



Overview of Future Accelerators

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CERN Accelerator School

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Daresbury

- **Future Accelerators for particle physics**
 - What is needed & why
 - The Large Hadron Collider (LHC)
 - The Linear Collider (LC)
 - The Muon Collider (MC)
 - The Neutrino Factory (NF)

Will not have time to discuss this

- EURISOL and Beta Beams
- 'Factories' (e.g. c)

• **Future Accelerators for other sciences**

- What is needed and why
- X-ray sources

... sorry, because it is important

- Nuclear Physics

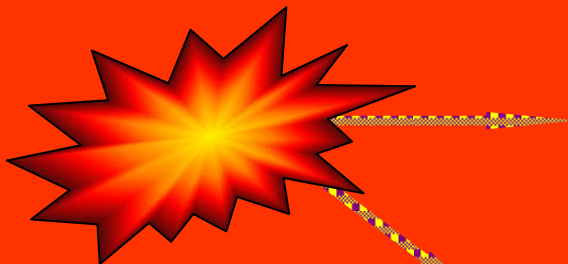
- **Other applications**
 - Accelerators in Medicine
- **Summary**

What is needed, and why

2 routes to new knowledge about the fundamental structure of the matter

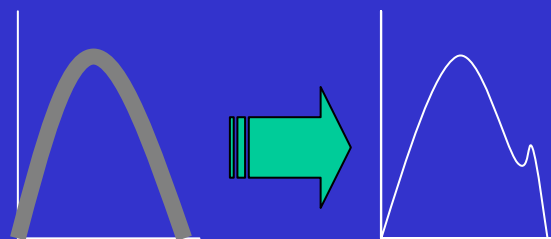
High Energy Frontier

New phenomena
(new particles)
created when the
“usable” energy $> mc^2$ [x2]



High Precision Frontier

Known phenomena studied
with high precision *may* show
inconsistencies with theory



The Standard Model

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + i\bar{\psi}D\psi$$

$$+ \psi_i \lambda_{ij} \psi_j h + h.c.$$

$$+ |D_\mu h|^2 - V(h)$$

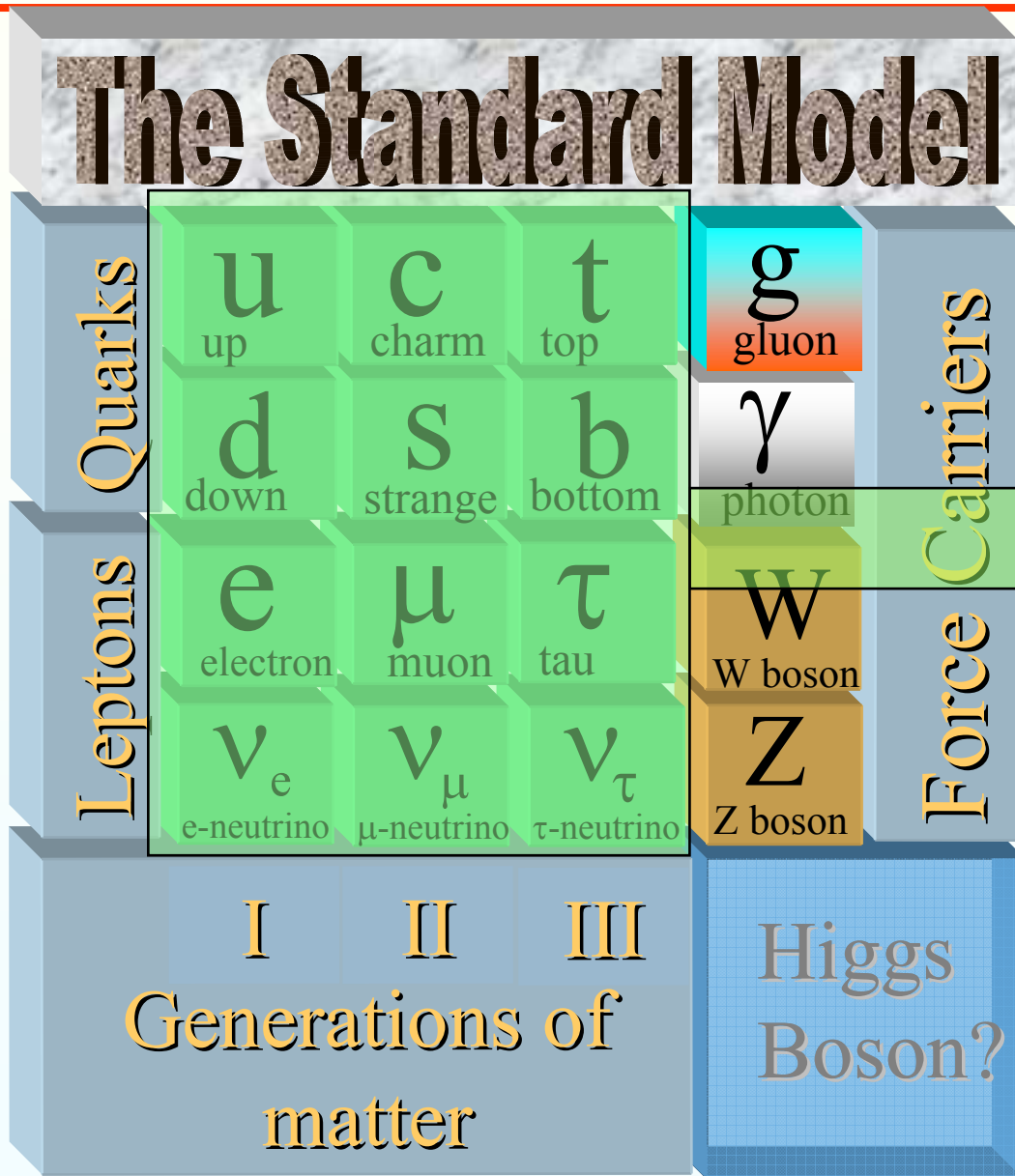
$$+ \frac{1}{M} L_i \lambda_{ij}^\nu L_j h^2 \text{ or } L_i \lambda_{ij}^\nu N_j$$

The gauge sector (1)

The flavor sector (2)

The EWSB sector (3)

The ν -mass sector (4)



Each with its own 'antiparticle'

The Standard Model

The Standard Model Effective Lagrangean

$$\mathcal{L}_{(\text{Standard Model})} =$$

$[W^\pm]$	$-\frac{1}{2}(\partial_\mu W_\nu - \partial_\nu W_\mu)(\partial^\mu W^{\nu\mu} - \partial^{\nu\mu} W^{\mu\nu}) + M_W^2 W_\mu W^{\mu\nu}$
[Photon]	$-\frac{1}{4}F_{\mu\nu}^A F^{\mu\nu A}$
$[Z^0]$	$-\frac{1}{2}F_{\mu\nu}^Z F^{\mu\nu Z} + \frac{1}{2}M_Z^2 Z_\mu Z^\mu$
$[\ell, \nu_\ell]$	$+i\bar{L}_\ell \not{\partial} L_\ell + i\bar{R}_\ell \not{\partial} R_\ell - m_\ell \bar{\ell}\ell$
$[W\ell\nu]$	$-\frac{g}{\sqrt{2}}\bar{L}_\ell(\tau_+ W + \tau_- W)L_\ell$
$[\gamma\ell^+\ell^-]$	$+e_{\ell/m}\bar{\ell}\not{A}\ell$
$[Z\ell^+\ell^-, Z\nu\bar{\nu}]$	$-\frac{g}{\cos\theta_w}\bar{L}_\ell\left(\frac{\tau_3}{2}\cos^2\theta_w + \frac{1}{2}\sin^2\theta_w\right)\not{Z}L_\ell - \frac{g\sin^2\theta_w}{\cos\theta_w}\bar{R}_\ell\not{Z}R_\ell$
[H]	$+\frac{1}{2}\partial_\mu H\partial^\mu H - \frac{1}{2}\mu^2 H^2 - \frac{1}{2}\lambda\mu H^3 - \frac{1}{8}\lambda^2 H^4$
[HH&H W^+W^-]	$+\frac{g^2}{8}\left(H^2 + \frac{2\mu}{\lambda}H\right)(2W_\mu W^{\mu\nu})$
[HH&H ZZ]	$+\frac{g^2}{8}\left(H^2 + \frac{2\mu}{\lambda}H\right)\left(\frac{1}{\cos^2\theta_w}Z_\mu Z^\mu\right)$
[H $\ell^+\ell^-$]	$-m_\ell\sqrt{\sqrt{2}G_F}\bar{\ell}\ell H$
[quark γ]	$+\bar{Q}q\not{A}q$
[quark Z]	$-\frac{g}{\cos\theta_w}\bar{L}_q\left(\frac{\tau_3}{2}\cos^2\theta_w + \frac{\sin^2\theta_w}{2}\right)\not{Z}L_q$
[quark W]	$-\frac{g}{\sqrt{2}}\bar{U}V_{CKM}(\tau_+ W + \tau_- W)\mathcal{D}$
[quark H]	$-m_q\sqrt{\sqrt{2}G_F}\bar{q}q H$
[gluons]	$-\frac{1}{4}F_{\mu\nu}^a F^{\mu\nu a}$
[quarks]	$+\bar{U}(i\not{\partial} - m_U)U + \bar{D}(i\not{\partial} - m_D)D$
[quark gluon]	$+igT^a(\bar{U}\not{A}^a U + \bar{D}\not{A}^a D)$
[3 gluons]	$+\frac{g}{2}(\partial_\mu A_\nu^a - \partial_\nu A_\mu^a)f^{abc}A^{b\mu}A^{c\nu}$
[4 gluons]	$-\frac{g^2}{4}f^{abc}f^{abd}A_\mu^b A_\nu^c A^{\mu\nu} A^{ad}$

The Higgs Sector

excluding GRAVITY

The Parameters

- 6 quark masses
 - m_u, m_c, m_t
 - m_d, m_s, m_b
- 3 lepton masses
 - m_e, m_μ, m_τ
- 2 vector boson masses
 - M_W, M_Z
 - $(m_\gamma, m_g=0)$
- 1 Higgs mass
 - M_h
- 3 coupling constants
 - G_F, α, α_s
- 3 quark mixing angles
 - $\theta_{12}, \theta_{23}, \theta_{13}$
- 1 quark phase
 - δ

The Standard Model in action

The Standard Model Effective Lagrangean

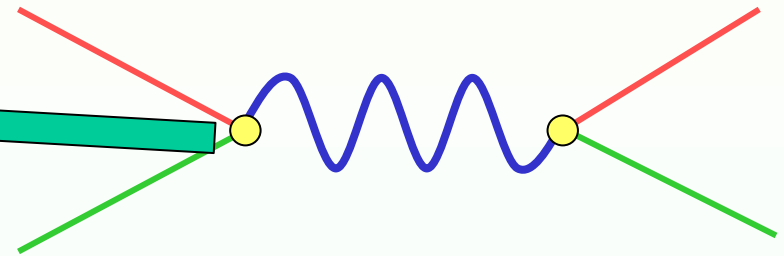
$$\mathcal{L}_{(\text{Standard Model})} =$$

- [W^\pm] $-\frac{1}{2}(\partial_\mu W_\nu - \partial_\nu W_\mu)(\partial^\mu W^{\nu\mu} - \partial^{\nu\mu} W^{\mu\nu}) + M_W^2 W_\mu W^{\mu\nu}$
- [Photon] $-\frac{1}{4}F_{\mu\nu}^A F^{A\mu\nu}$
- [Z^0] $-\frac{1}{2}F_{\mu\nu}^Z F^{Z\mu\nu} + \frac{1}{2}M_Z^2 Z_\mu Z^\mu$
- [ℓ, ν_ℓ] $+\bar{\ell}i\not{\partial}L_\ell + \bar{\nu}_\ell i\not{\partial}R_\ell - m_\ell\bar{\ell}\ell$
- [$W\ell\nu$] $-\frac{g}{\sqrt{2}}\bar{L}_\ell(\tau_+ W + \tau_- W)L_\ell$
- [$\gamma\ell^+\ell^-$] $+e_e/m\bar{\ell}A\ell$
- [$Z\ell^+\ell^-, Z\nu\bar{\nu}$] $-\frac{g}{\cos\theta_w}\bar{L}_\ell\left(\frac{\tau_3}{2}\cos\theta_w + \frac{1}{2}\sin^2\theta_w\right)\not{Z}L_\ell - \frac{g\sin^2\theta_w}{\cos\theta_w}\bar{R}_\ell\not{Z}R_\ell$
- [H] $+\frac{1}{2}\partial_\mu H\partial^\mu H - \frac{1}{2}\mu^2 H^2 - \frac{\lambda}{4}H^4$
- [HH&H W^+W^-] $+\frac{g^2}{8}\left(H^2 + \frac{2\mu}{\lambda}H\right)(2W_\mu W^{\mu\nu})$
- [HH&H ZZ] $+\frac{g^2}{8}\left(H^2 + \frac{2\mu}{\lambda}H\right)\left(\frac{1}{\cos^2\theta_w}Z_\mu Z^\mu\right)$
- [H $\ell^+\ell^-$] $-m_\ell\sqrt{2G_F}\bar{\ell}\ell H$
- [quark γ] $+Q\bar{q}Aq$
- [quark Z] $-\frac{g}{\cos\theta_w}\bar{L}_q\left(\frac{\tau_3}{2}\cos^2\theta_w + \frac{\sin^2\theta_w}{2}\right)\not{Z}L_q$
- [quark W] $-\frac{g}{\sqrt{2}}\bar{U}V_{CKM}(\tau_+ W + \tau_- W)D$
- [quark H] $-m_q\sqrt{2G_F}\bar{q}qH$
- [gluons] $-\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu}$
- [quarks] $+\bar{U}(i\not{\partial} - m_U)U + \bar{D}(i\not{\partial} - m_D)D$
- [quark gluon] $+\bar{q}gT^a(\bar{U}A^aU + \bar{D}A^aD)$
- [3 gluons] $+\frac{g}{2}(\partial_\mu A_\nu^a - \partial_\nu A_\mu^a)f^{abc}A^{b\mu}A^{c\nu}$
- [4 gluons] $-\frac{g^2}{4}f^{abc}f^{abd}A_\mu^b A_\nu^c A^{\mu\nu}A^d$

excluding GRAVITY

• Take a process

$$e^+e^- \rightarrow \mu^+\mu^-$$



$$4\pi\alpha^2/3s$$

α is the fine structure constant
 s is the (C.of.M Energy)²

(neglecting masses and $\sqrt{s} \ll M_Z$)

How good is the Standard Model?

The Standard Model Effective Lagrangian

$$\mathcal{L}(\text{Standard Model}) =$$

- [W[±]] $-\frac{1}{2}(\partial_\mu W_\nu - \partial_\nu W_\mu)(\partial^\mu W^{\nu\dagger} - \partial^\nu W^{\mu\dagger}) + M_W^2 W_\mu W^{\mu\dagger}$
- [Photon] $-\frac{1}{4}F_{\mu\nu}^A F^{A\mu\nu}$
- [Z⁰] $-\frac{F_{\mu\nu}^Z F^{Z\mu\nu}}{2} + \frac{1}{2}M_Z^2 Z_\mu Z^\mu$
- [ℓ, ν_ℓ] $+i\bar{L}_\ell \not{\partial} L_\ell + i\bar{R}_\ell \not{\partial} R_\ell - m_\ell \bar{\ell}\ell$
- [Wℓν] $-\frac{g}{\sqrt{2}}\bar{L}_\ell(\tau_+ W + \tau_- W)L_\ell$
- [γℓ⁺ℓ⁻] $+e_{\ell/m}\bar{\ell}\not{A}\ell$
- [Zℓ⁺ℓ⁻, Zνν̄] $-\frac{g}{\cos\theta_w}\bar{L}_\ell\left(\frac{\tau_3}{2}\cos^2\theta_w + \frac{1}{2}\sin^2\theta_w\right)\not{Z}L_\ell - \frac{g\sin^2\theta_w}{\cos\theta_w}\bar{R}_\ell\not{Z}R_\ell$
- [H] $+\frac{1}{2}\partial_\mu H\partial^\mu H - \frac{1}{2}\mu^2 H^2 - \frac{1}{2}\lambda\mu H^3 - \frac{1}{8}\lambda^2 H^4$
- [HH&H W⁺W⁻] $+\frac{g^2}{8}\left(H^2 + \frac{2\mu}{\lambda}H\right)(2W_\mu W^{\mu\dagger})$
- [HH&H ZZ] $+\frac{g^2}{8}\left(H^2 + \frac{2\mu}{\lambda}H\right)\left(\frac{1}{\cos^2\theta_w}Z_\mu Z^\mu\right)$
- [H ℓ⁺ℓ⁻] $-m_\ell\sqrt{2}G_F\bar{\ell}\ell H$
- [quark γ] $+Q\bar{q}\not{A}q$
- [quark Z] $-\frac{g}{\cos\theta_w}\bar{L}_q\left(\frac{\tau_3}{2}\cos^2\theta_w + \frac{\sin^2\theta_w}{2}\right)\not{Z}L_q$
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- [gluons] $-\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu}$
- [quarks] $+\bar{U}(i\not{\partial} - m_U)U + \bar{D}(i\not{\partial} - m_D)\mathcal{D}$
- [quark gluon] $+igT^a(\bar{U}\not{A}^a U + \bar{D}\not{A}^a \mathcal{D})$
- [3 gluons] $+\frac{g}{2}(\partial_\mu A_\nu^a - \partial_\nu A_\mu^a)f^{abc}A^{b\mu}A^{c\nu}$
- [4 gluons] $-\frac{g^2}{4}f^{abc}f^{abd}A_\mu^b A_\nu^c A^{\mu\nu} A^d$

18 measurements
5 free parameters
 $\chi^2 = 18.1/13$ d.o.f.
 $3 > 1\sigma$
 $1 > 2\sigma$
Almost too good!

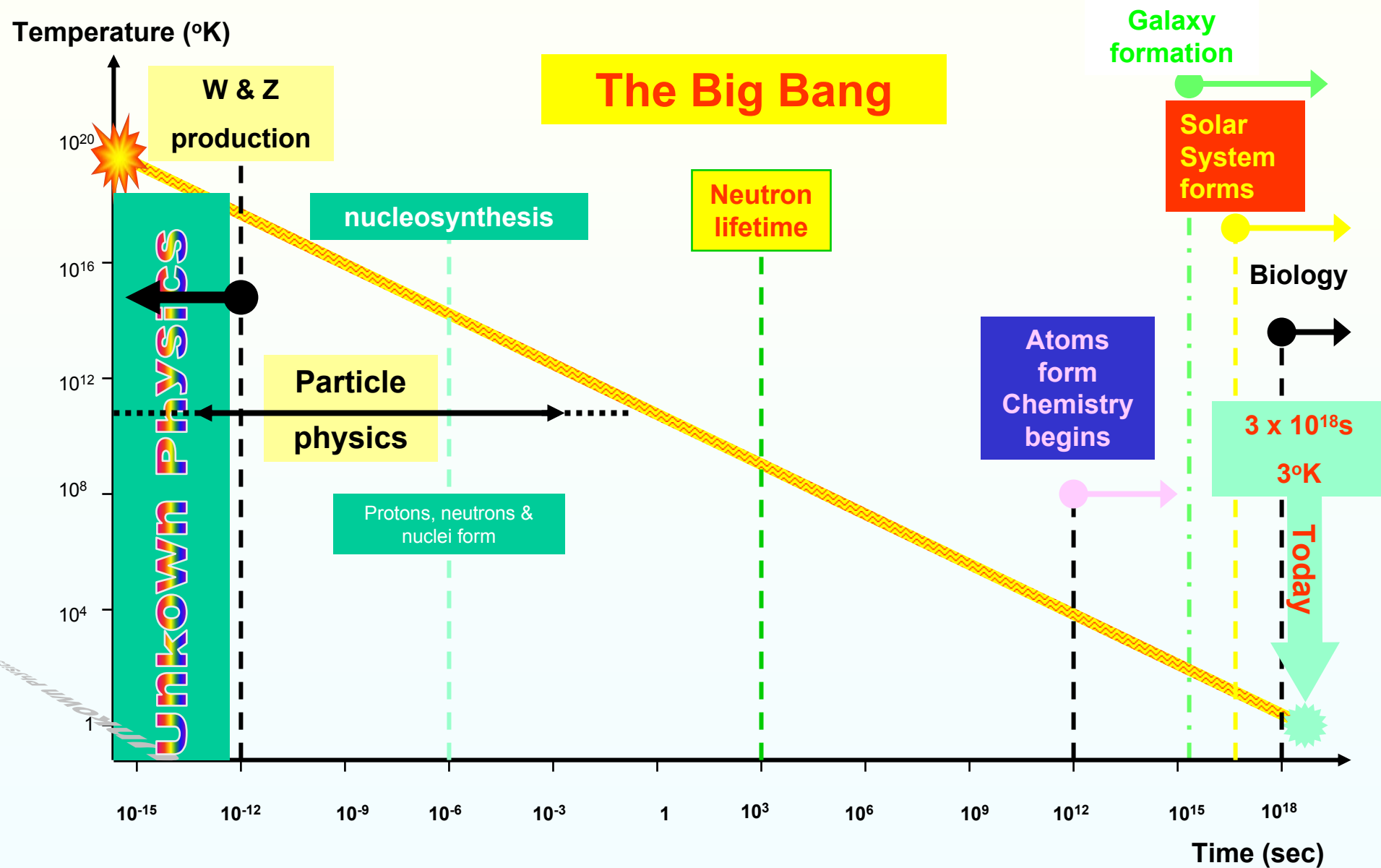
$\Delta\alpha_{\text{had}}^{(5)}(m_Z)$	0.02768
m_Z [GeV]	91.1875
Γ_Z [GeV]	2.4957
σ_{had}^0 [nb]	41.477
R_l	20.744
$A_{\text{fb}}^{0,l}$	0.01645
$A_l(P_\tau)$	0.1481
R_b	0.21586
R_c	0.1722
$A_{\text{fb}}^{0,b}$	0.1038
$A_{\text{fb}}^{0,c}$	0.0742
A_b	0.935
A_c	0.668
$A_l(\text{SLD})$	0.1481
$\sin^2\theta_{\text{eff}}^{\text{lept}}(Q_{\text{fb}})$	0.2314
m_W [GeV]	80.374
Γ_W [GeV]	2.091
m_t [GeV]	171.3

excluding GRAVITY

What remains to be done?

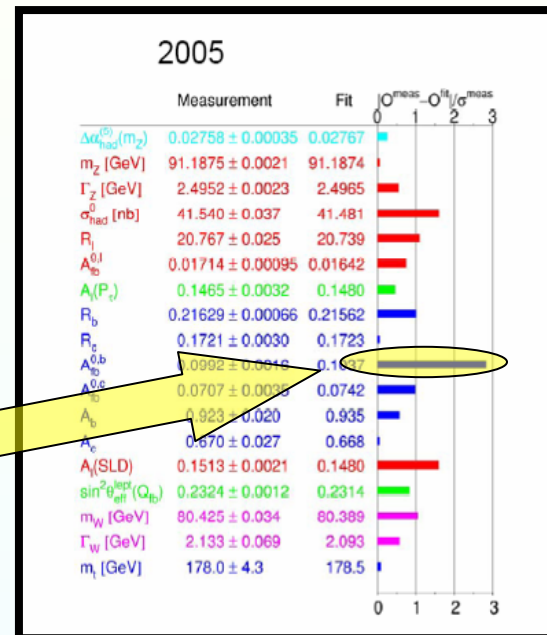
- The Standard Model is a very good *description* of the Universe at the particle scale ($\sim 2M_W$)
 - But does not *explain* many things
 - Why so many particles?
 - Why so many forces?
 - What is mass?
 - Why do particles have the masses they have?
 - How do neutrinos get mass?
 - Are neutrinos different? How do they fit in?
 - What is Dark Matter? Dark Energy?
 - Why is matter different from antimatter?
 - (Where did all the antimatter go?)
 - Where does gravity fit in?

The state of the Universe



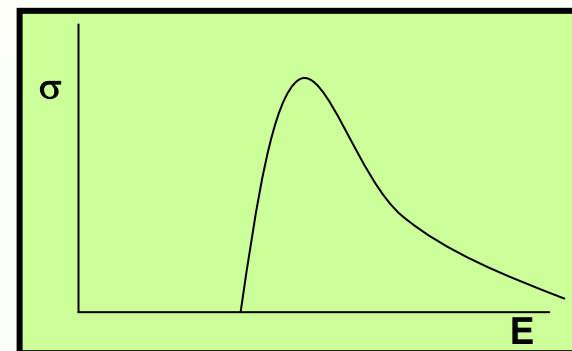
What do we need to make progress?

- **To reach higher energy**
 - To take us beyond the LEP/Tevatron energy scale
 - $\sim 100\text{-}200\text{GeV}$
- **To reach higher precision**
 - $10 \times$ statistics would make this effect (if real) 8σ
- **New types of accelerator**
 - Neutrino factories
 - Beta beams
 - Muon colliders ...



- We can accelerate *stable* particles
 - “*Stable*” means “with a lifetime long enough to capture and accelerate them”
 - in practice, $> \sim \mu\text{-second}$
- Hadrons
 - p, d, t, α, \dots nuclei (up to Pb) & antiprotons
 - Hadrons contain “partons” (quarks, gluons...)
- Leptons
 - e^\pm, μ^\pm
 - Leptons are “point-like”
 - (at our present energy scales)

- The *Energy* must be sufficiently high that the process of interest can occur



- The *Luminosity* must be sufficiently high that a sufficient number of events are obtained in a “reasonable” time
 – (a few years)

For fixed target (esp. neutrino experiments) the equivalent parameter is

Beam Power or **Protons on Target (POT)**

$$N_{ev} = L \times \sigma \times t$$

$$t \sim 10^7 \text{ s/year}$$

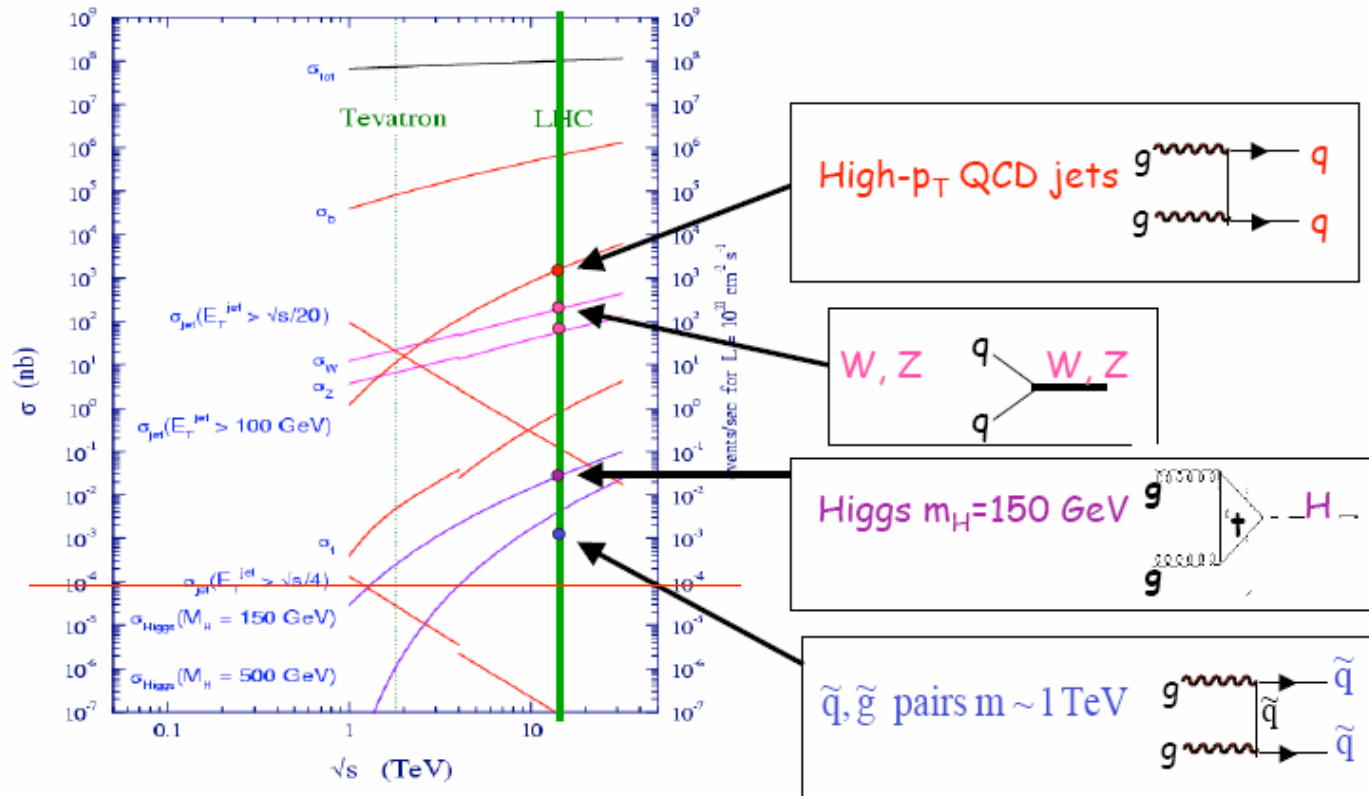
$$\sigma \sim \text{pb} (10^{-36} \text{ cm}^2)$$

For 1000 events in 1 year requires

$$L \sim 10^{32} \text{ cm}^2\text{s}^{-1}$$

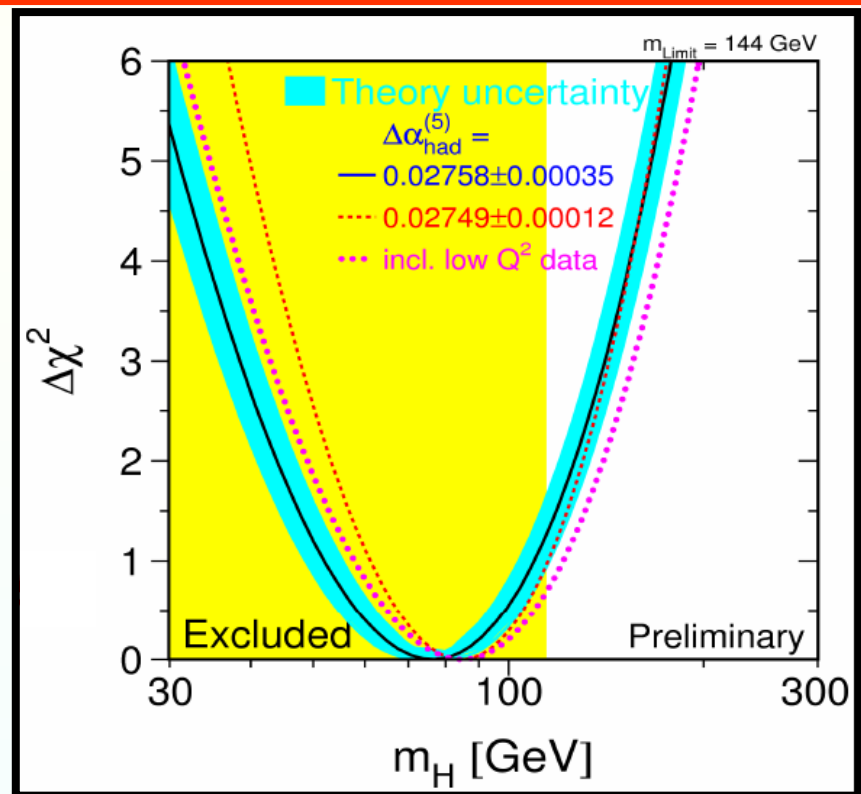
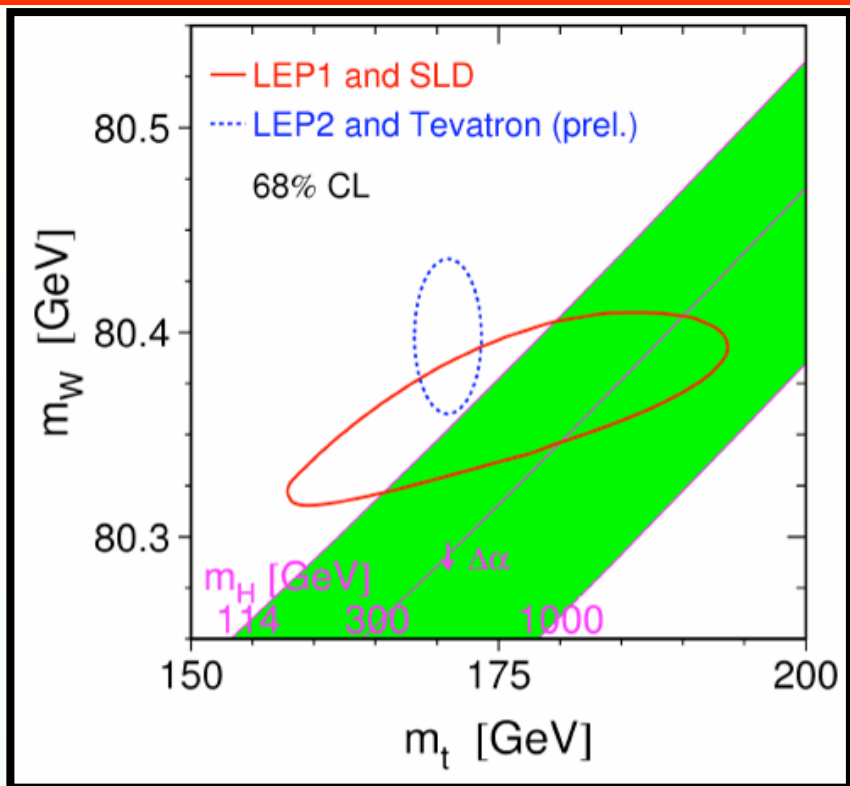
An example - the LHC

③ Huge (QCD) backgrounds (consequence of high energy ..)



- No hope to observe light objects ($W, Z, H ?$) in fully-hadronic final states \rightarrow rely on l, γ
- Fully-hadronic final states (e.g. $q^* \rightarrow qg$) can be extracted from backgrounds only with hard $O(100 \text{ GeV})$ p_T cuts \rightarrow works only for heavy objects
- Mass resolutions of $\sim 1\%$ (10%) needed for l, γ (jets) to extract tiny signals from backgrounds
- Excellent particle identification: e.g. e/jet separation

What are the big issues?



$M_H = 76^{+33}_{-24}$ GeV
 Incl. theory uncertainty:
 $M_H < 144$ GeV (95%CL)

Direct search limit (LEP-2):
 $M_H > 114$ GeV (95%CL)

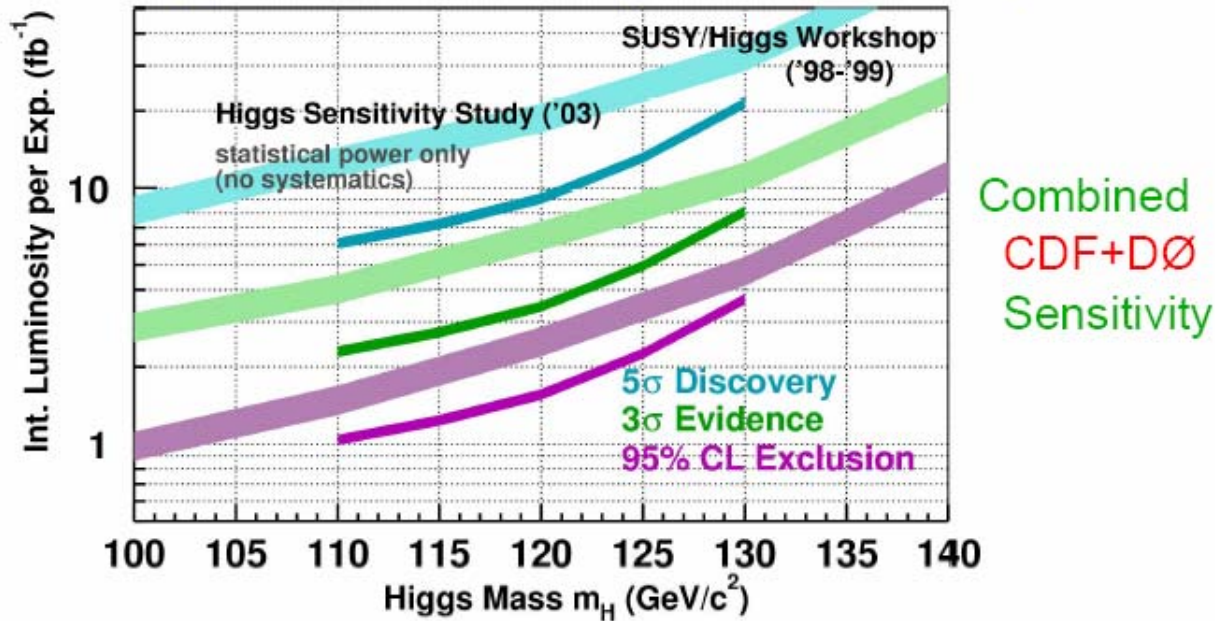
Probability $M_H > 114$ GeV: 15%

Renormalise probability
 for $M_H > 114$ GeV to 100%:
 $M_H < 182$ GeV (95%CL)

Higgs at the Tevatron

Standard Model Higgs Search

Combining production and decay channels and experiments:



Expectations:

- With 2/fb exclusion up to 123 GeV
- With 10/fb discovery up to 121 GeV

Currently: ~1/fb on tape

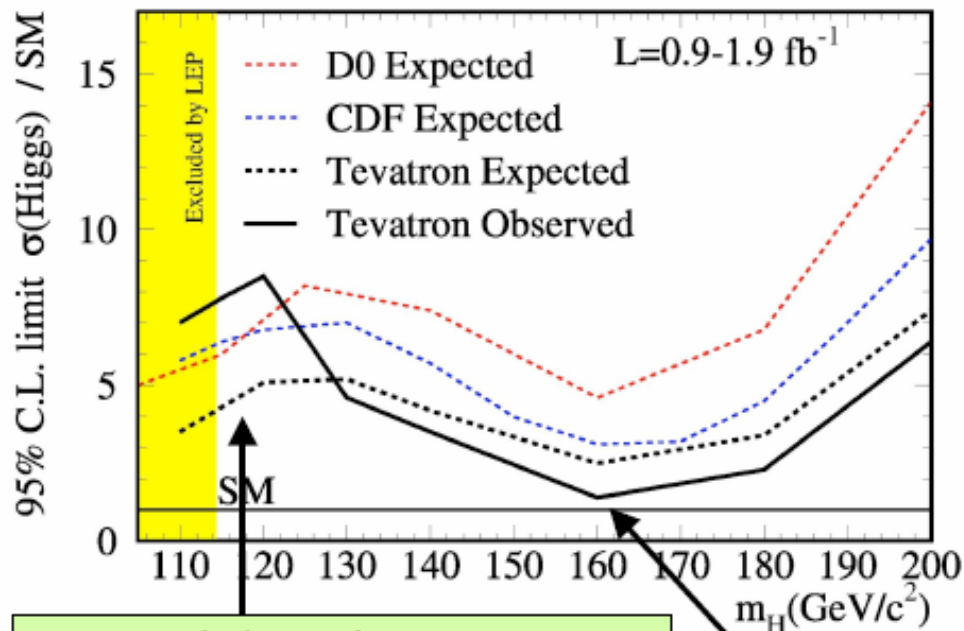
This was 3 years ago!!!

Where are they now?

The Tevatron Search today

What about the "competition" with Tevatron ?

Tevatron Run II Preliminary



Today : $\sim 2.8 \text{ fb}^{-1}$ /expt recorded
 End 2009: expect $6-7 \text{ fb}^{-1}$ /expt
 Operation beyond 2009 being discussed

With 7 fb^{-1} :

- 95% C.L. exclusion 150-180 GeV and $< 135 \text{ GeV}$ (if ~ 4 analysis improvement)
- 2.5σ evidence 155-170 GeV
- 3σ evidence up to 128 GeV (if ~ 10 analysis improvement)
- no 5σ sensitivity

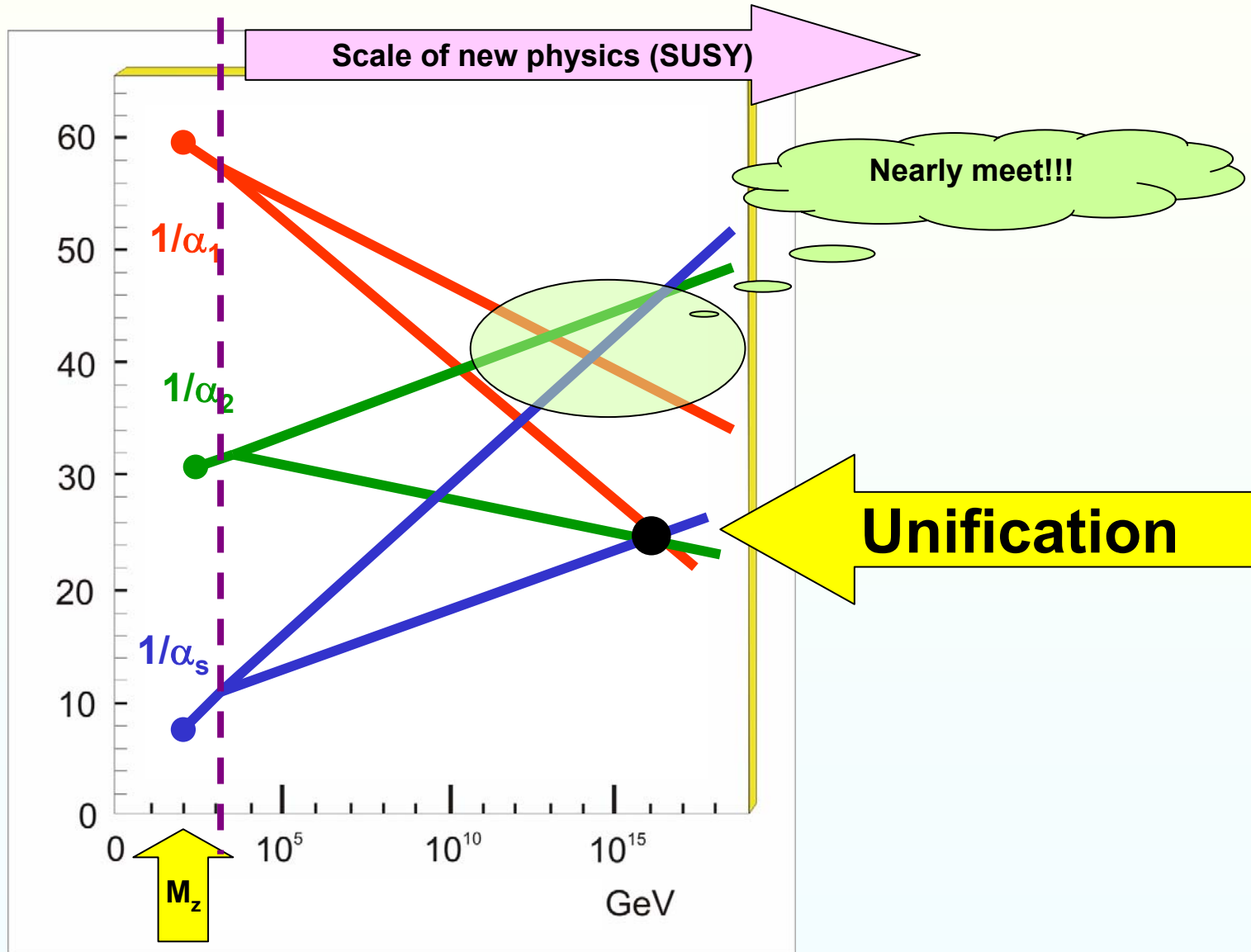
Note: big difference in statistics to go from exclusion to evidence (sophisticated cross-checks required ...)

- Several channels: $WH \rightarrow l\nu b\bar{b}, ZH \rightarrow \nu\nu b\bar{b}, \text{ etc.}$
- Expect analysis improvements (b-tagging, mass resolution, ..)
- With 7 fb^{-1} need: ~ 4 (10) improvement for 95% C.L. (3σ)

- 1 dominant channel: $H \rightarrow WW \rightarrow l\nu l\nu$ (counting channel)
- 3.8 fb^{-1} /expt for 95% C.L. exclusion (mid 2008 ?)
- end 2009: 2.5σ (6 fb^{-1}) to 3σ (8.5 fb^{-1}) evidence

After Gianotti, 07; Plot from Kim LP07

Unification of the forces?





The Large Hadron Collider

The Linear (e^+e^-) Collider

The Muon Collider

The Neutrino Factory

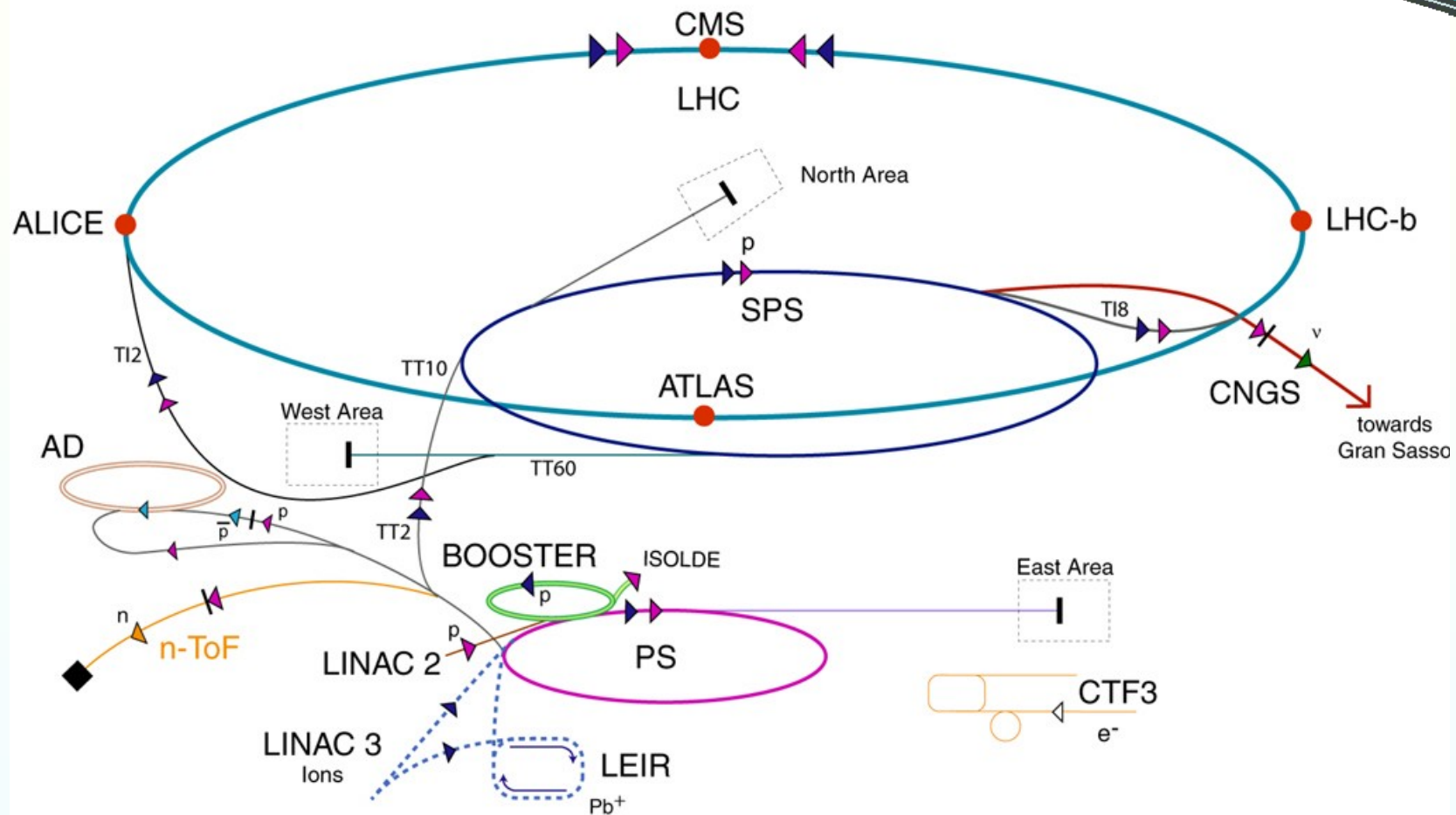
EURISOL and Beta Beams

“Factories” (ϕ , τ , c , b)



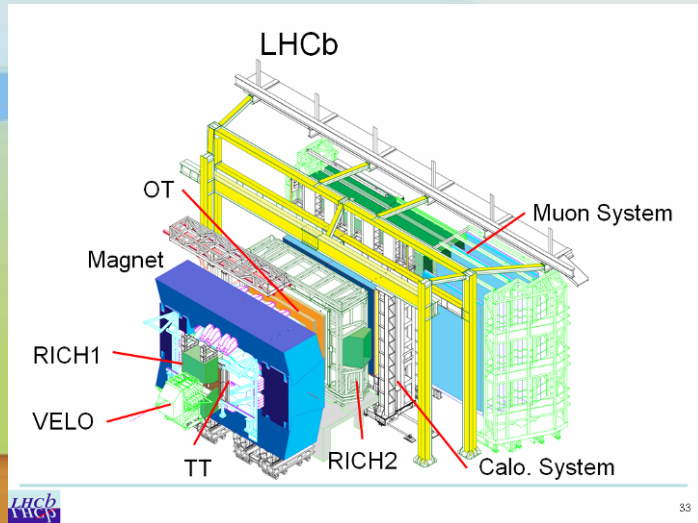
The Large Hadron Collider

- **The two main goals are:**
 - **Find the Higgs**
 - *If it exists!!!*
 - **Find the new physics**
 - *If it exists!!!*
- **We know ~ the energy scales**
 - $M_H < 250\text{GeV}$; $E_{NP} < 1\text{TeV}$
- **pp collisions at high energy**
 - **Collision energy ~10% of total energy**
 - Need a total collision energy $> 10\text{TeV}$
 - **Can calculate the cross-sections**
 - Need a luminosity $> 10^{33}\text{cm}^2/\text{s}$
- **The Large Hadron Collider (LHC) @ CERN**
 - $E \sim 14\text{TeV}$; $L \sim 10^{34}\text{cm}^2/\text{s}$

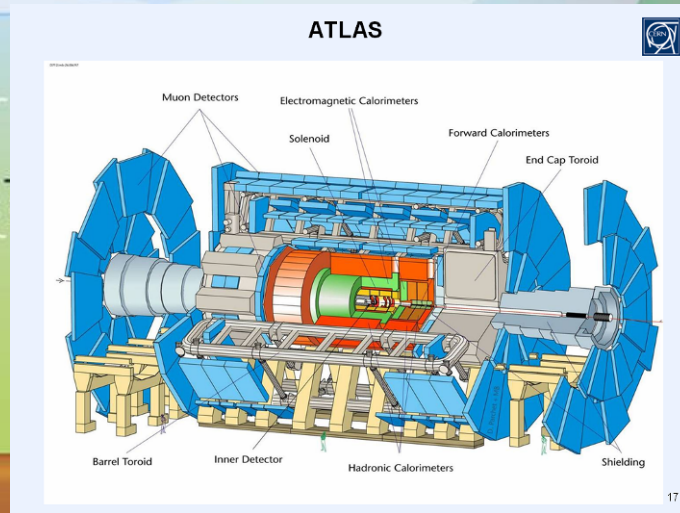


- | | | | |
|------------|---------------|------------------------------|--------------------------------|
| ▶ protons | ▶ antiprotons | AD Antiproton Decelerator | LHC Large Hadron Collider |
| ▶ ions | ▶ electrons | PS Proton Synchrotron | n-ToF Neutron Time of Flight |
| ▶ neutrons | ▶ neutrinos | SPS Super Proton Synchrotron | CNGS CERN Neutrinos Gran Sasso |
| | | | CTF3 CLIC Test Facility 3 |

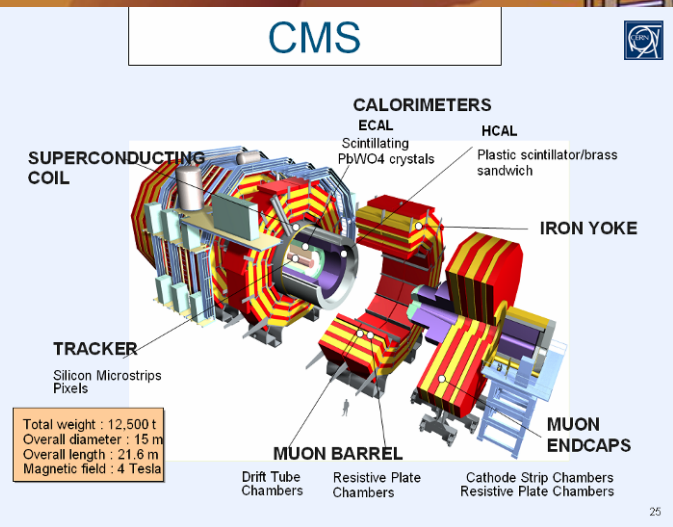
The Large Hadron Collider



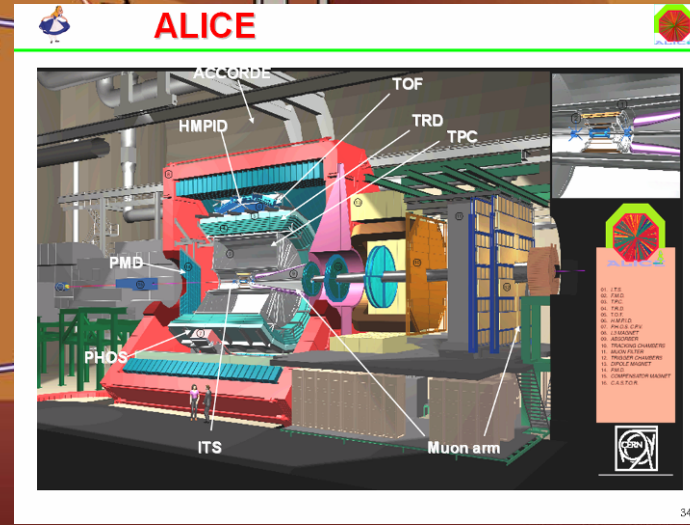
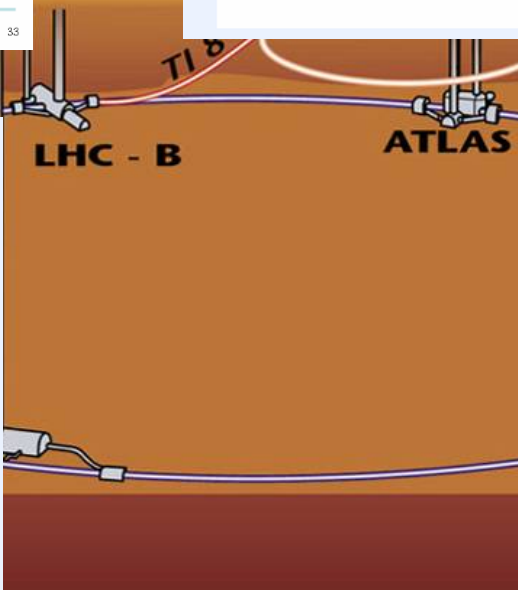
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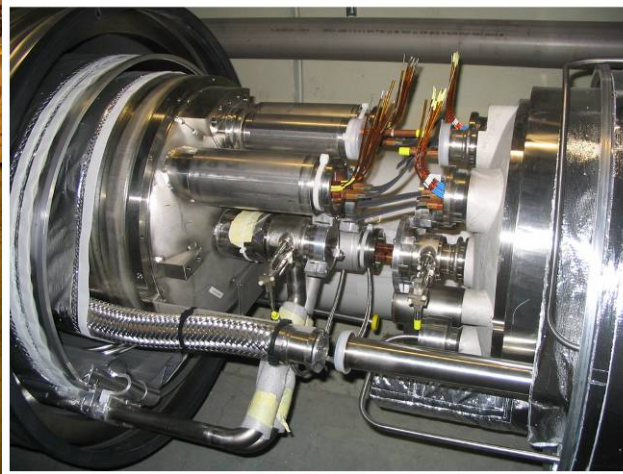


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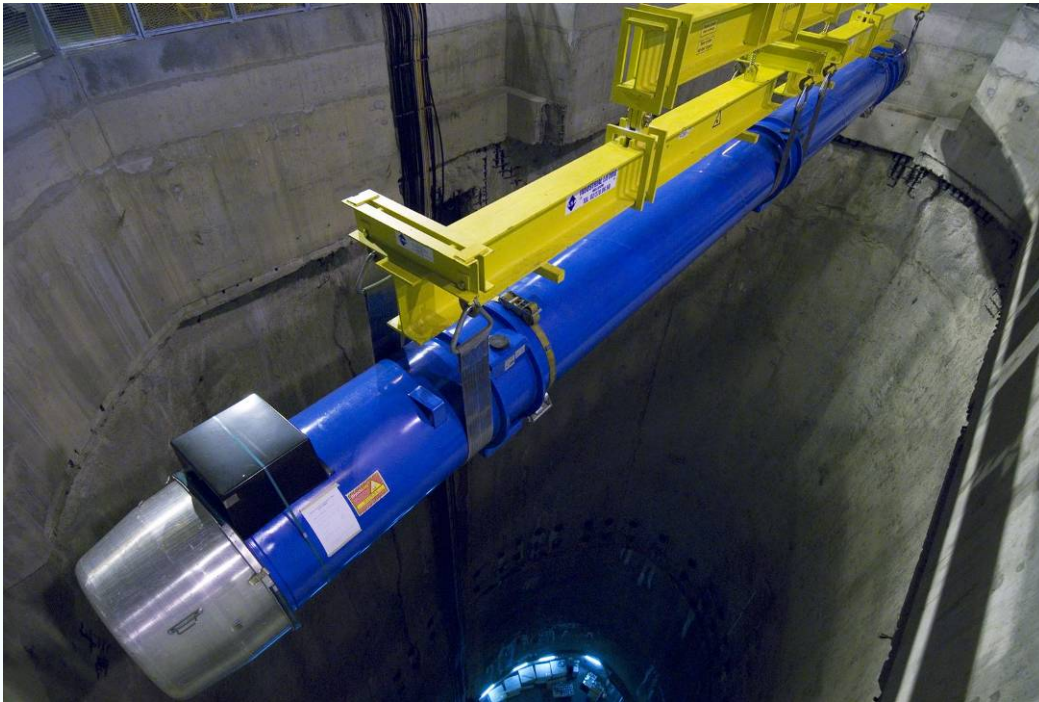
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The LHC installation





Descent of the last magnet, 26 April 2007

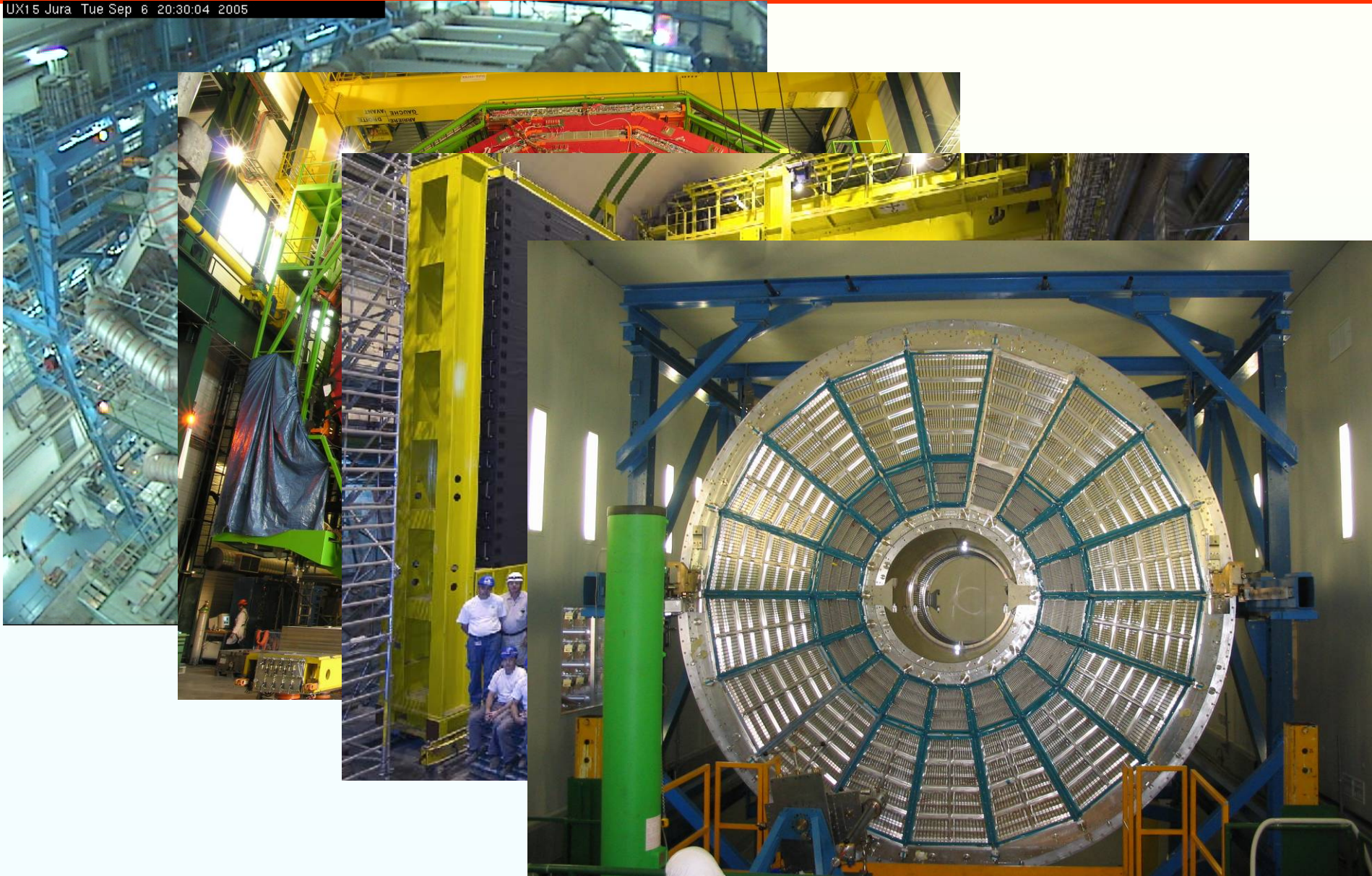


30'000 km underground at 2 km/h!



ATLAS, CMS, LHCb, ALICE

UX15 Jura Tue Sep 6 20:30:04 2005



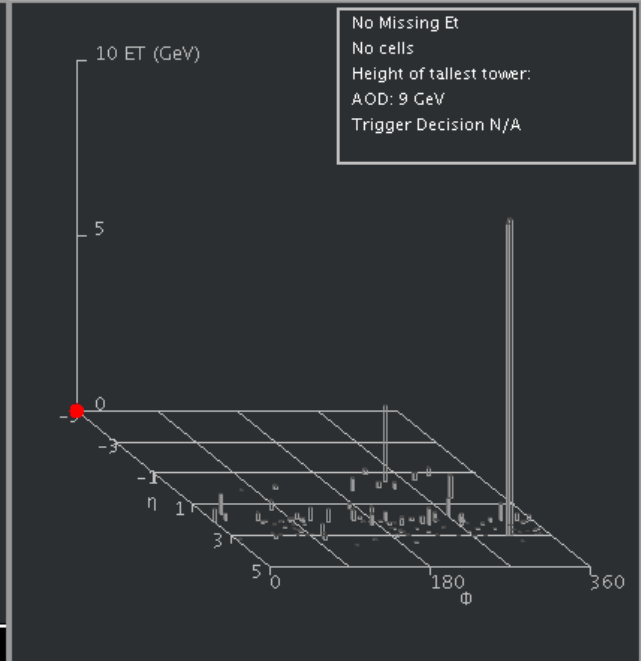
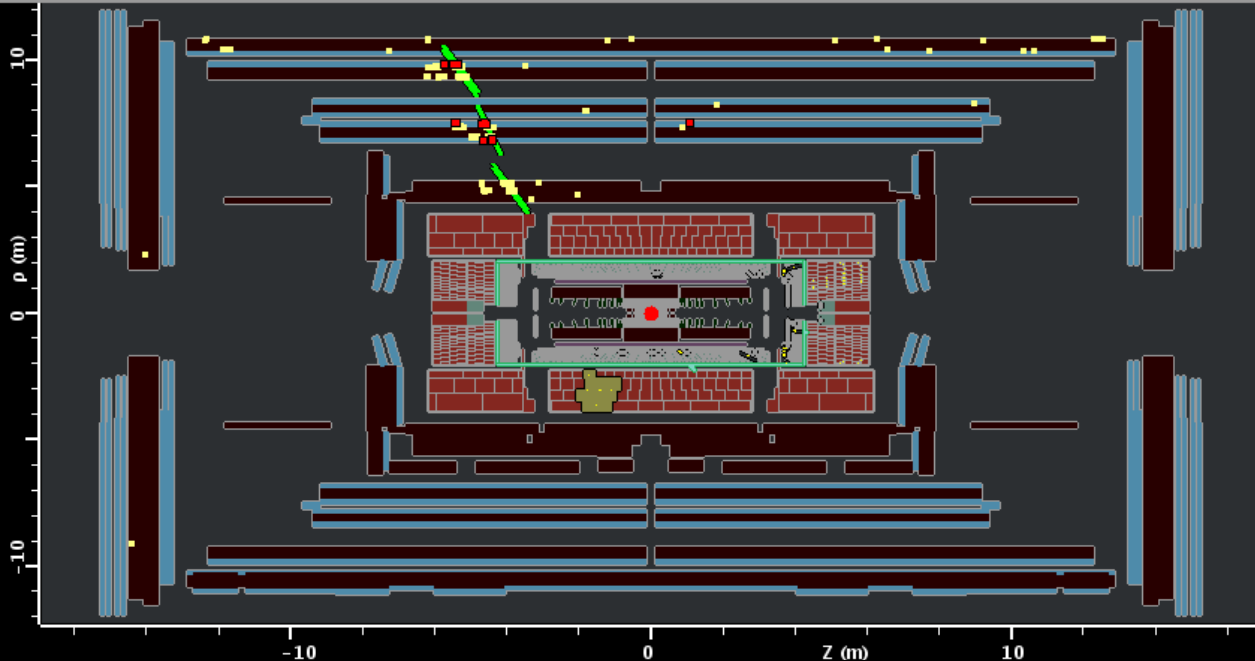
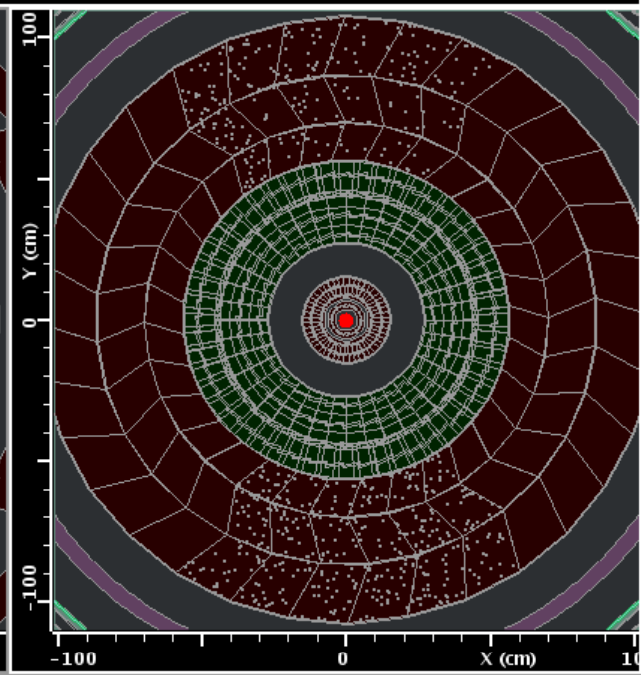
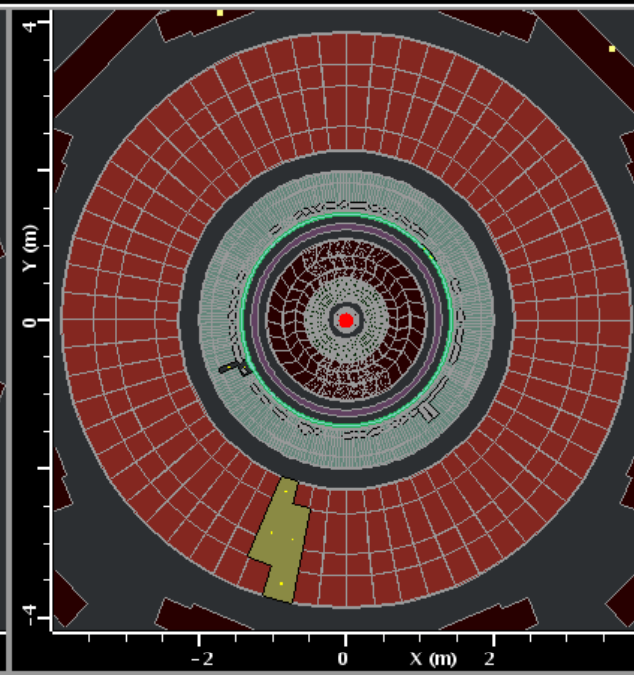
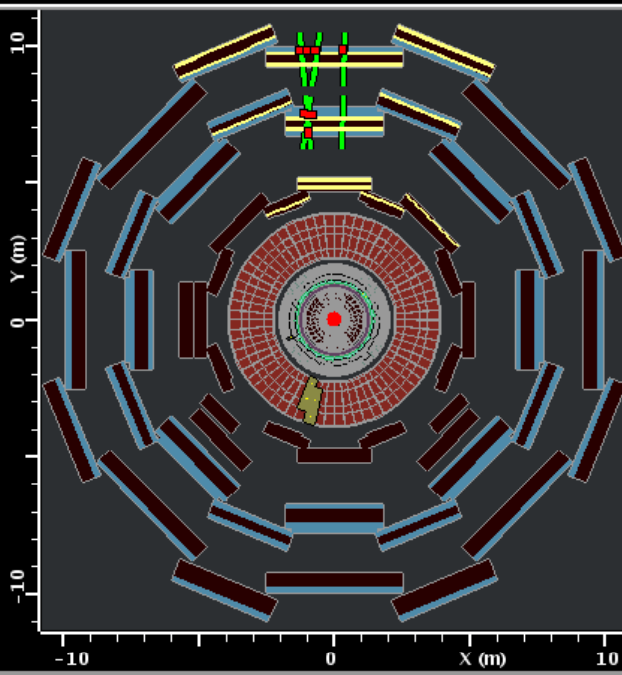
CMS Cavern



September 2007

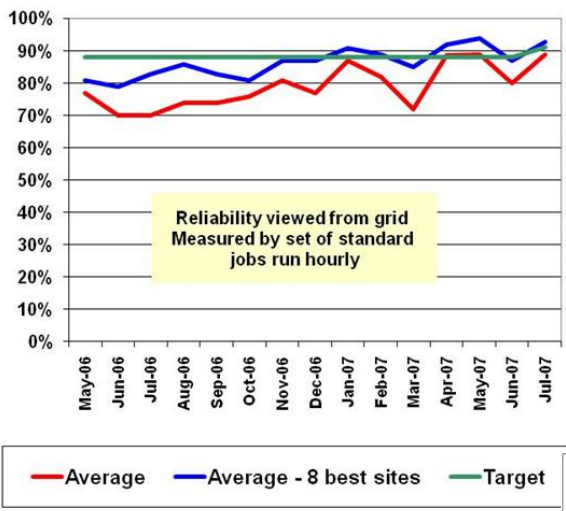
UX15 Geneva Wed Sep 26 10:00:03 2007



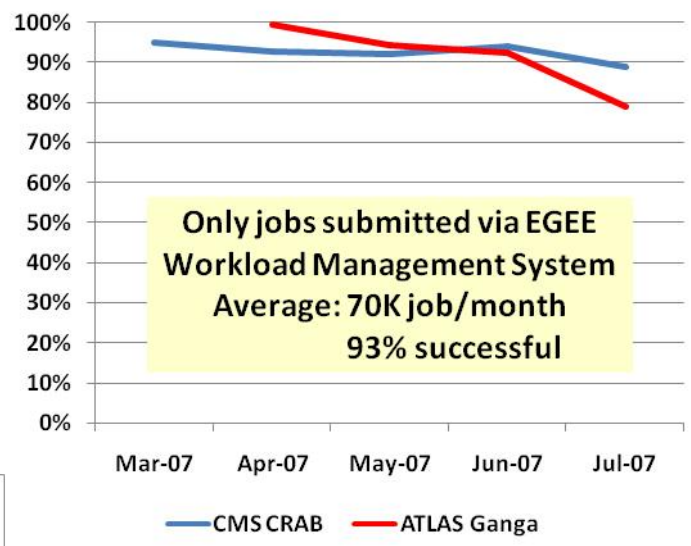




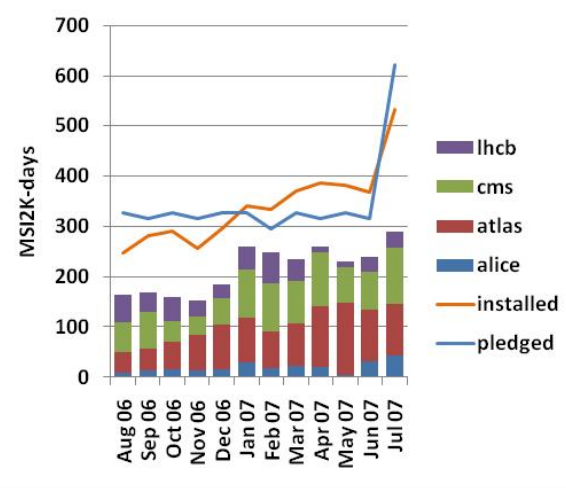
Site Reliability CERN + Tier-1s



User Job Efficiency



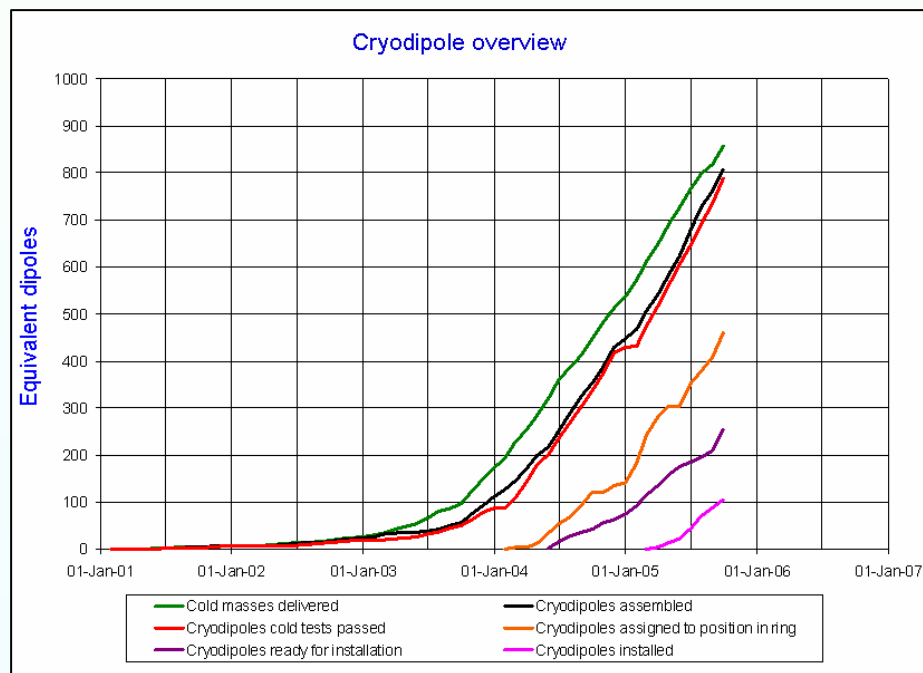
CERN + Tier-1s - CPU accounted to LHC VOs



Efficiency

Capacity

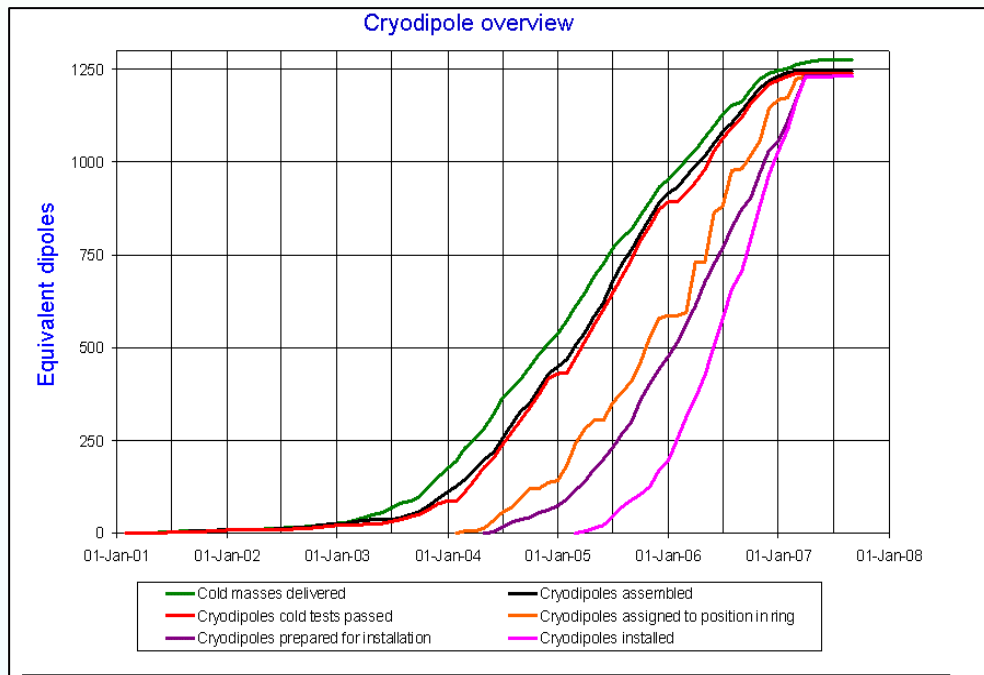
- Machine, experiments now installing
- Due for completion in summer 2007
- First collisions end summer 2007
- First results 2008
 - Higgs, SUSY
 - or something else



Updated 30 Sep 2005

Data provided by D. Tommasini AT-MAS, L. Bottura AT-MTM

- Machine installed, commissioning
- Experiments nearly installed, commissioning
- Due for completion in ~~summer 2007~~ **May 08**
- First collisions end ~~summer 2007~~ **8**
- First results ~~2008~~ **9**
 - **Higgs, SUSY**
 - or something else



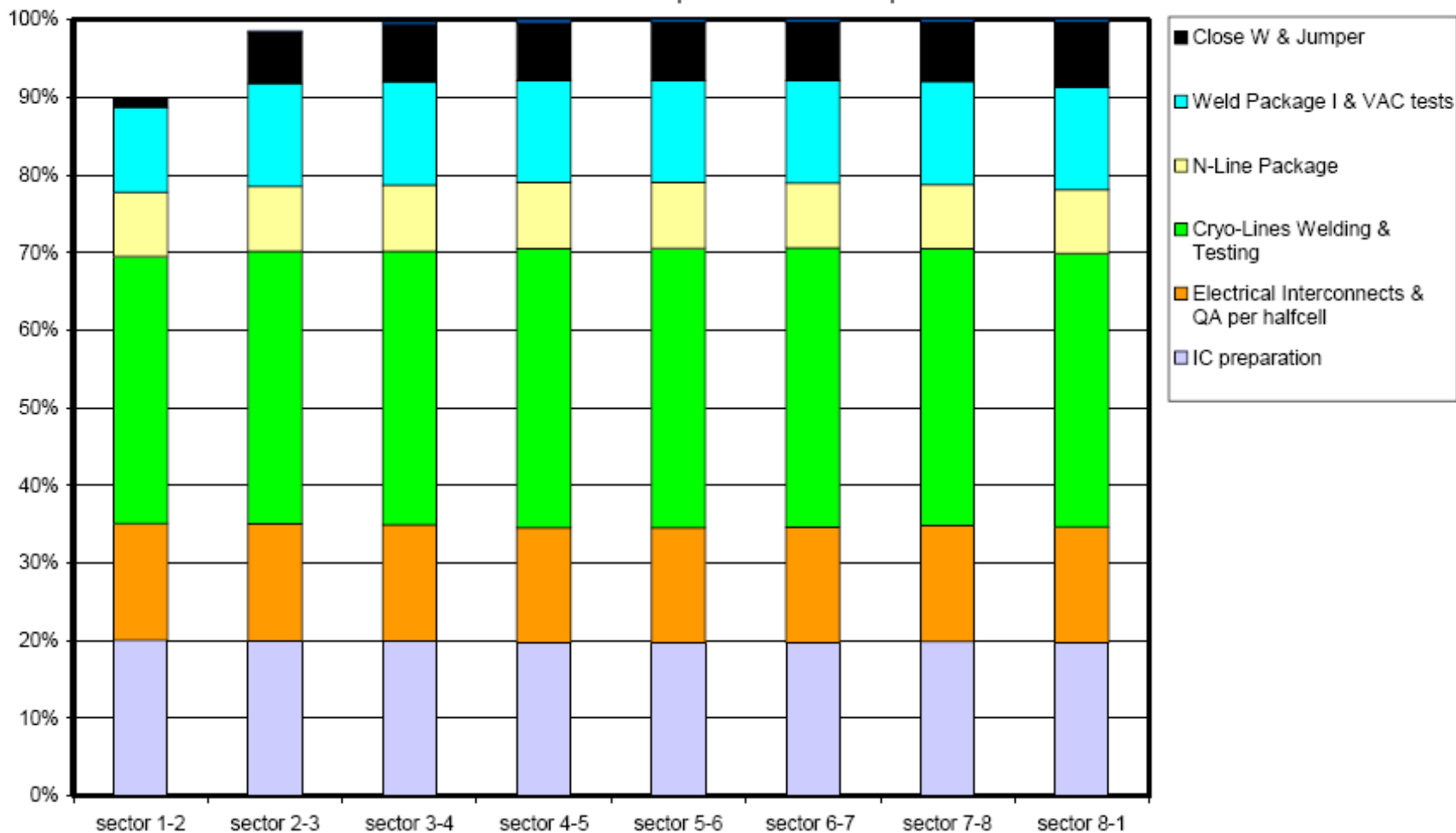
Updated 31 August 2007

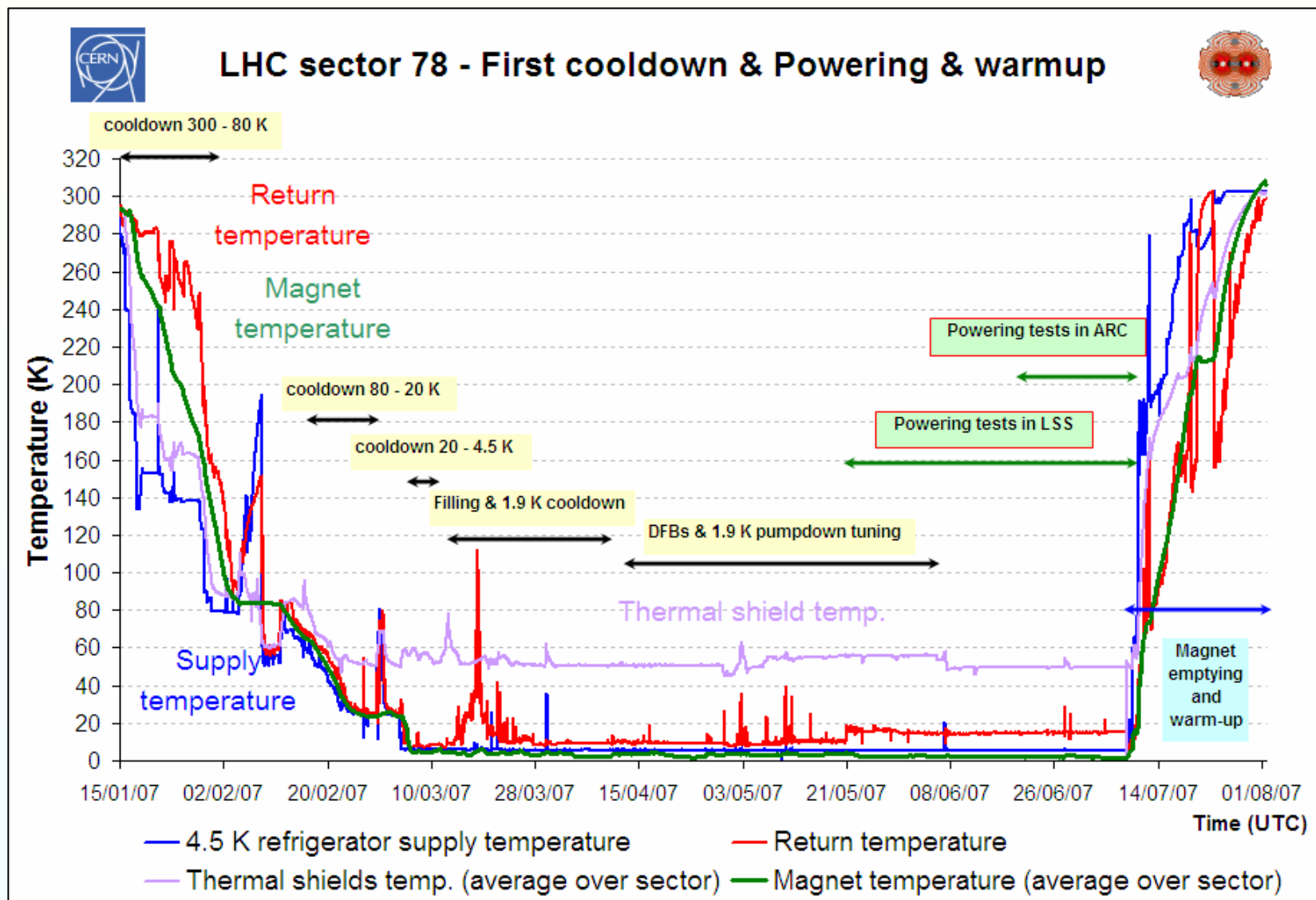
Data provided by D. Tommasini AT-MCS, L. Bottura AT-MTM



Magnet interconnections

General Advancement of Interconnects per Sector 10-September-2007

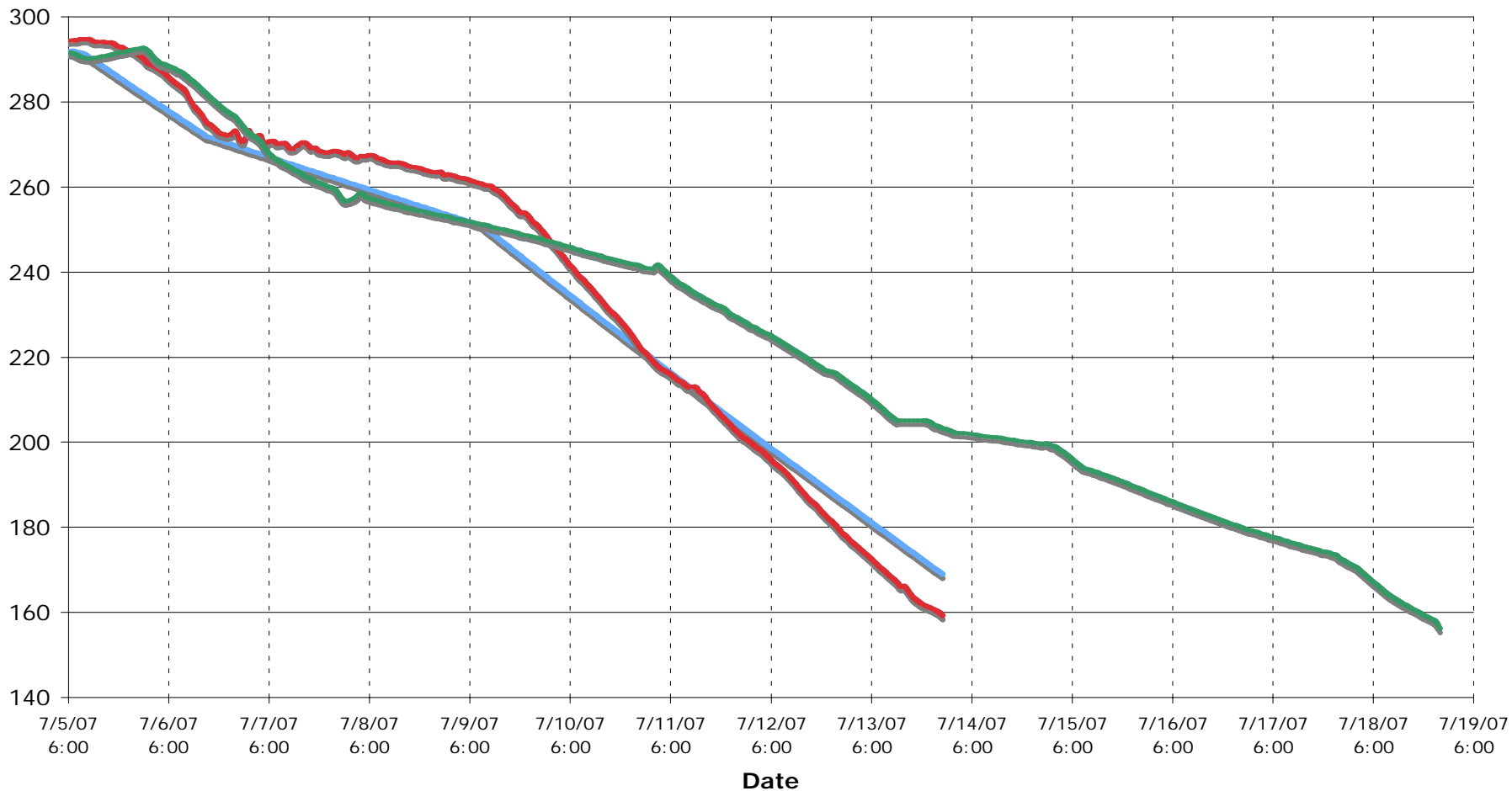






Cooldown until Saturday 14-7-7 - Cooldown speed

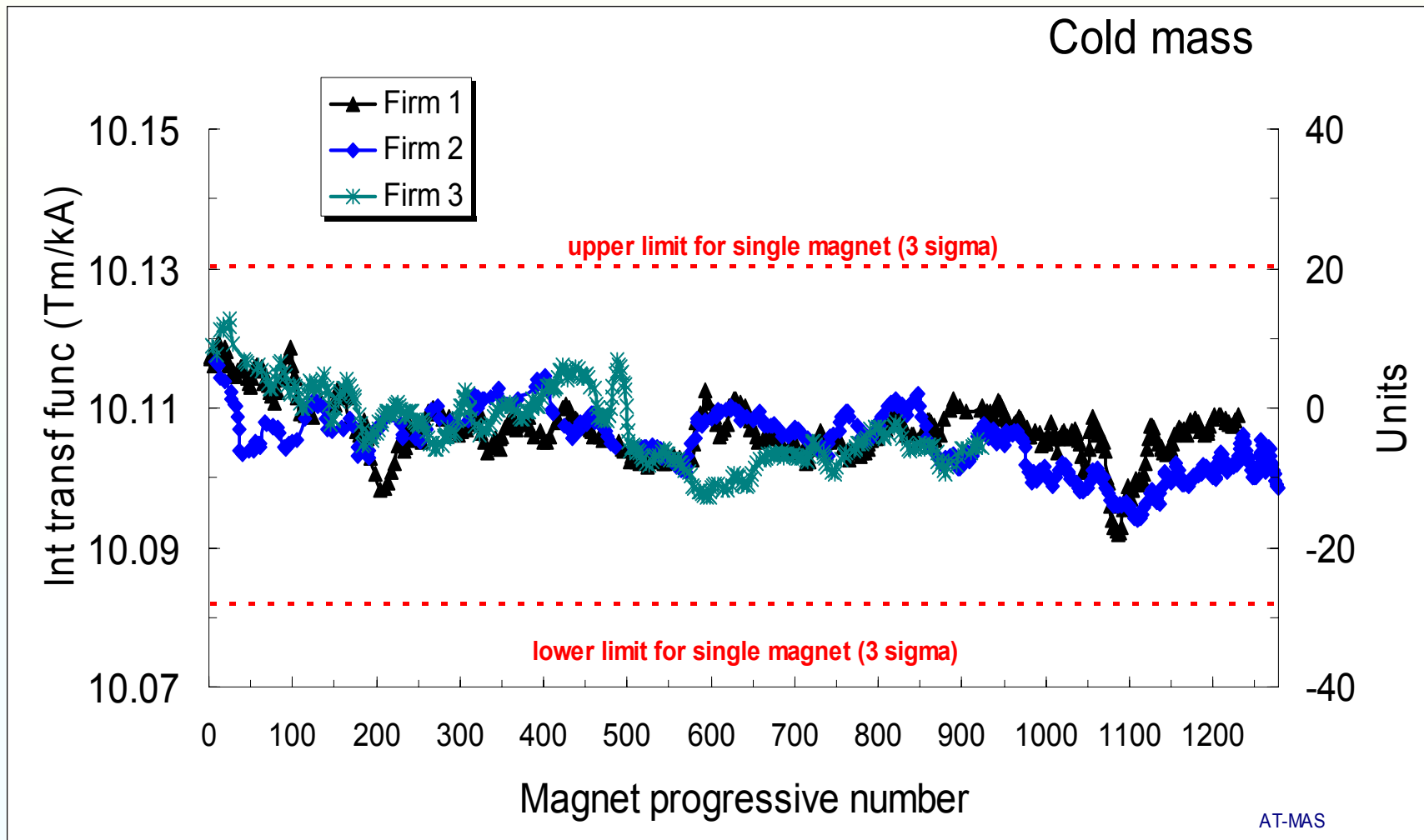
Avg. Cold Mass Temperature



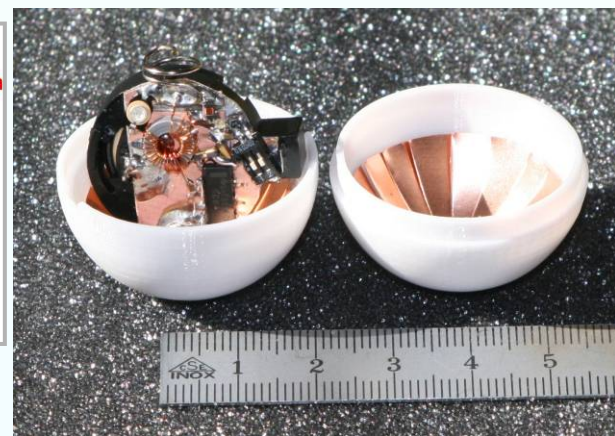
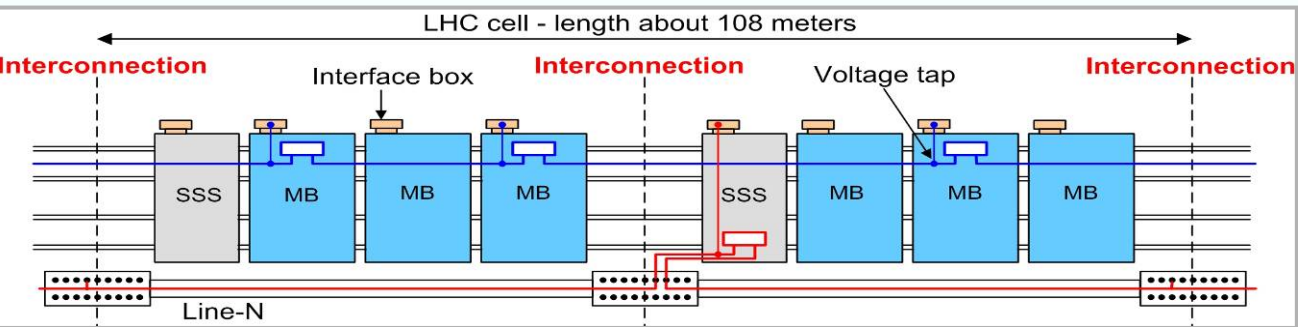
— planned [K] — measured [K] — data S7-8 [K]



Bending strength of dipoles



Module at room temperature ...

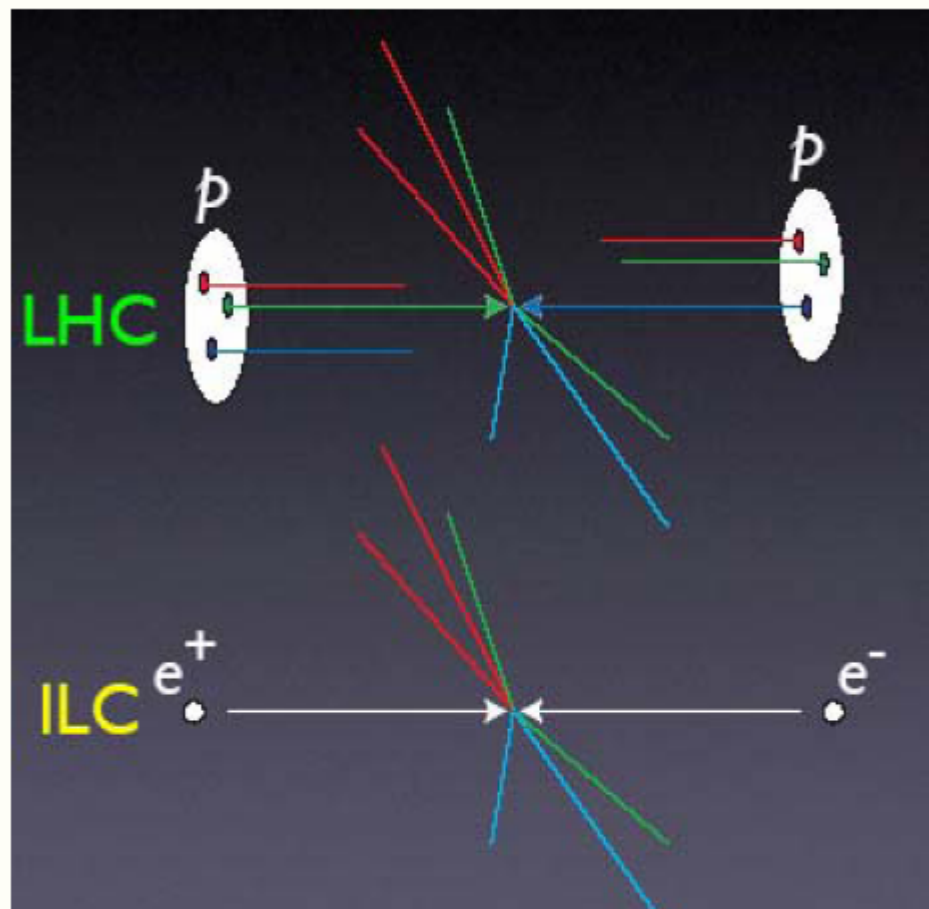




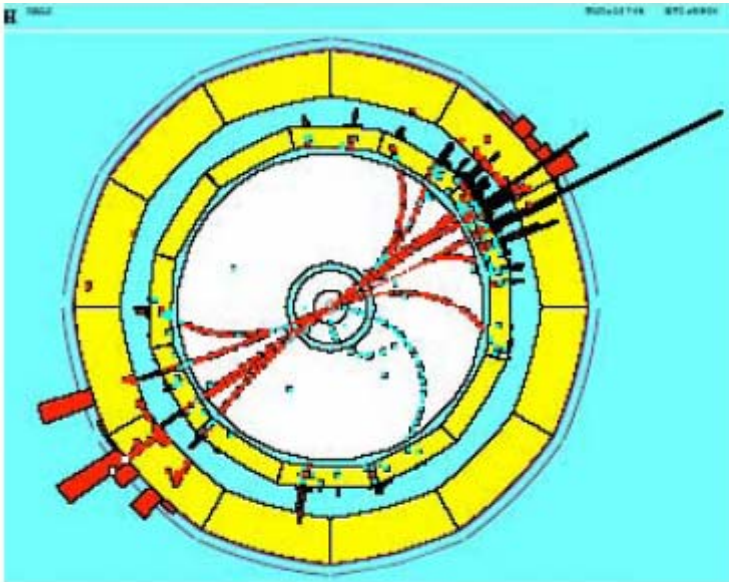
The Linear (e^+e^-) Collider

Why an e+e- collider?

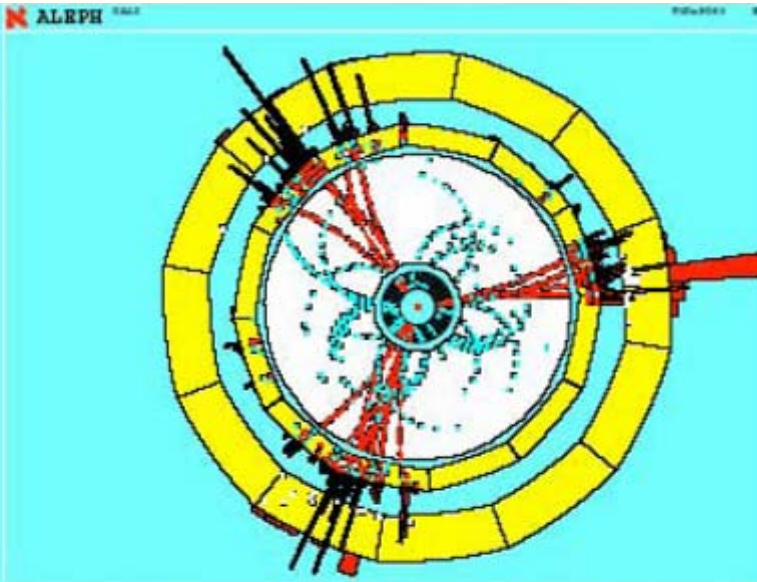
- elementary particles
- well-defined
 - energy,
 - angular momentum
- uses full COM energy
- produces particles democratically
- can mostly fully reconstruct events



Why an e+e- collider?



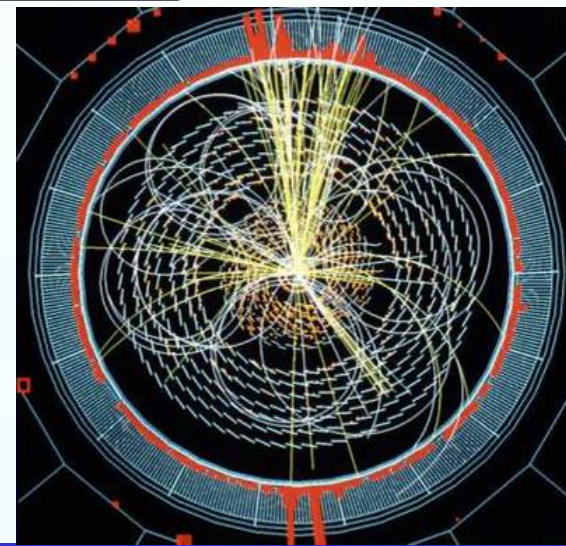
can see quarks



and a gluon ~1980

← LEP

LHC →



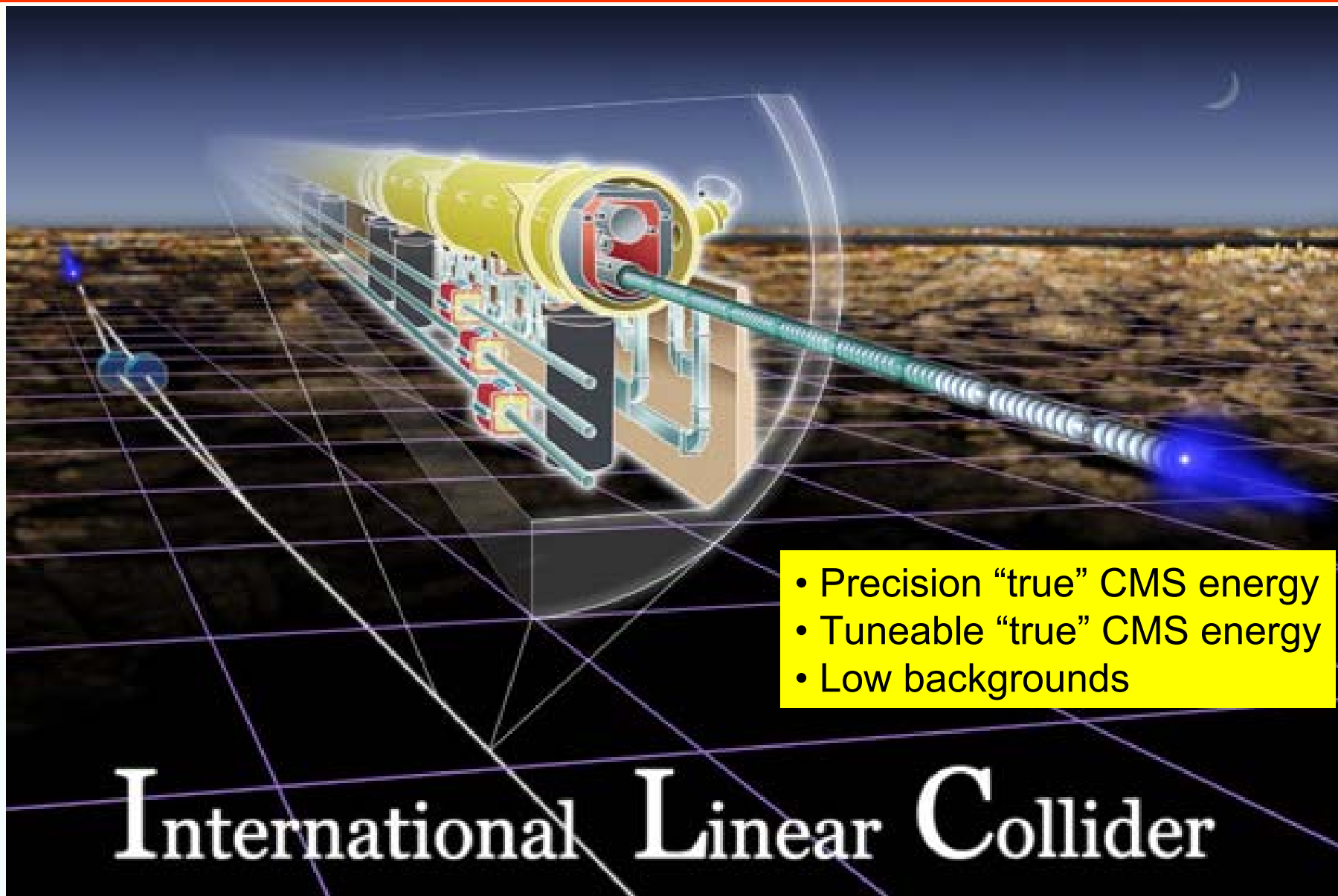
After Barry Barish

Synchrotron Radiation!

or rather

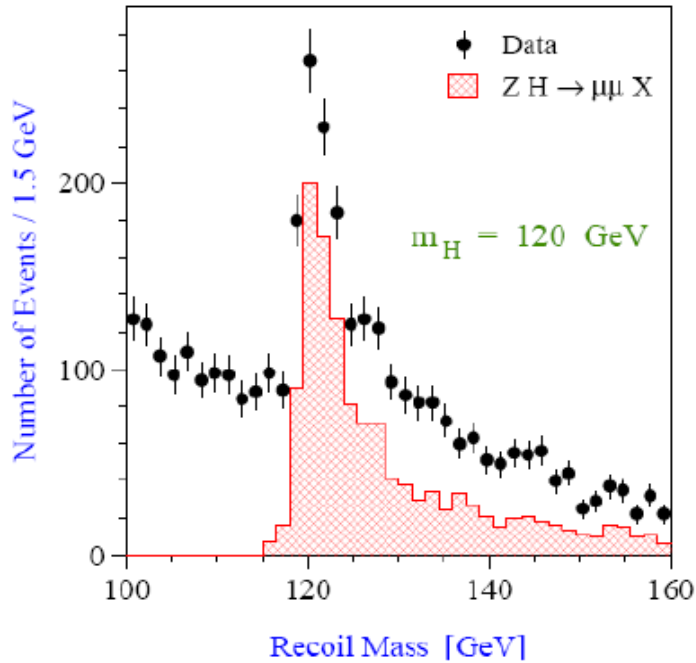
the lack of it in a linear machine

Key ILC Properties

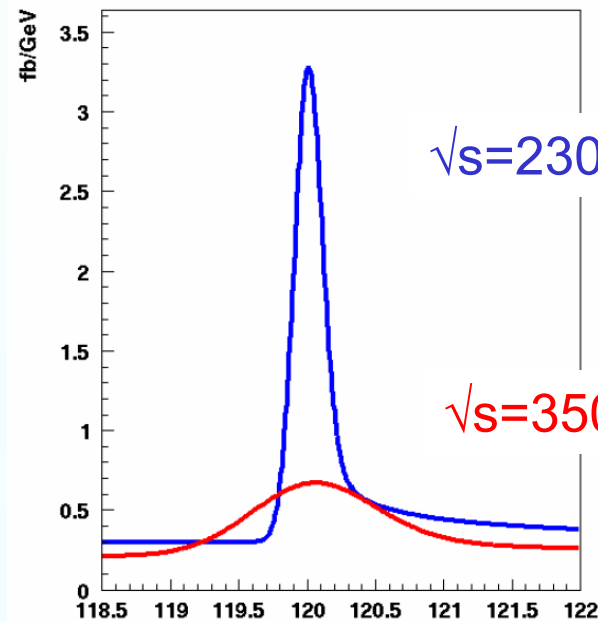
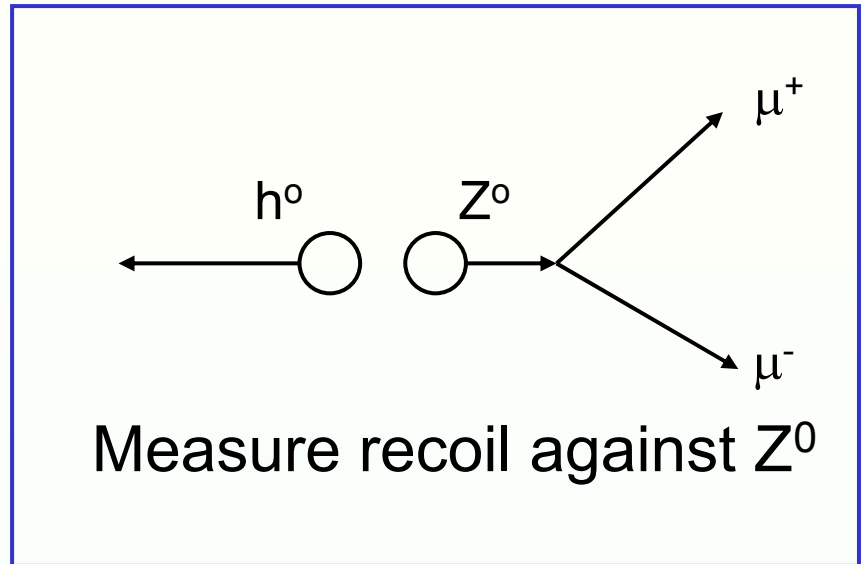


- Precision “true” CMS energy
- Tuneable “true” CMS energy
- Low backgrounds

International Linear Collider



120 GeV Higgs:
Advantages of
running at lower than
top threshold:



Bambade et al.



RDR vs ILC Physics Goals

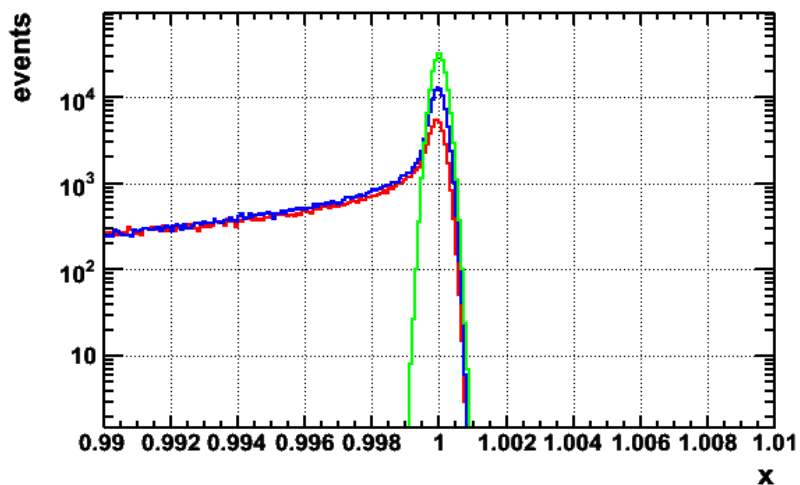
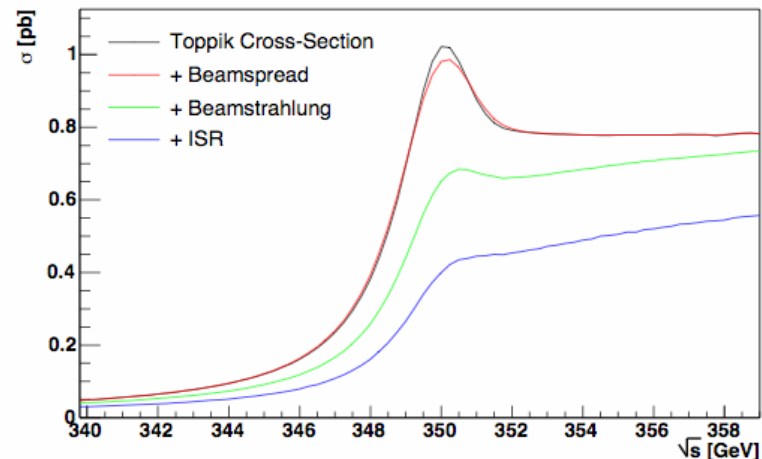
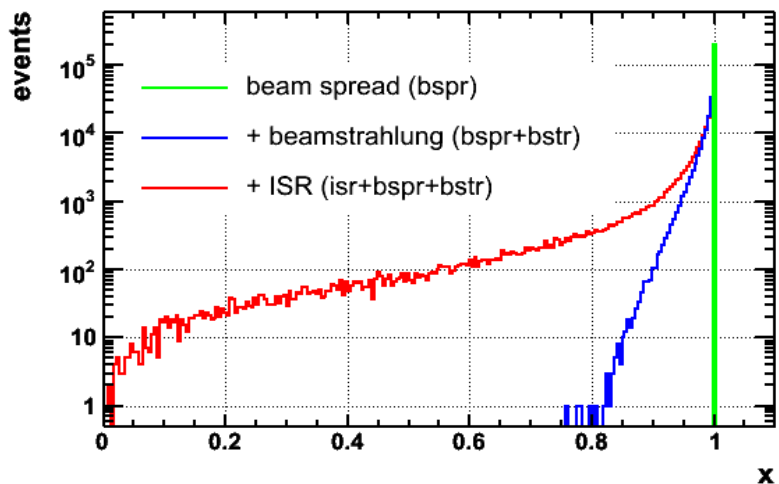
- E_{cm} adjustable from 200 – 500 GeV
- Luminosity $\rightarrow \int L dt = 500 \text{ fb}^{-1}$ in 4 years
- Ability to scan between 200 and 500 GeV
- Energy stability and precision below 0.1%
- Electron polarization of at least 80%

- The machine must be upgradeable to 1 TeV

The RDR Design meets these “requirements,” including the recent update and clarifications of the reconvened ILCSC Parameters group!



ILC energy spectrum; impacts physics output:



Need for:

- Energy measurement accuracy 10^{-4}
- Stability and ease of operation
- Minimal impact on physics data taking

Luminosity

Efficiency

$$L = \frac{\eta P_{\text{electrical}}}{E_{\text{CM}}} \sqrt{\frac{\delta}{\epsilon_{n,y}}} H_D$$

Beamstrahlung parameter

$$\delta \propto E_{\text{CM}} \frac{N^2}{\sigma_z \sigma_x^2}$$

Disruption parameter

Trade-off between

- Luminosity
- beam energy precision (beamstrahlung δ)
- backgrounds (related to H_D)
- running cost



ILC Parameter Plane

500 GeV CMS

	Nominal	Low Q	Large Y	Low P	
Number of Particles	2	1	2	2	10^{10}
Number of bunches	2625	5120	2625	1320	
Bunch interval (buckets)	369(480)	189(246)	369(480)	480(624)	ns()
Average current	9.0	9.0	9.0	6.8	mA
Norm.emittance at IP x/y	10/0.04	10/0.03	10/0.08	10(0.036)	μm
Beta at IP x/y	20/0.4	11/0.2	11/0.6	11/0.2	mm
Rms beamsize at IP x/y	639/5.7	474/3.5	474/9.9	474/3.8	nm
Rms bunch length	300	200	500	200	μm
Disruption param x/7	0.174/19.4	0.108/14.6	0.520/24.9	0.211/26.1	
Beamstrahlung param Y	0.048	0.050	0.038	0.097	
Energy loss by beamstr.	2.4	1.7	2.7	5.5	%
# of photons of beamstr.	1.32	0.91	1.77	1.72	
Pinch enhancement	1.71	1.48	2.18	1.64	
Geometric luminosity	1.20	1.35	0.935	1.21	10^{34}
Luminosity	2.0	2.0	2.0	2.0	10^{34}



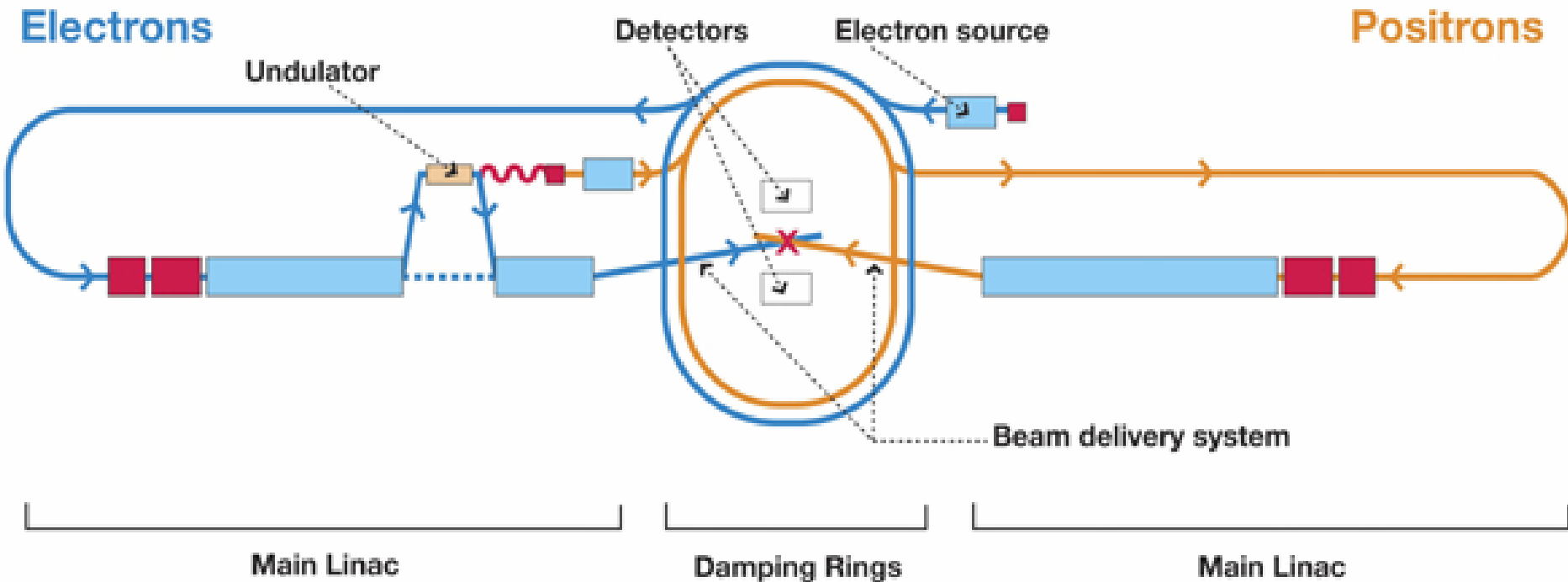
ILC Parameter Range

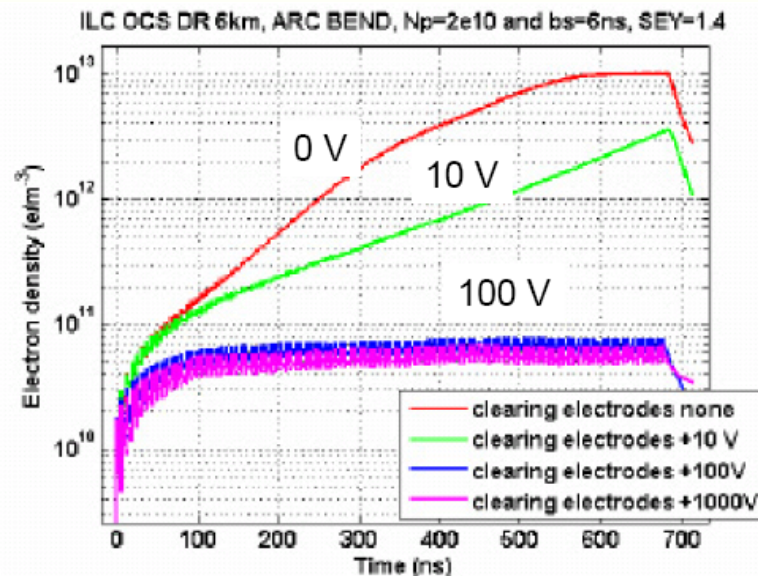
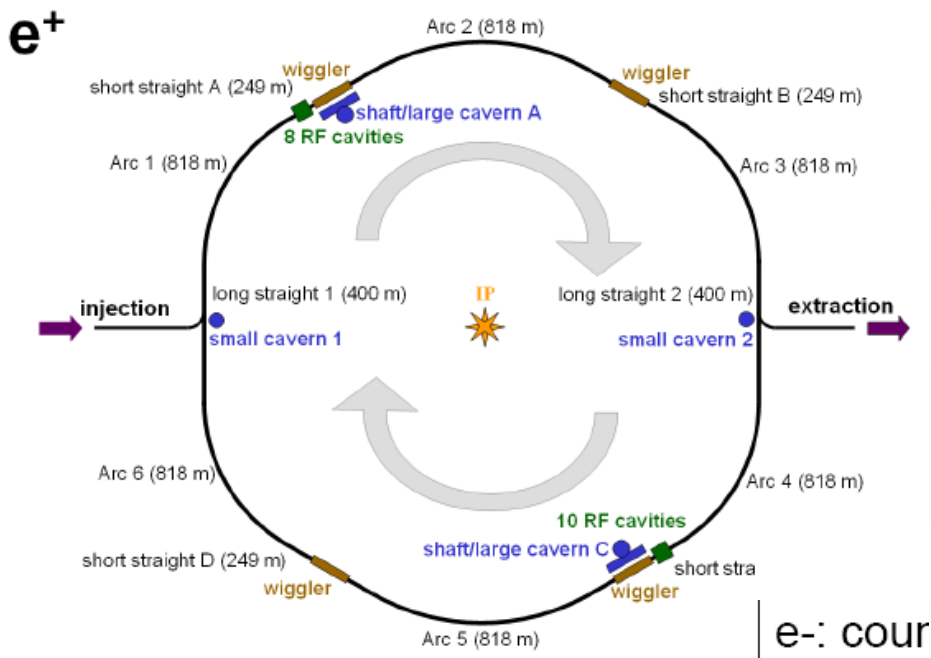
	min	-	nominal	-	max	
Number of particles	1	-	2	-	2	10^{10}
Number of bunches	1320	-	2625	-	5120	
Linac bunch interval	189	-	369	-	480	ns
DR bunch interval	3.08	-	6.15	-	12.3	ns
Bunch length	200	-	300	-	500	μm
Vertical emittance	0.03	-	0.04	-	0.08	μm
Beta at IP (x)	11	-	11	-	20	mm
Beta at IP (y)	0.2	-	0.4	-	0.6	mm

Special challenges in red

ILC Layout

500 GeV machine



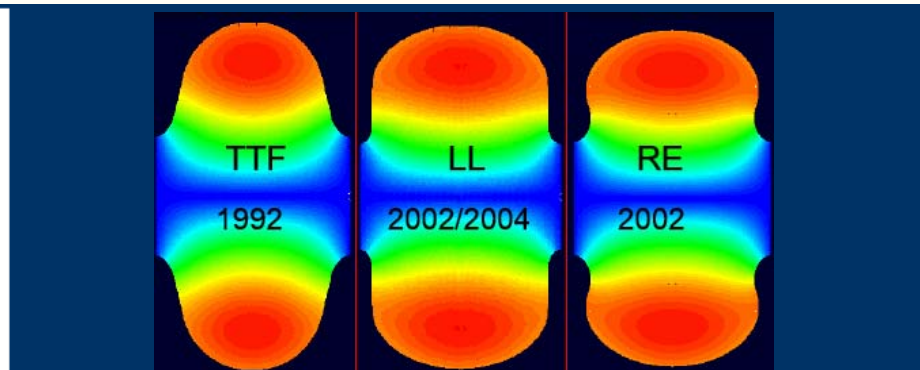
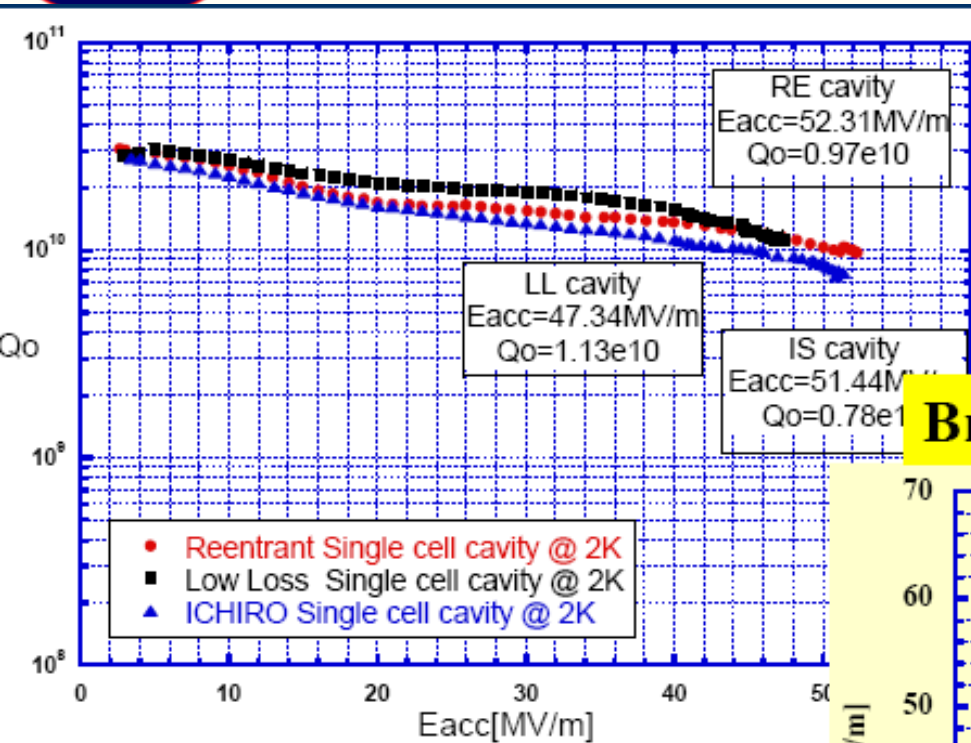


Electron cloud: suppression with clearing electrodes?

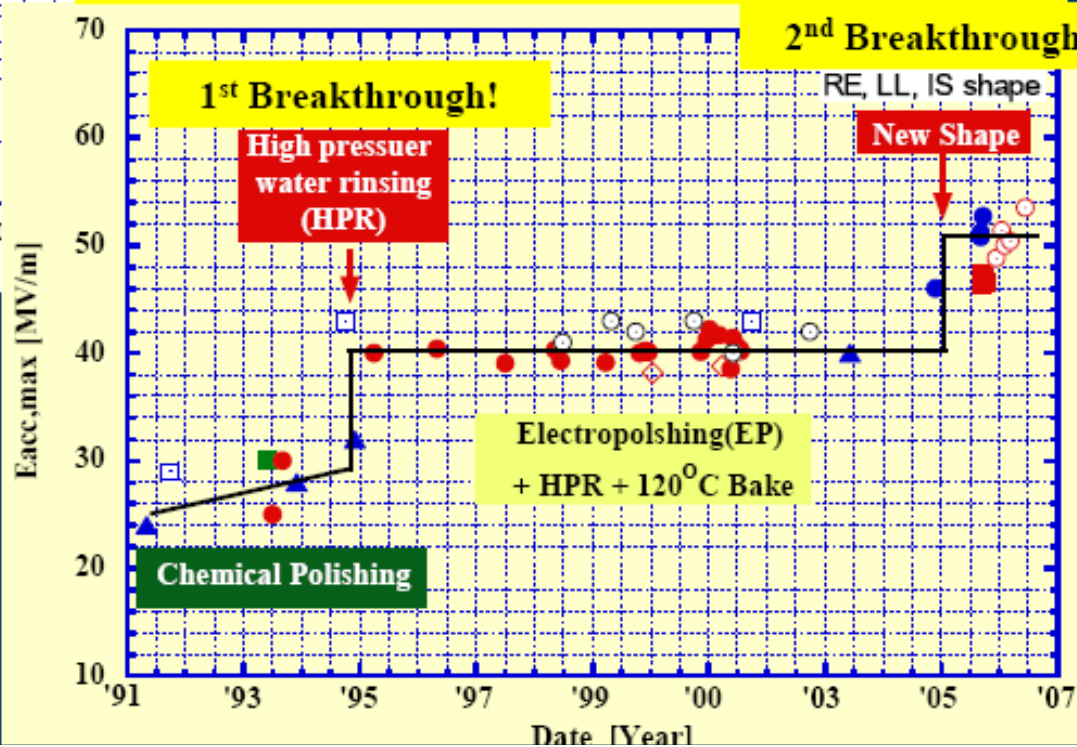
- Fast ion effects in electron DR:
feedback, vacuum design (1nTorr), train gaps?
- Long-range wake fields can drive multi-bunch instabilities,
- Short-range wake fields can drive single-bunch instabilities
- Requires: Fast kicker: 5ns rise time, 30 ns fall time...
(one train of $2625 \times 369ns \rightarrow 290$ km !)

“The DR have more accelerator physics than the rest of the accelerator...”

Higher Gradients?



Break-through by new-shapes



RE, LL and ICHIRO Single-cell (IS) achieved 45MV/m

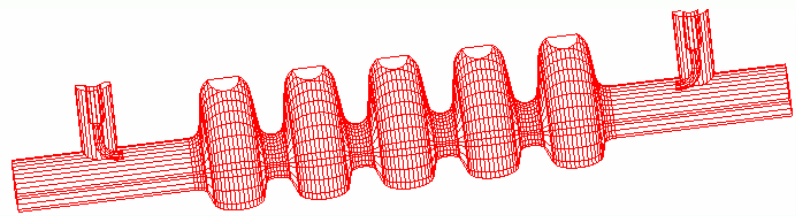
Crab crossing

$$\sigma_{x,projected} \approx \sqrt{\sigma_x^2 + \phi_c^2 \sigma_z^2}$$

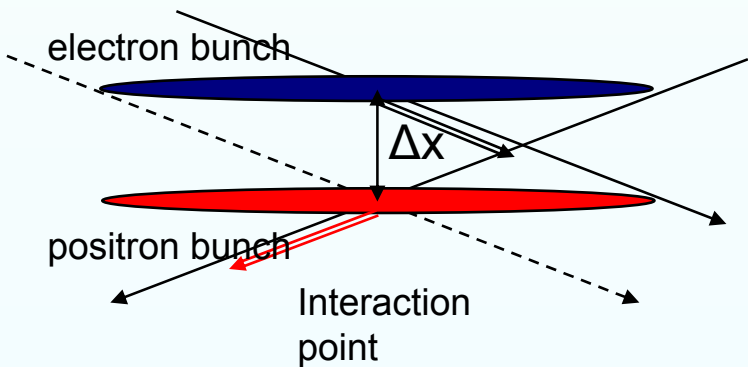
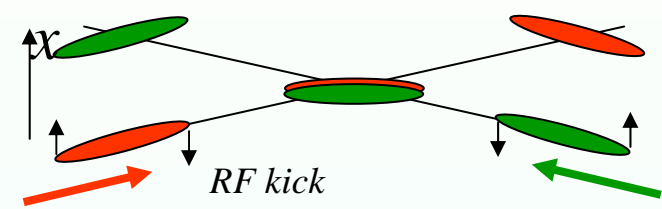
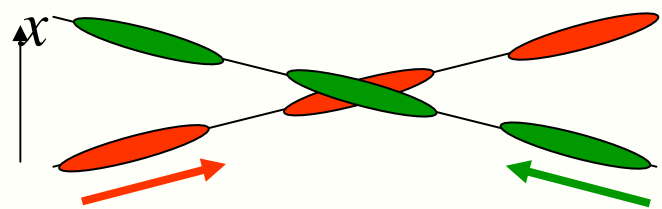
$$\approx \phi_c \sigma_z$$

$$= 20\text{mr} \times 100\mu\text{m} \approx 2\mu\text{m}$$

→ factor 10 reduction in Lumi



need one or two multi-cell cavities
~15m from IP

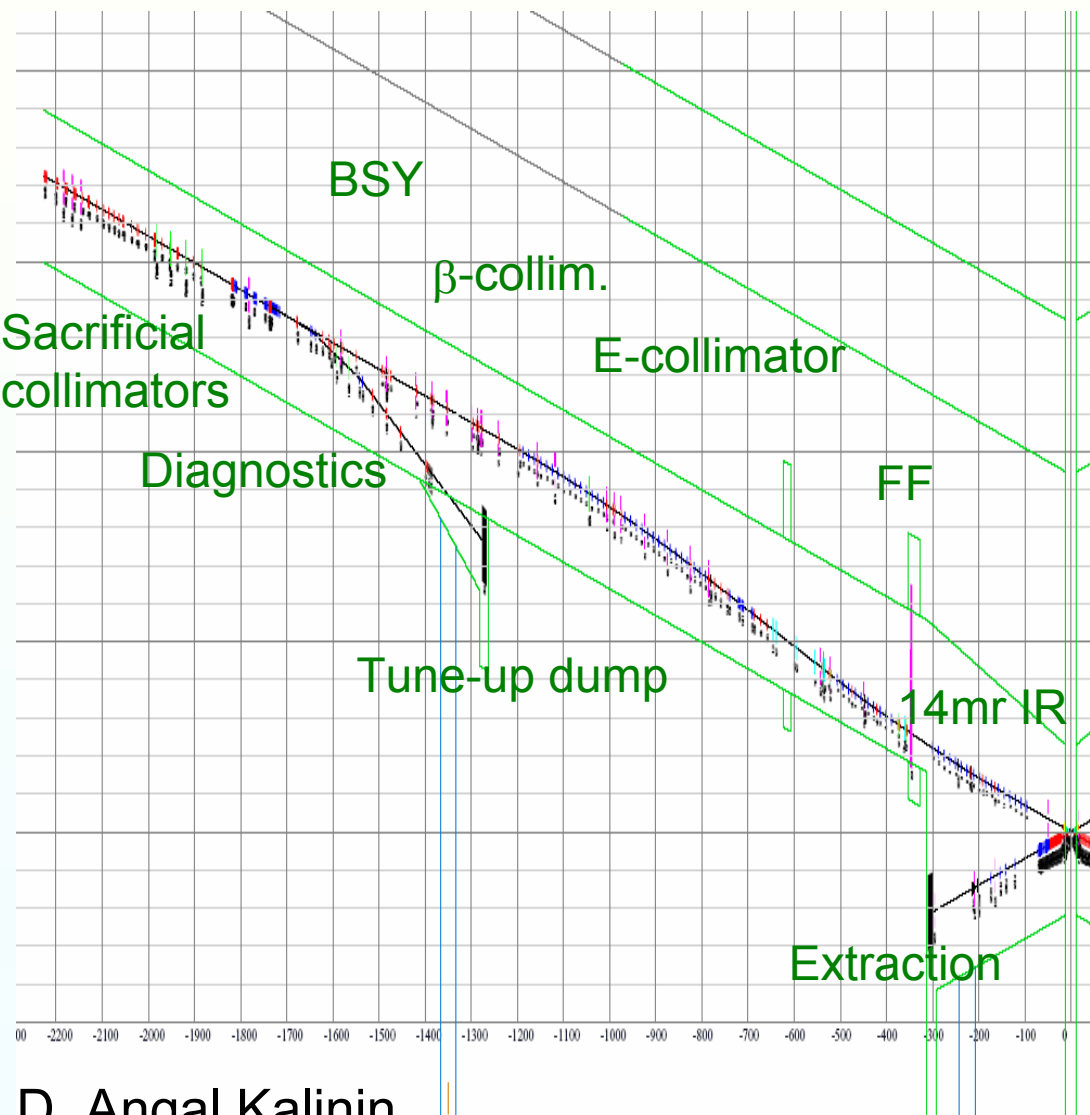


Burt et al.

Crossing angle	Phase error (degrees)	
	1.3GHz	3.9GHz
2mrad	0.222	0.665
10mrad	0.044	0.133
20mrad	0.022	0.066



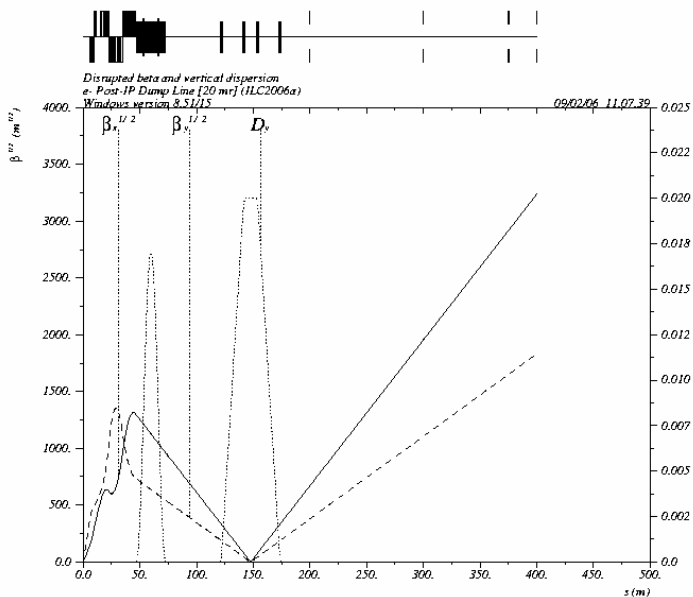
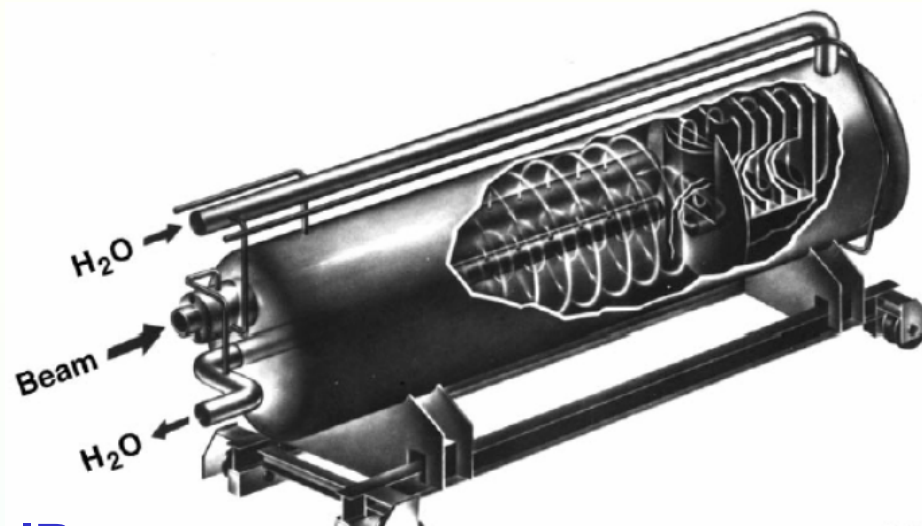
BDS : RDR configuration



Length (linac exit to IP distance)/side	m	2226
Length of main (tune-up) extraction line	m	300 (467)
Max Energy/beam (with more magnets)	GeV	250 (500)
Distance from IP to first quad, L^*	m	3.5-(4.5)
Crossing angle at the IP	mrad	14
Nominal beam size at IP, σ^* , x/y	nm	639/5.7
Nominal beam divergence at IP, θ^* , x/y	μ rad	32/14
Nominal beta-function at IP, β^* , x/y	mm	20/0.4
Nominal bunch length, σ_z	μ m	300
Nominal disruption parameters, x/y		0.17/19.4
Nominal bunch population, N		2.05×10^{10}
Beam power in each beam	MW	11.3
Preferred entrance train to train jitter	σ	< 0.5
Preferred entrance bunch to bunch jitter	σ	< 0.1
Typical nominal collimation depth, x/y		8-10/60-80
Vacuum pressure level, near/far from IP	nTorr	1/50

D. Angal Kalinin

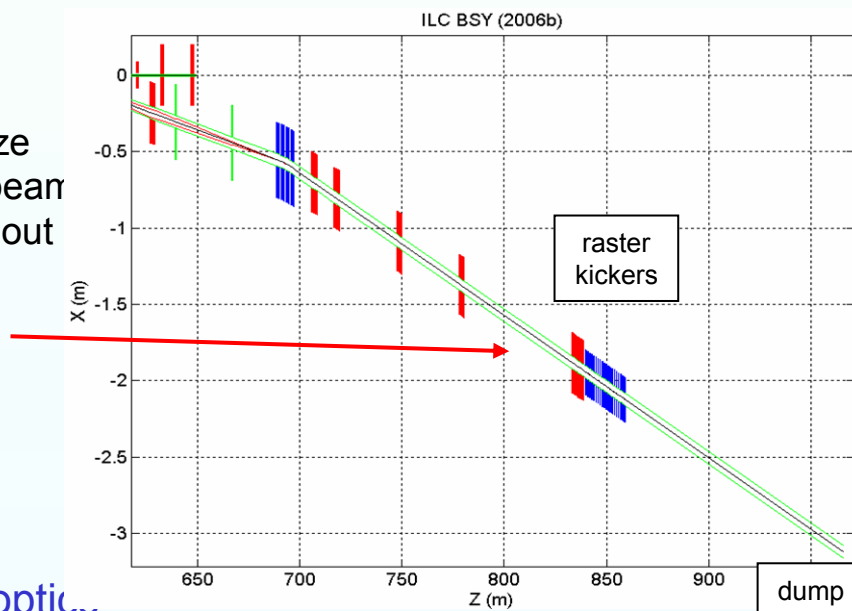
- Water vortex
- Window, 1mm thin, ~30cm diameter hemisphere
- Raster beam with dipole coils to avoid water boiling
- Deal with H, O, catalytic recombination
- Gas dump also being studied
- 3MW beamstrahlung dumps near IR

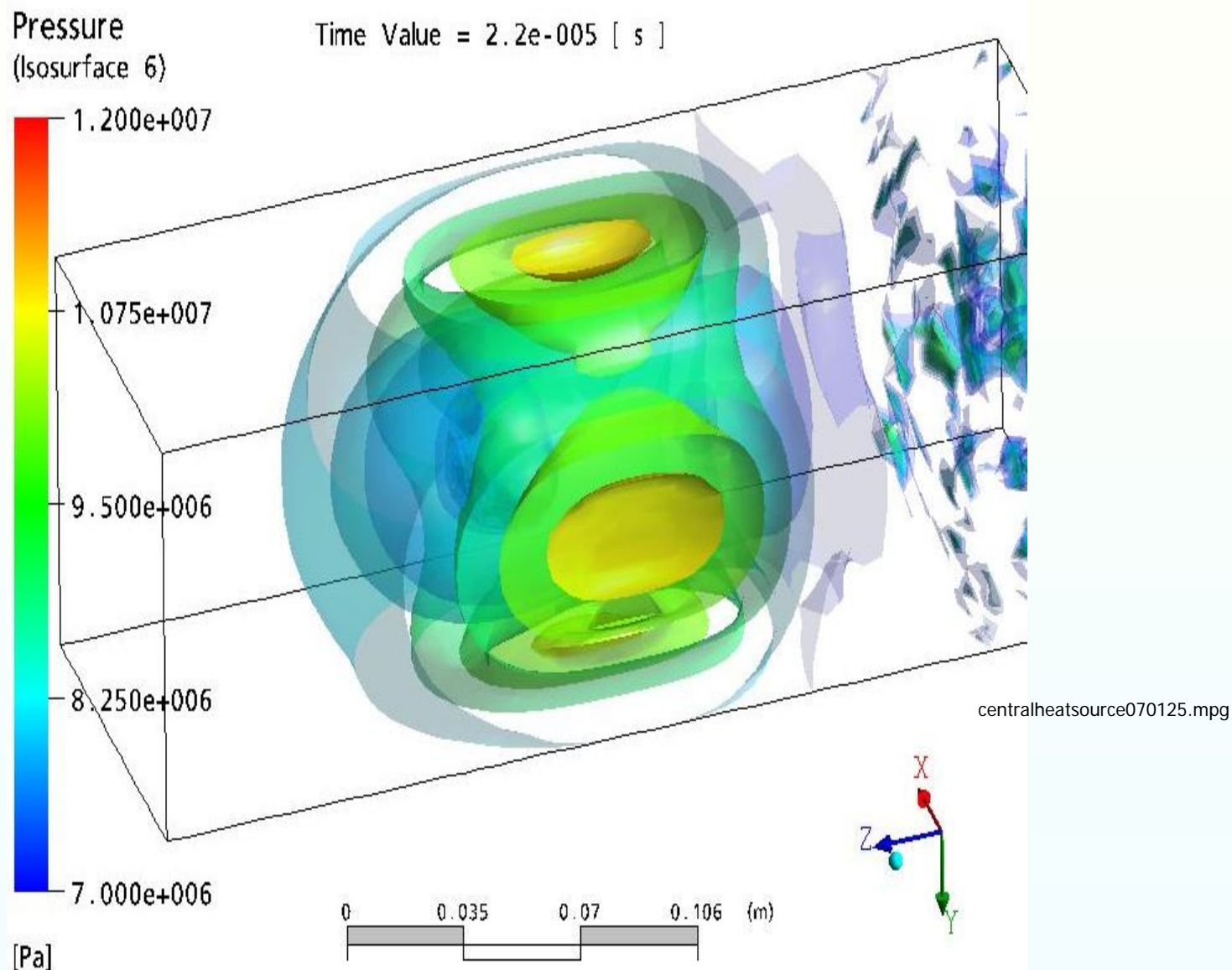


undisrupted or disrupted beam size does not destroy beam dump window without rastering.

Rastering to avoid boiling of water

20mr extraction optics





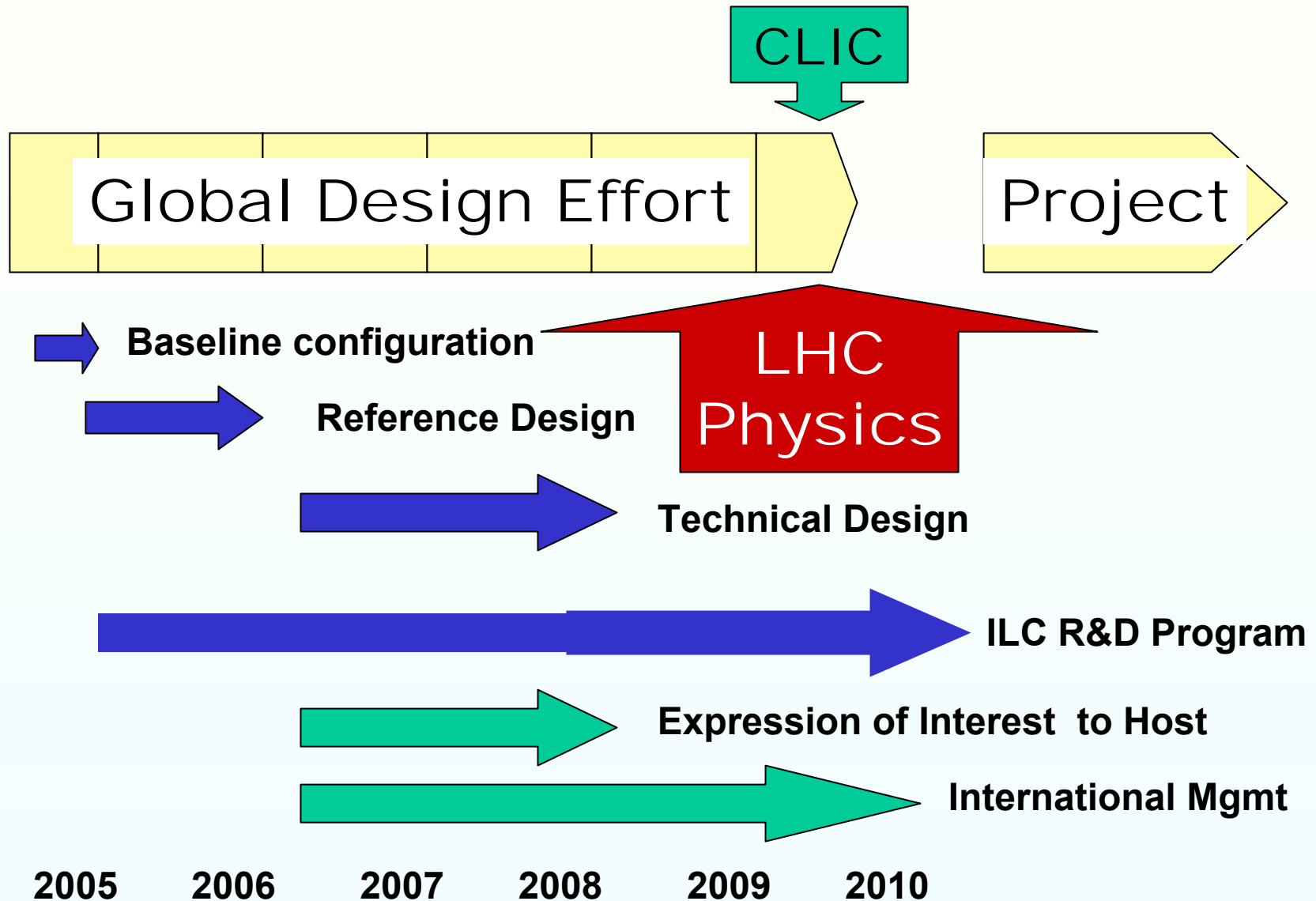


Final Focus

	SLC	FFTB	ATF2	ILC
E_{beam} (GeV)	45.6	46.6	1.3	250
σ_E/E (%)	0.25	0.25	0.1	0.1
N_{e^-} ($\times 10^{10}$)	4.2	1	1-2	2
σ_y (nm)	800	60	37	5.7
$\gamma\epsilon_y$ (m-rad)	1×10^{-5}	3×10^{-6}	3×10^{-8}	4×10^{-8}
Asp. ratio x/y	2.5	16	13	115
σ_z (mm)	~1	~1	~5	0.3

The ILC Plan and Schedule

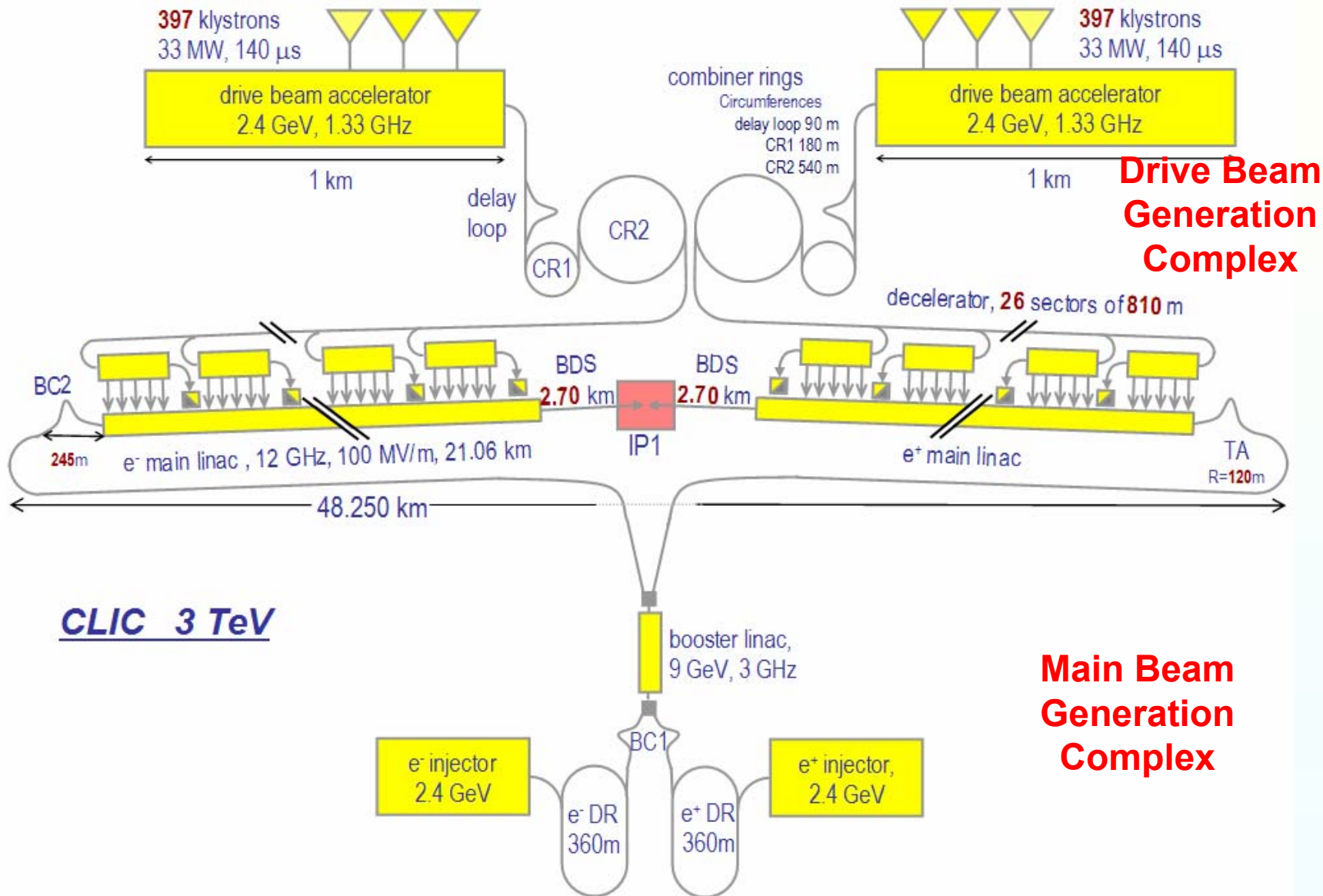
(B.Barish/CERN/SPC 050913)



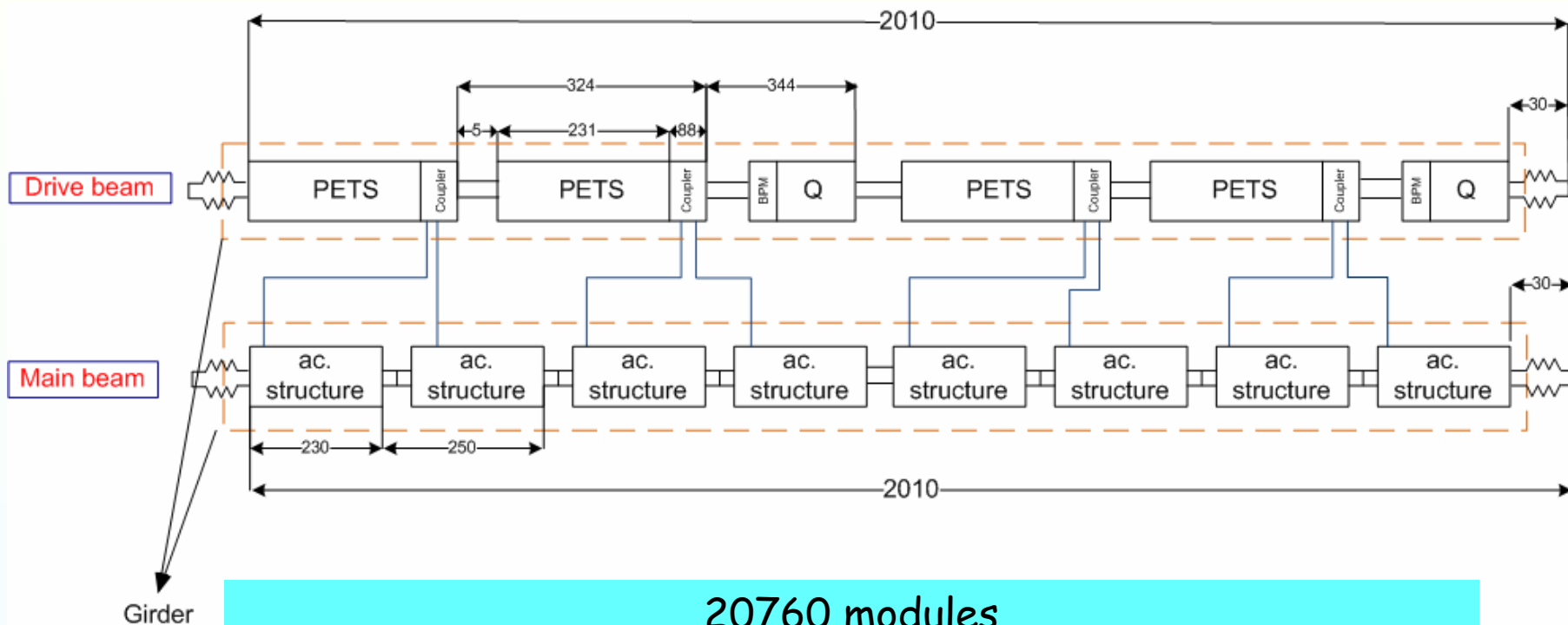
Compact Linea Collider



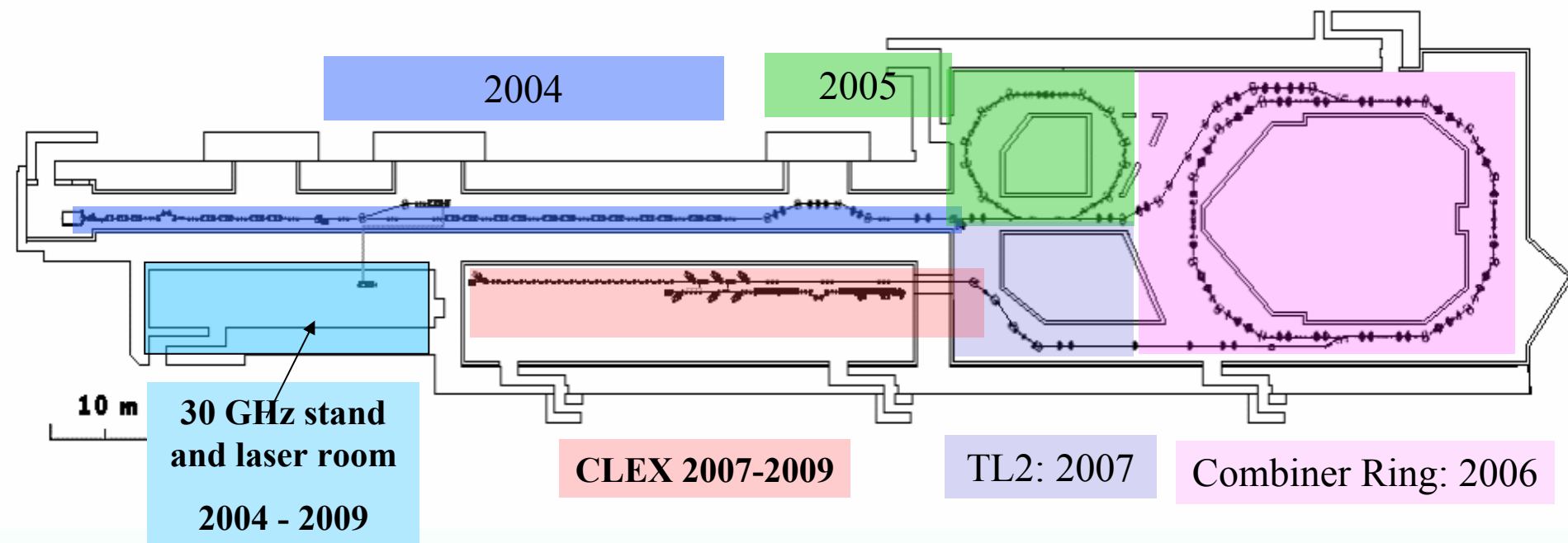
CLIC - overall layout



Two Beam Module



20760 modules
 71460 power production structures PETS (drive beam)
 143010 accelerating structures (main beam)



Key issues

From 2005: Accelerating structures Development & Tests (R2.1)

2007- 2008: Drive beam generation scheme (R1.2)

2008- 2009: Damped accelerating structure with nominal parameters (R1.1)
ON/OFF Power Extraction Structure (R1.3)
Drive beam stability bench marking (R2.2)
CLIC sub-unit (R2.3)

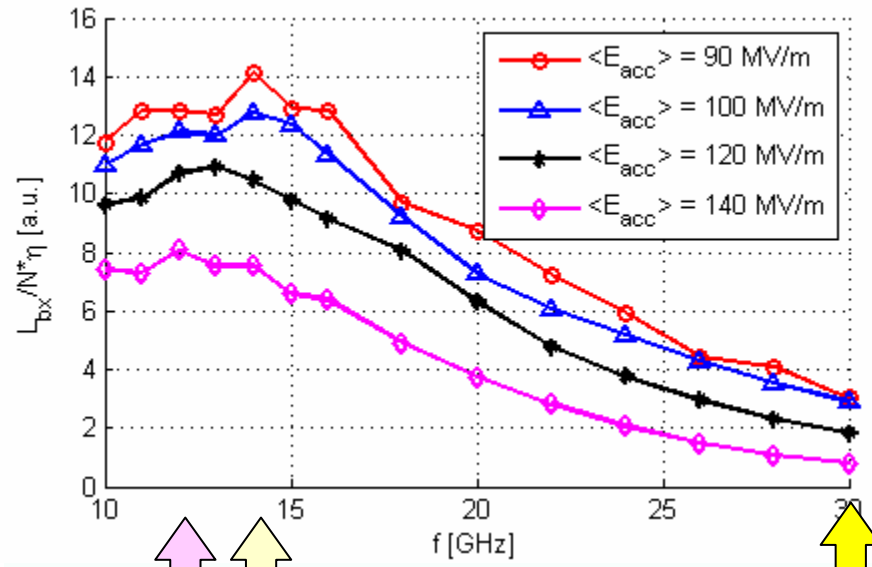
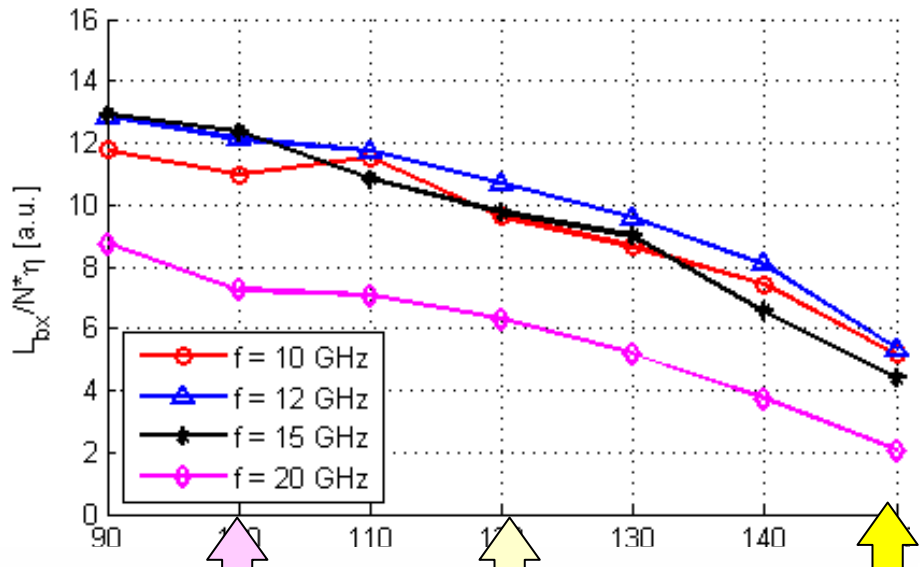


CLIC performance & cost optimisation

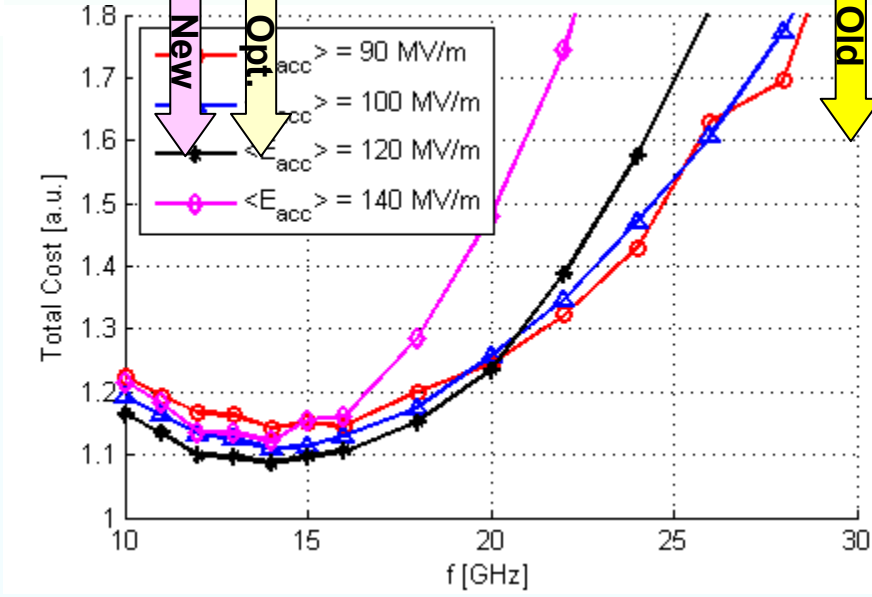
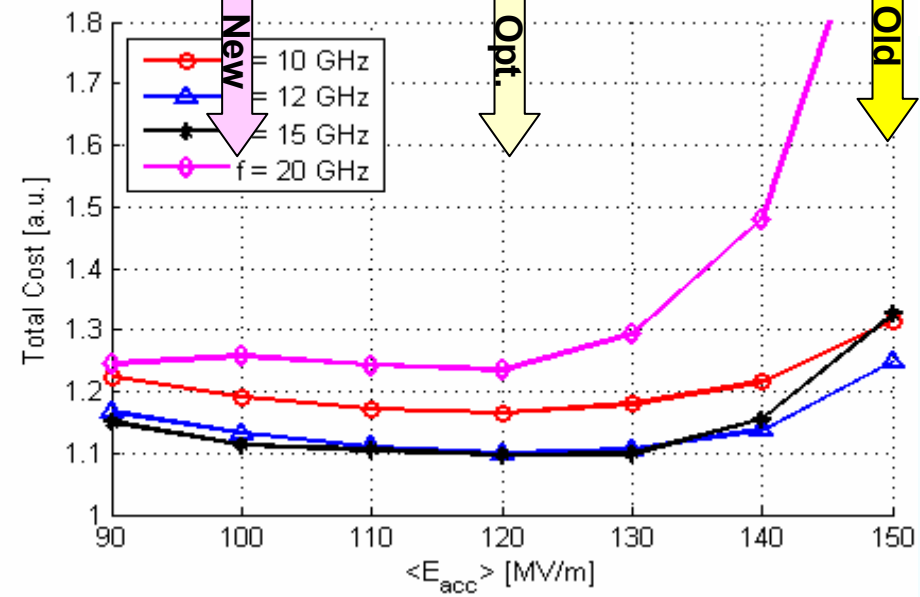
$E_{\text{cms}} = 3 \text{ TeV}$

$L_{(1\%)} = 2.0 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$

Performance



Cost



<i>Main Linac RF frequency</i>	30 GHz ⇒ 12 GHz
<i>Accelerating field</i>	150 MV/m ⇒ 100 MV/m
<i>Overall length @ $E_{CMS} = 3 \text{ TeV}$</i>	33.6 km ⇒ 48.2 km

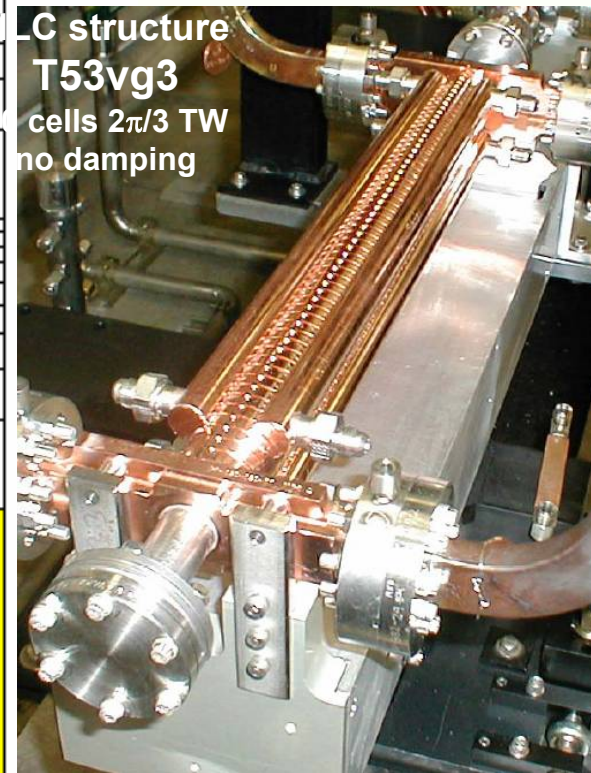
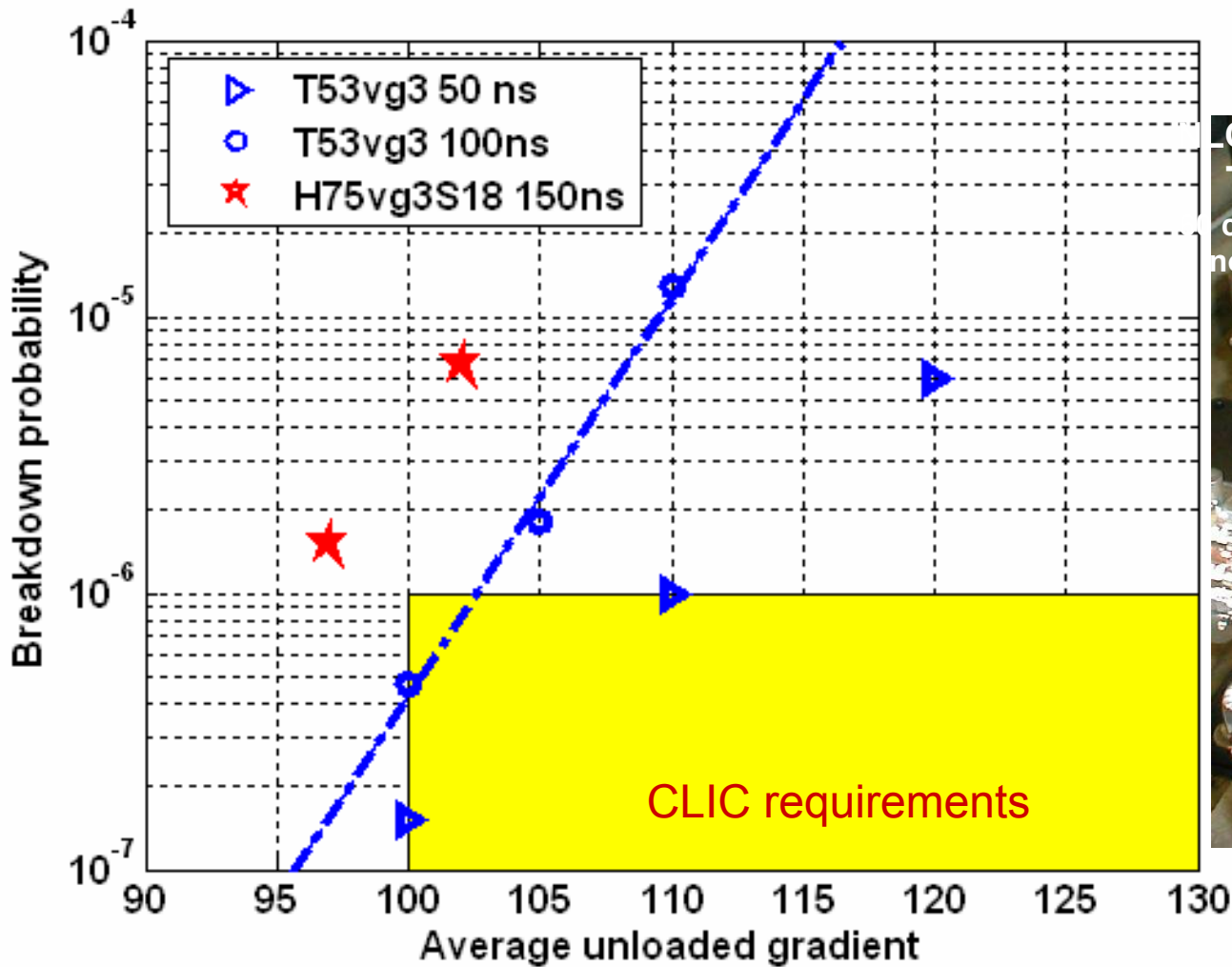
- **Substantial cost savings and performance improvements for 12 GHz / 100 MV/m indicated by parametric model (flat optimum in parameter range)**
- **Promising results already achieved with structures in test conditions close to LC requirements (low breakdown rate) but still to be demonstrated with long RF pulses and fully equipped structures with HOM damping.**
- **Realistic feasibility demonstration by 2010**



New CLIC main parameters

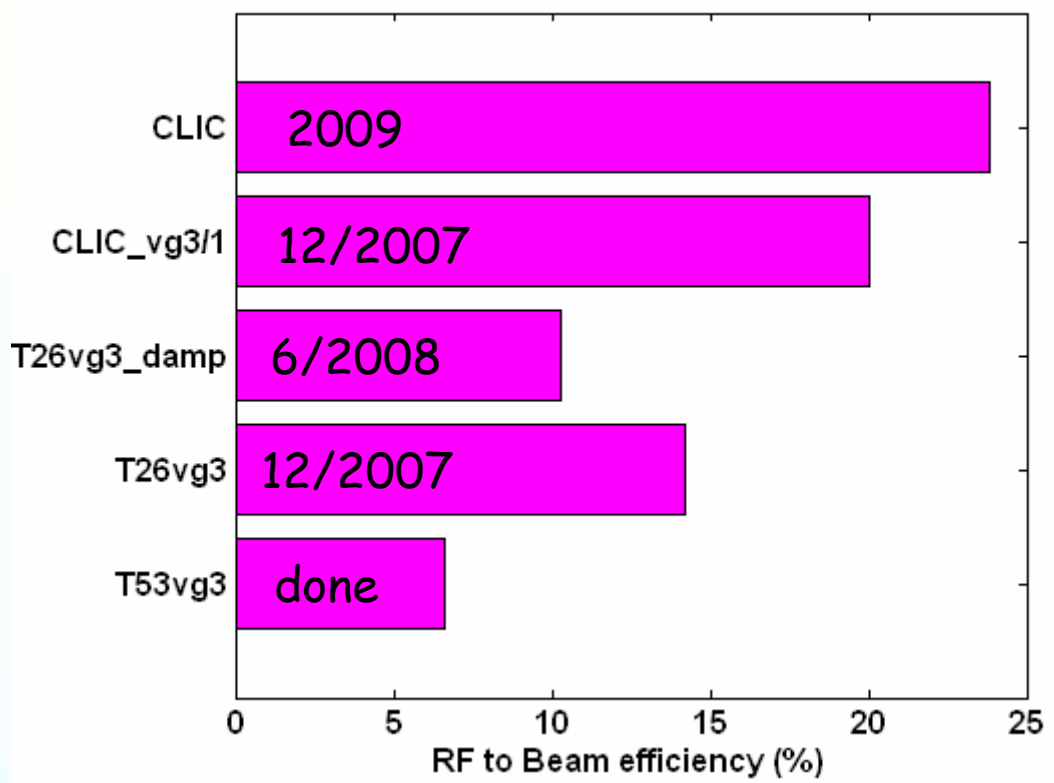
Center-of-mass energy	3 TeV
Peak Luminosity	$7 \cdot 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
Peak luminosity (in 1% of energy)	$2 \cdot 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
Repetition rate	50 Hz
Loaded accelerating gradient	100 MV/m
Main linac RF frequency	12 GHz
Overall two-linac length	41.7 km
Bunch charge	$4 \cdot 10^9$
Beam pulse length	200 ns
Average current in pulse	1 A
Hor./vert. normalized emittance	660 / 20 nm rad
Hor./vert. IP beam size bef. pinch	53 / ~1 nm
Total site length	48.25 km
Total power consumption	390 MW

Provisional values





RF to Beam Efficiency milestones



P = 65 MW; 297 ns \Leftrightarrow nb = 311

P = 70 MW; 295 ns \Leftrightarrow nb = 359

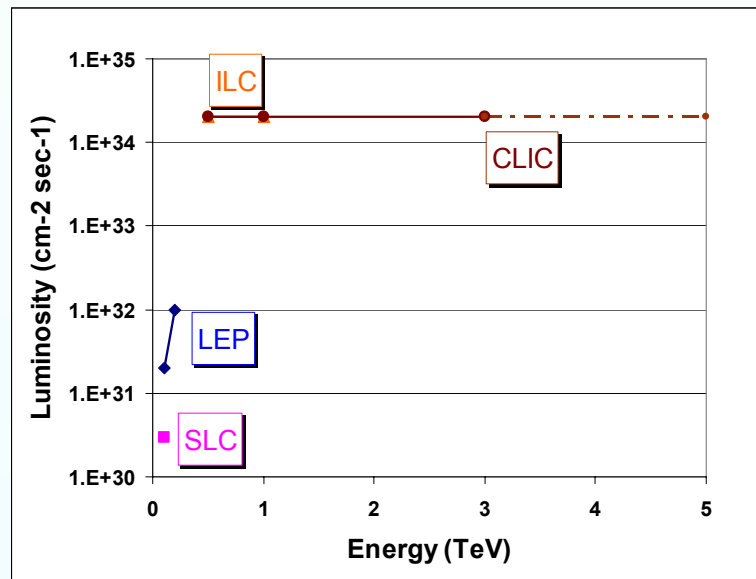
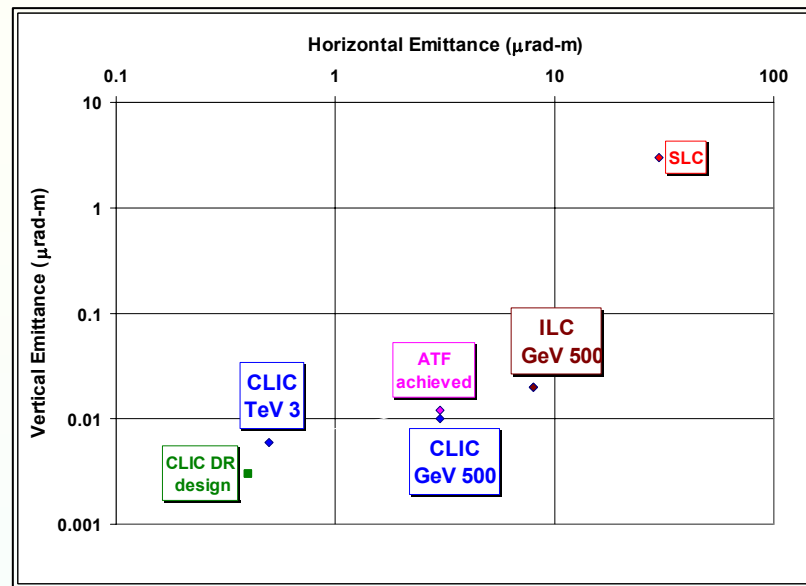
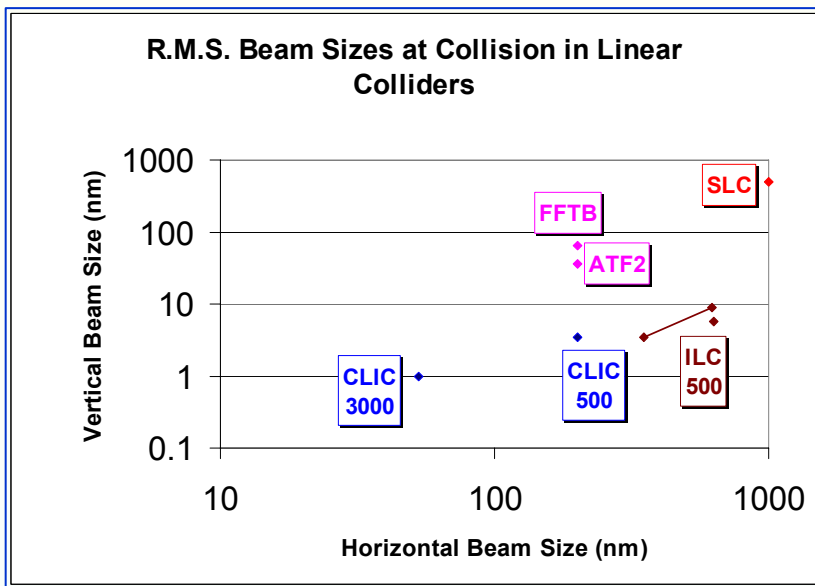
P = 111 MW; 102 ns \Leftrightarrow nb = 66

P = 102 MW; 113 ns \Leftrightarrow nb = 93

P = 134 MW; 104 ns \Leftrightarrow nb = 27

100 MV/m loaded, 10^{-6} break down rate, $q_b=4 \cdot 10^9$,
8 rf period bunch spacing, $P \cdot pl/C = 18 \text{ Wue}$

Performances of Lepton Colliders





The Muon Collider

- **Why?**

1. For some processes, muons are better than electrons

- $\sigma(\mu^+\mu^-\rightarrow H)$ is $(m_\mu/m_e)^2 \times \sigma(e^+e^-\rightarrow H)$
 - (40000 times larger)

2. If CLIC does *not* work, this may be the only route to multi-TeV collisions under clean conditions

- **Why not?**

- Muon lifetime is only 2 μ sec!
- Need to produce, collect, cool and accelerate large numbers ($\gg 10^{13}$) muons per second



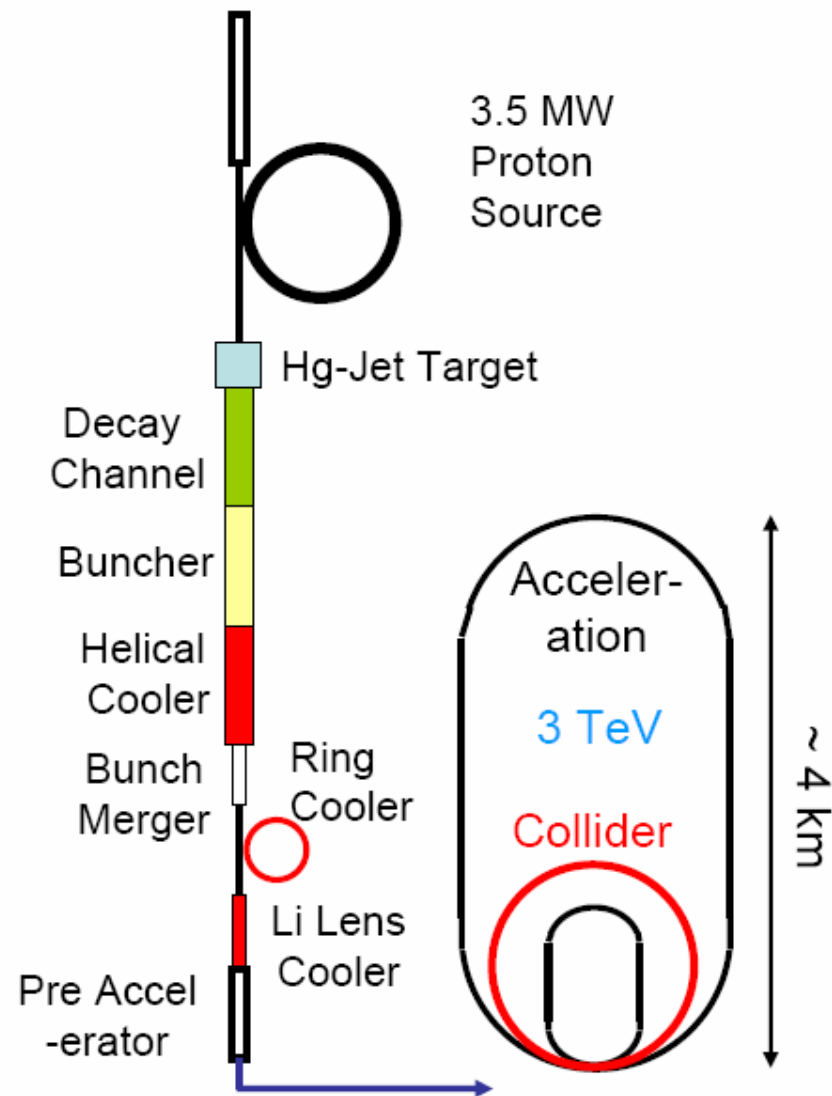
Introduction



- Muon Collider concept is attractive because muons are point-like particles that do not radiate as readily as electrons ($m_\mu / m_e \sim 207$):
 - Circular (compact) multi-TeV Lepton Collider that would fit on an existing laboratory site.
 - Very small beam energy spread enabling precise scans and width measurements.
- To produce sufficient luminosity for an interesting physics program will require very bright muon beams. This is challenging because:
 - Muons produced as a tertiary beam that occupies a large longitudinal & transverse phase space. The beam must be cooled by a large factor.
 - Muons decay ($t_0 = 2\mu\text{s}$). Beam manipulation & acceleration must be rapid.
- Much of the R&D required for Muon Colliders is in common with the R&D for Neutrino Factories, & is being pursued in the US by the Neutrino Factory & Muon Collider Collaboration (NFMCC: 135 Scientists & Engineers, 4 US Labs, 17 US Universities, 14 non-US Institutions).

Main Components of a Muon Collider

- Proton Driver
 - primary beam on production target
- Target, Capture, and Decay
 - create π ; decay into μ
- Bunching & Phase Rotation
 - reduce ΔE of bunch
- Cooling
 - reduce 6D emittance
- Acceleration
 - 130 MeV \rightarrow up to 1.5 TeV
- Storage Ring
 - store for ~ 1000 turns
 - One IP



Steve Geer



Parameters

Energy	0.1 TeV	3 TeV
Proton Srce Power	4 MW	3.5 MW
Rate	15 Hz	30 Hz
Muons / bunch	4×10^{12}	2×10^{12}
Bunches	1×1	2×2
Circumference	0.35 km	6 km
Effective turns	450	900
β_{\perp}	9.4 mm	3 mm
ϵ_{\perp} (mm radians)	0.195	0.05
$\epsilon_{//}$	5 mm	72 mm
Δv	0.022	0.044
$\Delta p/p$	0.01 %	0.16 %
Luminosity ($\text{cm}^{-2} \text{s}^{-1}$)	2.2×10^{34}	7×10^{34}

Requires $1 \times 10^{21} \mu^+ / \text{year}$

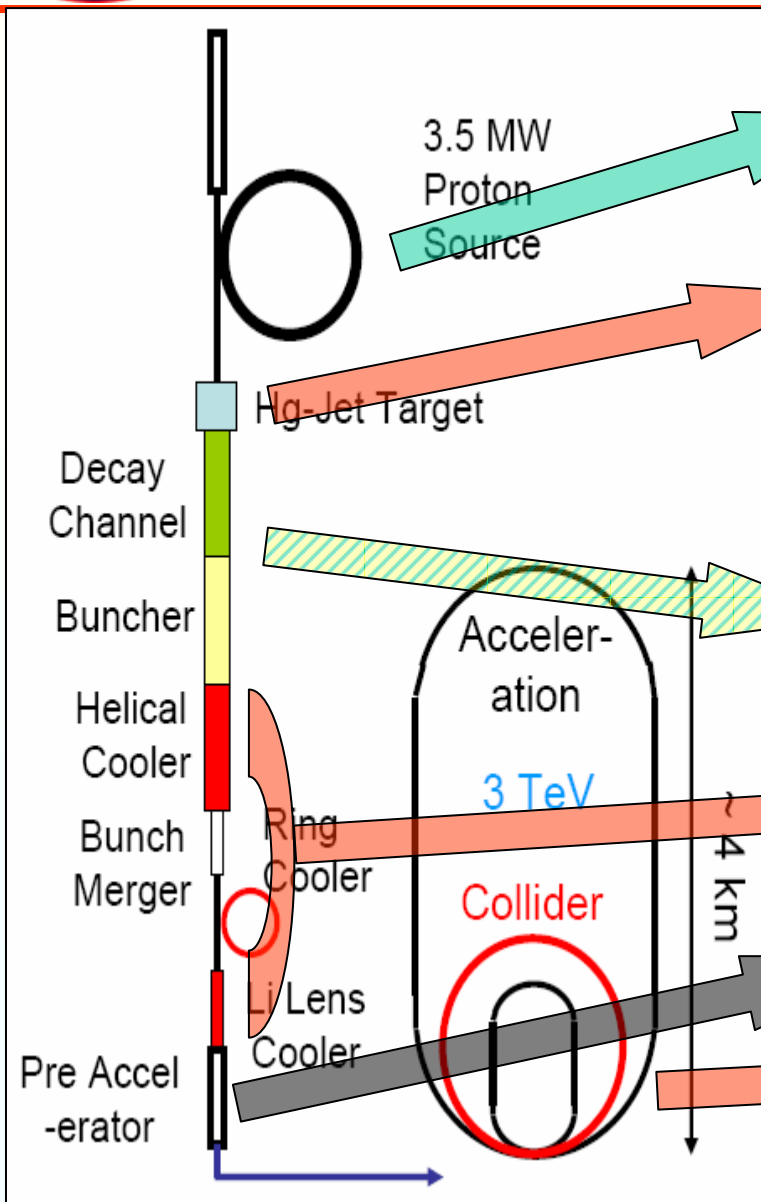
Similar to (a little more than) number of muons expected from Neutrino Factory front-end

Collider energy limited by “neutrino radiation” – high energy muon decays in straight sections produce a collimated beam of neutrinos that interact in the Earth to produce a radiation field at their exit point.

Lower emittance beams (more cooling) would enable use of fewer muons → higher energy colliders become plausible.

Steve Geer

Challenges for the Muon Collider



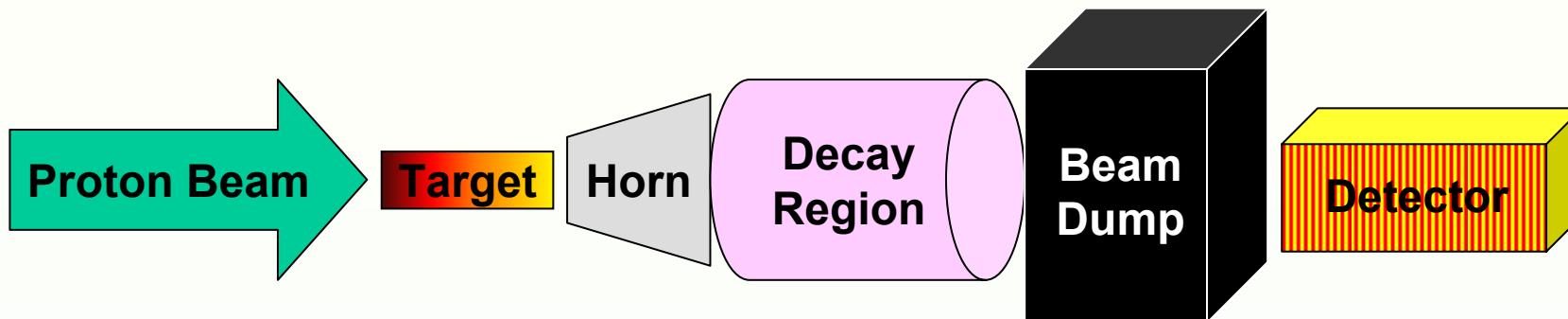
- **Proton Driver**
 - ~4MW, 1 ns bunches
- **Target**
 - ~4MW, 1 ns bunches
 - “open” geometry
- **Muon Collection and bunching**
- **Bunch cooling & compression (Pre-acceleration)**
- **Collider**

All are needed for a Neutrino Factory



Neutrino Factory

Conventional Neutrino Beams



• Main components

- Proton Beam
 - Energy, Intensity, frequency
- Target
- Horn (focussing)
- Decay Region
- Beam Dump
- Detector

Note

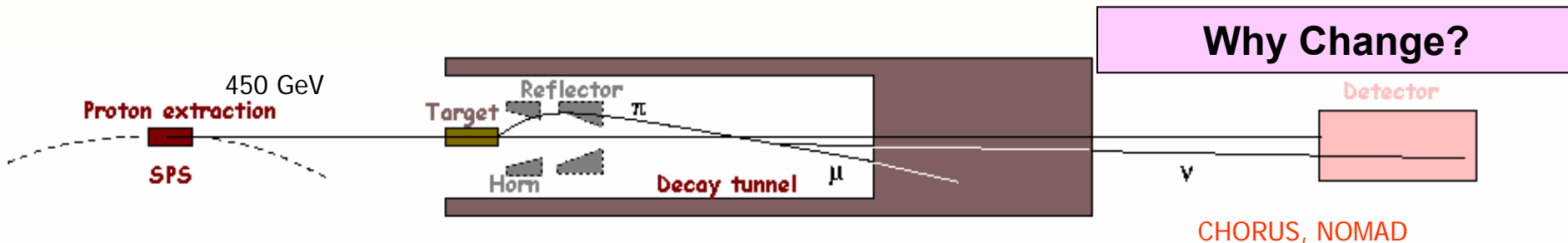
For any (class of) experiment

$$N_{\text{ev}} \propto \mathbf{P} \times \mathbf{M} \quad (\times E_{\nu})$$

Beam Power	Target Mass	Neutrino Energy
-----------------------	------------------------	----------------------------

Example of a Neutrino Beam

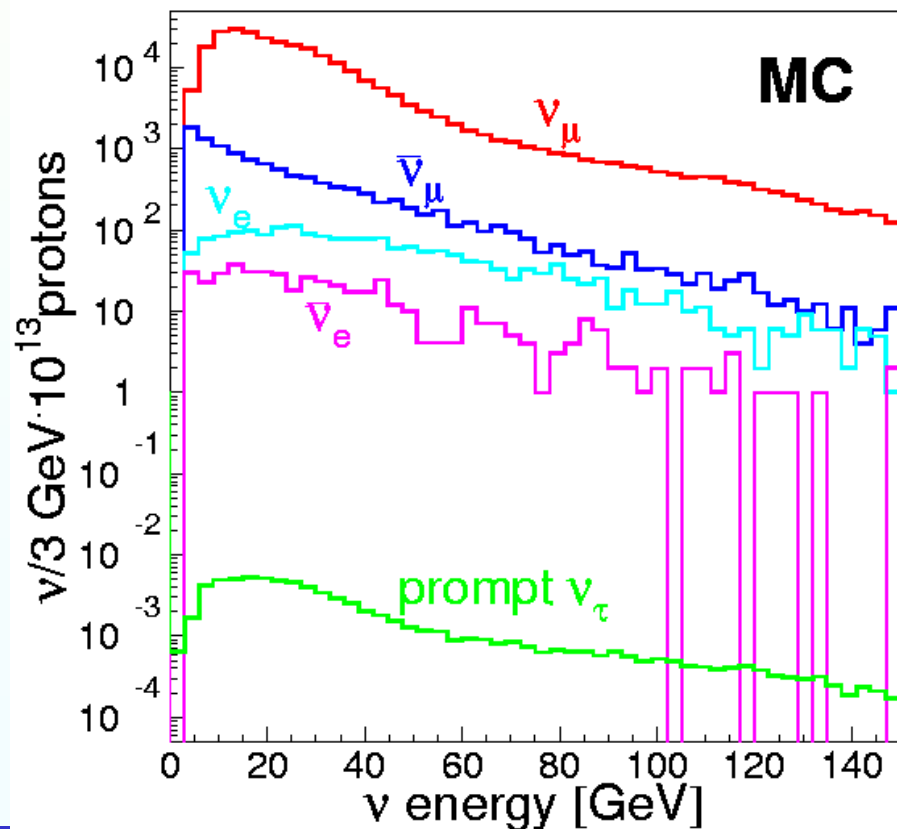
West Area Neutrino Facility at CERN SPS



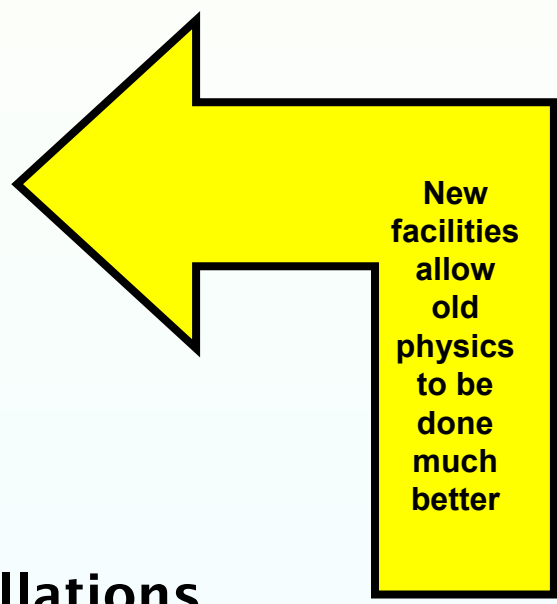
Wide Band Beam

- 5.06×10^{19} POTs (1994-1997)
- $\langle E_{\nu_{\mu}} \rangle \sim 27 \text{ GeV}$
- $\langle L \rangle \sim 0.6 \text{ km}$
 $\langle L \rangle / \langle E \rangle \sim 2 \times 10^{-2} \text{ km/GeV}$
 $\rightarrow \Delta m^2 > 1 \text{ eV}^2$
- Prompt ν_{τ} : negligible

$\sim 10^{12}$ neutrinos



- 1950's and early 60's
 - Nature (and existence) of the neutrino
 - (Reines & Cowan, Lederman, Schwartz and Steinberger)
- Late 1960s, 1970s, 1980s
 - Structure of the nucleon
 - F_2 , xF_3 etc
 - Structure of the weak current
 - Neutral currents, $\sin_2\theta_w$ etc
- Now, and future
 - Nature of the neutrino
 - Neutrino Mass and Neutrino Oscillations
 - Standard Model assumption of massless neutrinos is *wrong!*
 - Note: difficult to add neutrino mass to SM *a la Higgs*
 - Lack of Charge \rightarrow additional mass-like (Majorana) terms



Atmospheric

3G

solar

Majorana

$$U_{MNS} = \begin{bmatrix} 1 & & & & & \\ & c_{23} & s_{23} & & & \\ & -s_{23} & c_{23} & & & \\ & & & 1 & & \\ & & & & c_{13} & s_{13}e^{-i\delta} \\ & & & & -s_{13}e^{i\delta} & c_{13} \end{bmatrix} \otimes \begin{bmatrix} c_{12} & s_{12} & \\ -s_{12} & c_{12} & \\ & & 1 \end{bmatrix} \otimes \begin{bmatrix} 1 & & \\ & e^{i\alpha} & \\ & & e^{i\beta} \end{bmatrix} \\
 = \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{i\delta} \\ -c_{23}s_{12} - s_{13}s_{23}c_{12}e^{i\delta} & c_{23}c_{12} - s_{13}s_{23}s_{12}e^{i\delta} & c_{13}s_{23} \\ s_{23}s_{12} - s_{13}c_{23}c_{12}e^{i\delta} & -s_{23}c_{12} - s_{13}c_{23}s_{12}e^{i\delta} & c_{13}c_{23} \end{bmatrix} \otimes \begin{bmatrix} 1 & & \\ & e^{i\alpha} & \\ & & e^{i\beta} \end{bmatrix}$$

$$|\Delta m_{32}^2| = 2.4(1_{-0.26}^{+0.21}) \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{23} = 0.44(1_{-0.51}^{+0.41})$$

$$\Delta m_{21}^2 = 7.92(1 \pm 0.09) \times 10^{-5} \text{ eV}^2$$

$$\sin^2 \theta_{12} = 0.314(1_{-0.15}^{+0.18})$$

$$\sin^2 \theta_{13} < 3.2 \times 10^{-2}$$

Sign(Δm_{32}^2) unknown

δ, α, β unknown

masses $< O(1\text{eV})$

2 σ

Lisi, NuFACT05

Parameters of neutrino oscillation

1 absolute mass scale

2 squared mass diffs

3 mixing angles

1 phase

2 Majorana phases

m_{ν_e}

$$\Delta m_{12}^2, \Delta m_{23}^2 \begin{cases} \Delta m_{ji}^2 = m_j^2 - m_i^2 \\ \Delta m_{31}^2 = \Delta m_{32}^2 + \Delta m_{21}^2 \end{cases}$$

$\theta_{12}, \theta_{23}, \theta_{13}$

δ (always $\sin\theta_{13}e^{i\delta}$)

α, β

$$c_{ij} = \cos\theta_{ij}$$

$$s_{ij} = \sin\theta_{ij}$$

Measuring the Parameters

$$\begin{aligned}
 P(\nu_\mu \Rightarrow \nu_e) = & 4c_{13}^2 s_{12}^2 (c_{12}^2 c_{23}^2) \sin^2 \left(\frac{\Delta m_{21}^2 L}{4E} \right) \\
 & 4c_{13}^2 s_{12}^2 \left(c_{12}^2 c_{23}^2 - s_{12}^2 s_{13}^2 s_{23}^2 - 2c_{12} c_{23} s_{12} s_{23} s_{13} \boxed{\cos \delta} \right) \sin^2 \left(\frac{\Delta m_{21}^2 L}{4E} \right) \\
 & + 8c_{13}^2 s_{12} s_{13} s_{23} \left(c_{12} c_{23} \boxed{\cos \delta} - s_{12} s_{13} s_{23} \right) \cos \left(\frac{\Delta m_{32}^2 L}{4E} \right) \sin \left(\frac{\Delta m_{31}^2 L}{4E} \right) \sin \left(\frac{\Delta m_{21}^2 L}{4E} \right) \\
 & + 4c_{13}^2 s_{13}^2 s_{23}^2 \sin^2 \left(\frac{\Delta m_{13}^2 L}{4E} \right) \left(1 + \left(1 - 2s_{13}^2 \right) \frac{2a}{\Delta m_{31}^2} \right) \nu_\mu \Rightarrow \bar{\nu}_\mu \Leftrightarrow a \Rightarrow -a \\
 & - 8c_{13}^2 c_{12} c_{23} s_{12} s_{13} s_{23} \sin \delta \sin \left(\frac{\Delta m_{32}^2 L}{4E} \right) \sin \left(\frac{\Delta m_{31}^2 L}{4E} \right) \sin \left(\frac{\Delta m_{21}^2 L}{4E} \right) \\
 & - 8c_{13}^2 s_{13}^2 s_{23}^2 \cos \left(\frac{\Delta m_{32}^2 L}{4E} \right) \sin \left(\frac{\Delta m_{31}^2 L}{4E} \right) \sin \left(\frac{\Delta m_{21}^2 L}{4E} \right) \left(1 - 2s_{13}^2 \right) \frac{aL}{4E}
 \end{aligned}$$

$$a = 2\sqrt{2} G_F n_e E_\nu = 7.6 \cdot 10^{-5} \rho E$$

Where n_e is the electron density ; ρ is the density (g/cm³) ; E is the neutrino energy (GeV)

$$c_{ij} = \cos \theta_{ij}, \quad s_{ij} = \sin \theta_{ij}$$

(Richter: hep-ph/0008222)

Neutrinos

- ν_e disappearance
- $\nu_e \rightarrow \nu_\mu$ appearance
- $\nu_e \rightarrow \nu_\tau$ appearance

- ν_μ disappearance
- $\nu_\mu \rightarrow \nu_e$ appearance
- $\nu_\mu \rightarrow \nu_\tau$ appearance

... and the corresponding antineutrino interactions

Note: the beam requirements for these experiments are:

high intensity	known flux
known spectrum	known composition (preferably no background)

The need for long baselines

$$P_{Oscillation} \approx (\text{Sines \& Cosines}) \sin^2 \left(\frac{\Delta m_{ij}^2 L}{4E} \right)$$

$$\Rightarrow \frac{1}{4} \lambda = \frac{2\pi E}{\Delta m_{ij}^2}$$

$$\approx \frac{0.8E}{\Delta m_{ij}^2}$$

[E in eV; Δm_{ij}^2 in eV^2 ; L in km]

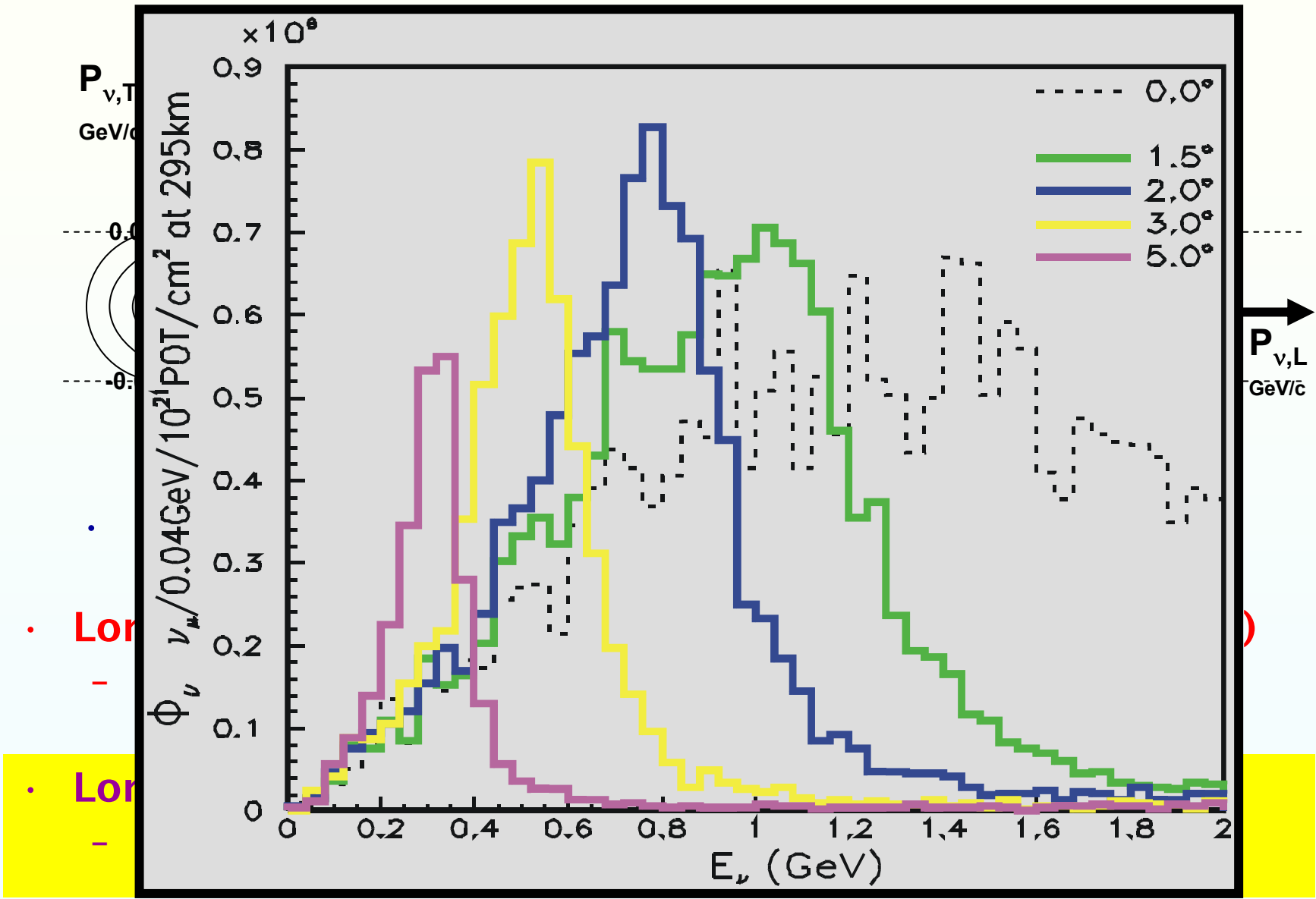
Means more intense neutrino beams!

$$\frac{1}{4} \lambda_{12} \approx 11 \times 10^3 E \text{ km}$$

$$\frac{1}{4} \lambda_{23} \approx 0.3 \times 10^3 E \text{ km}$$

Well matched to the dimensions of the earth!

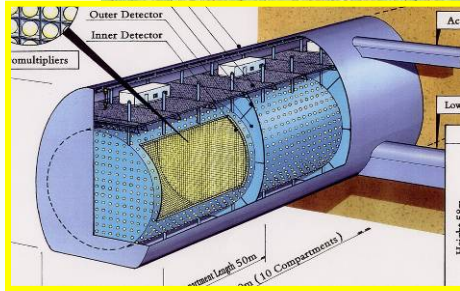
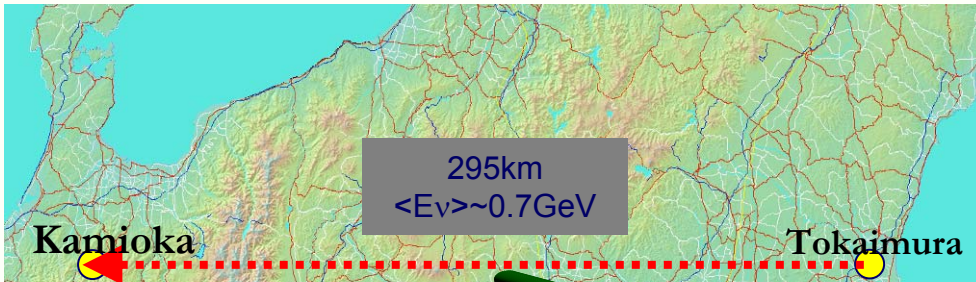
The "Off Axis" trick



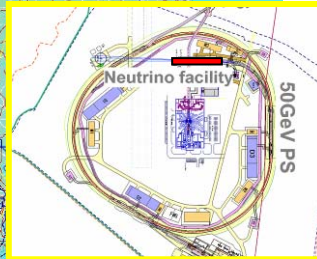
• Long

• Long

T2K & Nova?

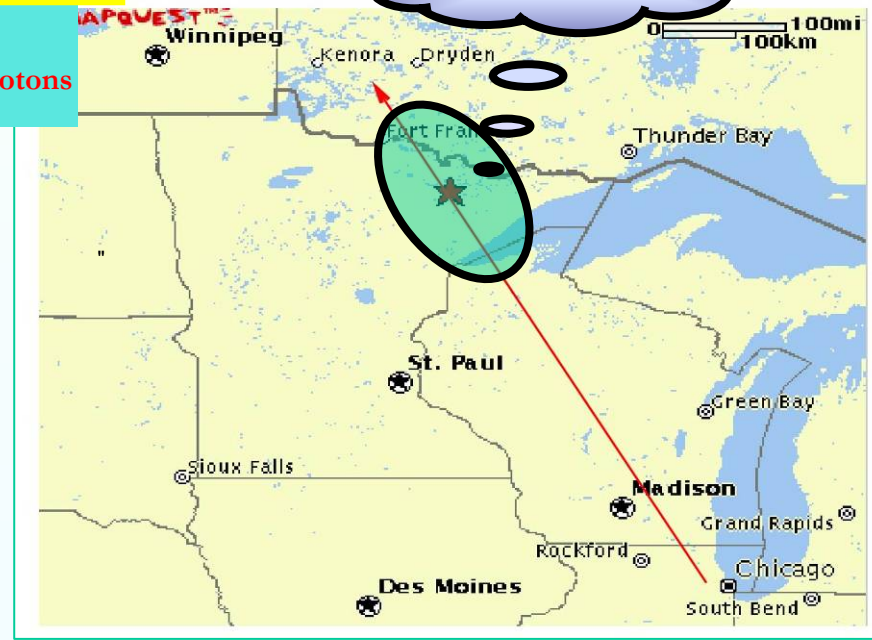


0.54Mton Kamiokande

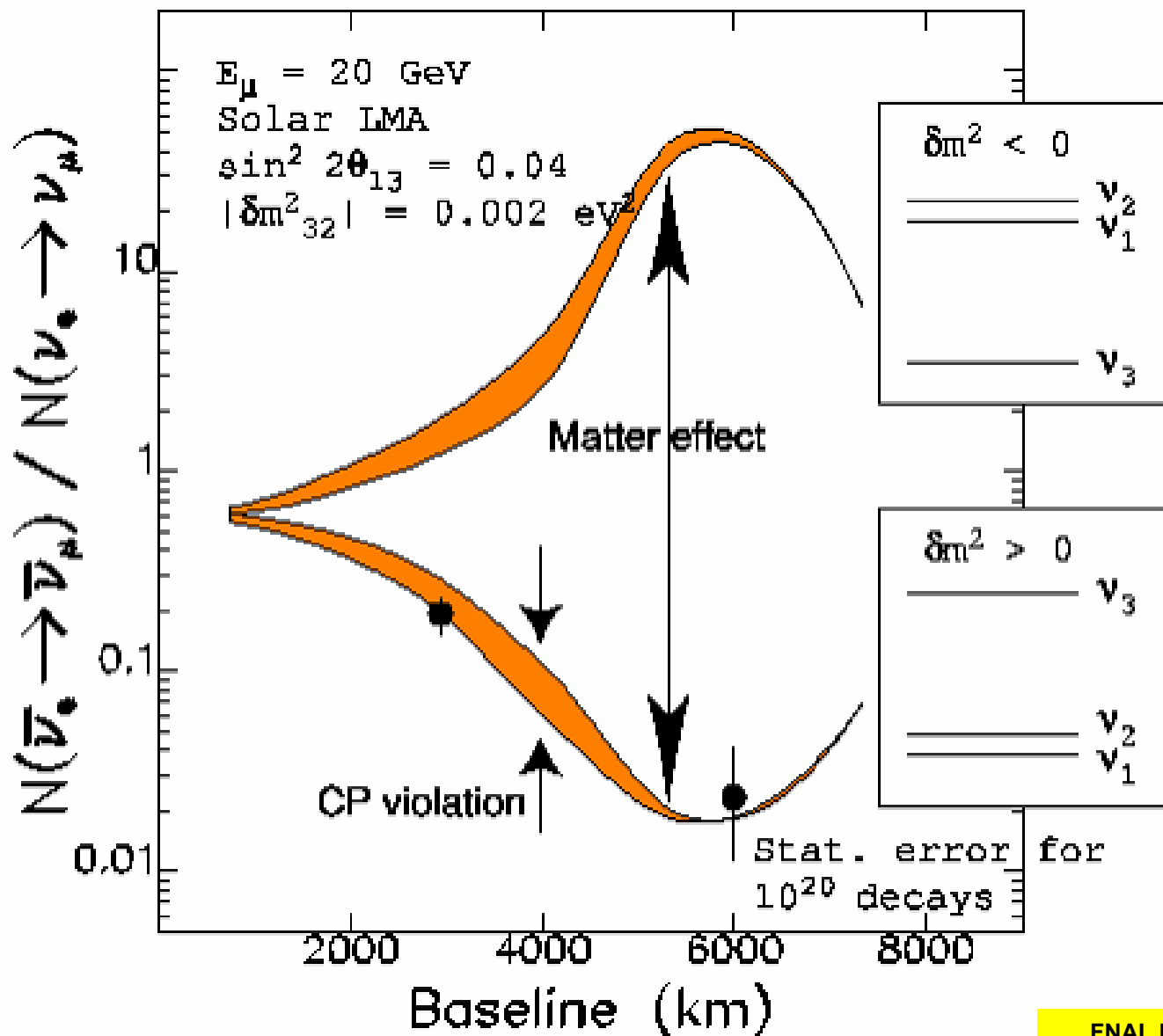


4MW 50GeV Protons

M Shiozawa NuFACT03



CP-violation





The Neutrino Factory

CPV: $> 10^{20}$ muon decays

Conventional ν beams

- π, μ & K decay
- Some flavour selectivity
- Contamination

Reactor ν beams

- Pure ν_e
- Huge Fluxes
- Very low energy (MeV)

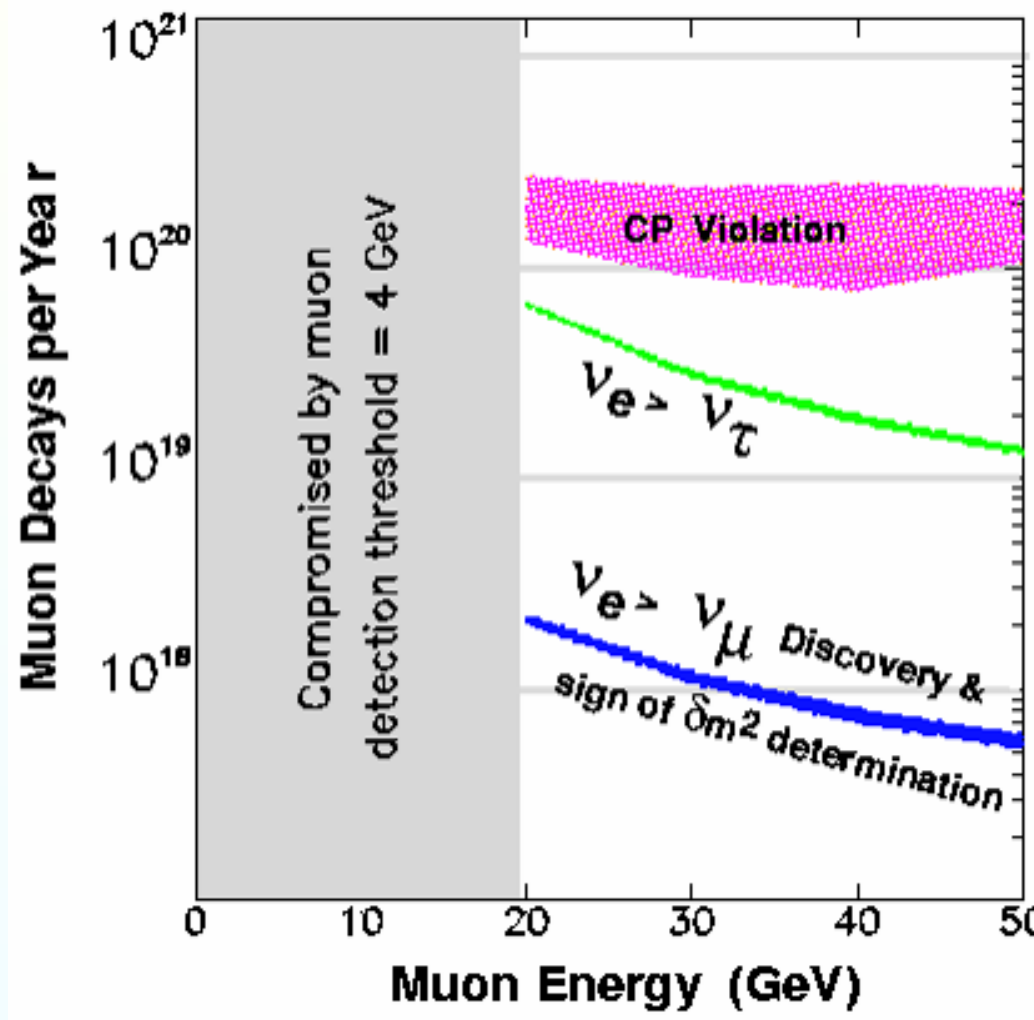
Super Conventional ν beams

- π , (& some μ) decay
- Flavour selectivity (ν_μ)
- Low Contamination at $E < 200 \text{ MeV}$

The Neutrino Factory

β beams

$L = 2800 \text{ km}, \sin^2 2\theta_{13} = 0.04$

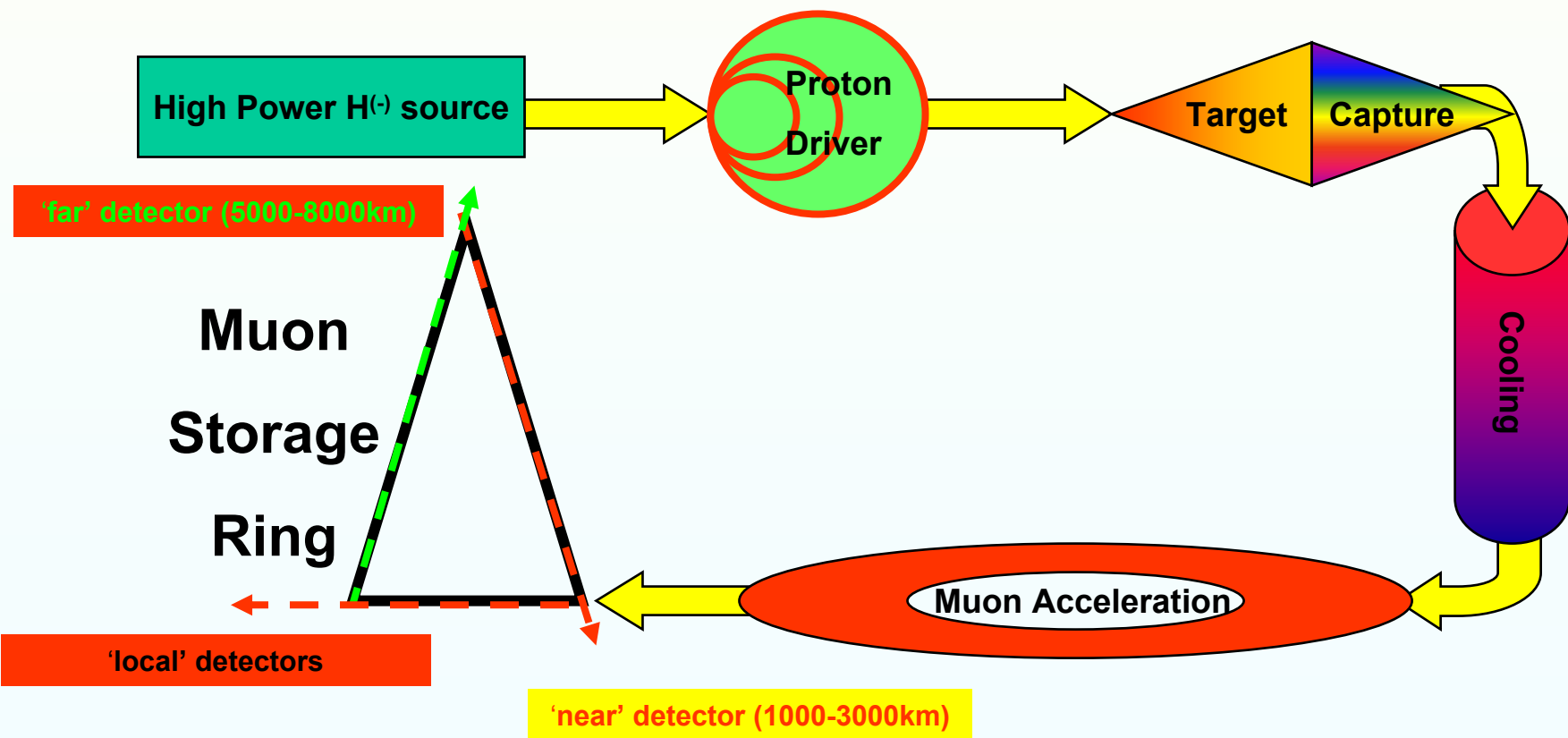


FNAL Feasibility Study 1

A Neutrino Factory is ...

... an accelerator **complex** designed to produce $>10^{20}$ muon decays per year directed at a detector thousands of km away

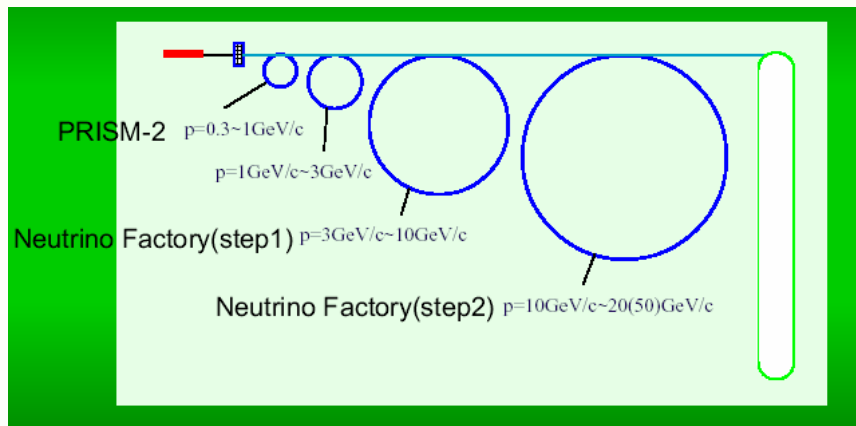
Principal Components



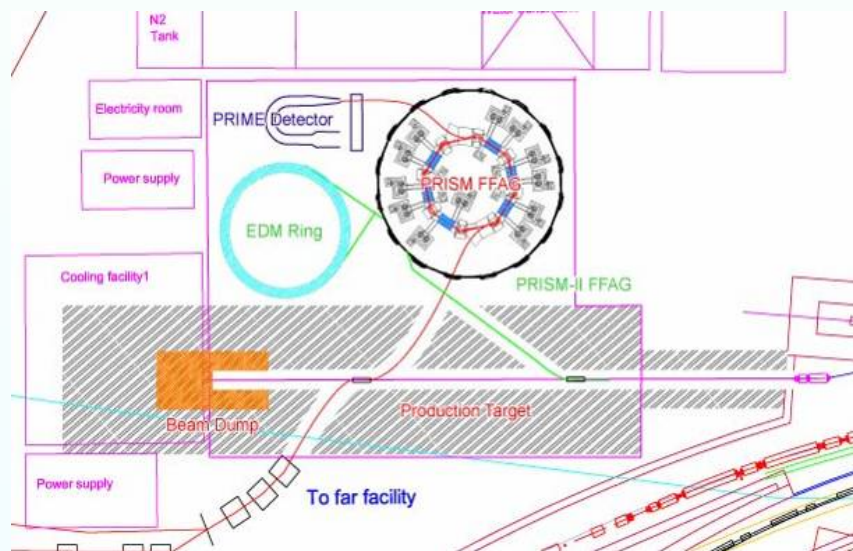
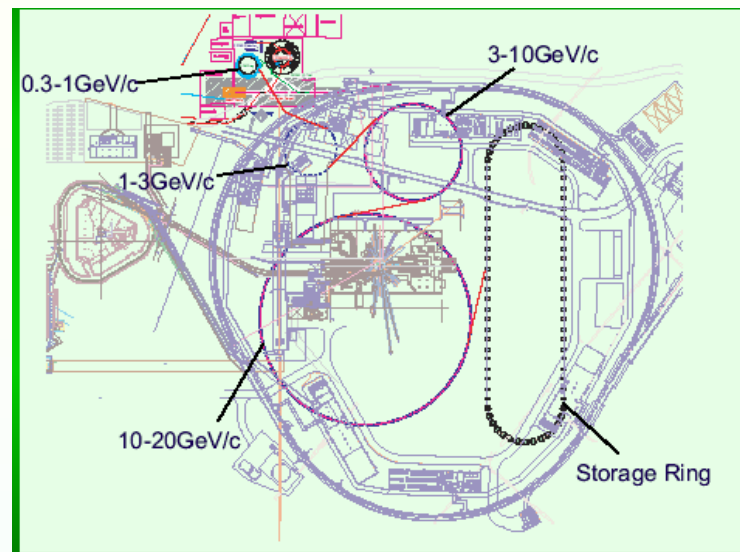


Neutrino Factory Challenges

- **Parameters**
 - **Need to know that θ_{13} is not zero**
 - Other parameters well known to fix (E_μ, L)
- **Technology**
 - **Proton driver**
 - RCS or LINAC?
 - **Proton energy?**
 - HARP, E910, MIPP
 - **Target**
 - MW beam power
 - Mercury, solid, liquid-cooled, pellet, ...
 - **Pion/muon collection and/or cooling**
 - Magnetic Horns or Solenoids?
 - Phase Rotators, FFAG's, cooling?
 - **RF and acceleration**
 - RLA's or FFAG's?
 - **Muon Storage Ring**
 - Racetrack, triangular or bow-tie
 - Conventional or FFAG?
- **Other uses of high power protons & muons?**



- **High Power Proton Driver**
 - Muon g-2
- **Muon Factory (PRISM)**
 - Muon LFV
- **Muon Factory-II (PRISM-II)**
 - Muon EDM
- **Neutrino Factory**
 - Based on 1 MW proton beam
- **Neutrino Factory-II**
 - Based on 4.4 MW proton beam



Targets

- ~ same power as SNS targets
 - Open
 - Small
 - Environmental protection?

Muon Cooling

- Certainly needed for a muon collider
- Almost certainly needed for a neutrino factory
 - (combined FFAG/cooling or ring-coolers?)

- **Requirements**

- Pulsed Beam
- Small Beams
- Short High-Z Material
- Embedded Targets
- High Beam Power

- **Tests**

- **Stress = $E\alpha\Delta T$**
- **Need tough materials to withstand stress**
- **Materials with small α (CTE) to reduce stress**
- **Beam test (short and longer term) are mandatory**

- **Options**

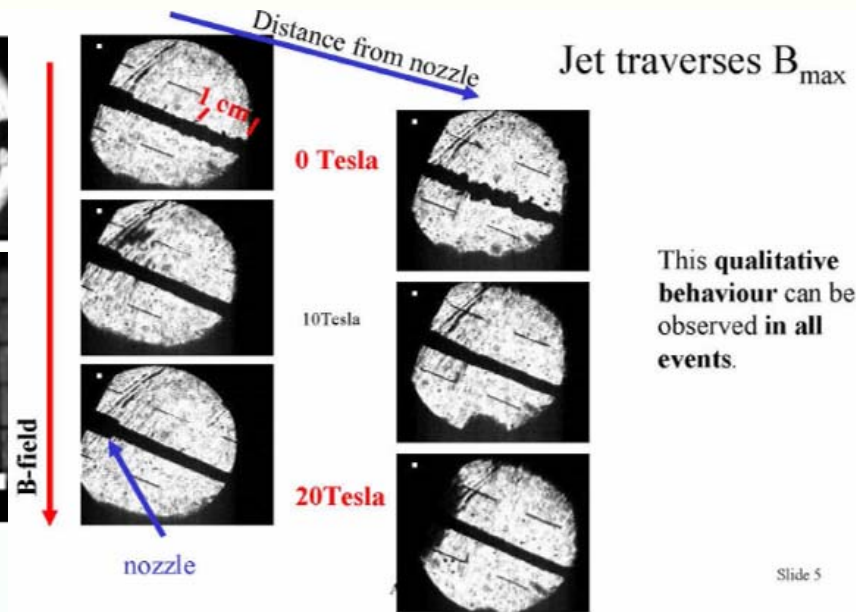
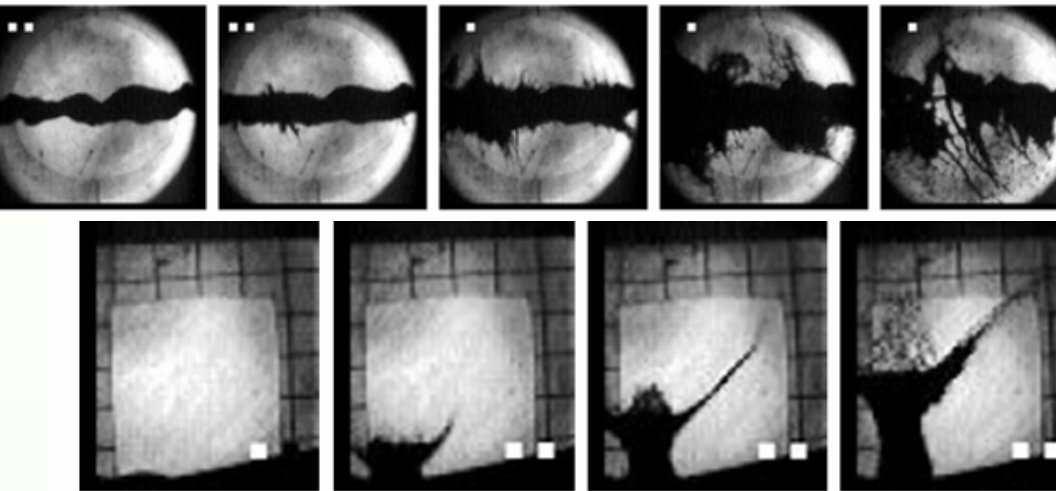
- **Solid**

- Radiation cooled
- Liquid cooled
- Continuous
- pellets

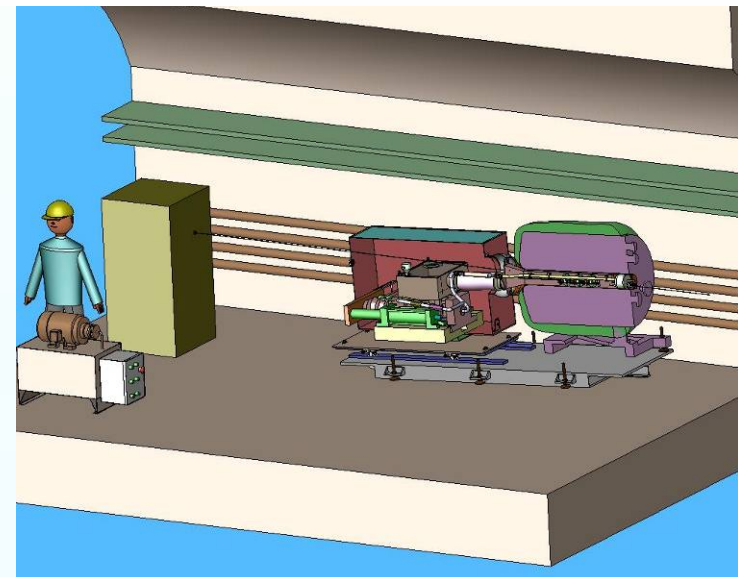
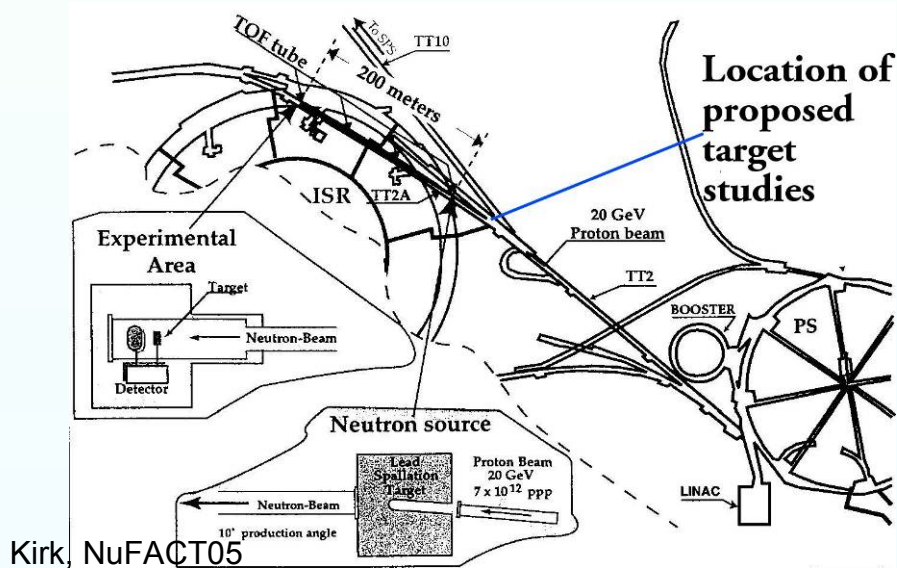
- **Liquid**

- Mercury?

Target Studies

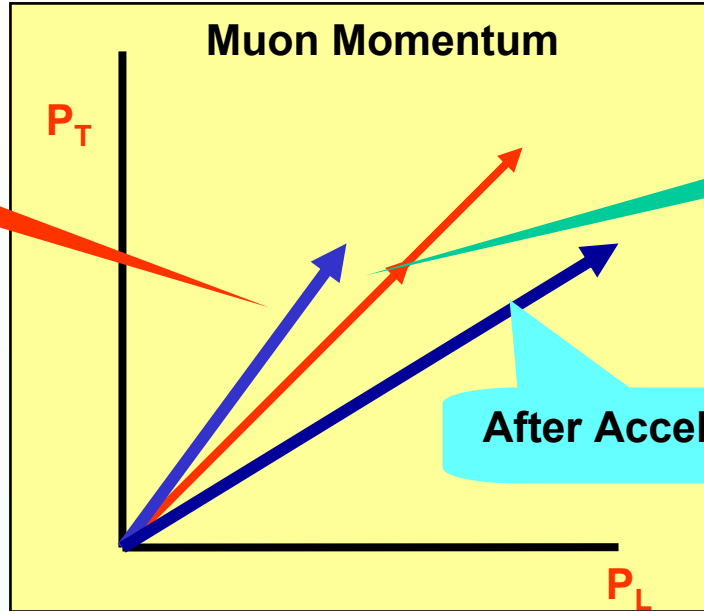


Slide 5



Ionization Cooling

After Multiple Scattering

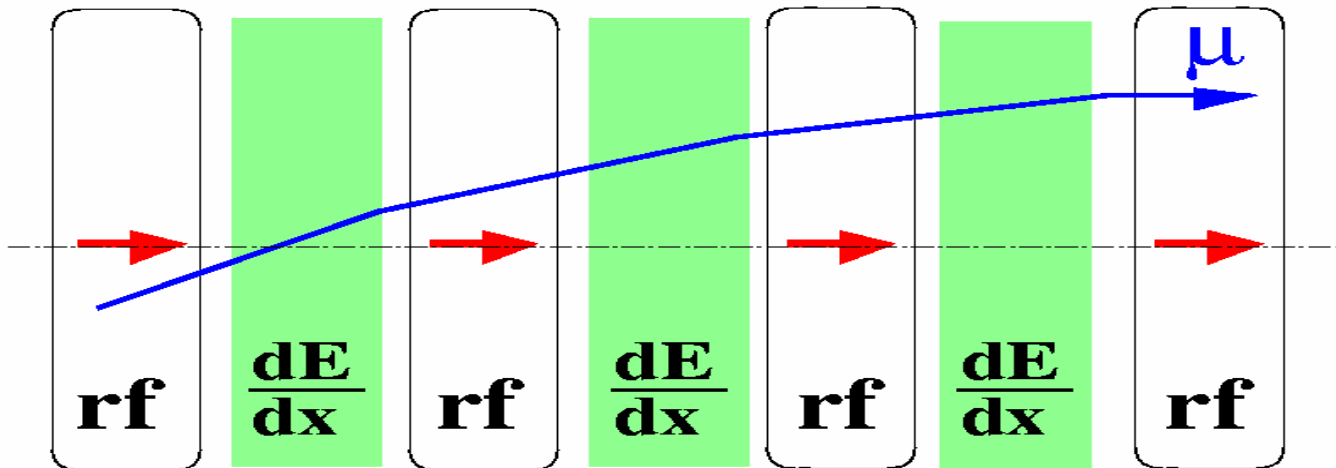


After ionisation energy loss

$$\frac{d\varepsilon_n}{ds} = -\frac{\varepsilon_n}{\beta^2 E_\mu} \frac{dE_\mu}{ds}$$

After Acceleration

$$\frac{d\varepsilon_n}{ds} = \frac{\beta_\perp}{\beta^3 m_\mu c^2 X_0 E_\mu} \frac{dE_\mu}{ds}$$

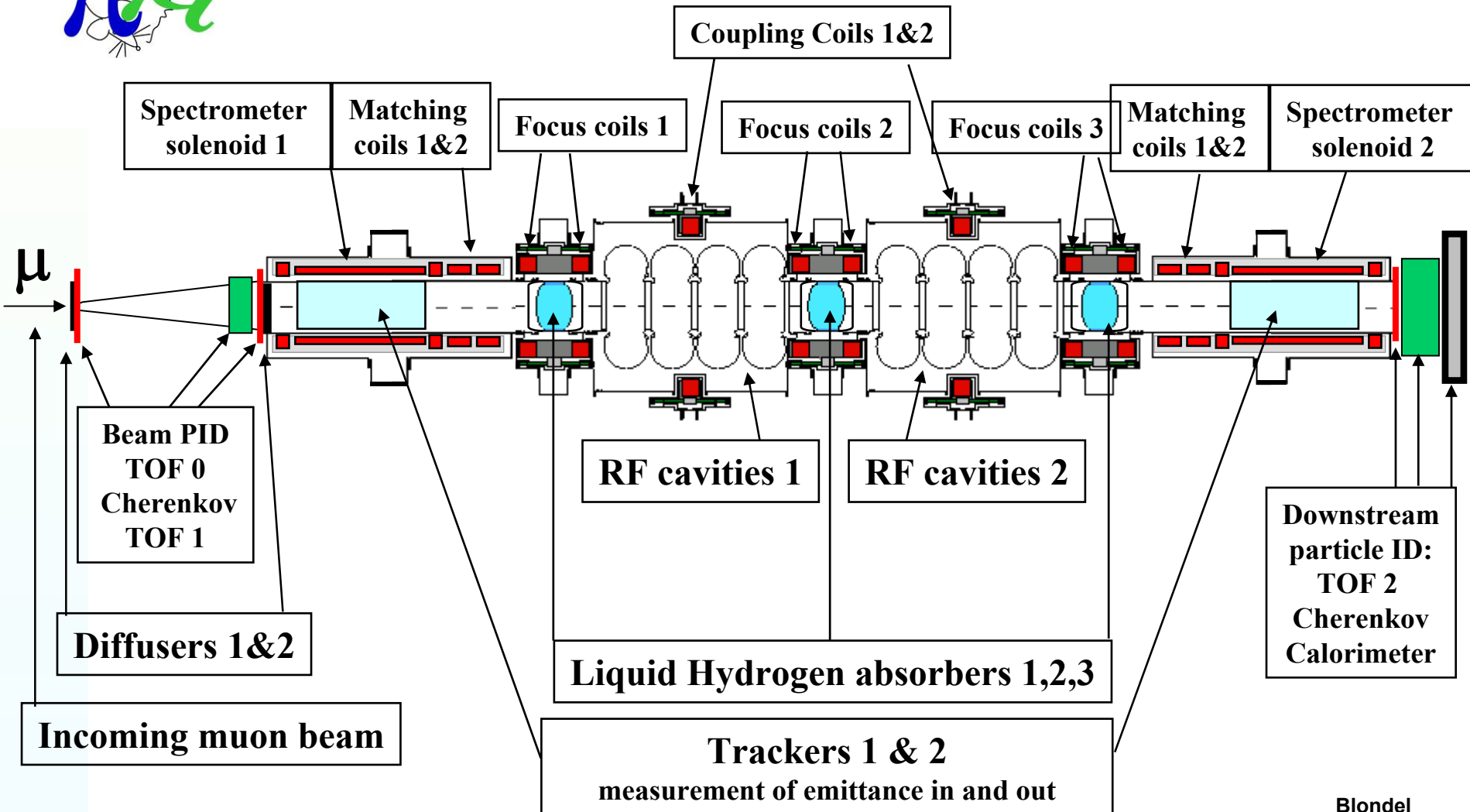
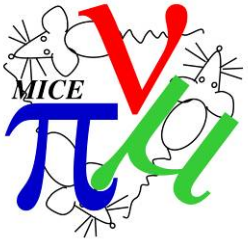


10% cooling of 200 MeV/c muons requires ~ 20 MV of RF

single particle measurements =>

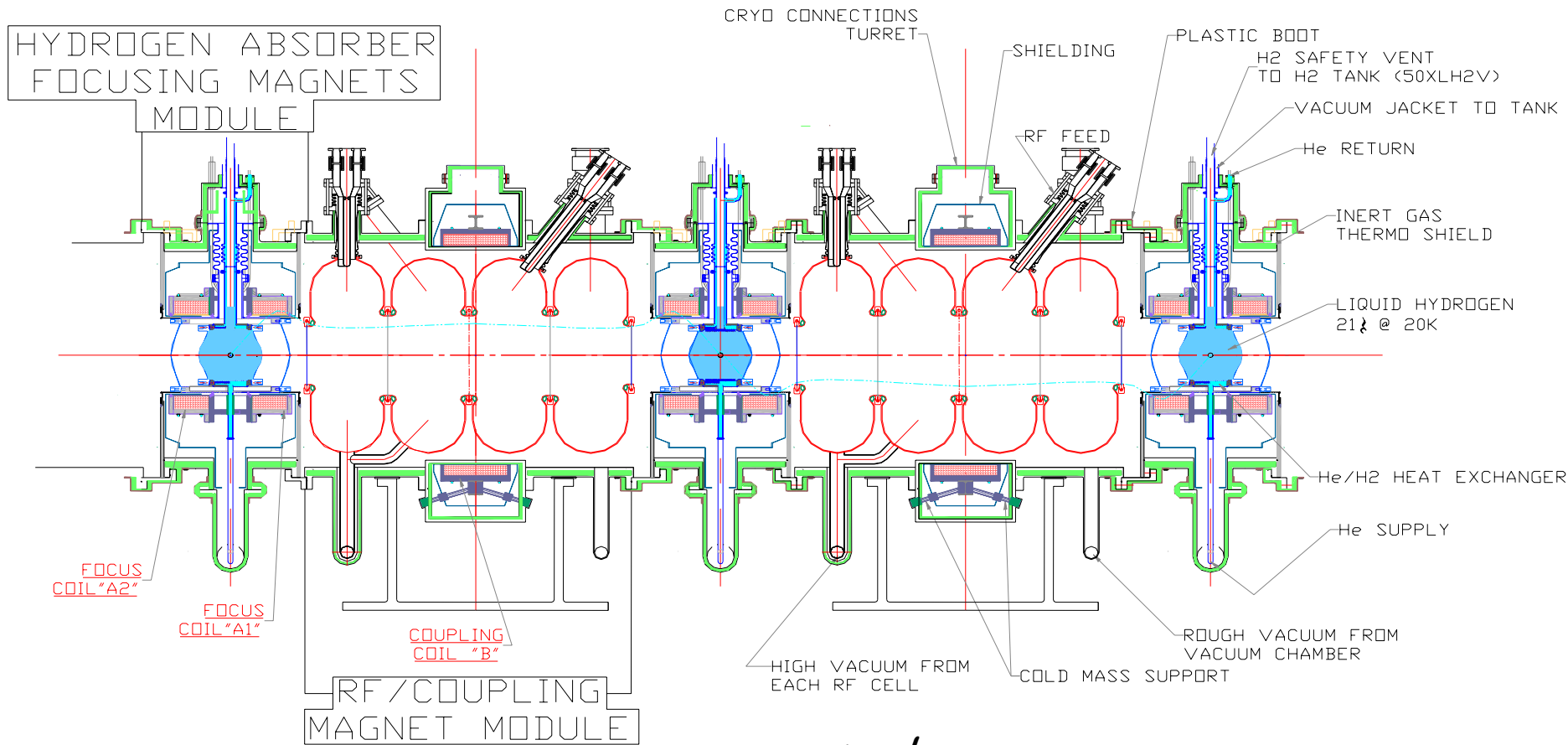
measurement precision can be as good as $\Delta(\epsilon_{\text{out}}/\epsilon_{\text{in}}) = 10^{-3}$

never done before either....

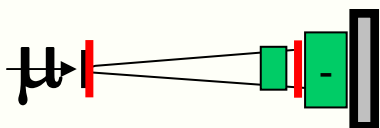


... after engineering ...

reality (simplified)

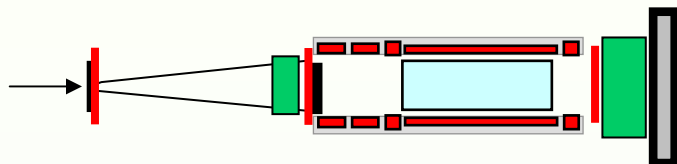


... maybe ...

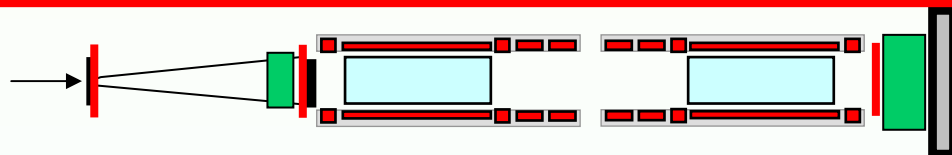


STEP I: spring 2007

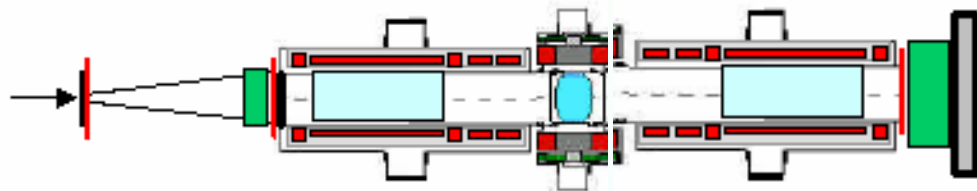
PHASE I
approved



STEP II: fall 2007

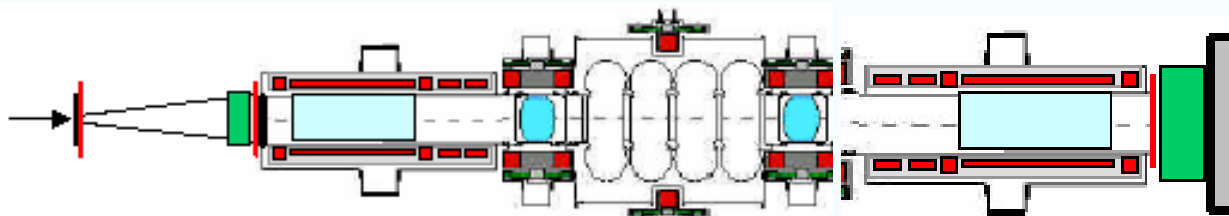


STEP III: 2008

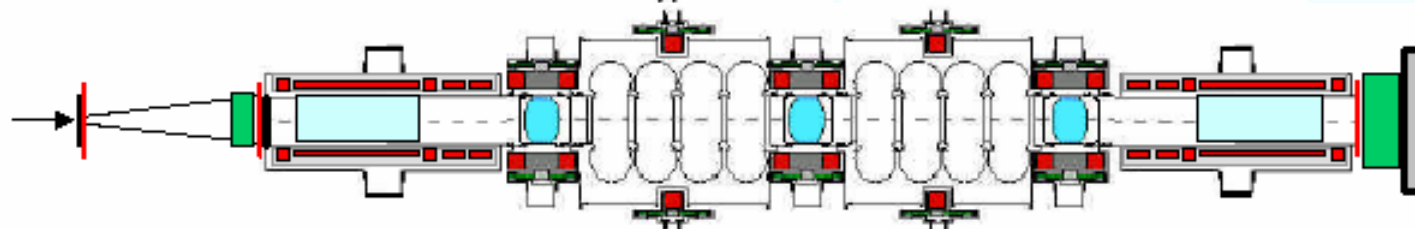


STEP IV: 2008

PHASE II
in preparation

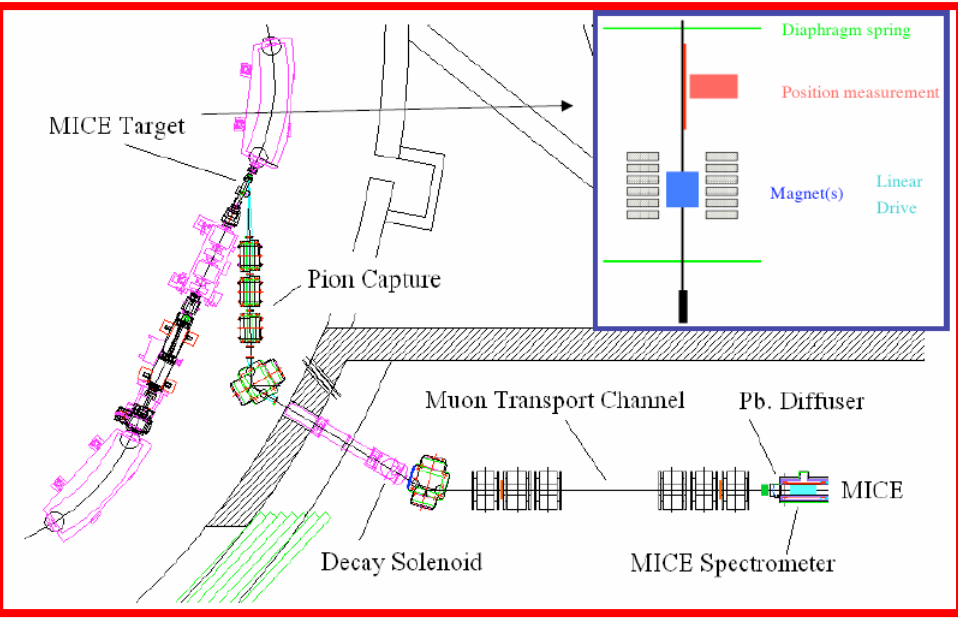
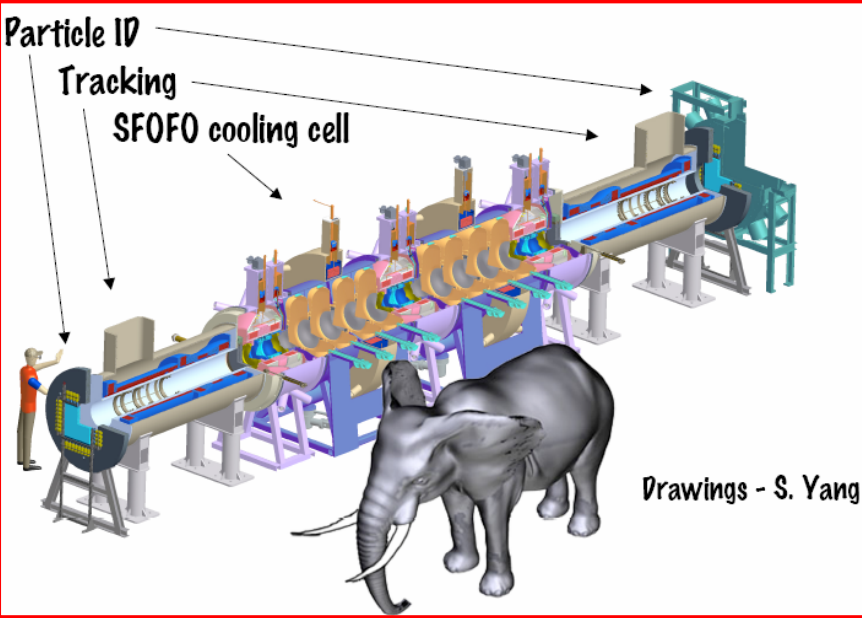


STEP V: fall 2008?

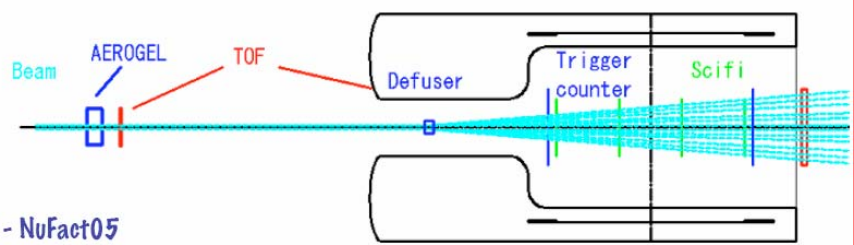
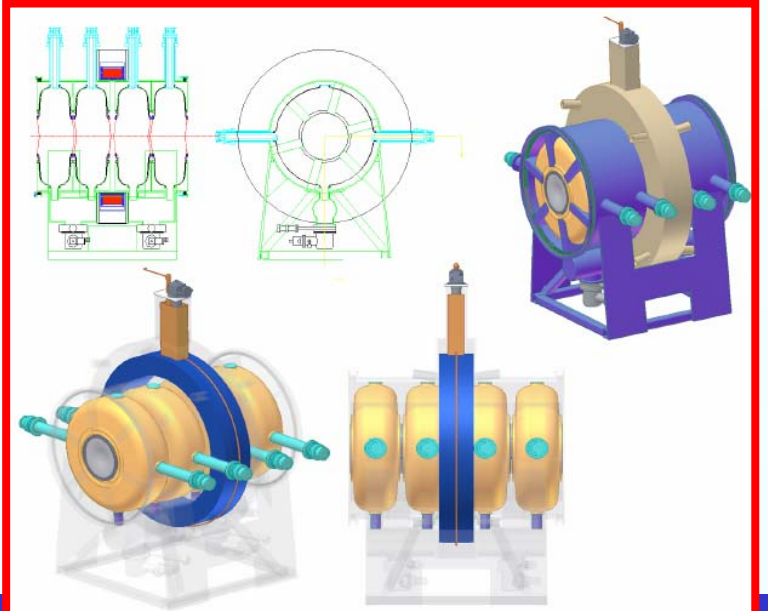
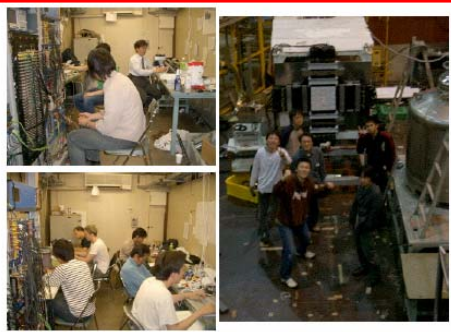


STEP VI:
2009?

MICE



- First part completed June 2 at KEK T2 beamline
- Data taken for
 - TOF calibration
 - Aerogel Cerenkov performance
 - Beam survey
 - DAQ test
- Ready to test tracker prototype



Neutrino Factory Challenges

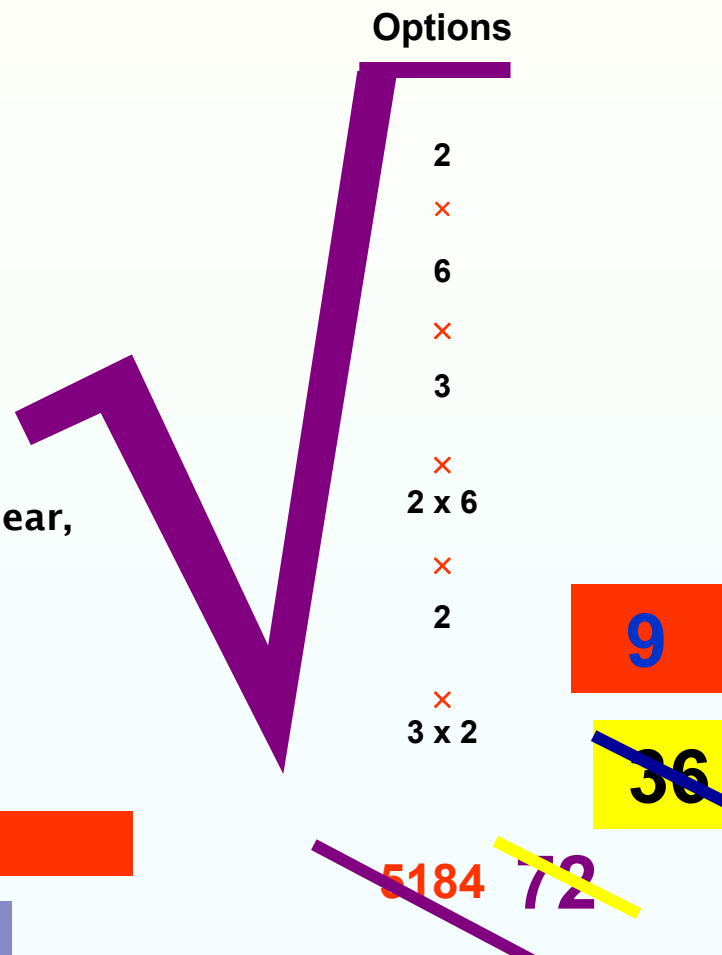
Not all match efficiently

Technology

- **Proton driver**
 - RCS or LINAC?
- **Proton energy?**
 - HARP, E910, MIPP
- **Target**
 - MW beam power
 - Mercury, solid, liquid-cooled, pellet, ...
- **Pion/muon collection and/or cooling**
 - Magnetic Horns or Solenoids?
 - Phase Rotators, FFAG's, cooling (ring, linear, ...)?
- **RF and acceleration**
 - RLA's or FFAG's?
- **Muon Storage Ring**
 - Racetrack, triangular or bow-tie
 - Conventional or FFAG?

¾ don't make any sense

~ 9 feasible NF designs!!!!

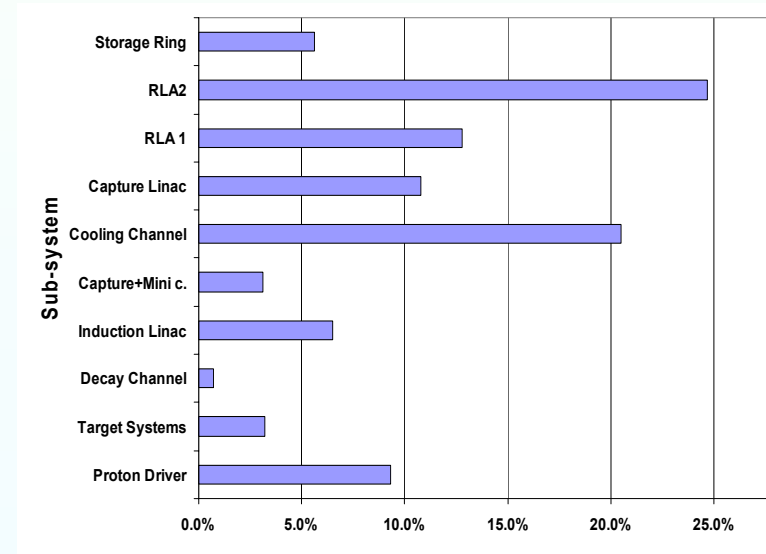
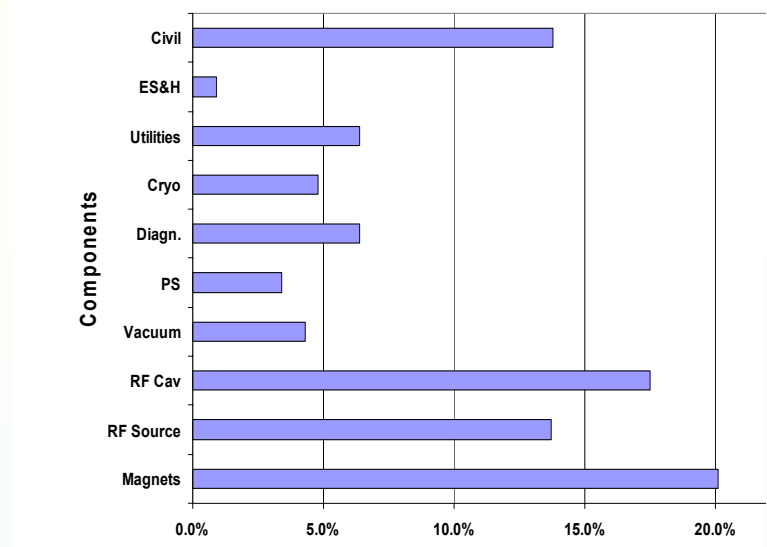


Some don't work - divide by 2



Neutrino Factory R&D

- **High Power proton drivers**
 - MW power, ns pulses
- **RF**
 - 30% of the cost?
- **Cooling**
 - How much? (20% of the cost?)
- **RLA or FFAG?**
 - Which is cheaper?





Laser-Plasma accelerators



Plasma accelerators driven by TW lasers

Tajima & Dawson *Phys Rev. Lett.* **43** 267 (1979)

VOLUME 43, NUMBER 4

PHYSICAL REVIEW LETTERS

23 JULY 1979

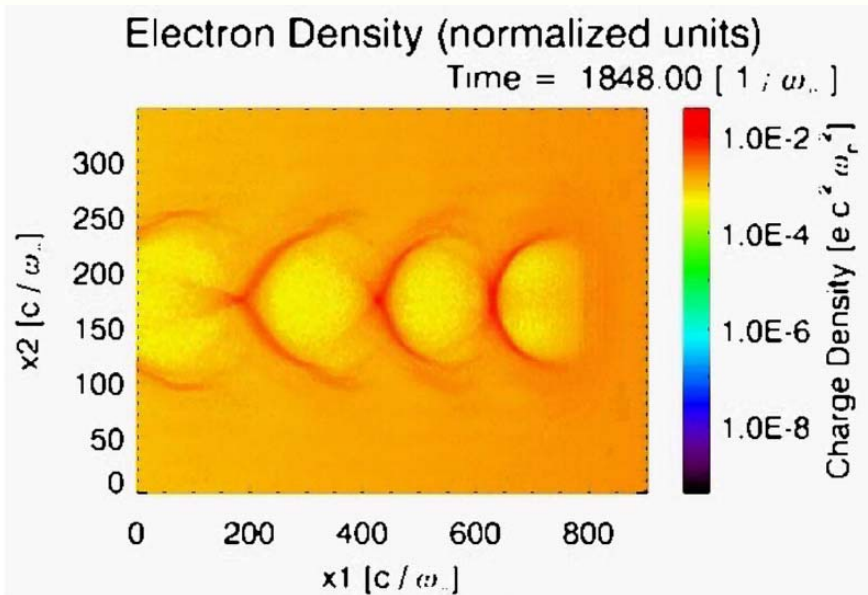
Laser Electron Accelerator

T. Tajima and J. M. Dawson

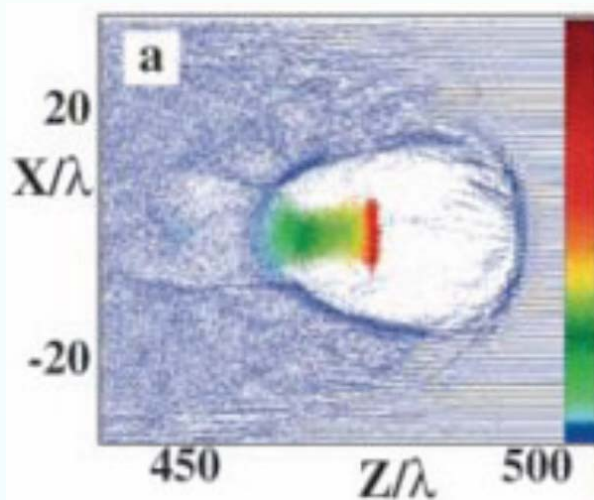
Department of Physics, University of California, Los Angeles, California 90024

(Received 9 March 1979)

An intense electromagnetic pulse can create a weak of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density $10^{18}\text{W}/\text{cm}^2$ shone on plasmas of densities 10^{18}cm^{-3} can yield gigaelectronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.



- Plasma frequency decreases with intensity.
- Wavefronts of plasma wave become curved.
- At very high intensities reach the “blow-out” or “bubble” regime.



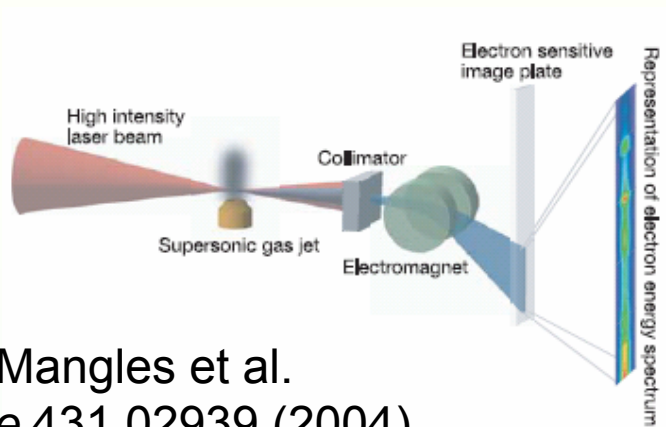
Pukhov *et al.* Appl. Phys. Lett. **74** 355 (2002)

Generation of quasi-monoenergetic beams

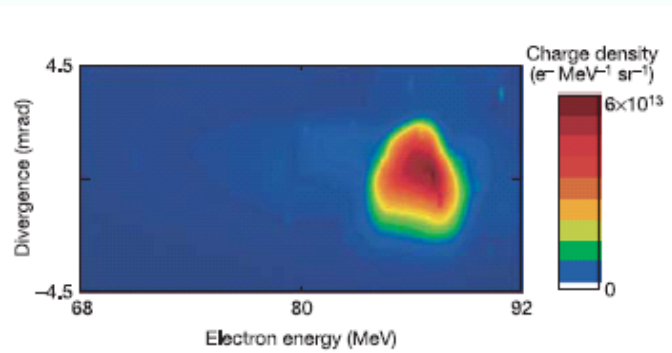
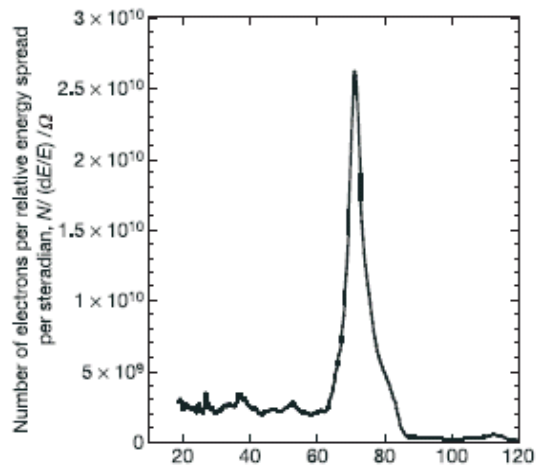


- Three milestone results published in Nature at end of 2004 by:
 - Karl Krushelnick (Imperial College, UK)
 - Victor Malka (LOA, France)
 - Wim Leemans (Lawrence Berkeley, USA)
- These showed evidence of *quasi-monoenergetic* electron beams for the first time.

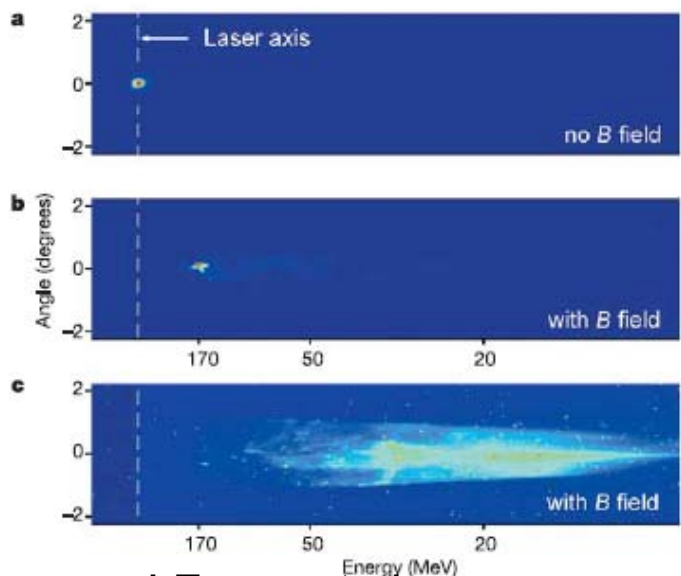
Generation of quasi-monoenergetic beams



S. D. Mangles et al.
Nature 431 02939 (2004)

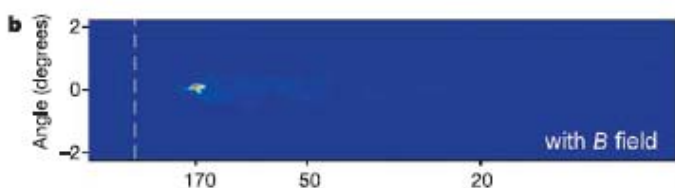
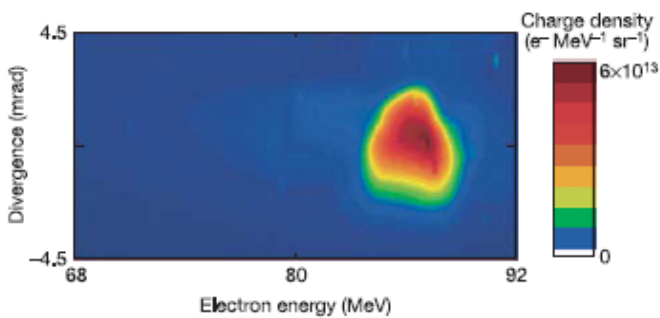
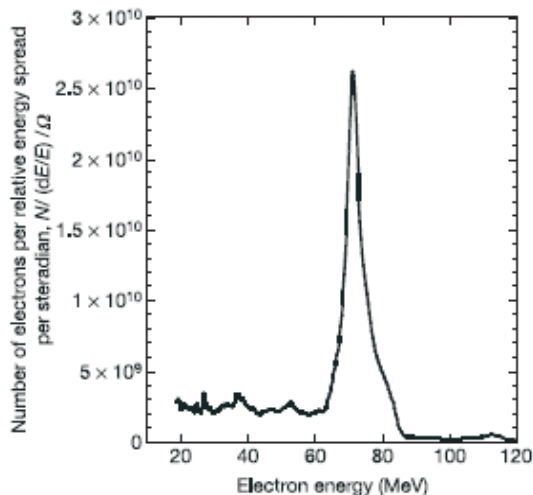


C. G. R. Geddes et al.
Nature 431 02900 (2004)



J Faure et al.
Nature 431 02963 (2004) Hooker, Oxford

Generation of quasi-monoenergetic beams



- Typical output parameters:

- Output energy: 100 - 170 MeV
- Energy spread: 2.5 - 8%
- Bunch charge: 20 - 500 pC
- Normalized emittance: 1-2 π mm mrad

- **In the last 5 years laser-driven plasma accelerators have made enormous progress:**
 - **Demonstration of quasi-mononergetic beams**
 - **Increase of output energy to 1 GeV**
 - **Demonstration of controlled injection**



Future Accelerators for other sciences

What is needed and why

X-ray sources
No time to discuss this

Neutron sources

Nuclear Physics
... sorry, because it is important

Future Accelerators for other applications

Accelerators in Medicine

Development of NS-FFAG EMMA & PAMELA

(After Takeichiro Yokoi)

Introduction ...

- FFAG(Fixed Field Alternating Gradient) Accelerator
 - rapid particle acceleration
 - large beam acceptance
 - wide variety of applications
 - fundamental science (e.g. Neutrino factory)
 - practical applications (e.g. cancer therapy)
- Compared to existing fixed field accelerators, a new approach to FFAG, **Non-Scaling FFAG**, has advantages such as small beam excursion and flexibility in machine design and operation and variable energy beam extraction
- However, no NS-FFAG has ever been realized up to now.



Real working machine is needed !!

CONFORM (Construction of a Non-scaling FFAG for Oncology, Research and Medicine) aims to develop the Non-scaling FFAG as a versatile accelerator. (Project HP: www.conform.ac.uk)

Scaling FFAG & Non-Scaling FFAG

Both have large acceptance and ability of fast acceleration

* Acceleration speed of fixed field accelerator is unlimited

(in synchrotron, ramping speed of magnet limits the repetition rate)

- **Scaling FFAG**

- Similar orbit shape

- Large beam excursion

- Stable betatron tune

- Combined function($B \propto r^k$)

- **Non-Scaling FFAG**

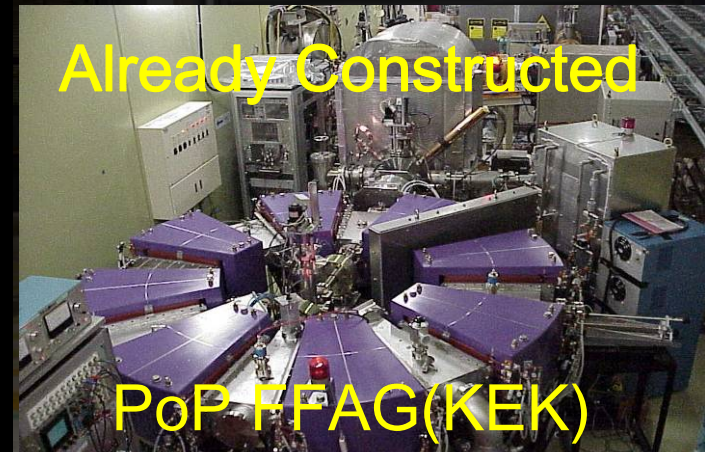
- Non-similar orbit shape

- Small beam excursion

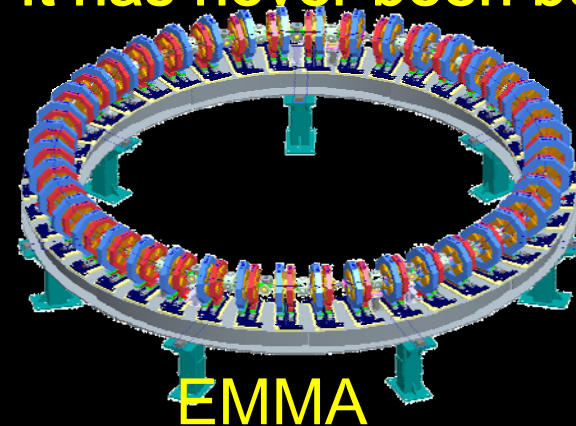
- (small path length variation)

- Large tune change

- Linear lattice (quadrupole etc)



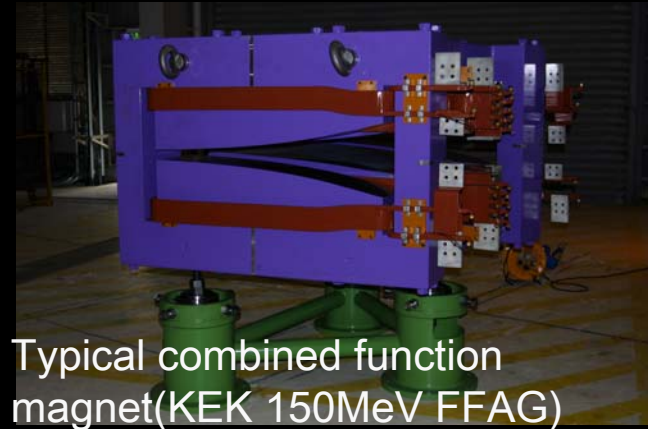
It has never been built



How to realize NS-FFAG lattice...

Element of NS-FFAG provides bending and focusing field in a magnet

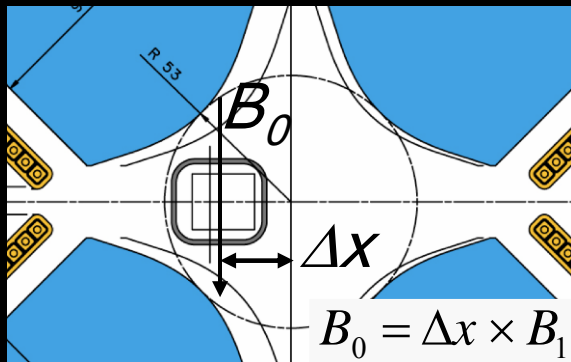
⇒ “**Combined function magnet**”



Typical combined function magnet(KEK 150MeV FFAG)

In NS-FFAG, focusing force is **linear** (quadrupole field) ...

⇒ **Shifted quadrupole** magnet works as a combined function magnet

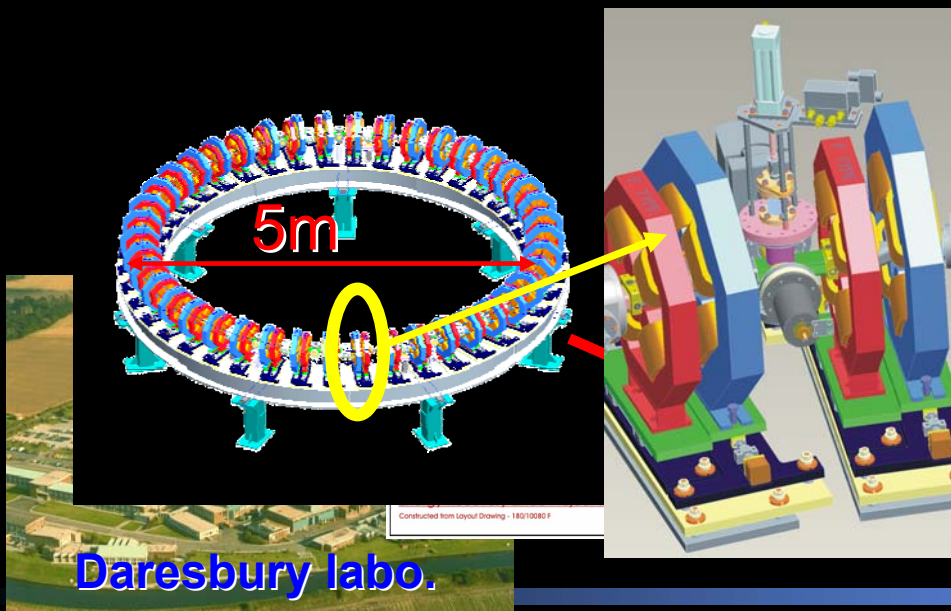


By changing the horizontal position of magnet, bending power and focusing power can be changed independently.

⇒ “**Separated function FFAG**”

EMMA: Electron Model for Many Applications

- Electron NS-FFAG as a **proof of principle** is to be built as 3-year project.(host lab: Daresbury lab.)
- It is also **a scaled-down model of muon accelerator** for neutrino factory.
- Research items are . . .
 - (1) Research of beam dynamics of NS-FFAG
 - (2) Demonstration of NS-FFAG as a practical accelerator
 - (3) Demonstration of fast acceleration with fixed frequency RF



Number of Cell	42 (doublet Q)
Circumference	16.57m
Injection energy	10~20MeV(variable)
Extraction energy	10~20MeV(variable)
RF	1.3GHz
Acceptance	3mm(normalized)

Daresbury labo.

Beam acceleration : EMMA

Resonance is a **coherent** effect

⇒ Fast acceleration can circumvent the problem

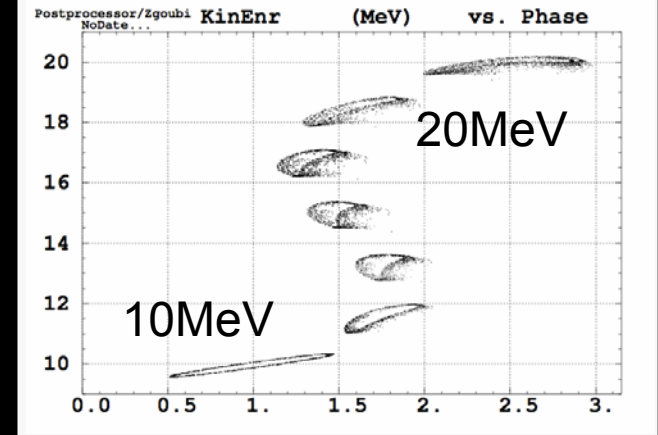
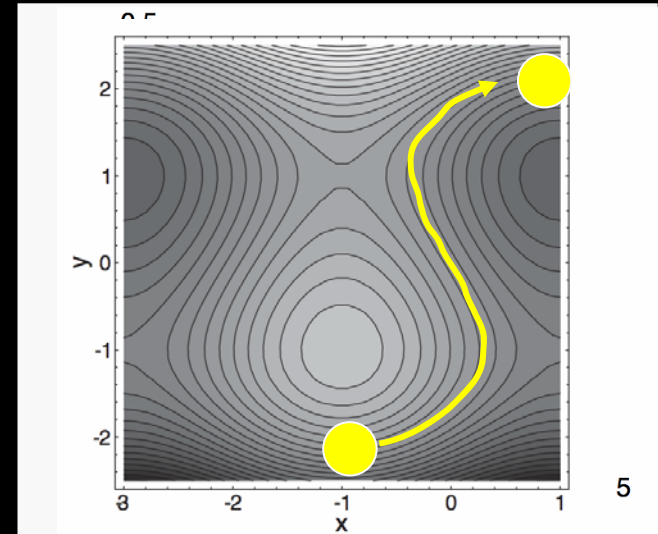
Resonant crossing acceleration

Small variation of path length makes it possible to adopt **fixed frequency rf** for relativistic particle

Fast asynchronous acceleration

* In EMMA, Acceleration completes within 10 turns (~500ns)

EMMA is a unique system to observe transient process of resonance precisely.
⇒ Unique playground for nonlinear dynamics !!

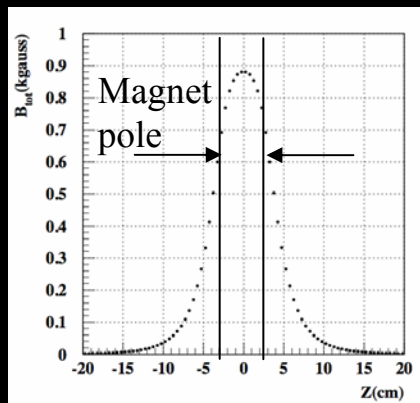


Linear Model, Nonlinear Reality

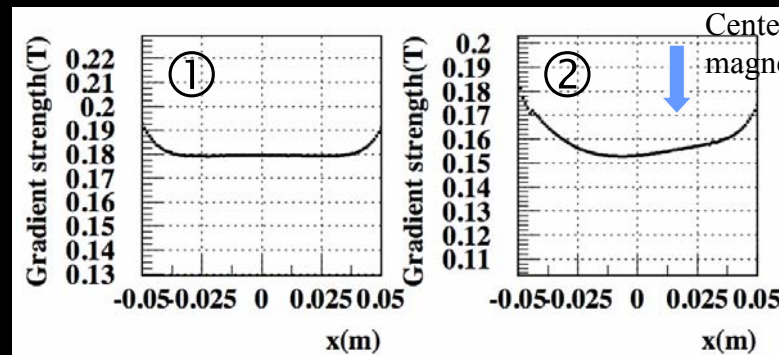
in the actual lattice of EMMA ...

Magnet aperture \sim Magnet length \sim Magnet distance

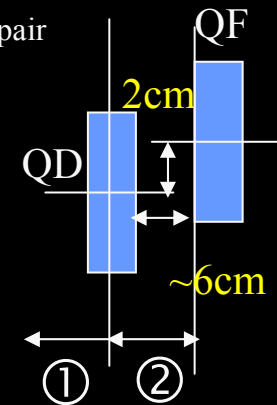
\Rightarrow Severe nonlinearity arises due to coupling and fringing field



Fringing field is dominant!!



Inter-magnet coupling introduces strong nonlinearity

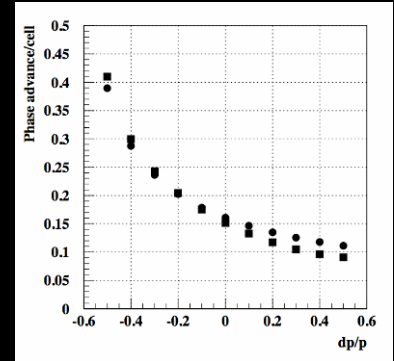


Tracking study with realistic 3D field is indispensable in machine design

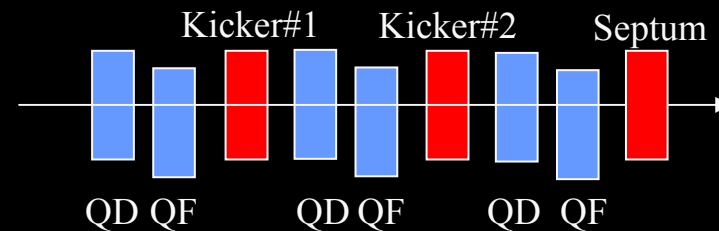
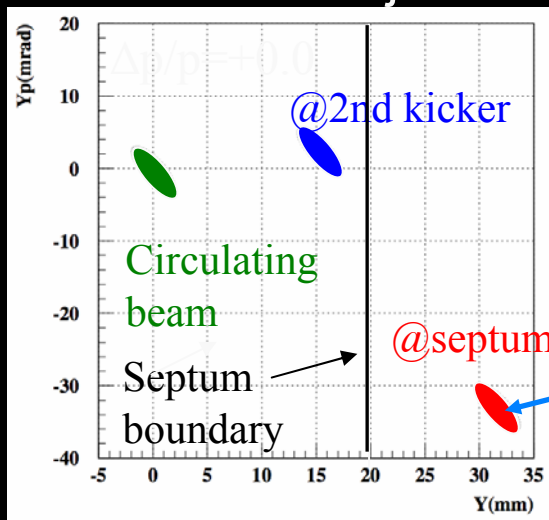
Beam Injection & extraction

Small beam excursion of NS-FFAG makes energy variable beam extraction easier

⇐ Unique feature for fixed field accelerator



However, large tune change requires phase adjustment mechanism in injection & extraction \Rightarrow **multi-kicker system**

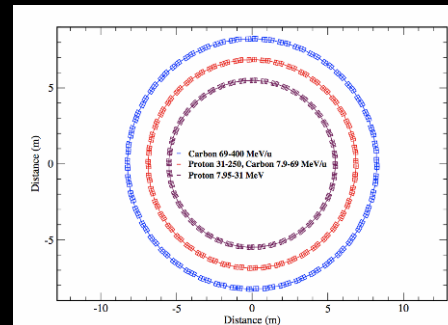


By changing the field strength and direction, beam position in phase space can be adjusted

Example of beam extraction (PAMELA)

PAMELA : Particle Accelerator for Medical Applications

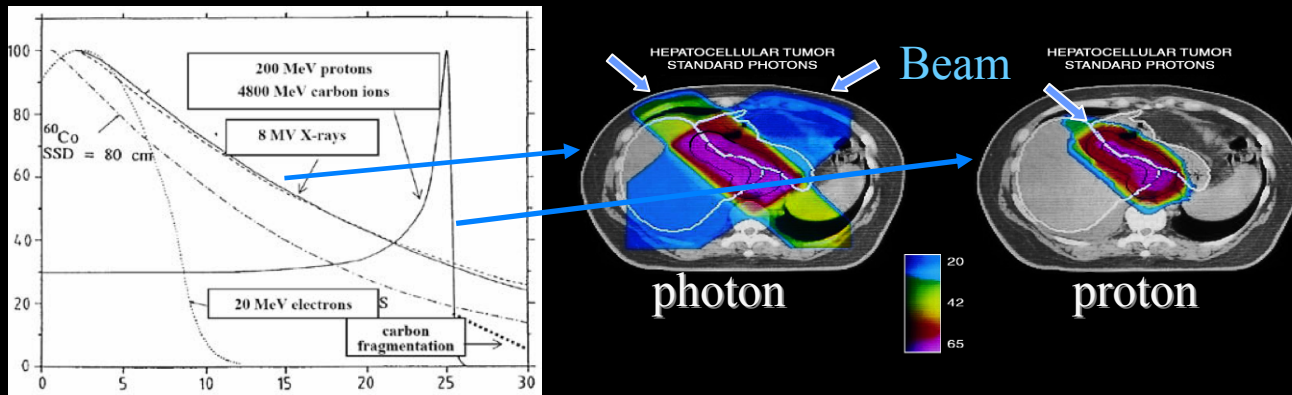
- Particle therapy has advantages in cancer therapy compared to X-ray therapy due to **good dose concentration and better biological effectiveness** (especially HI therapy).
- As an accelerator for particle therapy, the advantage of FFAG is **higher intensity** compared to ordinary synchrotron, **flexible machine operation** compared to cyclotron, and **simultaneous(multi-port) beam extraction**
- PAMELA aims to design particle therapy accelerator facility for proton and carbon using NS-FFAG
- It also aims to design a smaller machine for biological study as a prototype.



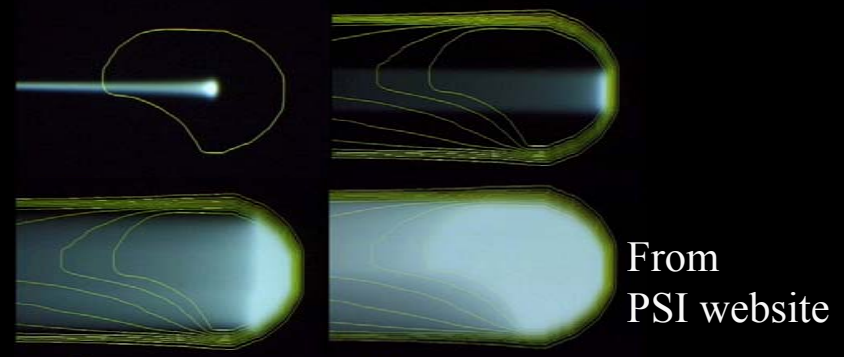
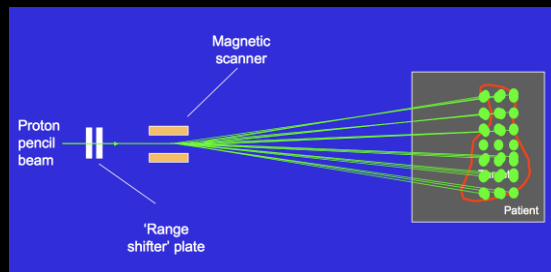
3-ring scheme by E.Keil,
A.Sessler, D. Trbojevic

Particle therapy

With the help of **Bragg peak**, proton and heavy ion beam can form sharp-edged irradiation field, and can minimize radiation damage to normal tissues.

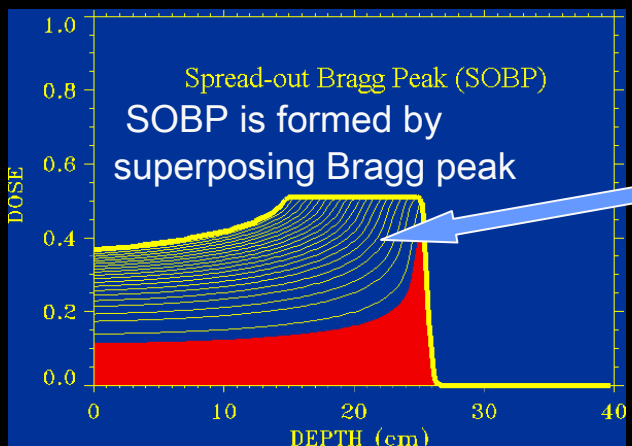


Spot scanning can fully exert the advantage of particle therapy and **pulsed beam** of FFAG matches well to the treatment



Monitoring and Control

:Key issues for medical applications



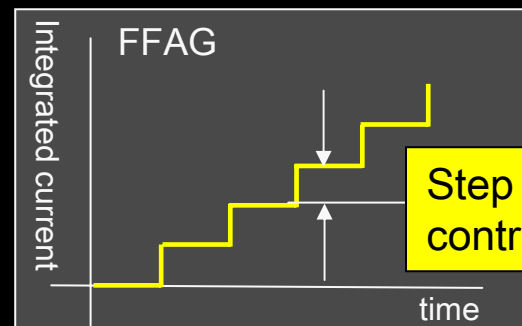
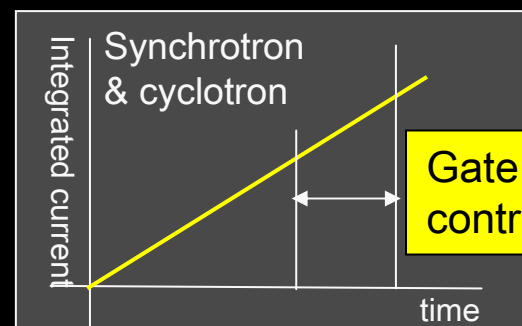
Dose uniformity should be $< \sim 2\%$
 \Rightarrow To achieve the uniformity, precise **intensity modulation** is a must

Beam of FFAG is quantized.

\Rightarrow Active intensity control at the injection level and precise loss control are indispensable.

New approach to medical accelerator control is required in PAMELA

(New postdoc is employed for the issue)



Beam acceleration : PAMELA

Two approaches in NS-FFAG
for non-relativistic beam acceleration.....

(1) Harmonic number jump (A. Ruggiero)

- Fixed frequency RF (high Q rf : high gradient)
- Amplitude modulation

← Can high Q cavity accommodate amplitude modulation ?

(2) Frequency modulation

- low Q rf (low gradient)
- no need of amplitude modulation

(adiabatic capture requires AM)

← Can beam be accelerated sufficiently fast?

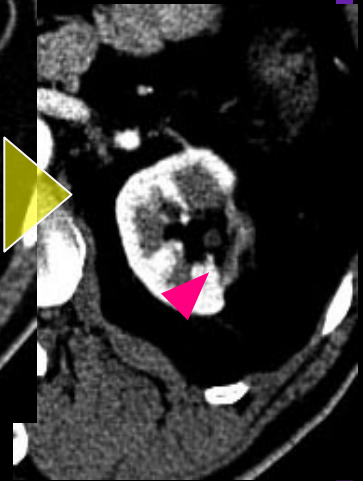
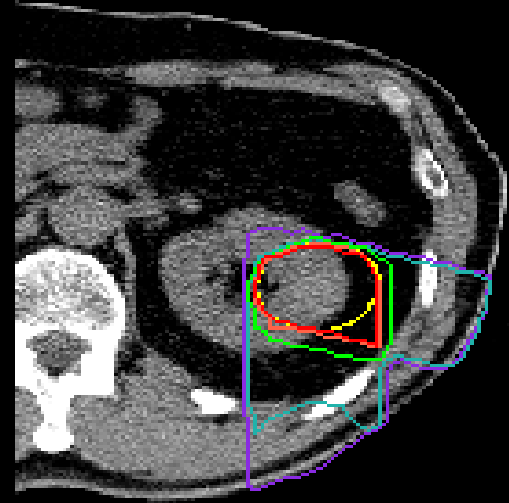
How fast beam should be accelerated in NS-FFAG ?

* Now, preparing for the study

Cancer of the Kidney

Stage I: T1a N0 M0

80GyE / 16fr. /4wks



Summary

- NS-FFAG is a novel accelerator and will open up new fields in accelerator science and its application
- COMFORM is to be engaged in two developments, EMMA(electron) and PAMELA(proton,HI)

Results of EMMA and design of PAMELA will come within 4 years

- **Particle Accelerators have an exciting future**
 - **In particle physics**
 - LHC, LC, CLIC, NF, factories ...
 - **In other sciences**
 - Light sources, FELs, spallation sources
 - **In society**
 - Medical accelerators (isotopes, hadron-therapy...)
- **And they are *fun too!***