

Measurements, Statistics and Errors

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- Autocorrelation
- Noise Propagation
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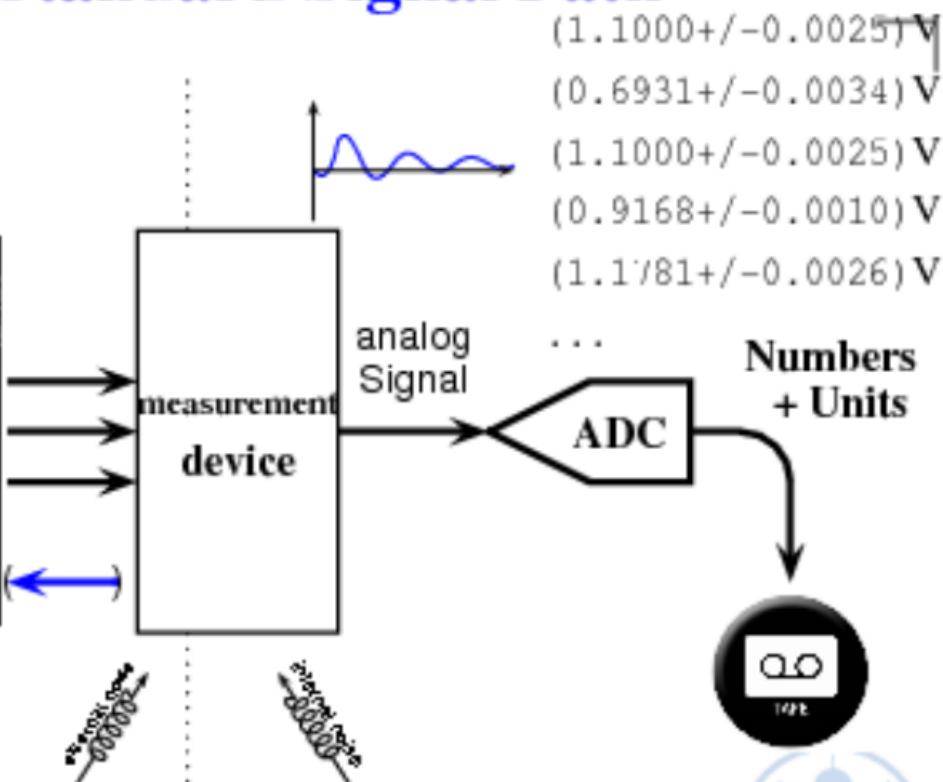
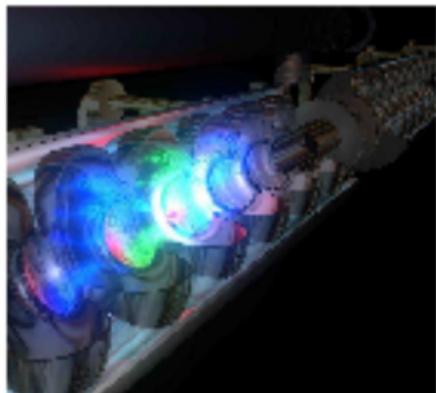
5. Applications



1. Measurement

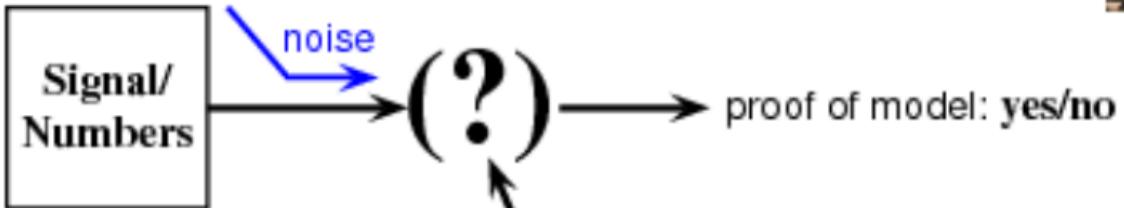


The Standard Signal Path



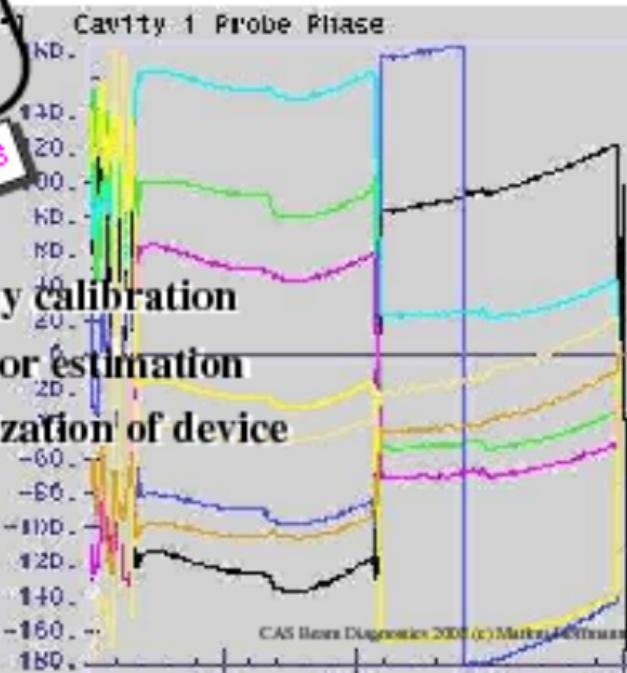
- The documentation about the measurement device → significance numbers + units.

Gnosis



Measurement means:

- Abstraction from raw data by calibration
 - Significance of values by Error estimation
 - Reduction of noise by optimization of device
 - Model + simulate the signals
 - proof or falsify model
 - adapt model



Signals

Signal =

Voltage as a function of time

$$U(t)$$

Definition:

- Time: $t \in \mathbb{R}$ (sometimes $\in \mathbb{R}_0^+$)
- Amplitude: $s(t) \in \mathbb{R}$
- Power: $s^2(t) \in \mathbb{R}_0^+$ (constants are renormalized to 1, e.g. $P = \frac{U^2}{R}$ or $P = I^2 R$.)

discrete	continuous
analog	digital
causal	non-causal

Energy-Signal:

$$\int_{-\infty}^{\infty} s^2(t) dt < \infty$$

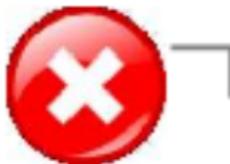
(Problems with $\sin()$, $\cos()$, $\text{rect}()$ and stochastic/noise signals)

Power-Signal:

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} s^2(t) dt < \infty$$



Errors



1. systematic (deterministic) error
2. unsystematic (statistic) error (noise)

systematic error

- due to characteristics of the measurement device.
(ADC/DAC: *offset, gain, linearity-errors*)
- Improved by improvements of apparatus.
- Limits: Quantum mechanics

statistic error

- unforeseen fluctuations, random, stochastic, noise
- it is possible to estimate the extend
- can be reduced by statistical methods (averaging), multiple measurements

accuracy

Definition is context dependent: accuracy of 100 devices can be a matter of precision !

precision



2. Noise



Noise

acoustic \Rightarrow electronics



Noise sources:

- Brownian movement of charges (*thermal noise*)
- Variations of the number of charges involved in conduction (*Shottky noise/shot noise, flicker noise*)
- quantum effects (*zero point fluctuations*)

Noise classes:

- **white noise** \rightarrow **flat** spectrum
- **pink noise** \rightarrow low-passed spectrum
- **blue noise** \rightarrow high-passed spectrum

Amplitude density distribution:

probability density



\rightarrow Gauß

if many indep.
sources contribute

\rightarrow stochastic signal



Sensor Noise

Sources:

- Thermal noise (Johnson-/Nyquist-noise)
- Shot noise (Shottky-Noise)
(quants of the elementary charge)

Spectrum:

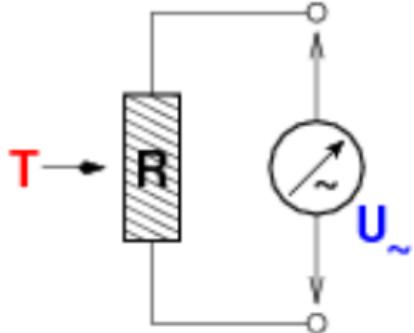


- $1/f^\alpha$ noise
(e.g. flicker noise)



Thermal Noise

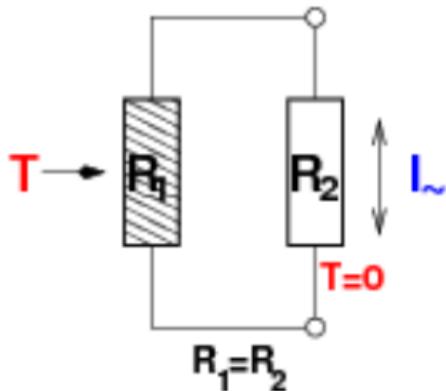
Brownian motion of charge carriers



effective value measurement bandwidth

Noise Voltage: $U_{\text{rms}} = \sqrt{4kT \cdot R \cdot B}$

Note: Thermal noise is only emitted by structures with electromagnetic losses (resistors, R), which also absorb power!



$$I_{\text{rms}} = \sqrt{4kT \frac{B}{R}}$$

Reileigh-Jeans approximation, high T , small ν

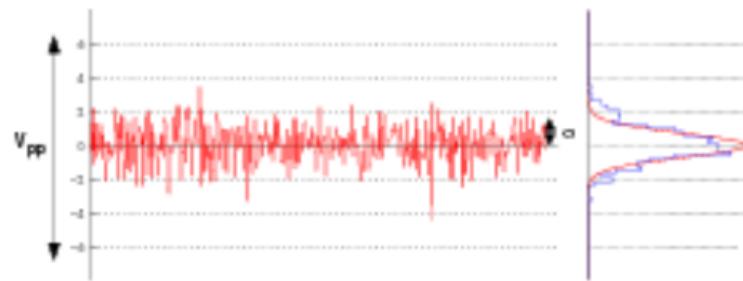
independent from R !

Noise Power: $P_n = \frac{1}{4} U_{\text{rms}} I_{\text{rms}} = kTB$

$$\sim \frac{U^2}{I^2}$$

received from R_1 (power matching)

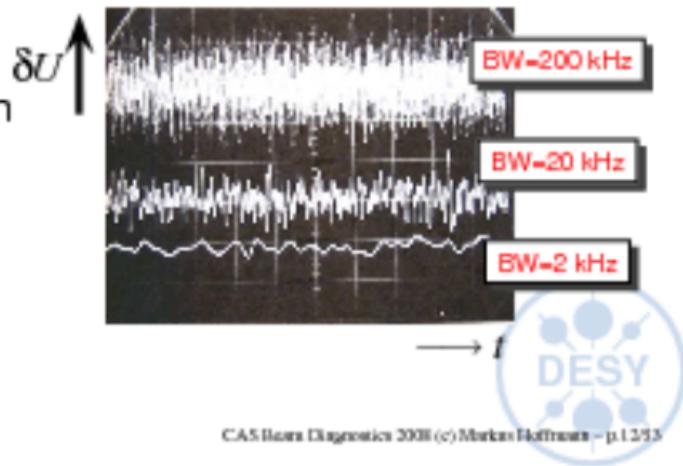
Amplitude Characteristics



$$U_\sigma = U_{\text{rms}} \sim \sqrt{\text{BW}}$$

$$\bar{U}_{\text{mean}} = 0$$

$$U_{\text{peak}} = \infty$$



The Physics behind

- Power Spectral Density $\rho(v)$:

$$\rho = kT \quad ; \quad P = \int_{\text{BW}} \rho dv$$

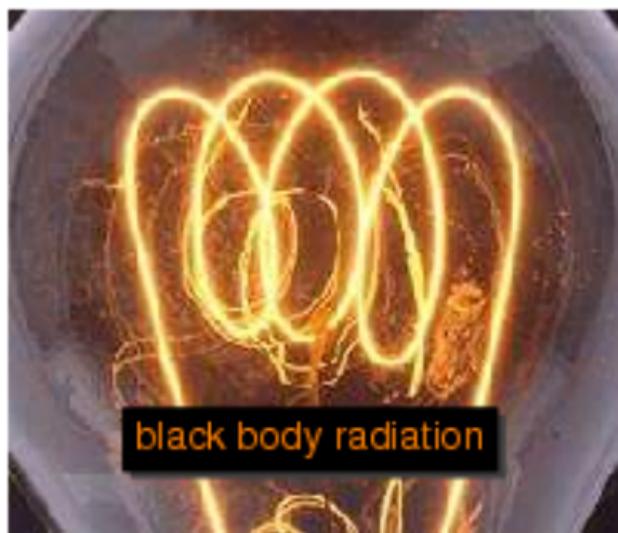
- derived from **Planck's law**:

$$\Rightarrow \rho = kT \frac{\frac{hv}{kT}}{e^{\frac{hv}{kT}} - 1} \approx kT$$

(+ zero point vacuum fluctuations $\frac{1}{2}hv$)

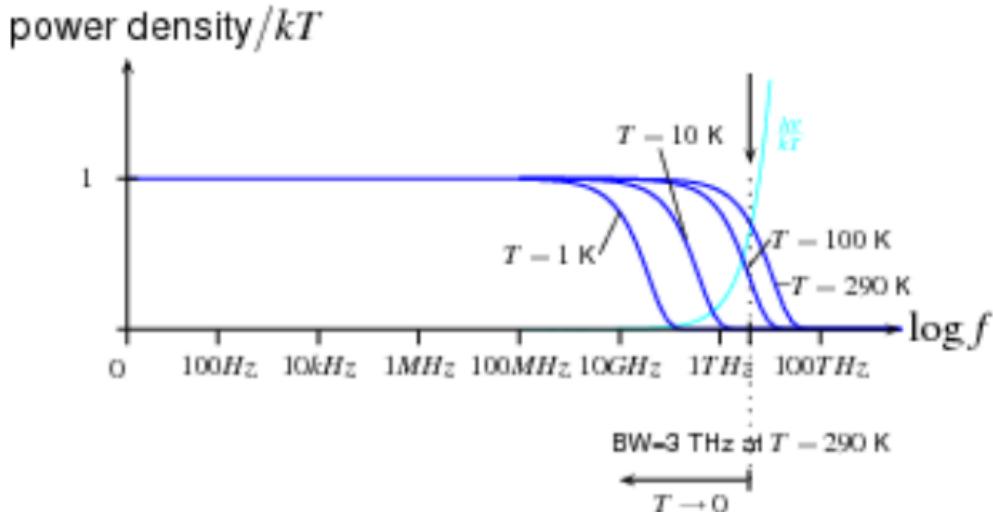
$$\left[\frac{\text{W}}{\text{Hz}} = \text{J} \right] \underset{T=290\text{K}}{=} 4 \cdot 10^{-21} \frac{\text{W}}{\text{Hz}}$$

Problem: total power = ∞ !



black body radiation

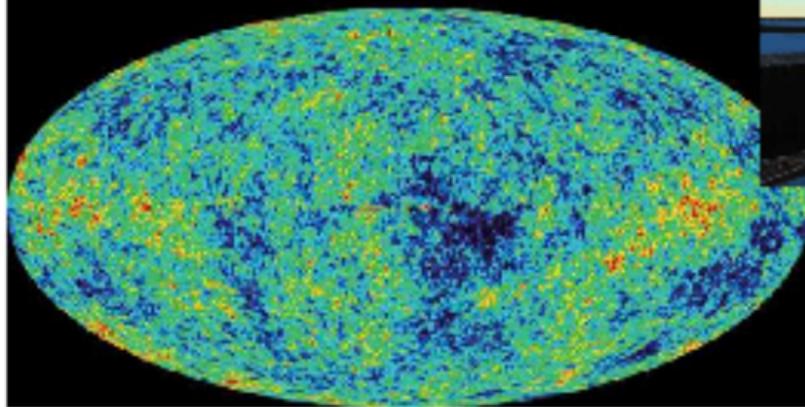
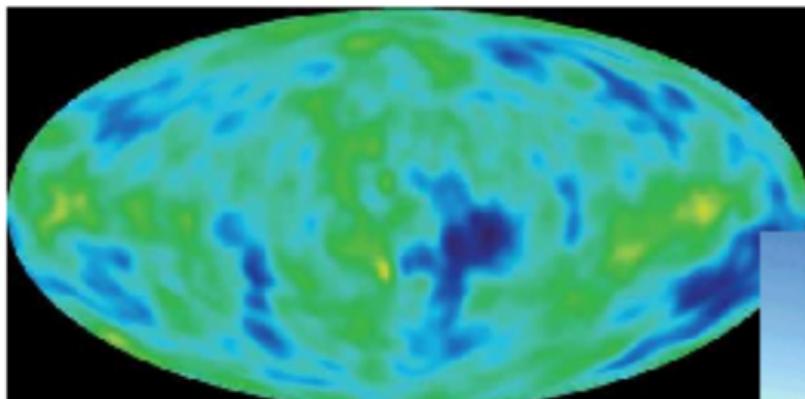
Thermal Noise Spectrum



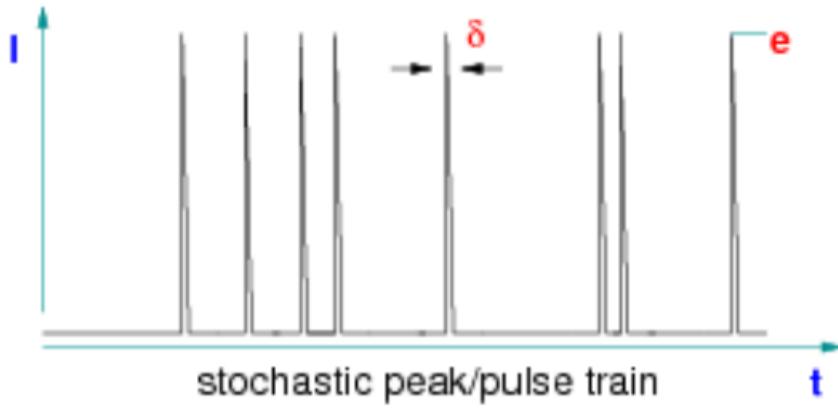
- For frequencies $>10\text{ GHz}$ you can gain more than proportionally by **cooling!**
- The **zero-point noise** *can not deliver power* to a load. But it can be amplified to a noticeable magnitude.



Cosmic Noise



Shot(tky) Noise

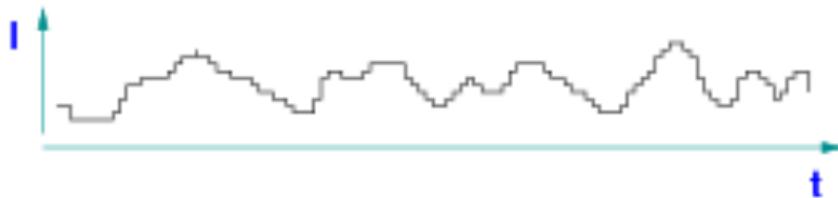
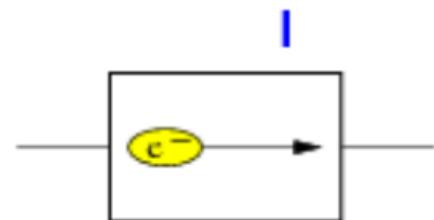
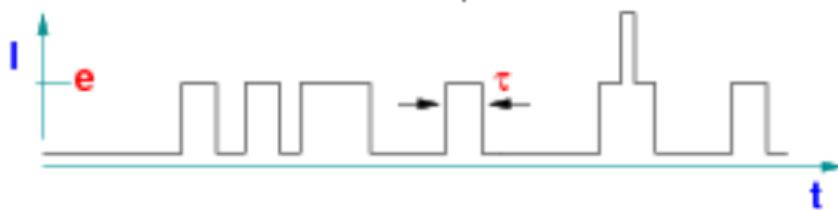


no current,
no shot noise!

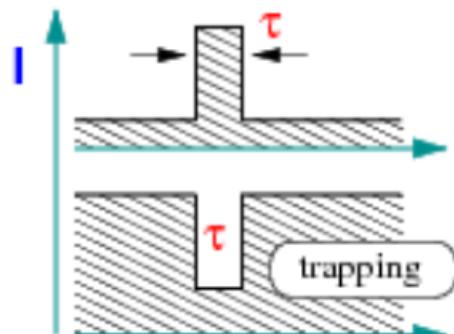
$$P_s \sim \langle I \rangle^2 = 2eI\Delta f$$

$$\delta I \sim \sqrt{2eI\Delta f}$$

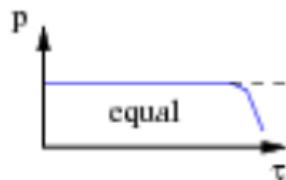
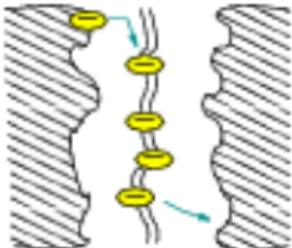
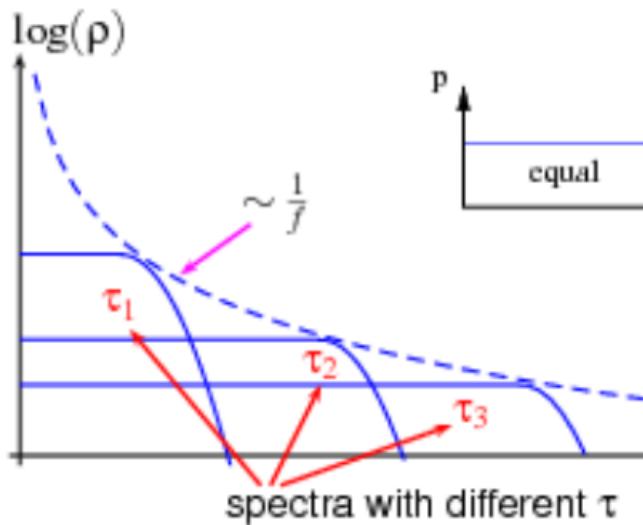
→ white spectrum



Flicker Noise

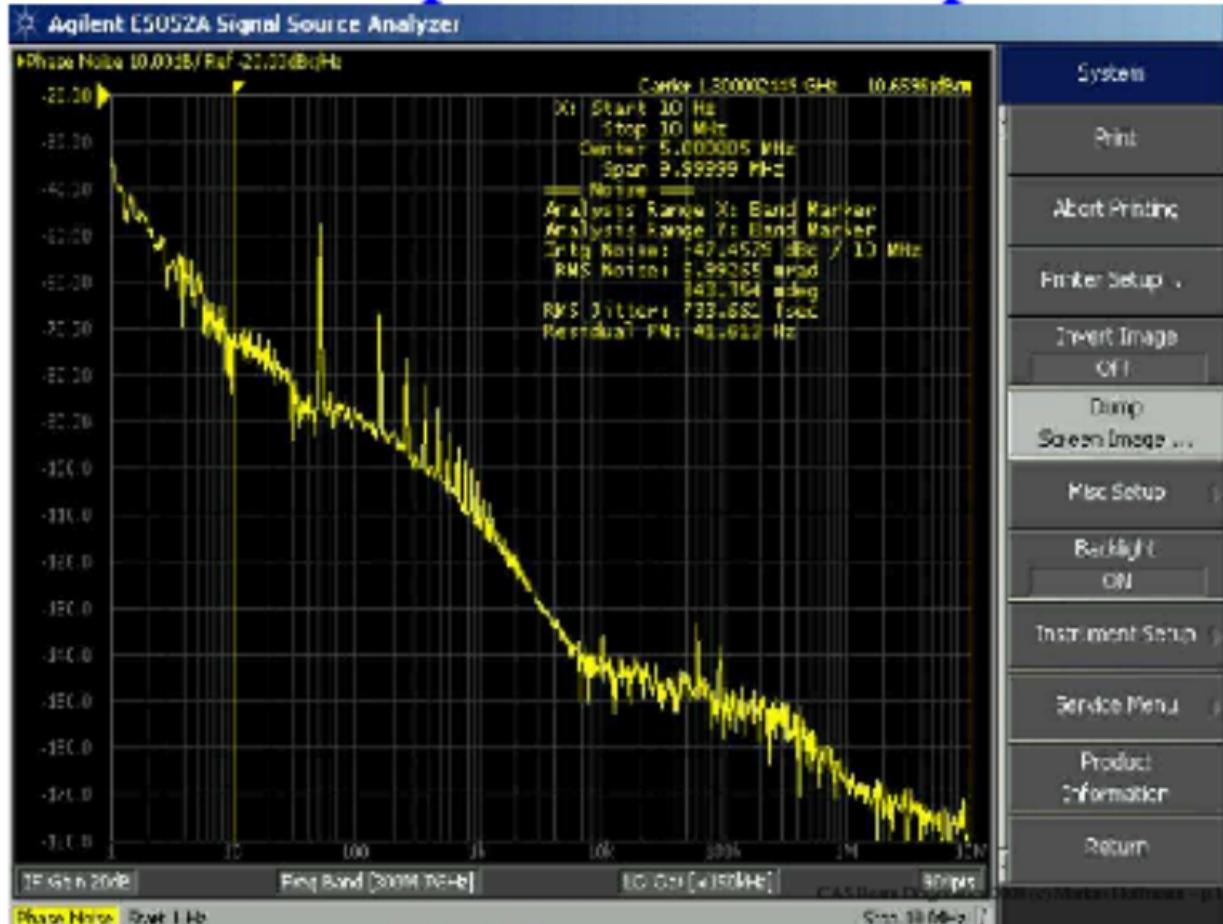


- Markow Process
- Lorentz spectrum



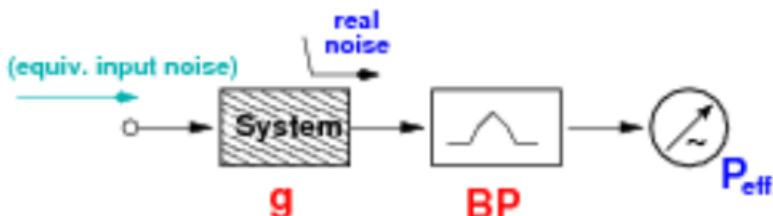
Phase noise spectrum of the FLASH Master Oscillator

Noise Spectrum – Example



System Noise Measurement

Input Noise:



$$\frac{P_{\text{eff}}}{\text{gain}} =: P_{\text{inputnoise}} \rightarrow \text{projected to the input}$$

- $\underbrace{\text{SNR} = 1}$ if input signal level = inputnoise level.
At the output



1. open input, measure $P_{\text{eff},1}$
2. close input, inject additional noise (P_r) so that $P_{\text{eff},2} = 2P_{\text{eff},1}$
3. read P_r from generator, $\rightarrow P_r = \text{input noise.}$



Signal to Noise Ratio

Power ratio, inside bandwidth!

$$\text{SNR} := \frac{\bar{P}_{\text{signal}}}{\bar{P}_{\text{noise}}} = \left(\frac{\hat{A}_{\text{signal,rms}}}{\hat{A}_{\text{noise,rms}}} \right)^2$$

Measurement:

$$\frac{S+N}{N} = \frac{S}{N} + 1$$

$$\bar{P} := \int_{\text{BW}} \rho(v) dv$$

$$\begin{aligned}\text{SNR(dB)} &:= 10 \log_{10} \left(\frac{\bar{P}_{\text{signal}}}{\bar{P}_{\text{noise}}} \right) = 20 \log_{10} \left(\frac{\hat{A}_{\text{signal,rms}}}{\hat{A}_{\text{noise,rms}}} \right) \\ &= P_{\text{signal}}[\text{dBm}] - P_{\text{noise}}[\text{dBm}]\end{aligned}$$

Units: [SNR(dB)] = dBc ("dB below carrier")



Noise Figures

noise factor:

$$F = \frac{P_r}{kT_0B} + 1 \quad \left(= \frac{SNR_{in}}{SNR_{out}} \right)$$

290 K

noise figure:

$$F' = 10 \log F \quad [\text{dB}]$$

- Relation of input noise + noise of the source to noise of the source only
- number of additional kTB units of noise necessary on the input to double the output noise

- is a measure of the system itself
- The figure is independent from T and B !

Examples

- the system adds no additional noise: $F = 1, F' = 0 \text{ dB}$ (The input noise equals thermal noise on the input. If the real temperature is lower than 290 K, F can be smaller.)
- a very noisy system reduces the SNR from 100 (20dB) to 10 (10dB):
 $F = 10, F' = 10 \text{ dB}$



A Data Sheet

For N8A models that either have a serial prefix less than 0B4446 or are fitted with the Noise Figure/BF board, N8973-60001.

Instrument's own noise figure

Frequency	Noise figure	Noise figure over a limited temperature range of $23^\circ\text{C} \pm 3^\circ\text{C}$
10 MHz to < 500 MHz	$< 4.9 + 3 + (0.0025 \times \text{freq in MHz})$	$< 4.4 \text{ dB} + (0.0025 \times \text{freq in MHz})$
500 MHz to < 2.3 GHz	$< 7.4 + 3 + (0.00135 \times \text{freq in MHz})$	$< 6.1 \text{ dB} + (0.00135 \times \text{freq in MHz})$
2.3 GHz to 3.0 GHz	$< 9.98 + (0.0115 \times \text{freq in MHz})$	$< 7.0 \text{ dB} + (0.0115 \times \text{freq in MHz})$
> 3.1 GHz to 13.2 GHz	< 12.0 dB	< 10.5 dB
> 13.2 GHz to 26.5 GHz	< 15.0 dB	< 12.5 dB

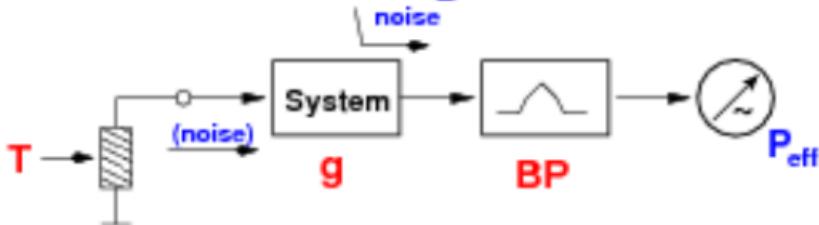
For N8A models that either have a serial prefix greater than 0B4446 or are fitted with the Noise Figure/BF board, N8973-60001.

Instrument's own noise figure

Frequency	Noise figure	Noise figure over a limited temperature range of $23^\circ\text{C} \pm 3^\circ\text{C}$
10 MHz to 3.0 GHz	$< 4.8 \text{ dB} + (0.00174 \times \text{freq in MHz})$	$< 4.4 \text{ dB} + (0.00174 \times \text{freq in MHz})$
> 3.0 GHz to 13.2 GHz	< 12.0 dB	< 10.5 dB
> 13.2 GHz to 26.5 GHz	< 16.0 dB	< 12.5 dB



Noise Temperature



Put a resistor T_{noise} so that P_{eff} doubles:

$$T_{\text{noise}} := \underbrace{(F - 1) \cdot 290 \text{ K}}_{\text{for white noise}} \quad [\text{K}]$$

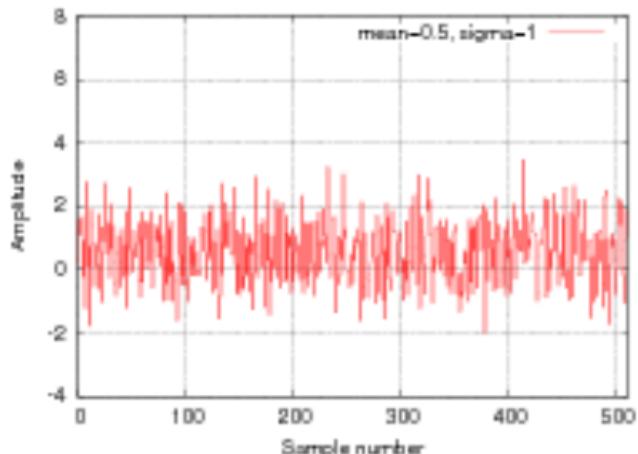
- for an ideal noise-free system $\rightarrow T_{\text{noise}} = 0 \text{ K}$.
- is not necessarily identical with the real temperature the system is on.
- A low-noise amplifier can have $T_{\text{noise}} = 20 \text{ K} \rightarrow F = 1.06, F' = 0.28 \text{ dB}$.
(This is really good!)



3. Statistics

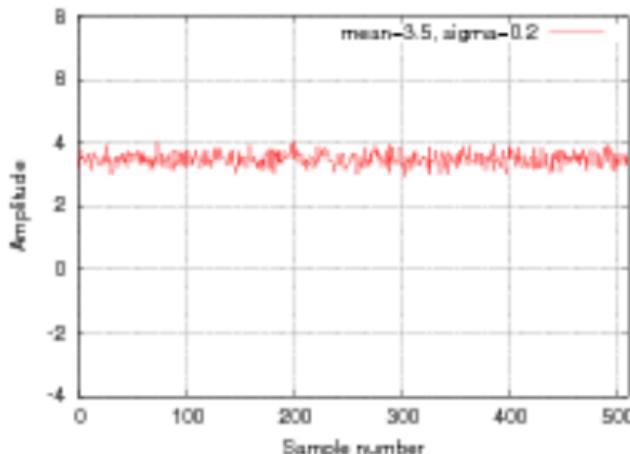


Mean and Standard Deviation



$$\hat{x} := \frac{1}{N} \sum_{i=0}^{N-1} x_i$$

mean over N samples.



$$\sigma^2 := \frac{1}{N-1} \sum_{i=0}^{N-1} (x_i - \hat{x})^2$$

variance = (standard deviation)²
"power of fluctuations"



RMS and the Running Statistics

$$\sigma_N^2 = \frac{1}{N-1} \left[\underbrace{\sum_{i=0}^{N-1} x_i^2}_{\text{sum of squares}} - \frac{1}{N} \underbrace{\left(\sum_{i=0}^{N-1} x_i \right)^2}_{\text{sum}^2} \right]$$

Signal to Noise Ratio (SNR): $\text{SNR} = \frac{\hat{x}^2}{\sigma^2}$

Coefficient of Variation (CV): $\text{CV} = \frac{\sigma}{\hat{x}} \cdot 100\%$

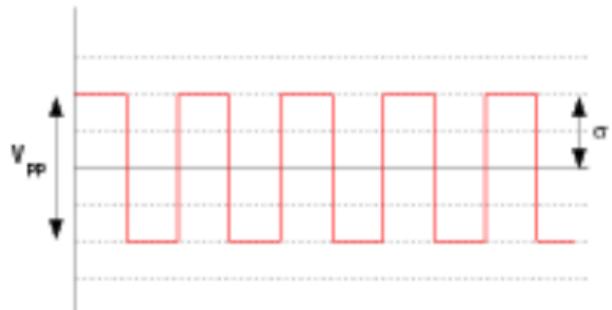
Root Mean Square (RMS): $x_{\text{rms}} := \sqrt{\frac{1}{N} \sum_{i=0}^{N-1} x_i^2}$

"Power of fluctuations plus power of DC component"

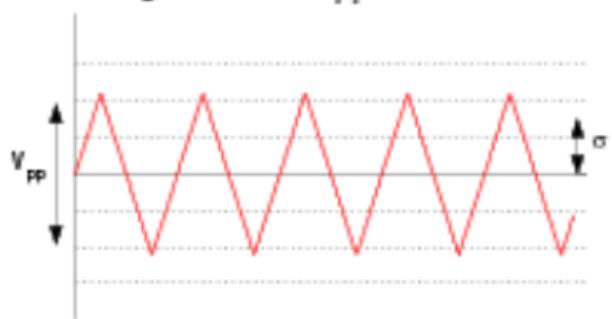


Standard Deviation of Common Waveforms

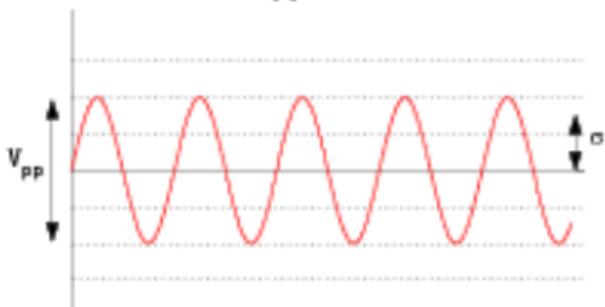
a. Square wave, $V_{pp} = 2\sigma$



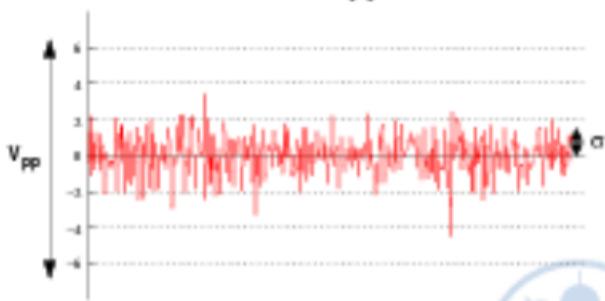
c. Triangle wave, $V_{pp} = \sqrt{12}\sigma$



b. Sine wave, $V_{pp} = 2\sqrt{2}\sigma$

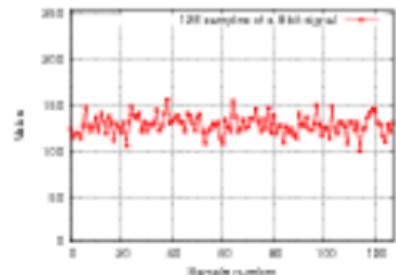


d. Random noise, $V_{pp} = 6 - 8\sigma$

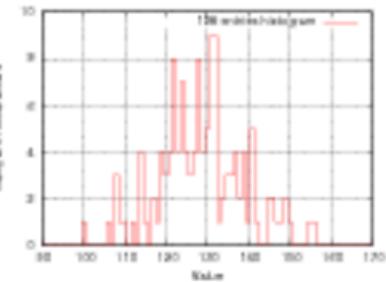


Histograms

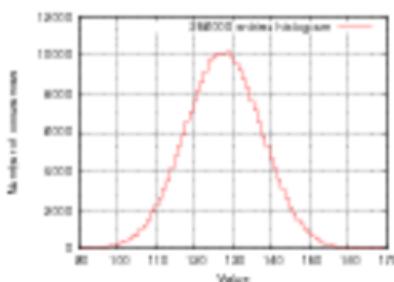
Snapshot of N samples



N small

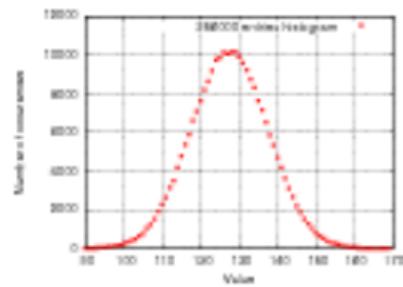


N large

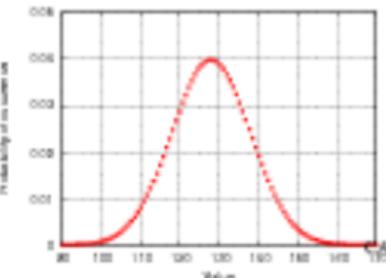


$$N = \sum_{i=0}^{M-1} H_i \quad \hat{x} := \frac{1}{N} \sum_{i=0}^{M-1} i \cdot H_i \quad \sigma^2 := \frac{1}{N-1} \sum_{i=0}^{M-1} (i - \hat{x})^2 H_i$$

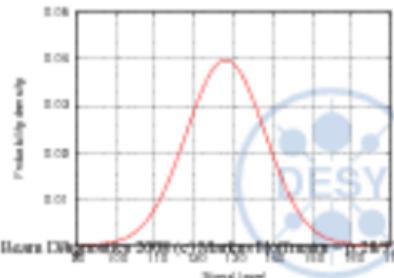
Histogram



Probability mass function



Probability density distribution



$N \rightarrow \infty$

Probability of a sample

Value

Probability density

Value

Probability density

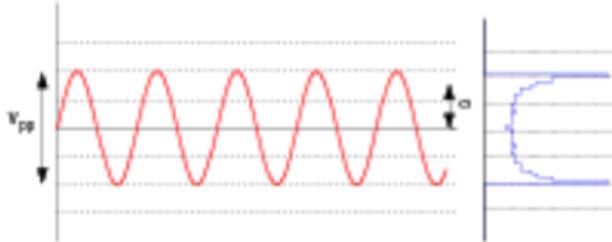
Signal level

Probability Distribution Functions

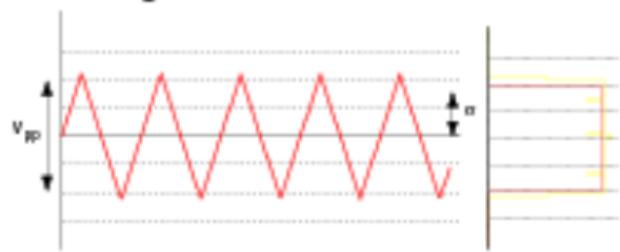
a. Square wave



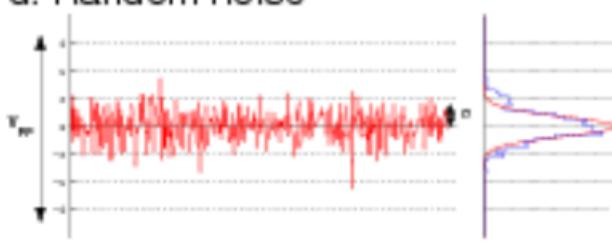
b. Sine wave



c. Triangle wave



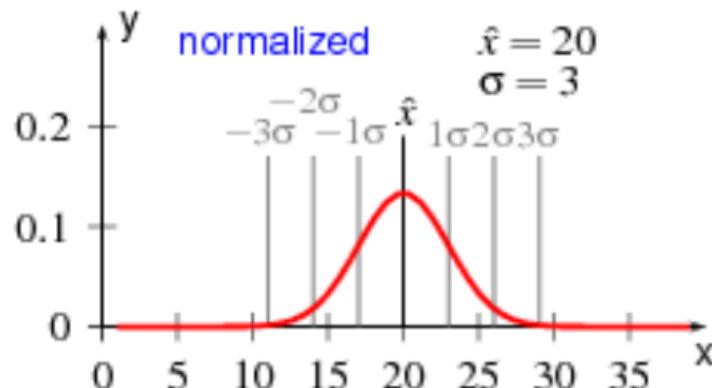
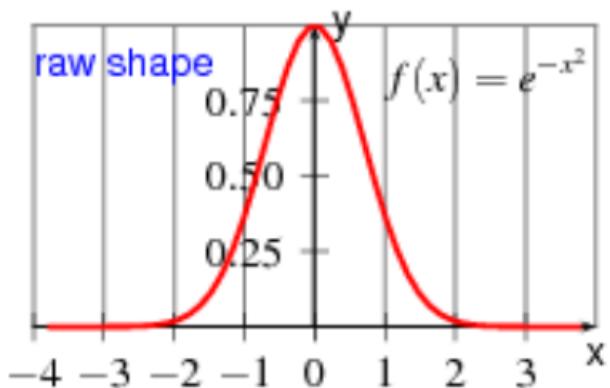
d. Random noise



The Normal Distribution

$$P(x) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-\hat{x})^2}{2\sigma^2}}$$

Gauß function



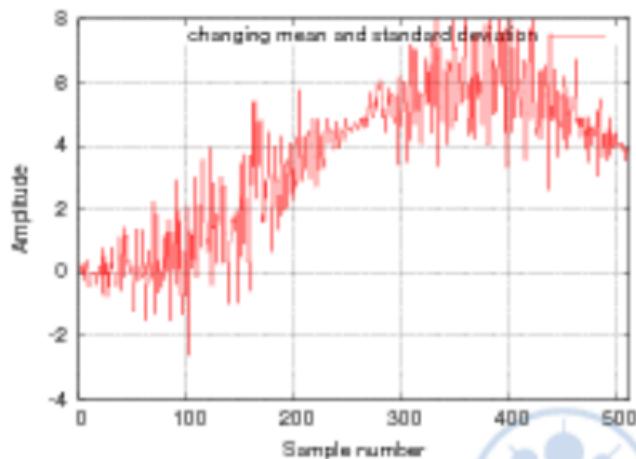
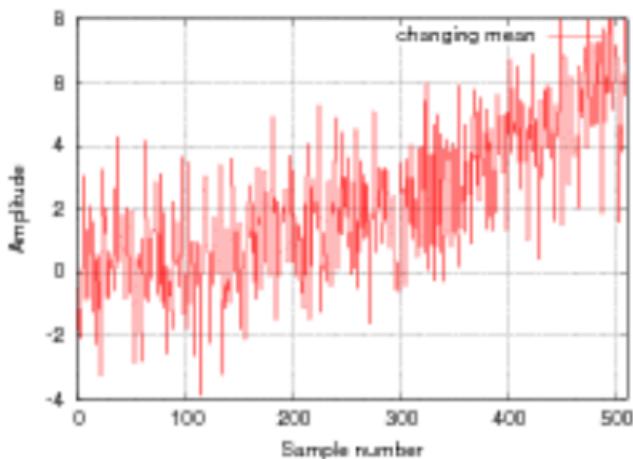
$$\int_{-\infty}^{+\infty} P(x) dx = 1$$



Underlying Process

$$\text{typical error: } \Delta A = \frac{\sigma_N}{\sqrt{N}}$$

- σ_N is an estimate of the standard deviation of the **underlying process** by N samples (e.g. extracted from the histogram).



- Extract the most information about the underlying process out of the sampled signal.

Propagation of Error

in general:

$$f = f(\alpha_1, \alpha_2, \dots, \alpha_n)$$

a function of (model) parameters α_i with corresponding errors $\Delta\alpha_i$.

$$\Rightarrow \Delta f = \sqrt{\sum_i \left(\frac{\partial f}{\partial \alpha_i} \Delta \alpha_i \right)^2}$$

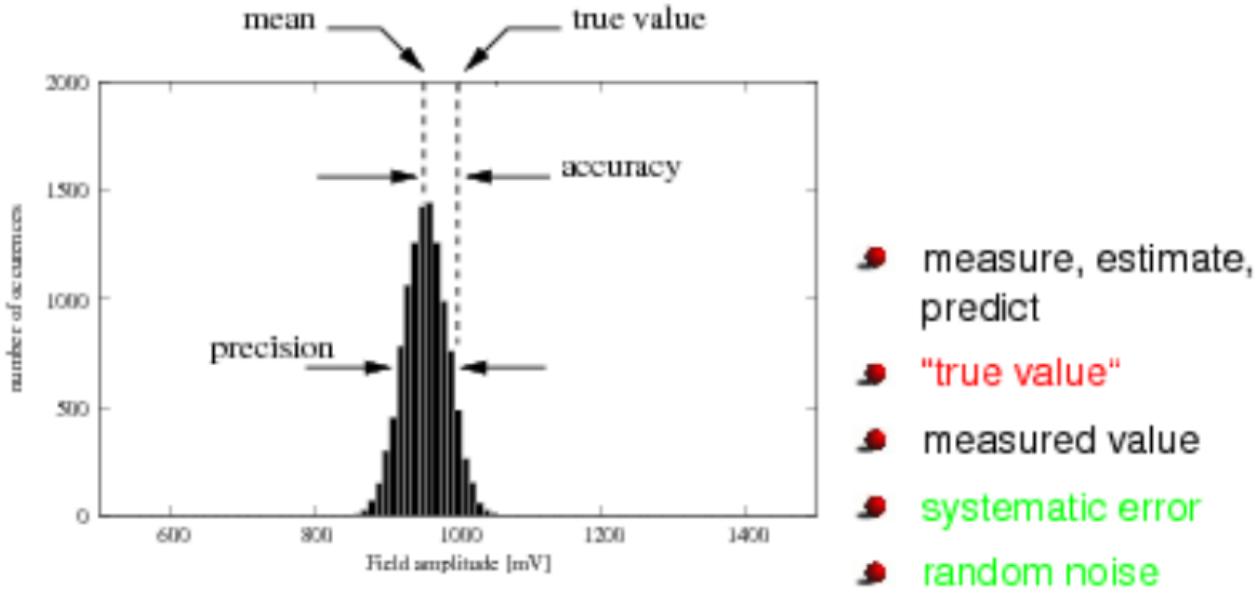
Example:

$$x(t) = v \cdot t + x_0$$

$$\begin{aligned}\Rightarrow \Delta x &= \sqrt{\left(\frac{\partial x}{\partial v} \Delta v \right)^2 + \left(\frac{\partial x}{\partial x_0} \Delta x_0 \right)^2 + \left(\frac{\partial x}{\partial t} \Delta t \right)^2} \\ &= \sqrt{(\Delta v \cdot t)^2 + (\Delta x_0)^2}\end{aligned}$$



Accuracy and Precision



- **accuracy** is a measure of calibration
- **precision** is a measure of statistics

Question: measuring with 100 different devices. → accuracy or precision?



The Central Limit Theorem

- generation of random numbers (noise)

- white noise $RND := [0; 1[$

- gaussian noise:

"the sum of random numbers (any distribution) becomes gaussian distributed"

better:

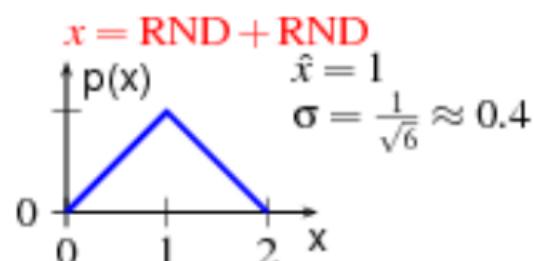
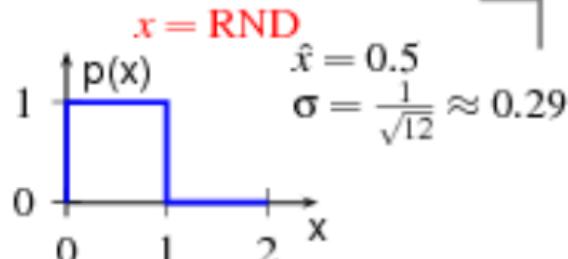
$$x = \sqrt{-2 \log_{10}(RND_1)} \cdot \cos(2\pi RND_2)$$

$$\hat{x} = 0, \quad \sigma = 1$$

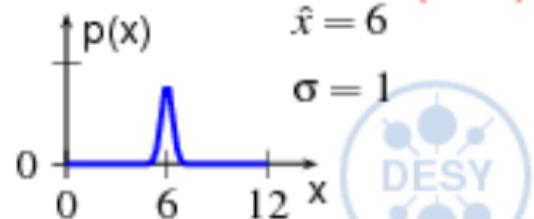
- pseudo-random:

$$RND = (as + b) \bmod c$$

seed



$x = RND + \dots + RND \text{ (12 x)}$



Chi Square Distribution

X_i are k independent, *normally distributed* random variables with $\bar{X}_i = 0$, $\sigma_{x_i}^2 = 1$:

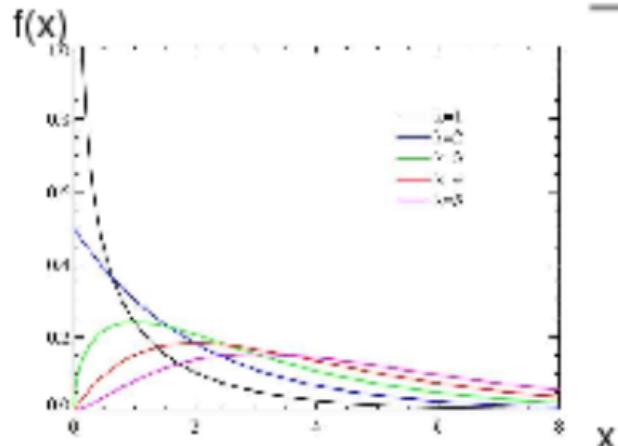
$$Q := \sum_{i=1}^k X_i^2$$

is distributed according to „Chi Square“

$$Q \sim \chi_k^2$$

k : number of degrees of freedom

For $k \gg 1$: Q becomes Gaussian distributed.



Probability distribution:

$$f(x; k) = \begin{cases} \frac{x^{\frac{k}{2}-1}}{2^{\frac{k}{2}} \Gamma(\frac{k}{2})} e^{-\frac{x}{2}} & \text{for } x > 0, \\ 0 & \text{for } x \leq 0. \end{cases}$$



Chi Square Test

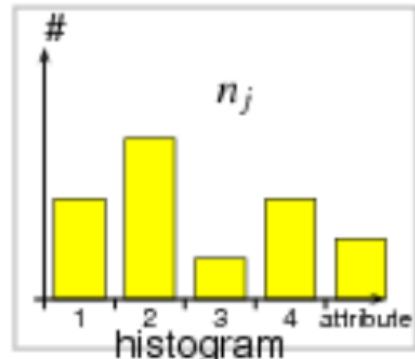
=your model

Null-hypothesis: is true if your alternative hypothesis cannot be supported.

Statistical Test: falsify the Null-hypothesis!

χ^2 -Test: the statistics of samples of data (derived from model) has a χ^2 distribution if the null-hypothesis is true.

Or gauss or any, if the probability distribution approximates large # of samples



probability, that x is in j

Theory: $n_{j0} = F_0(x)_j \cdot N$

$$\Rightarrow \chi^2 = \sum_{j=1}^m \frac{(n_j - n_{j0})^2}{n_{j0}}$$

χ^2 distributed,
 $m - 1$ degrees of freedom
if N large

If χ^2 is larger than a **significance level α** ,
the hypothesis will be rejected.

from tables
(out of the scope)



4. Advanced Concepts



Power Spectral Density

averaged power levels related to intervals of noise frequency

average over many time intervals T

$$\rho_x(v) := \lim_{T \rightarrow \infty} \frac{1}{T} \left\langle \left| \int_{-\frac{T}{2}}^{\frac{T}{2}} X(t) e^{2\pi i vt} dt \right|^2 \right\rangle \quad \text{power spectrum of } X(t)$$

Unit: $\left[\frac{W}{Hz} \right]$, $\left[\frac{dBm}{Hz} \right]$, $\left[\frac{dBc}{Hz} \right]$, or $\left[\frac{rad^2}{Hz} \right]$ (phase noise). Amplitude: $\left[\frac{rad}{\sqrt{Hz}} \right]$

Wiener-Khinchin theorem:

$\rho_x(v) = \text{Fourier Transformation of Autocorrelation function of } X(t)$

$$\rho_x(v) = \int_{-\infty}^{\infty} g_x(\tau) e^{2\pi i v \tau} d\tau \quad \text{fourier transform of}$$

$$g_x(t) := \left\langle \int_{-\infty}^{\infty} x(\tau) x(t + \tau) d\tau \right\rangle \quad \text{average autocorrelation function}$$



The Fourier Transformation

time domain

$$s(t)$$

signal

frequency domain

$$S(\omega)$$

spectrum

Given $f : D \rightarrow \mathbb{C}$, where $D \subseteq \mathbb{R}$, the

Fourier transformation of f is:

and the **back transformation**

$$F(\omega) := \int_D f(t) e^{-i\omega t} dt$$

$$f(t) := \int_{-\infty}^{\infty} F(\omega) e^{+i\omega t} d\omega .$$

- **Note:** We use $\omega = 2\pi\nu$ to get rid of the constants.



Fourier Transformation Examples

$s(t)$	time domain	$S(f)$	frequency domain $ S $
$\delta(t)$		$\delta(f)$	
$\varpi(t)$		$\varpi(f)$	
$e^{-\pi t^2}$		$e^{-\pi f^2}$	
$2 \cos(2\pi F t)$		$\delta(f+F) + \delta(f-F)$	

Calculation with Fourier Transforms

$$x(t) \xrightarrow{\text{FT}} X(\omega) \xrightarrow{\text{FT}} x(-t)$$

Symmetry:

$$\text{FT}^2\{x(t)\} = x(-t)$$

Linearity:

$$\text{FT}\{c_1x_1(t) + c_2x_2(t)\} = c_1X_1(\omega) + c_2X_2(\omega)$$

Scaling:

$$\text{FT}\{x(\lambda t)\} = \frac{1}{|\lambda|}X\left(\frac{\omega}{\lambda}\right)$$

Convolution:

$$\text{FT}\{x_1(t) * x_2(t)\} = X_1(\omega) \cdot X_2(\omega) ; \quad \text{FT}\{x_1(t) \cdot x_2(t)\} = X_1(\omega) * X_2(\omega)$$

For a **real** input, the transformation produces a **complex** spectrum which is symmetrical:

$$X(\omega) = X^*(-\omega)$$

complex conjugate

$$\begin{aligned} X(\cos\text{-like}) &= \text{real} \\ X(\sin\text{-like}) &= \text{imaginary} \end{aligned}$$

Time-Shift:

$$\text{FT}\{x(t + t_0)\} = e^{j\omega t_0}X(\omega)$$

Autocorrelation

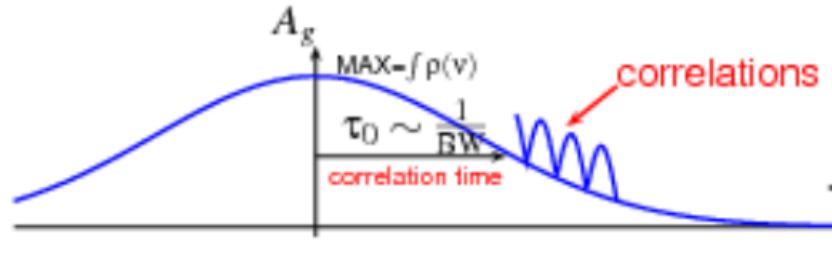
Given two functions $f, g : D \rightarrow \mathbb{C}$,
where $D \subseteq \mathbb{R}$, the **cross correlation**
of f with g :

$$(f \circ g)(t) := K \int_D f(\tau) g(t + \tau) d\tau$$

Autocorrelation

$$A_g(t) := g \circ g = K \int_D g(\tau) g(t + \tau) d\tau$$

Detect a known waveform in a noisy background, e.g. echoes.



Symmetry:

$$A_g(-t) = A_g^*(t)$$

Convolution:

$$f[n] \circ g[n] = f[n] * g[-n]$$

Peak:

$$|A_g(t)| \leq A_g(0)$$

Periodicity:

$$g(t) \text{ periodic} \Leftrightarrow A_g(t) \text{ periodic}$$



Power Spectral Density (2)

- Spectrum of the autocorrelation function:

$$s(t) * s(-t) \xrightarrow{\text{FT}} S(\omega) \cdot S^*(\omega) = |S(\omega)|^2 = \rho(\omega) \quad [\text{W/Hz} = \text{J}]$$

Energy spectrum

- we loose information about phase (or time/shift/location), *Time invariance*
- white noise part is contained in $g_x(0)$,
- relation to the **variance** and RMS of $x(t)$:

$$\langle x^2 \rangle = \int_{v_1}^{v_2} \rho_x(v) dv = P_{\text{BW}} = (\text{RMS}_x)^2 =: x_{\text{eff}}^2 =: \sigma_x^2|_{[v_1, v_2]}$$

Wiener-Khinchin:

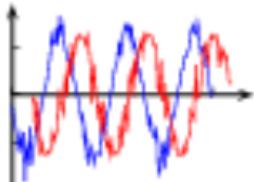
mean square value in time = mean square value in frequency

- analog and digital measurement techniques possible.

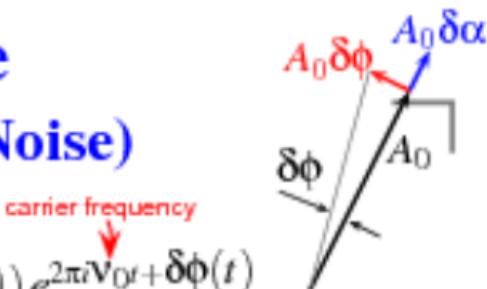


Complex Noise

(Amplitude & Phase Noise)



Harmonic Signal with Noise: $x(t) = A_0 (1 + \delta\alpha(t)) e^{2\pi i v_0 t + \delta\phi(t)}$

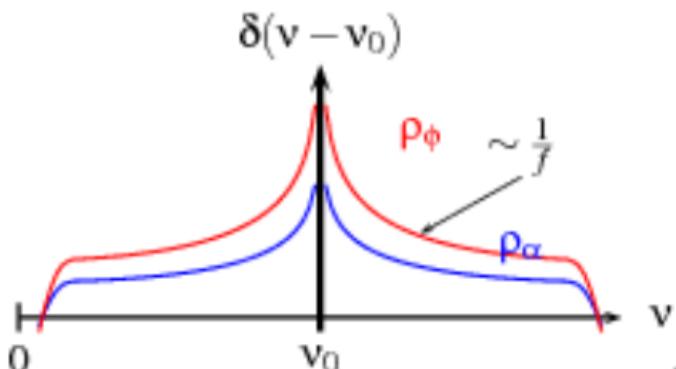


carrier δ peak

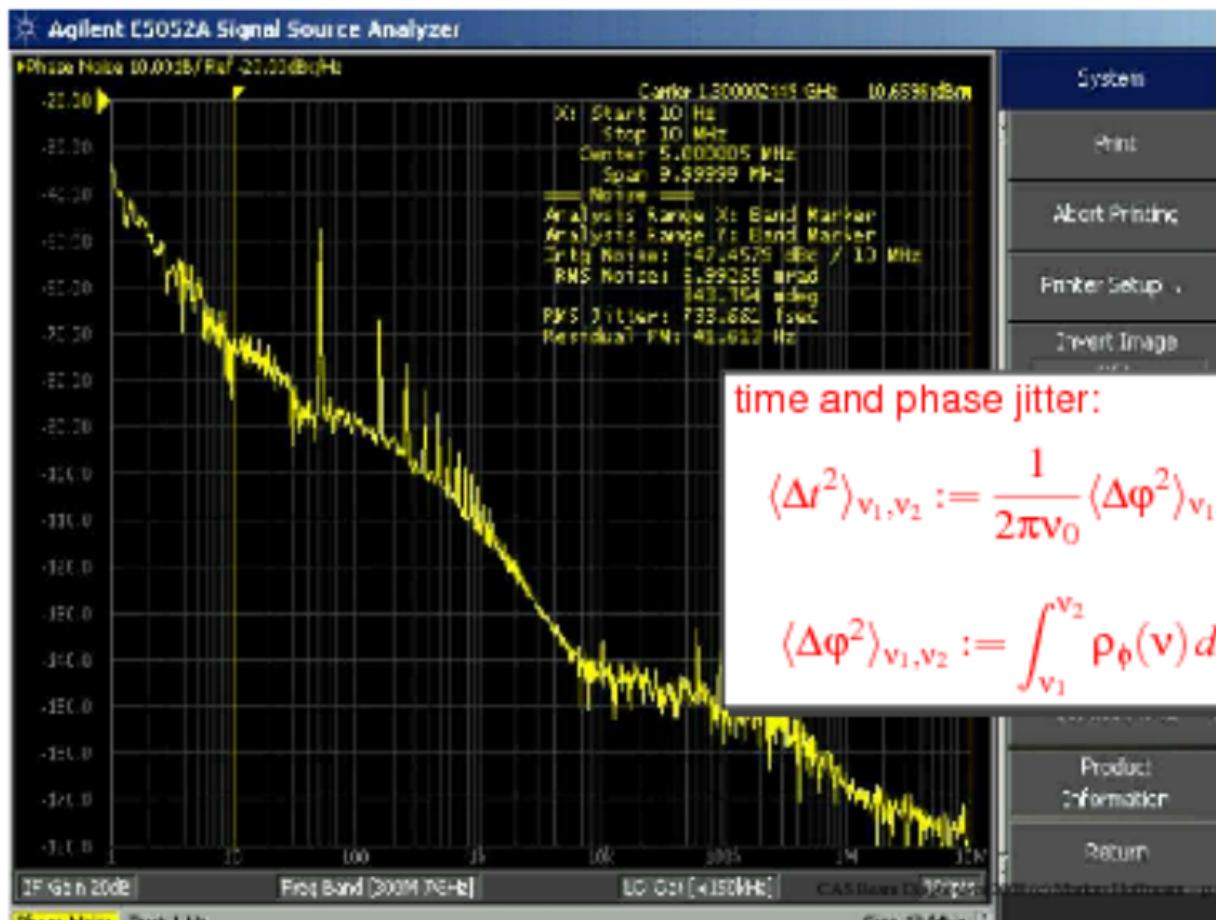
amplitude noise

phase noise

$$p(v) = A_0^2 (\delta(v - v_0) + p_\alpha(v - v_0) + p_\phi(v - v_0) + \mathcal{O}^n(p_\alpha, p_\phi))$$

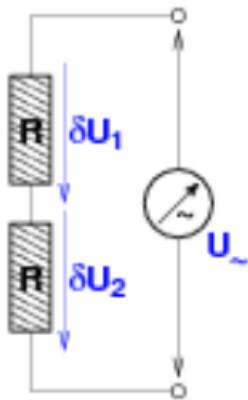


Phase noise spectrum of the FLASH Master Oscillator

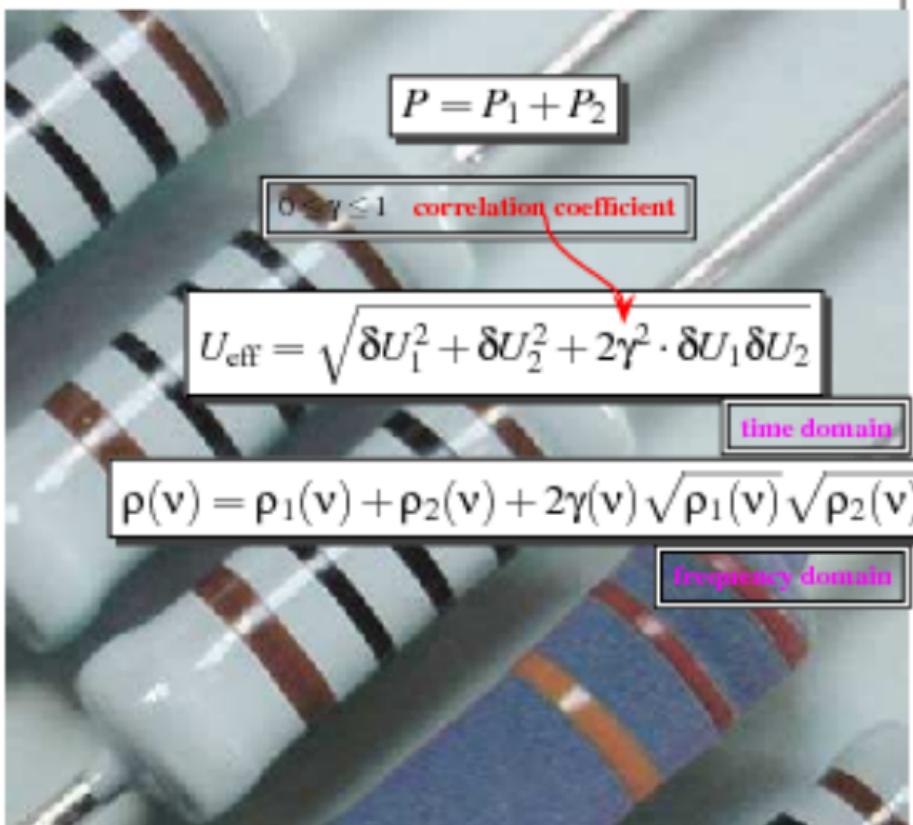
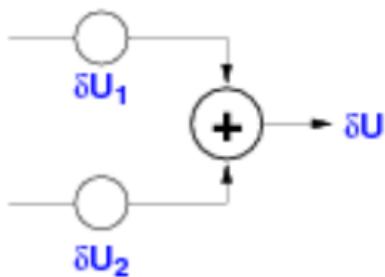


Noise Propagation

- Adding noise sources:

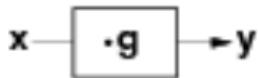


\Leftrightarrow



Noise Propagation (2)

- Amplifier/Attenuator:



$$y(t) = g x(t)$$

$$P_y = g^2 P_x$$

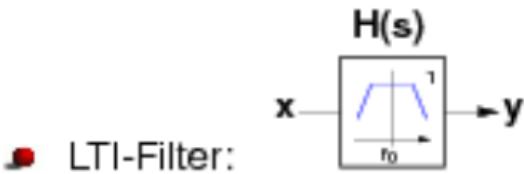
$$U_{\text{eff},y} = g \cdot \delta U_x$$

$$\rho_y(v) = g^2 \rho_x(v)$$

$$\rho_{\alpha,y}(v) = \rho_{\alpha,x}(v)$$

$$\rho_{\phi,y}(v) = \rho_{\phi,x}(v)$$

The power of the phase noise is not(!) amplified!



- LTI-Filter:

$$y(t) = h(t) * x(t)$$

transfer function

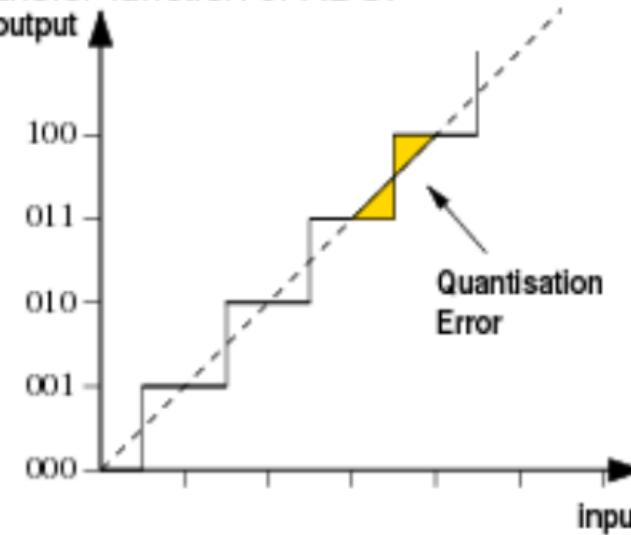
$$\rho_y(v) = |H(v)|^2 \rho_x(v)$$

$$\rho_{\alpha,y}(v) = \left| \frac{H(v + v_0)}{H(v_0)} \right|^2 \rho_{\alpha,x}(v)$$

$$\rho_{\phi,y}(v) = \left| \frac{H(v + v_0)}{H(v_0)} \right|^2 \rho_{\phi,x}(v)$$

Quantisation Noise

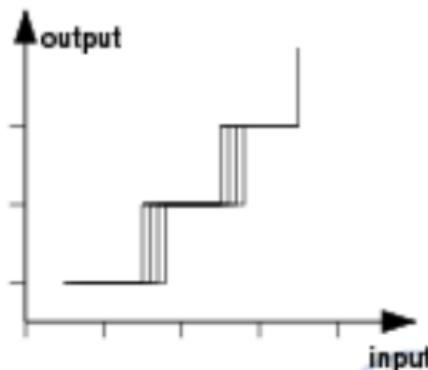
Transfer-function of ADC:



- Quantisation-Error
 $|A| < 0.5 \text{ LSB}$

$$\text{RMS}(\Delta A) \approx \sqrt{12} \text{ LSB}$$

- missing codes
- code transition noise



For a full scale $\sin(\cdot)$ -signal:

$$\text{SNR} = 6.02n + 1.76 \text{ dB} + 10 \log \left(\frac{f_s}{2 \text{ BW}} \right)$$

increases with lower BW.

→ doubling the sampling frequency increases SNR by 3dB (same signal BW)
→ "oversampling"

Quantisation Noise Spectra

$$\text{SNR} = 6.02n + 1.76\text{dB} + 10 \log \left(\frac{f_s}{2 \text{BW}} \right)$$

it is assumed that the noise is equally distributed over the full BW.

! This is often not the case !

Mostly the noise is **correlated** with the input signal!

- The lower the signal, the more correlation!
- In case of strong correlation the noise is concentrated at the various harmonics of the input signal, just where you dont want them.
- dithering and broad input signal spectrum randomizes the quantisation noise.



5. Applications



Stochastic Signals in Accelerator Diagnostics

Apply additional and/or artificial
(pseudo random) noise

Use the noise which is there anyway

- **Tune measurements** from shottky noise of the beam itself (without increasing the emittances), esp. hadron beams)
- Synchrotron light emmission → noise source, beam exitation

- dithering methods
- beam size blow ups (nbunch lengthening)
- decorrelate signals to avoid
 - ↳ systematic errors,
 - ↳ interference and
 - ↳ resonant excitations.
- **transfer function** measurements
- **Stochastic cooling**

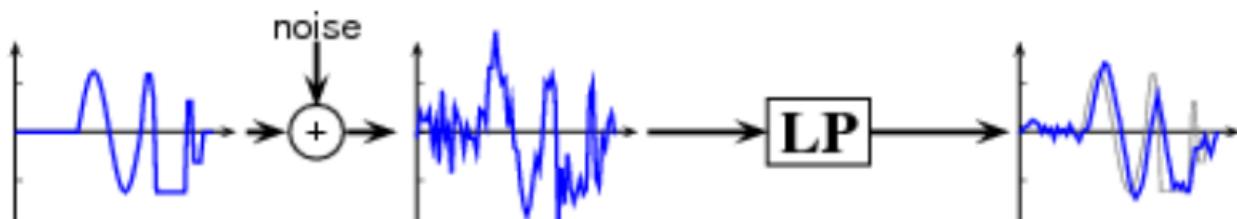
in many other cases the stochastic part of the signal is unwanted (noise)

→ Noise Filtering Techniques



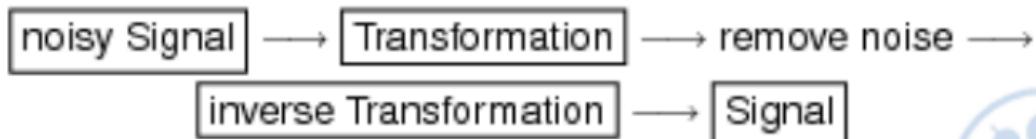
Fighting the Noise

- How can we reconstruct a signal to which noise has been added?
- One Idea is to use a Low-Pass Filter:



Problems are: latency, dissipation, high frequency cut-off and still noise in the low frequencies.

- solution to this: **the Kalman Filter**
- other solution: (non-causal)



Measurements, Statistics and Errors



The art of measurement is always also the art of error treatment!

The End

