{\beta_1} & 0 \\ -\frac{\alpha_1}{\sqrt{\beta_1}} & \frac{1}{\sqrt{\beta_1}} \end{pmatrix} \begin{pmatrix} \cos(\Delta\mu_i) & \sin(\Delta\mu_i) \\ -\sin(\Delta\mu_i) & \cos(\Delta\mu_i) \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{\beta_0}} & 0 \\ \frac{\alpha_0}{\sqrt{\beta_0}} & \sqrt{\beta_0} \end{pmatrix}$$

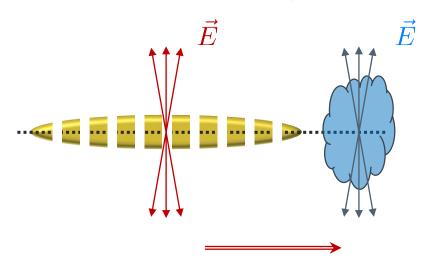
$$\Delta \vec{x}'[i] = -\frac{e^2}{m\gamma\beta^2c^2}\vec{E}_{e^-}C \qquad \qquad \Delta \dot{\vec{x}} = -\frac{e}{m}\left(\vec{E}_{p^+}[i] + \frac{\dot{\vec{x}}\times\vec{B}}{c}\right)\Delta t$$
 Particles in/fields from slice i



- Basic loop of tracking with electron clouds:
 - Transport beam along segment to interaction point
 - Apply e-cloud kick
 → get fields from PIC step

E-cloud beam system



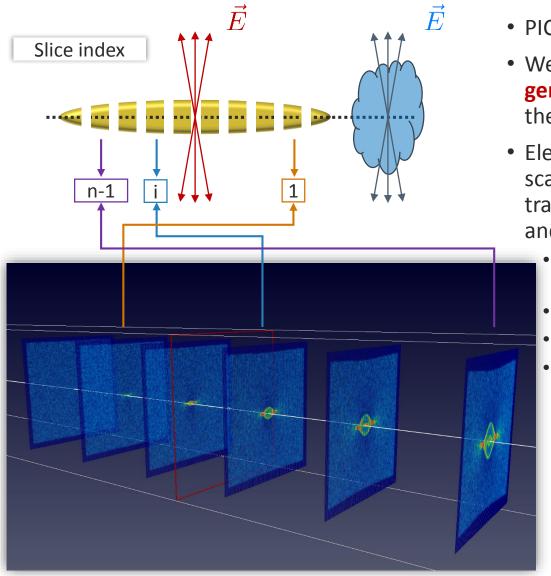


- PIC stands for Particle-In-Cell
- We use this method to compute fields generated by particles to solve e.g. the Poisson equation
- Electron motion occurs at the time scale of a slice of a bunch length → track single slices through the e-cloud and apply integrated kicks
 - Compute electric fields from one slice and from e-cloud
 - Apply kicks to protons and electrons
 - Push electrons by one slice length
 - Track next slice through e-cloud



E-cloud beam system





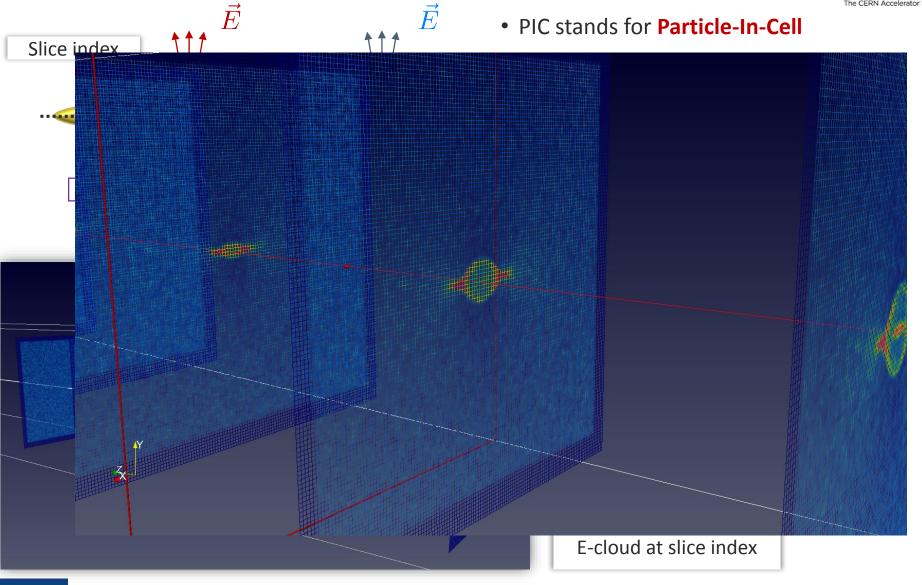
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 - Push electrons by one slice length
 - Track next slice through e-cloud

E-cloud at slice index



E-cloud beam system

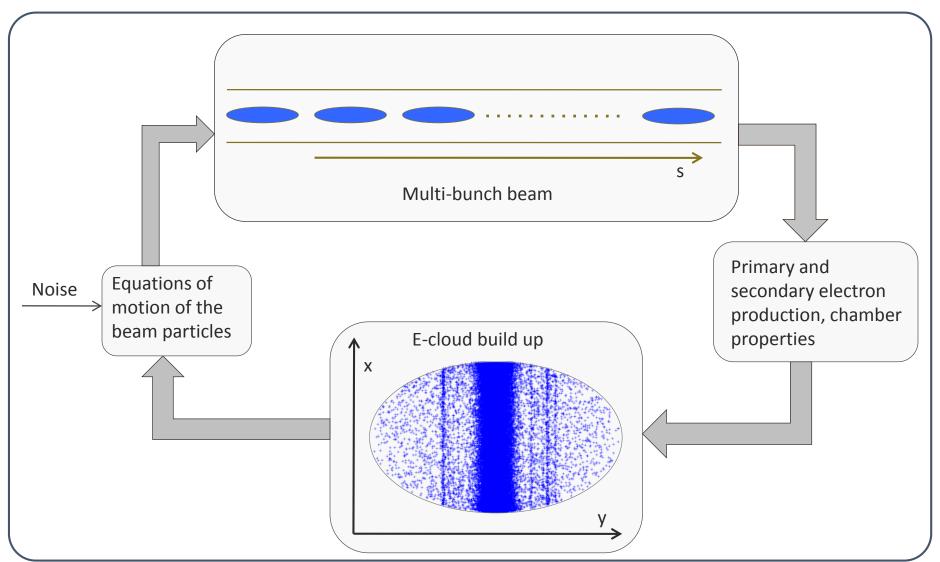






Numerical model of electron cloud effects

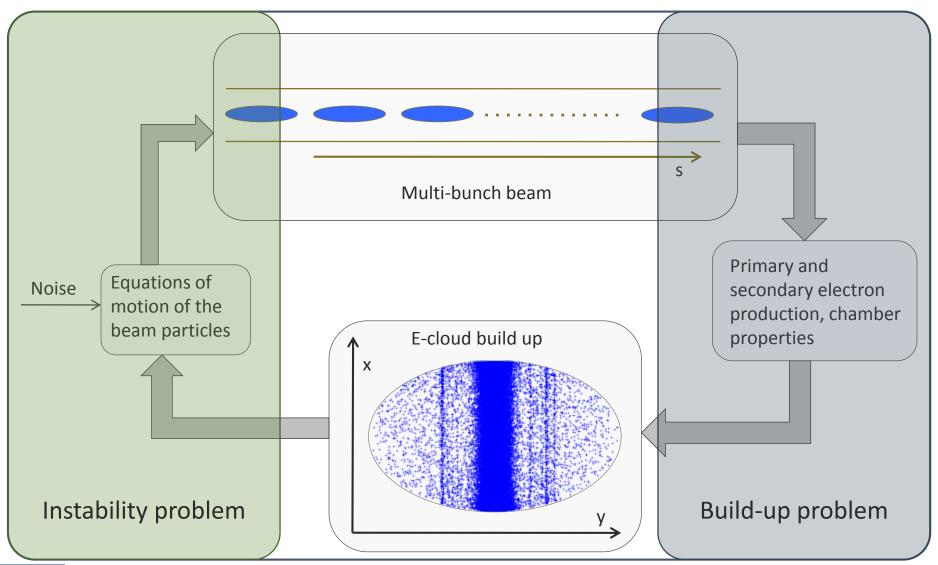






Numerical model of electron cloud effects







Electron cloud induced instabilities

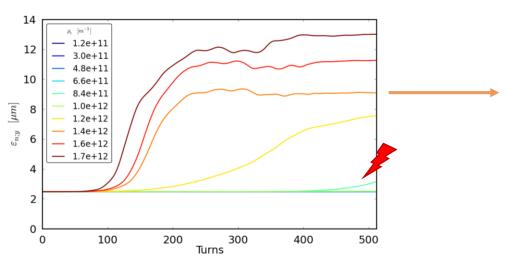


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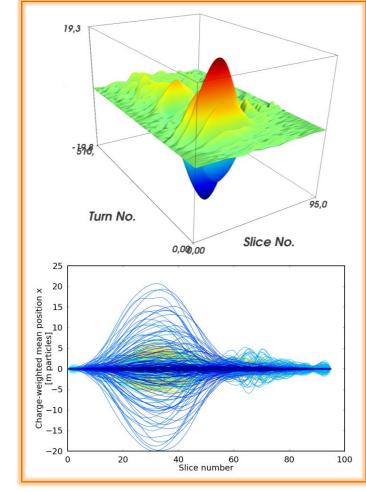
Typical e-cloud simulation try to identify the e-cloud central density threshold for an instability

• Scans in the central density are performed until an **exponential growth** can be observed in the

emittance



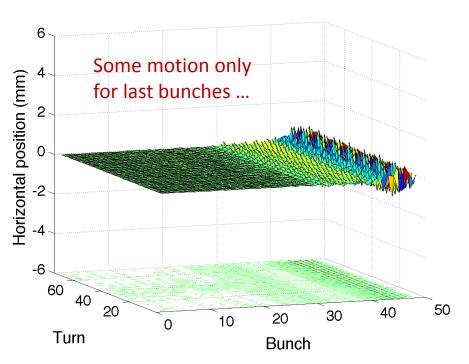
- Coherent instabilities occur when a certain central cloud density threshold is breached
- This leads to coherent intra bunch motion which grows exponentially
- A consequence is emittance blow-up and losses

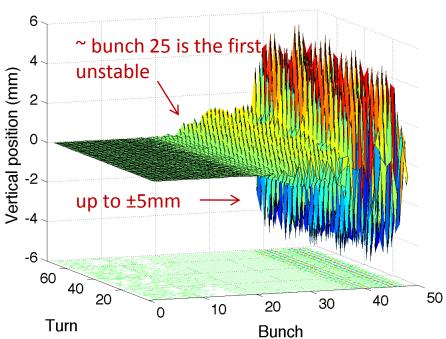


Ex. of coherent e-cloud effects in the LHC



- First injection of 48 bunches of 25 ns beam into the LHC in 2011
- Beam was dumped twice due to a violent instability in the vertical plane, causing losses above the interlock threshold

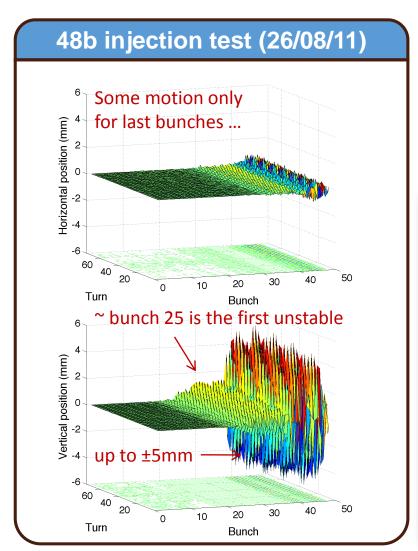


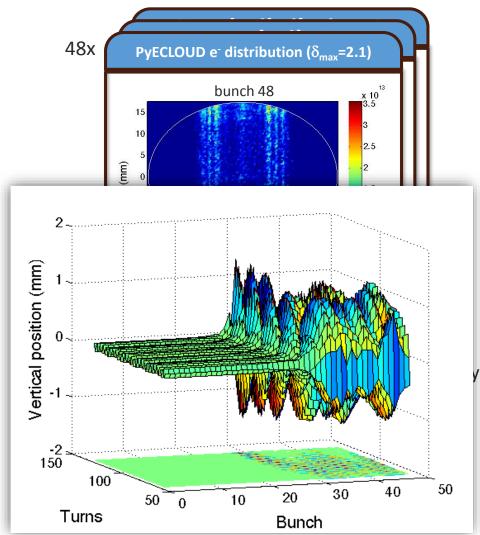




Ex. of coherent e-cloud effects in the LHC













- We have learned that **electron clouds** can build up in the vacuum chamber of an accelerator operating in a certain range of beam parameters.
- We have seen some of the **detrimental effects** of electron clouds on the machine.
- We have seen methods on how to **suppress or mitigate the build up** of electron clouds.
- We have seen how we can **conceptually model** the beam-electron cloud interaction and some **examples of electron cloud induced instabilities**.

Part 4: Electron cloud –

Build up and effects on beam dynamics

- Electron cloud build up
 - Electron production and multiplication
 - Observation in accelerator rings
- Scrubbing and other techniques of mitigation/suppression
- E-cloud induced instabilities and incoherent effects





The End





PIC solvers in brief



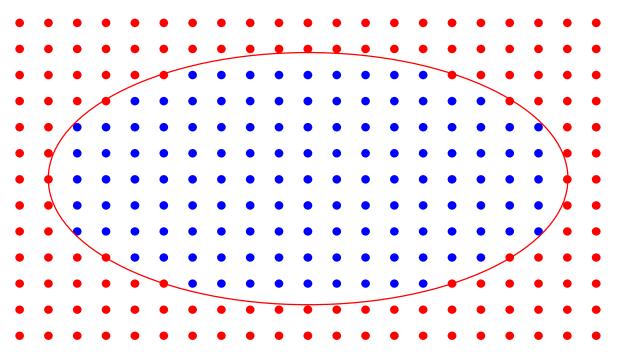
- In many of our codes, Particle in Cell (PIC) algorithms are used to compute the electric field generated by a set of charged particles in a set of discrete points (can be the locations of the particles themselves, or of another set of particles)
- The solution typically consists of 4 stages:
 - 1. Charge scatter from macroparticles (MPs) to grid (reduction of macroparticles)
 - 2. Calculation of the electrostatic potential at the nodes
 - 3. Calculation of the electric field at the nodes (gradient evaluation)
 - 4. Field gather from grid to MPs



PIC solvers in brief



- The solution typically consists of 4 stages:
 - 1. Charge scatter from macroparticles (MPs) to grid (reduction of macroparticles)
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Internal nodes

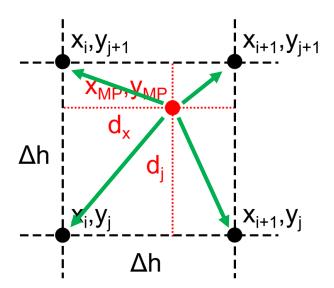
External nodes (optional)



Uniform square grid



- The solution typically consists of 4 stages:
 - 1. Charge scatter from macroparticles (MPs) to grid (reduction of macroparticles)
 - 2. Calculation of the electrostatic potential at the nodes
 - 3. Calculation of the electric field at the nodes (gradient evaluation)
 - 4. Field gather from grid to MPs



$$\rho_{i,j} = \rho_{i,j} + \frac{q \, n_{\text{MP}}}{\Delta h} \left(1 - \frac{d_x}{\Delta h} \right) \left(1 - \frac{d_y}{\Delta h} \right)$$

$$\rho_{i+1,j} = \rho_{i+1,j} + \frac{q \, n_{\text{MP}}}{\Delta h} \left(\frac{d_x}{\Delta h} \right) \left(1 - \frac{d_y}{\Delta h} \right)$$

$$\rho_{i,j+1} = \rho_{i,j+1} + \frac{q \, n_{\text{MP}}}{\Delta h} \left(1 - \frac{d_x}{\Delta h} \right) \left(\frac{d_y}{\Delta h} \right)$$

$$\rho_{i+1,j+1} = \rho_{i+1,j+1} + \frac{q \, n_{\text{MP}}}{\Delta h} \left(\frac{d_x}{\Delta h} \right) \left(\frac{d_y}{\Delta h} \right)$$





- The solution typically consists of 4 stages:
 - 1. Charge scatter from macroparticles (MPs) to grid (reduction of macroparticles)
 - 2. Calculation of the electrostatic potential at the nodes
 - 3. Calculation of the electric field at the nodes (gradient evaluation)
 - 4. Field gather from grid to MPs

$$\int \nabla^2 \phi(x,y) = -\frac{\rho(x,y)}{\varepsilon_0}$$

Boundary conditions (e.g., perfectly conducting, open, periodic)

- Different numerical approaches exist to solve these types of equations each with its own advantages and drawbacks:
 - Open space FFT solver (explicit, very fast but open boundaries)
 - Rectangular boundary FFT solver (explicit, very fast but only rectangular boundaries)
 - Finite Difference implicit Poisson solver (arbitrary chamber shape, sparse matrix, possibility to use Shortley Weller boundary refinement, KLU fast routines, computationally more demanding)
 - Dual or multi-grid in combination with direct or iterative solvers





- The solution typically consists of 4 stages:
 - 1. Charge scatter from macroparticles (MPs) to grid (reduction of macroparticles)
 - 2. Calculation of the electrostatic potential at the nodes
 - 3. Calculation of the electric field at the nodes (gradient evaluation)
 - 4. Field gather from grid to MPs

$$\mathbf{E} = -\nabla \phi$$



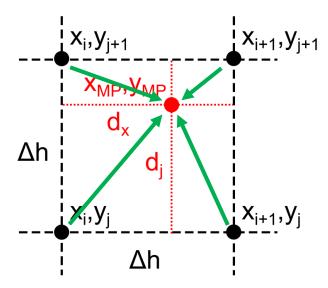
$$(E_x)_{i,j} = -\frac{\phi_{i+1,j} - \phi_{i-1,j}}{2\Delta h}$$

$$(E_y)_{i,j} = -\frac{\phi_{i,j+1} - \phi_{i,j-1}}{2\Delta h}$$





- The solution typically consists of 4 stages:
 - 1. Charge scatter from macroparticles (MPs) to grid (reduction of macroparticles)
 - 2. Calculation of the electrostatic potential at the nodes
 - 3. Calculation of the electric field at the nodes (gradient evaluation)
 - 4. Field gather from grid to MPs



$$\mathbf{E}(x_{\mathrm{MP}}, y_{\mathrm{MP}}) =$$

$$\mathbf{E}_{i,j} \left(1 - \frac{d_x}{\Delta h} \right) \left(1 - \frac{d_y}{\Delta h} \right) + \mathbf{E}_{i+1,j} \left(\frac{d_x}{\Delta h} \right) \left(1 - \frac{d_y}{\Delta h} \right)$$

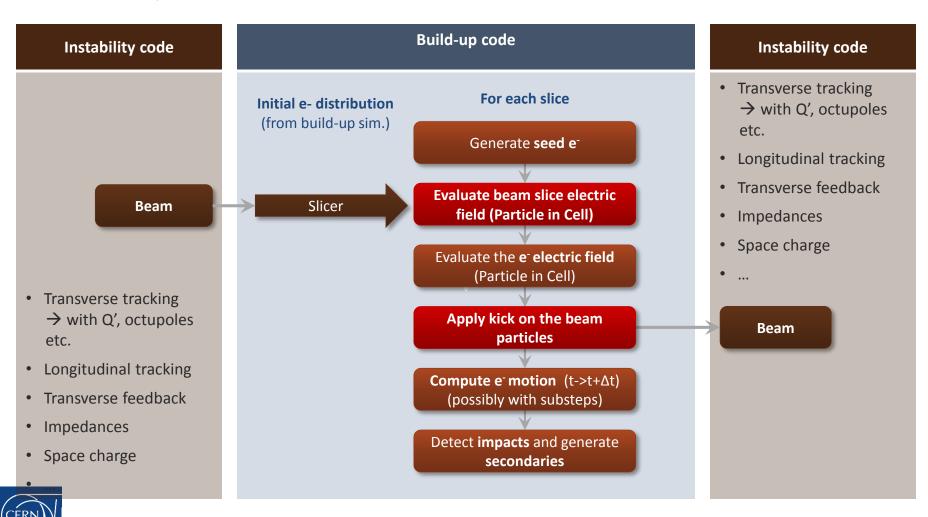
$$+ \mathbf{E}_{i,j+1} \left(1 - \frac{d_x}{\Delta h} \right) \left(\frac{d_y}{\Delta h} \right) + \mathbf{E}_{i+1,j+1} \left(\frac{d_x}{\Delta h} \right) \left(\frac{d_y}{\Delta h} \right)$$



Numerical model of electron cloud effects



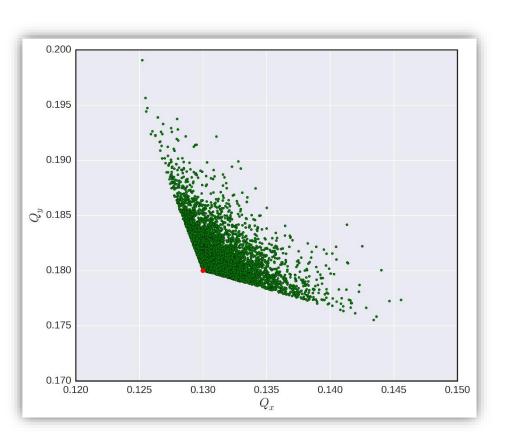
- A self-consistent treatment requires the combination of an instability and a build-up code
- Becomes easily possible with modular structure and good design of codes (e.g. object orientation)



<u>Legend</u>: From instability code – From build-up code – Interaction between the two codes

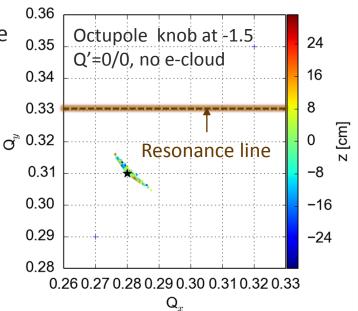


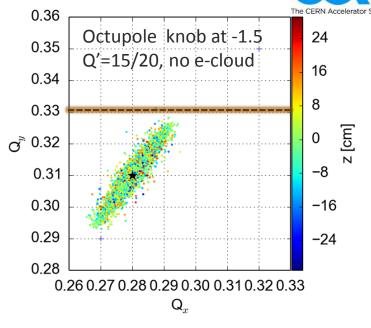
Remember tune footprint from octupoles in Part I



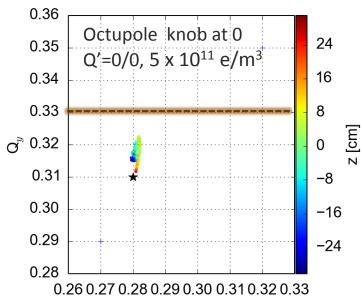


 Macroparticle simulations allow to obtain tune footprints from all effects separated





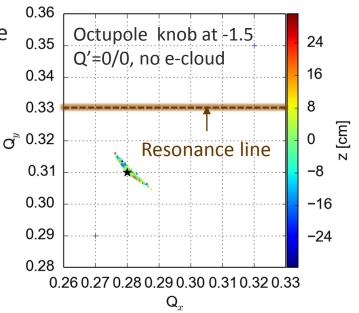
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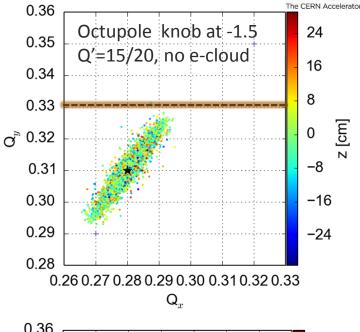


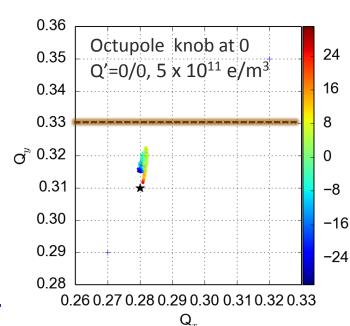


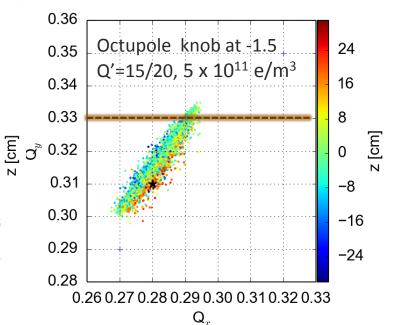
 ${f Q}_x$ 09/09/2017 Numerical Methods I - Kevin Li

- Macroparticle simulations allow to obtain tune footprint from all effects separated
- ... as well as from all effects combined





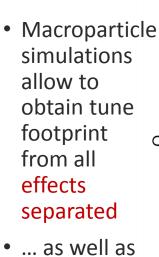




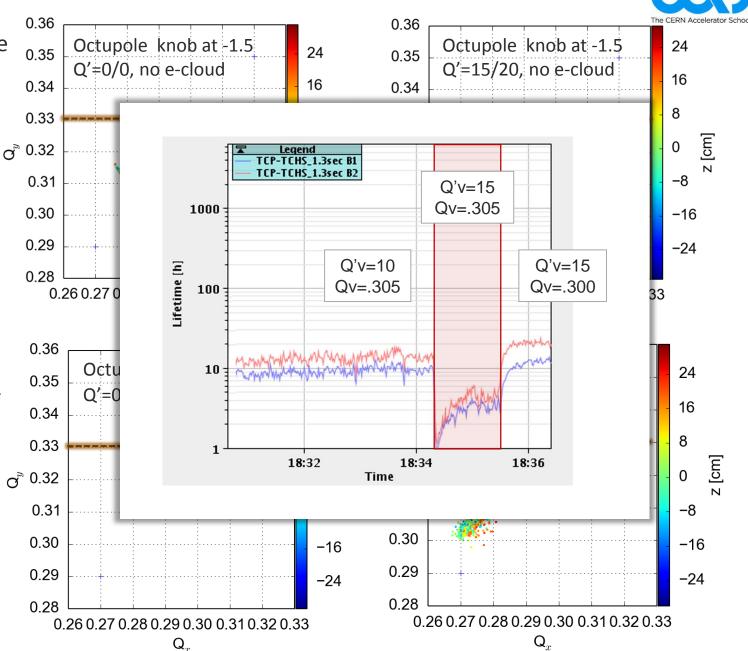
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- ... as well as from all effects combined
- ... to identify the source of incoherent losses in the LHC





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Backup - wakefields



Electron production



Generation of charged particles inside the vacuum chamber (primary, or seed, electrons)



- Acceleration of primary electrons in the beam field
- Secondary electron production when hitting the wall
 - Avalanche electron multiplication if SEY > 1



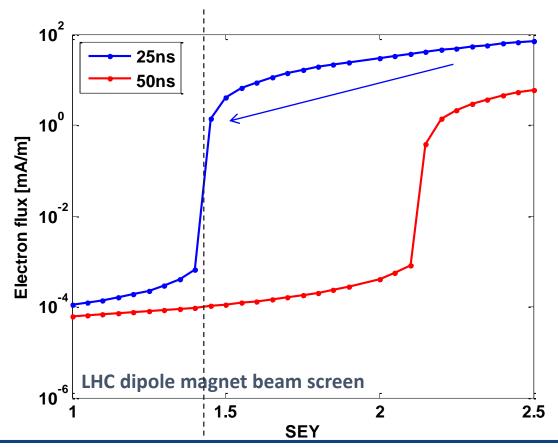
After the passage of several bunches, the electron distribution inside the chamber reaches a dynamic steady state (electron cloud)



Surface scrubbing



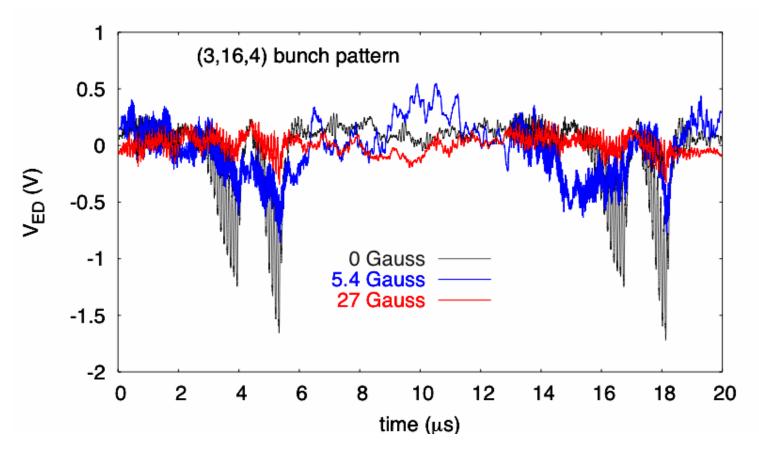
- If an accelerator can be run in **e-cloud regime**, scrubbing is expected to naturally occur (beam induced scrubbing)
 - Fortunately beam dynamics knobs exist to preserve beam stability, although lifetime might be poor in presence of significant e-cloud
 - Dedicated scrubbing runs can be used to lower the SEY





Examples: solenoids (I)



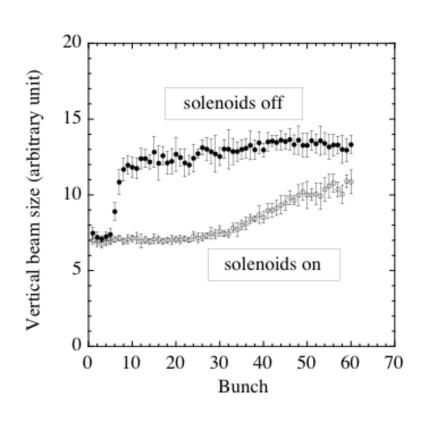


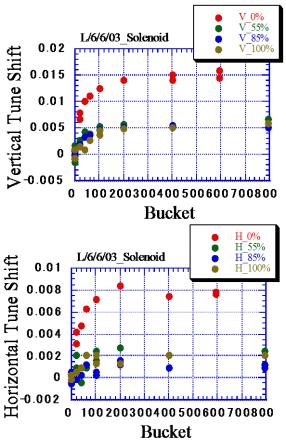
- Beneficial effect of the solenoids measured at the heavy ion collider RHIC (BNL)
- By changing the intensity of the magnetic field, the electron cloud was seen to be efficiently suppressed in a region equipped with an electron detector.



Examples: solenoids (II)





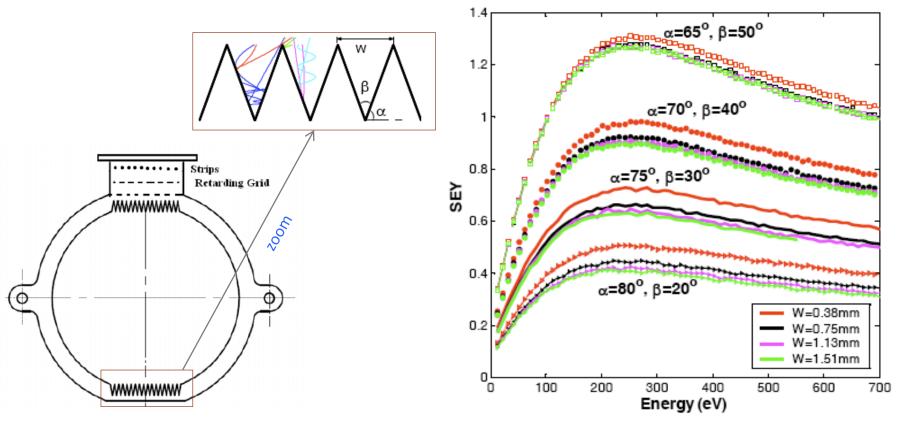


- Beneficial effect of solenoids measured at the LER of KEKB
- Drastic reduction of the beam size blow up as well as the tune shift along the batch



Examples: grooves



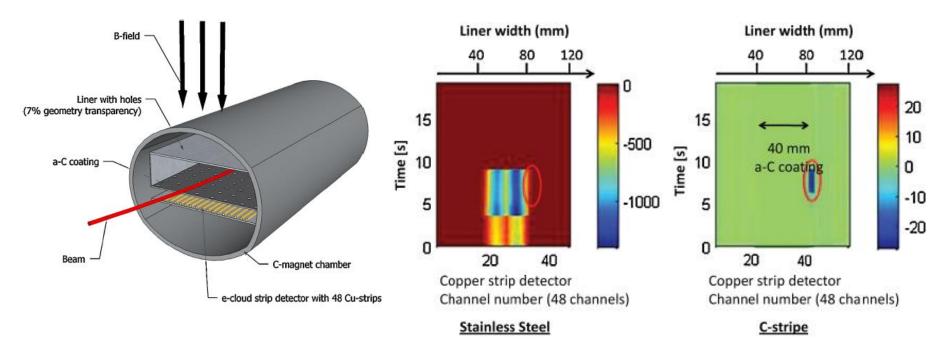


- To reduce the effective SEY, the inner surface of the beam pipe can be grooved, so that emitted electrons remain trapped
- Figure shows the effective SEY as a function of the groove angle and period, for a sample having δ_{max} =1.74 at E_{max}=330 eV



Examples: carbon coating





- To reduce the effective SEY, the inner surface of the beam pipe can be coated with amorphous carbon (a-C)
- It is possible to reach values of δ_{max} below 1, measured in the laboratory and also verified by measurements at an electron cloud detector in the SPS







- We have learned that **electron clouds** can build up in the vacuum chamber of an accelerator operating in a certain range of beam parameters.
- They are the origin of many detrimental effects, like pressure rise, additional heat load, beam degradation through instability and emittance growth
- They can be self-healing through beam induced scrubbing or they can be avoided by design (surface coating/treatment, solenoids, clearing electrodes)
- What is the mechanism through which an electron cloud degrades the beam?

Part 4: Electron cloud –

Build up and effects on beam dynamics

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- Scrubbing and other techniques of mitigation/suppression
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