Beam-beam effects

(an introduction)

Werner Herr CERN, AB Department

(http://cern.ch/lhc-beam-beam/talks/Daresbury_beambeam.pdf)

Werner Herr, beam-beam effects, CAS 2007, Daresbury

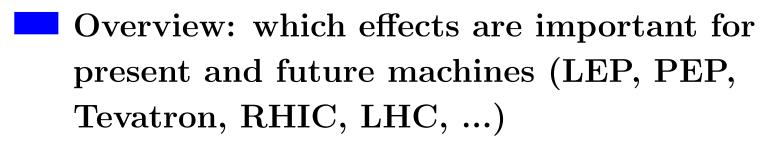
BEAMS: moving charges

- Beam is a collection of charges
- Represent electromagnetic potential for other charges
- ➡ Forces on itself (space charge) and opposing beam (beam-beam effects)
- \rightarrow Main limit in past, present and future colliders
- → Important for high density beams, i.e. high intensity and/or small beams: for high luminosity !

Beam-beam effects

Remember:

$$\implies \mathcal{L} = \frac{N_1 N_2 f n_B}{4\pi \sigma_x \sigma_y} = \frac{N_1 N_2 f n_B}{4\pi \cdot \sigma_x \sigma_y}$$



- Qualitative and physical picture of the effects
- Mathematical derivations in: Proceedings, Zeuthen 2003

Beam-beam effects

- A beam acts on particles like an electromagnetic lens, but:
 - Does not represent simple form, i.e. well defined multipoles
 - Very non-linear form of the forces, depending on distribution
 - Can change distribution as result of interaction (time dependent forces ..)
 - Results in many different effects and problems

Fields and Forces (I)

Need fields \vec{E} and \vec{B} of opposing beam.

In rest frame only electrostatic field: *E*', *B*' = 0
 Derive potential U(x, y, z) from Poisson equation:

$$\Delta U(x, y, z) = -\frac{1}{\epsilon_0}\rho(x, y, z)$$

The electrostatic fields become:

$$\vec{E}' = -\nabla U(x, y, z)$$

Fields and Forces (II)

Transform into moving frame and calculate Lorentz force \vec{F} on particle with charge $q = Z_2$ e

$$E_{\parallel} = E'_{\parallel}, \quad E_{\perp} = \gamma \cdot E'_{\perp} \text{ with }: \quad \vec{B} = \vec{\beta} \times \vec{E}/c$$

$$\vec{F} = q(\vec{E} + \vec{\beta} \times \vec{B})$$

Example Gaussian distribution:

$$\rho(x,y,z) = \frac{NZ_1e}{\sigma_x\sigma_y\sigma_z\sqrt{2\pi}^3} \exp\left(-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2} - \frac{z^2}{2\sigma_z^2}\right)$$

Simple example: Gaussian

For 2D case the potential becomes (see proceedings):

$$U(x, y, \sigma_x, \sigma_y) = \frac{NZ_1 e}{4\pi\epsilon_0} \int_0^\infty \frac{\exp(-\frac{x^2}{2\sigma_x^2 + q} - \frac{y^2}{2\sigma_y^2 + q})}{\sqrt{(2\sigma_x^2 + q)(2\sigma_y^2 + q)}} dq$$

Can derive \vec{E} and \vec{B} fields and therefore forces

For arbitrary distribution (non-Gaussian): difficult (or impossible, numerical solution required)

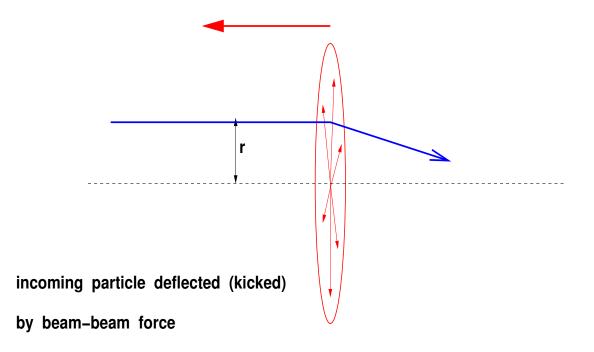
Simple example: Gaussian

Round beams: σ_x = σ_y = σ, Z₁ = −Z₂ = 1
Only components E_r and B_Φ are non-zero
→ Force has only radial component, i.e. depends only on distance r from bunch centre (where: r² = x² + y²) (see proceedings)

$$F_r(\mathbf{r}) = -\frac{Ne^2(1+\beta^2)}{2\pi\epsilon_0 \cdot \mathbf{r}} \left[1 - \exp(-\frac{\mathbf{r}^2}{2\sigma^2})\right]$$

Beam-beam kick:

• We need the deflection (kick) of the particle:



Beam-beam kick:

- → Kick $(\Delta r')$: angle by which the particle is deflected during the passage
- → Derived from force by integration over the collision (assume: $m_1=m_2$ and $Z_1=-Z_2=1$):

$$F_r(r, s, t) = -\frac{Ne^2(1+\beta^2)}{\sqrt{(2\pi)^3}\epsilon_0 r\sigma_s} \left[1 - \exp(-\frac{r^2}{2\sigma^2})\right] \cdot \left[\exp(-\frac{(s+vt)^2}{2\sigma_s^2})\right]$$

with Newton's law :
$$\Delta r' = \frac{1}{mc\beta\gamma} \int_{\infty}^{\infty} F_r(r,s,t)dt$$

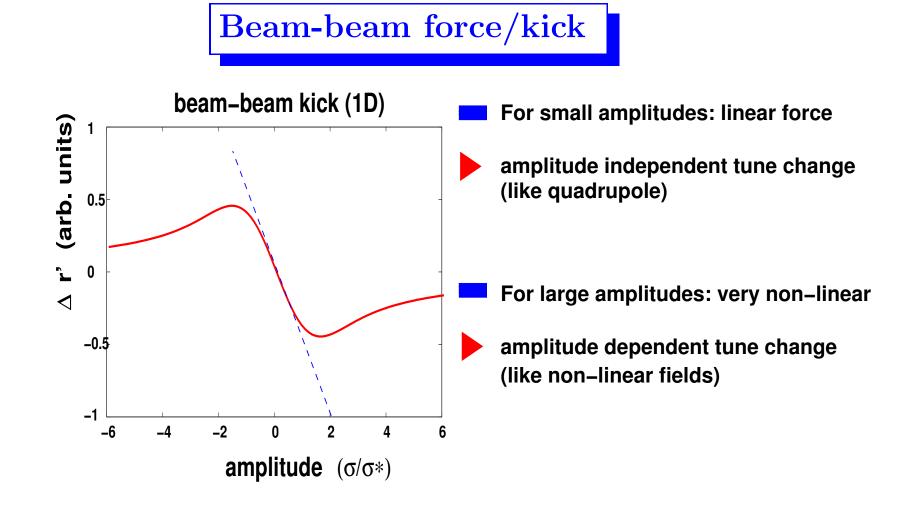
Beam-beam kick:

 \rightarrow Using the classical particle radius (implies $Z_1 = \pm Z_2$):

$$r_0 = e^2 / 4\pi\epsilon_0 mc^2$$

we have (radial kick and in Cartesian coordinates):

$$\Delta r' = -\frac{2Nr_0}{\gamma} \cdot \frac{r}{r^2} \cdot \left[1 - \exp(-\frac{r^2}{2\sigma^2})\right]$$
$$\Delta x' = -\frac{2Nr_0}{\gamma} \cdot \frac{x}{r^2} \cdot \left[1 - \exp(-\frac{r^2}{2\sigma^2})\right]$$
$$\Delta y' = -\frac{2Nr_0}{\gamma} \cdot \frac{y}{r^2} \cdot \left[1 - \exp(-\frac{r^2}{2\sigma^2})\right]$$



Can we quantify the beam-beam strength ?

- Try the slope of force (kick $\Delta r'$) at zero amplitude
 - This defines: beam-beam parameter ξ
 - For head-on interactions and round beams $(\beta^* = \beta_x^* = \beta_y^*)$ we get:

$$\boldsymbol{\xi} = \frac{\beta^*}{4\pi} \cdot \frac{\delta(\Delta r')}{\delta r} = \frac{N \cdot r_o \cdot \beta^*}{4\pi \gamma \sigma^2}$$

LEP - LHC

	$ m LEP~(e^+e^-)$	LHC (pp)
Beam sizes	160 - 200 $\mu {f m}$ \cdot 2 - 4 $\mu {f m}$	$16.6 \mu \mathbf{m} + 16.6 \mu \mathbf{m}$
Intensity N	$4.0 \cdot 10^{11}/\mathrm{bunch}$	$1.15 \cdot 10^{11}/\mathrm{bunch}$
Energy	$100 { m GeV}$	$7000 { m GeV}$
$egin{array}{ccc} eta_x^* & \cdot & eta_y^* \end{array}$	$1.25~\mathrm{m}~\cdot~0.05~\mathrm{m}$	$0.55~\mathrm{m}~\cdot~0.55~\mathrm{m}$
Crossing angle	0.0	${\bf 285}\mu{\bf rad}$
Beam-beam		
$parameter(\xi)$	$(+) \ 0.0700$	$(-) \ 0.0034$

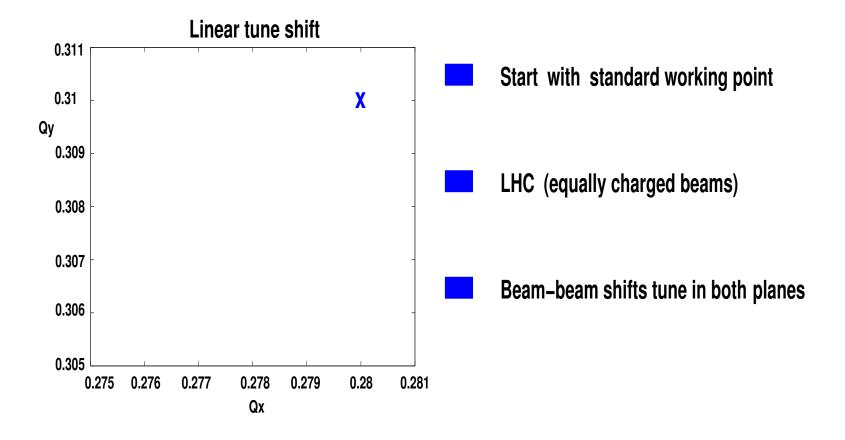
Can we quantify the beam-beam strength ?

In general for non-round beams $(\beta_x^* \neq \beta_y^*)$:

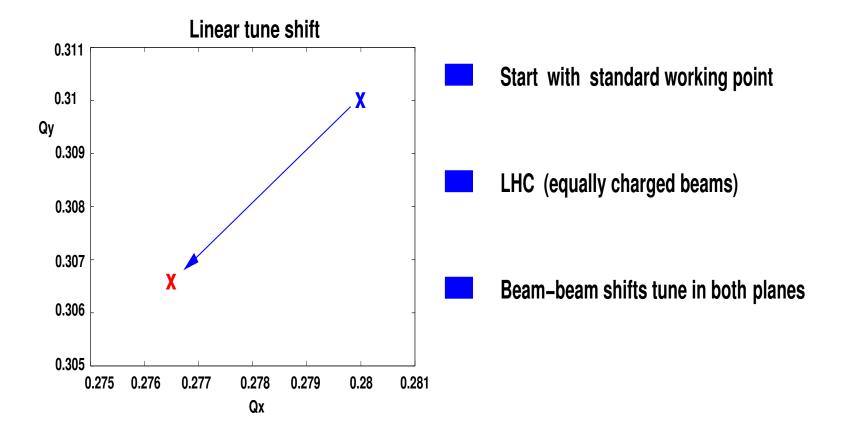
$$\xi_{x,y} = rac{N \cdot r_o \cdot eta_{x,y}^*}{2 \pi \gamma \sigma_{x,y} (\sigma_x + \sigma_y)}$$

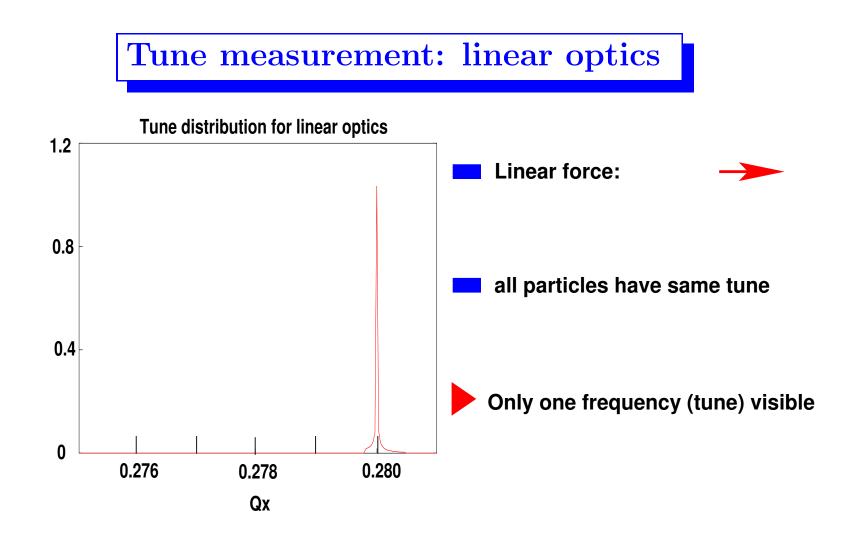
- Proportional to (linear) tune shift $\Delta \mathbf{Q}_{bb}$ from beam-beam interaction: $\Delta Q_{bb} \propto \pm \xi$
 - Good measure for strength of beam-beam interaction
 - BUT: does not describe
 - \rightarrow changes to optical functions
 - non-linear part of beam-beam force

Linear tune shift - two dimensions

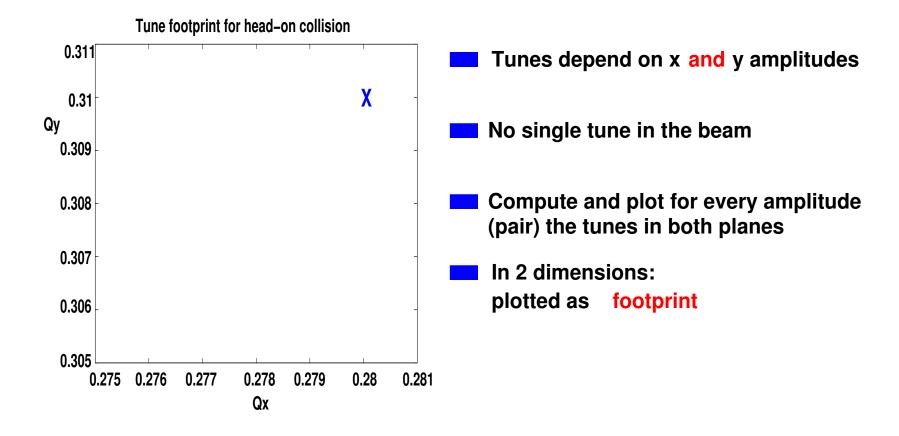


Linear tune shift - two dimensions

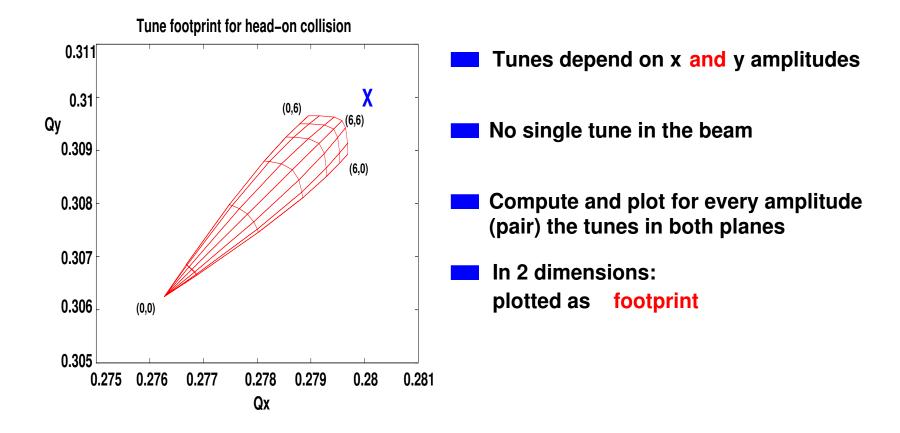




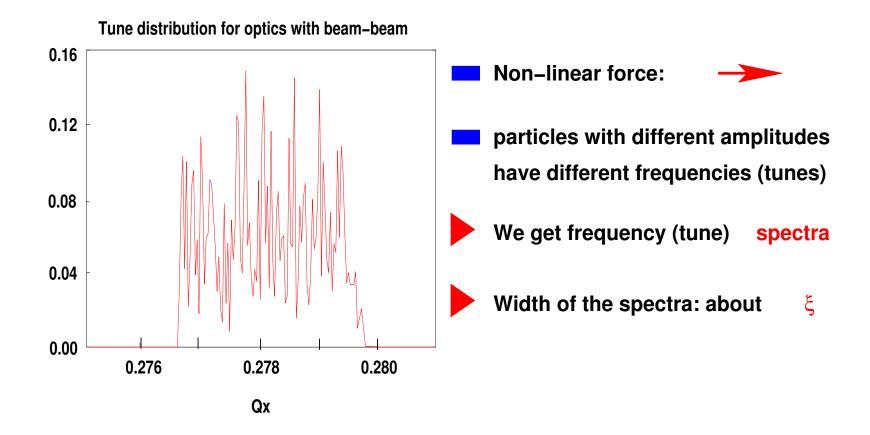
Non-linear tune shift - two dimensions



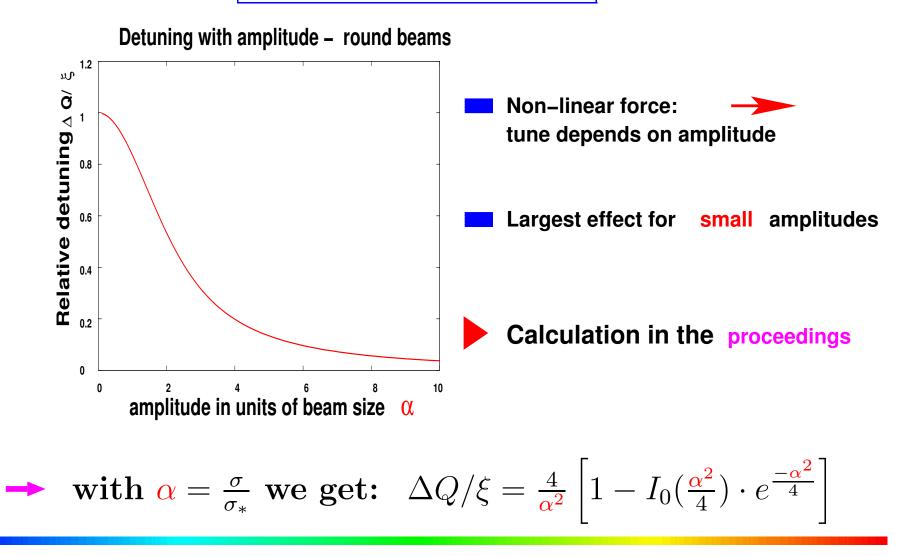
Non-linear tune shift - two dimensions



Tune measurement: with beam-beam







Weak-strong and strong-strong

- Both beams are very strong (strong-strong):
 - \rightarrow Both beam are affected and change due to beam-beam interaction
 - \rightarrow Examples: LHC, LEP, RHIC, ...
- One beam much stronger (weak-strong):
 - → Only the weak beam is affected and changed due to beam-beam interaction
 - \rightarrow Examples: SPS collider, Tevatron, ...

Incoherent effects

(single particle effects)

Single particle dynamics: treat as a particle through a static electromagnetic lens
Basically non-linear dynamics
All single particle effects observed:

→ Unstable and/or irregular motion
→ beam blow up or bad lifetime

Observations hadrons

- Non-linear motion can become chaotic

 → reduction of "dynamic aperture"
 → particle loss and bad lifetime

 Strong effects in the presence of noise or ripple
 Very bad: unequal beam sizes (studied at SPS, HERA)
 - **Evaluation is done by simulation**

Observations leptons

Remember:

$$\implies \mathcal{L} = \frac{N_1 N_2 f n_B}{4\pi \sigma_x \sigma_y}$$

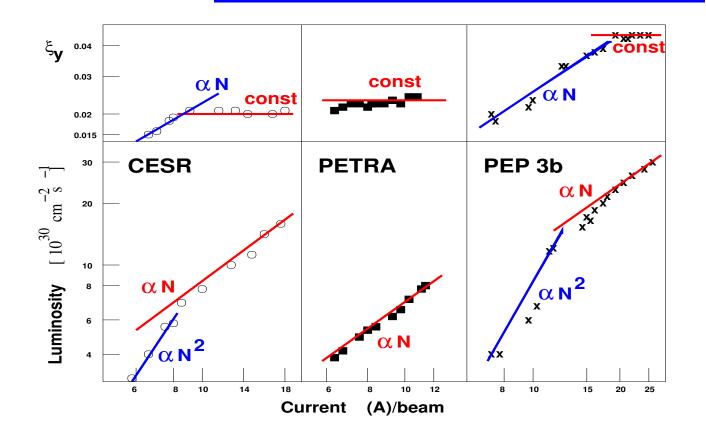
Luminosity should increase $\propto N_1 N_2$

 \rightarrow for: $N_1 = N_2 = N \rightarrow \infty N^2$

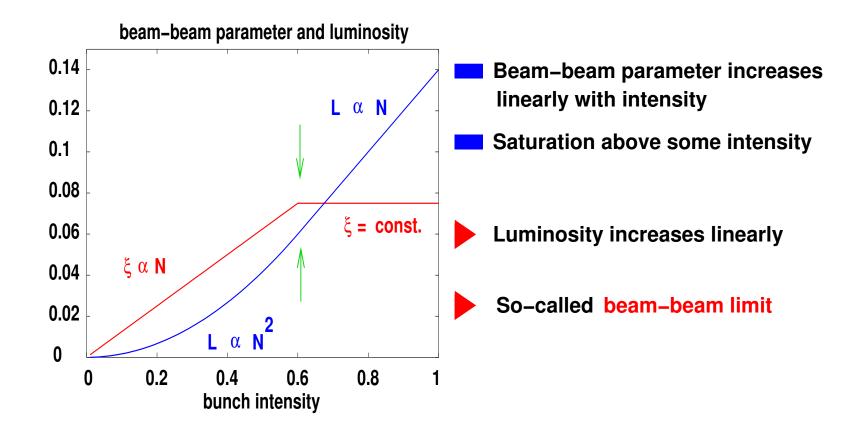
Beam-beam parameter should increase $\propto N$

But:

Examples: beam-beam limit



Beam-beam limit (schematic)



What is happening ?

we have
$$\xi_y = \frac{Nr_0\beta_y}{2\pi\gamma\sigma_y(\sigma_x + \sigma_y)} \stackrel{(\sigma_x \gg \sigma_y)}{\approx} \frac{r_0\beta_y}{2\pi\gamma(\sigma_x)} \cdot \frac{N}{\sigma_y}$$
$$\frac{N^2 f n_D}{N} = \frac{N f n_D}{N} \frac{N}{\sigma_y}$$

and
$$\mathcal{L} = \frac{N^2 f n_B}{4\pi \sigma_x \sigma_y} = \frac{N f n_B}{4\pi \sigma_x} \cdot \frac{N}{\sigma_y}$$

Above beam-beam limit: σ_y increases when N increases to keep ξ constant \rightarrow equilibrium emittance !

Therefore: $\mathcal{L} \propto N$ and $\xi \approx \text{constant}$

- ξ_{limit} is NOT a universal constant !
 - Difficult to predict

The next problem

Remember:

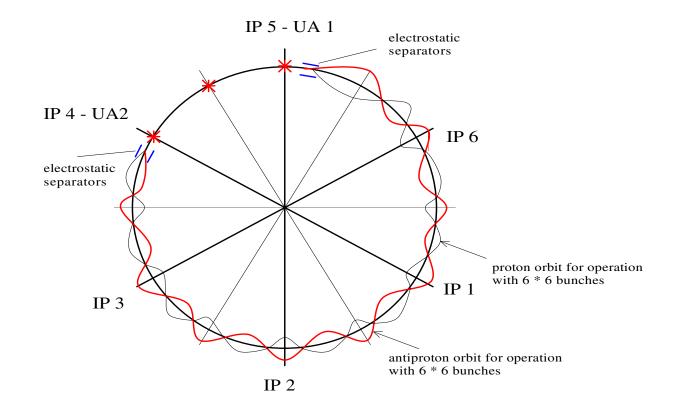
$$\implies \mathcal{L} = \frac{N_1 N_2 f \cdot n_B}{4\pi \sigma_x \sigma_y}$$

How to collide many bunches (for high L) ??
Must avoid unwanted collisions !!
Separation of the beams:
Pretzel scheme (SPS,LEP,Tevatron)
Bunch trains (LEP,PEP)
Crossing angle (LHC)

Separation: SPS

- $\Rightarrow Few equidistant bunches$ (6 against 6)
- Beams travel in same beam pipe (12 collision points !)
 - Two experimental areas
- Need global separation
- Horizontal pretzel around most of the circumference

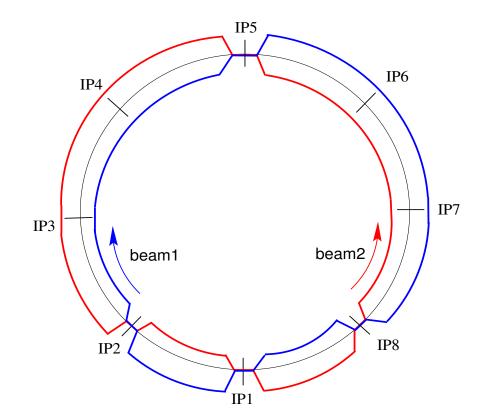
Separation: SPS

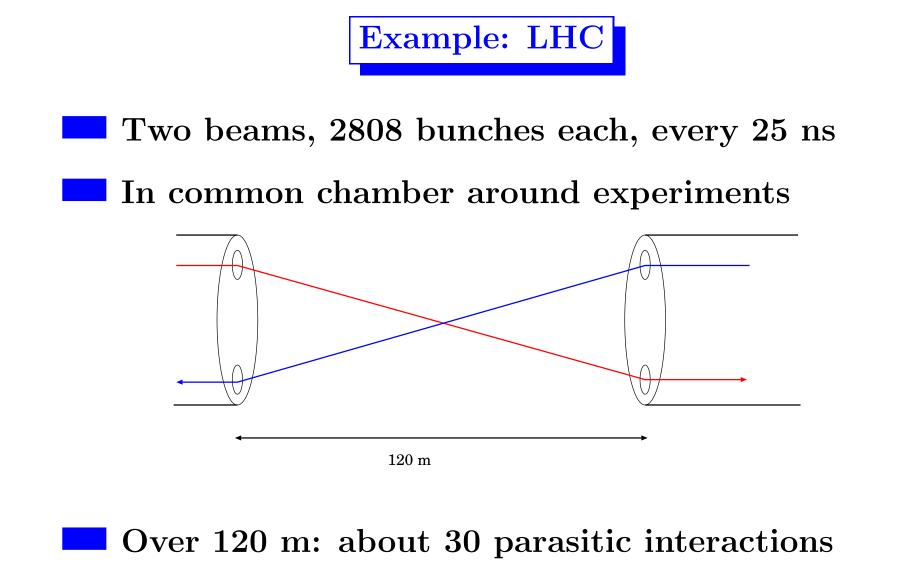


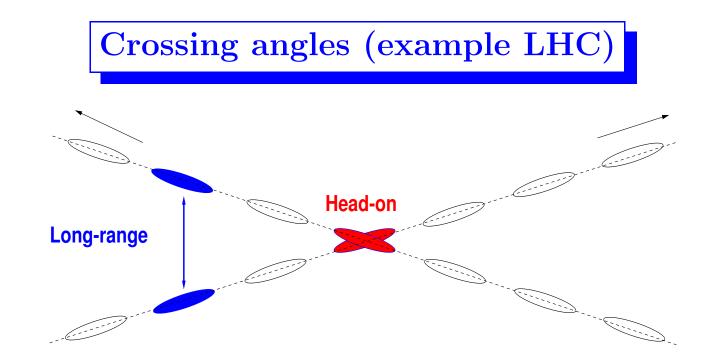
Separation: LHC

- $\bullet \Rightarrow$ Many equidistant bunches
- **Two beams in separate beam pipes except:**
 - Four experimental areas
 - Need local separation
- **Two horizontal and two vertical crossing angles**









Particles experience distant (weak) forces

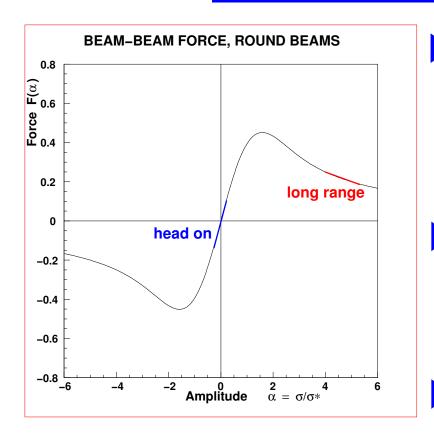
Separation typically 6 - 12 σ

→ We get so-called long range interactions

What is special about them ?

- Break symmetry between planes, also odd resonances
 - Mostly affect particles at large amplitudes
 - Cause effects on closed orbit
 - **PACMAN** effects
 - Tune shift has opposite sign in plane of separation

Why opposite tuneshift ???

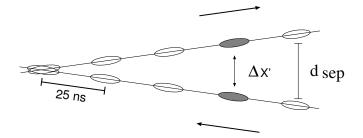


Local slope of force has opposite sign for large separation

Opposite sign for focusing

Used for partial compensation

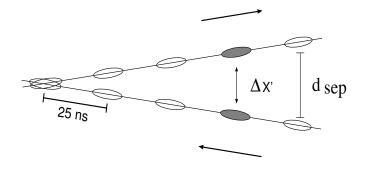
Long range interactions (LHC)



\rightarrow For horizontal separation d:

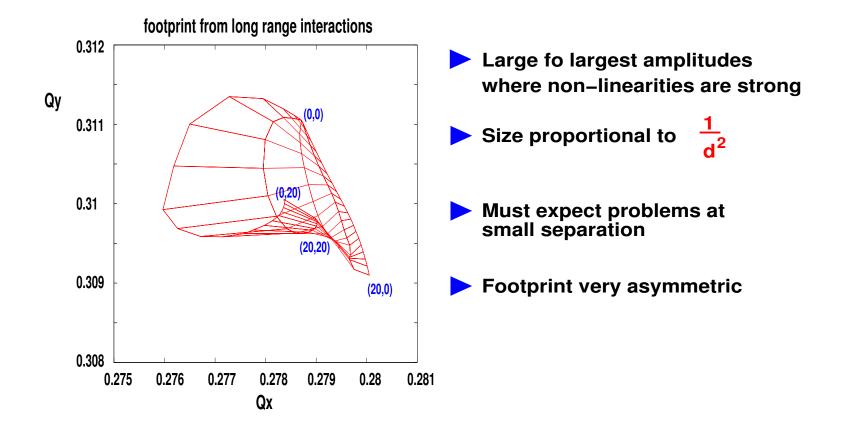
$$\Delta x'(x+d, y, r) = -\frac{2Nr_0}{\gamma} \cdot \frac{(x+d)}{r^2} \left[1 - \exp(-\frac{r^2}{2\sigma^2}) \right]$$
(with: $r^2 = (x+d)^2 + y^2$)

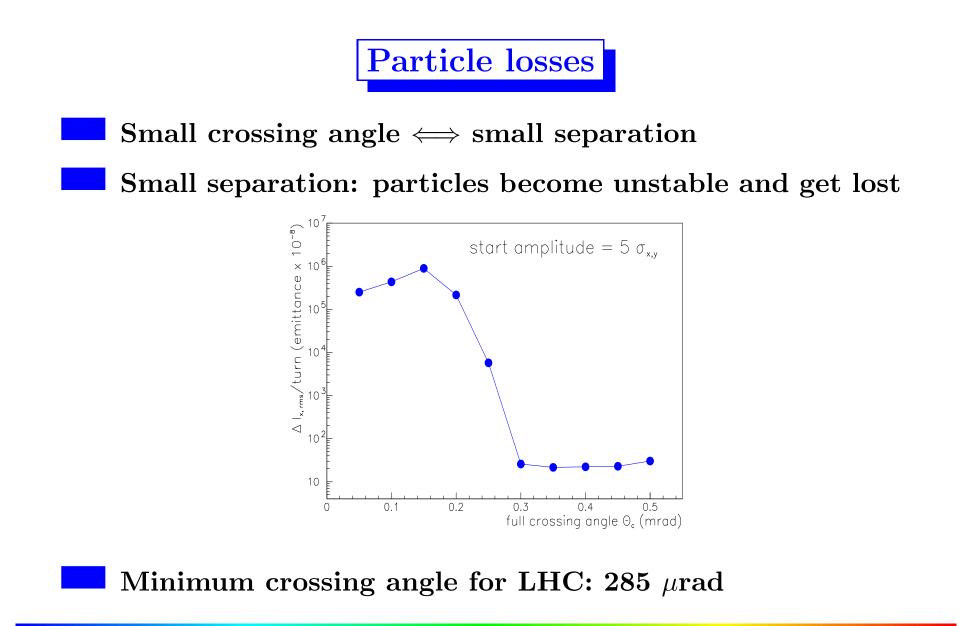
Long range interactions (LHC)



Number of long range interactions depends on spacing and length of common part
 In LHC 15 collisions on each side, 120 in total !
 Effects depend on separation: ΔQ ∝ - ^N/_{d²} (for large enough d !) footprints ??

Footprints





Closed orbit effects

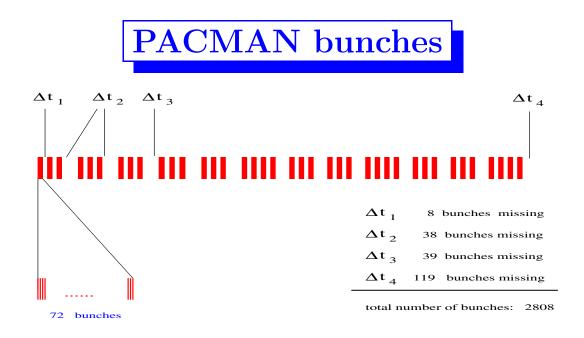
$$\Delta x'(\mathbf{x} + \mathbf{d}, y, r) = -\frac{2Nr_0}{\gamma} \cdot \frac{(\mathbf{x} + \mathbf{d})}{r^2} \left[1 - \exp(-\frac{r^2}{2\sigma^2}) \right]$$

For well separated beams $(d \gg \sigma)$ the force (kick) has an amplitude independent contribution: \rightarrow orbit kick

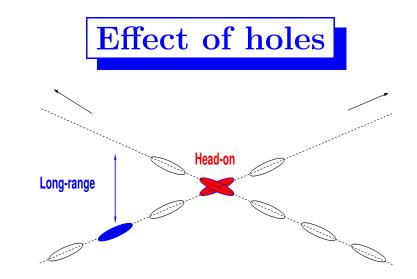
$$\Delta x' = \frac{const.}{d} \cdot [1 - \frac{x}{d} + O\left(\frac{x^2}{d^2}\right) + \dots$$

Closed orbit effects

- → Beam-beam kick from long range interactions changes the orbit
 - Has been observed in LEP with bunch trains
 - Self-consistent calculation necessary
 - Effects can add up and become important
- → Orbit can be corrected, but:



- LHC bunch filling not continuous: holes for injection, extraction, dump ..
 - **2808** of 3564 possible bunches \rightarrow 1756 "holes"
 - "Holes" meet "holes" at the interaction point
 - But not always ...

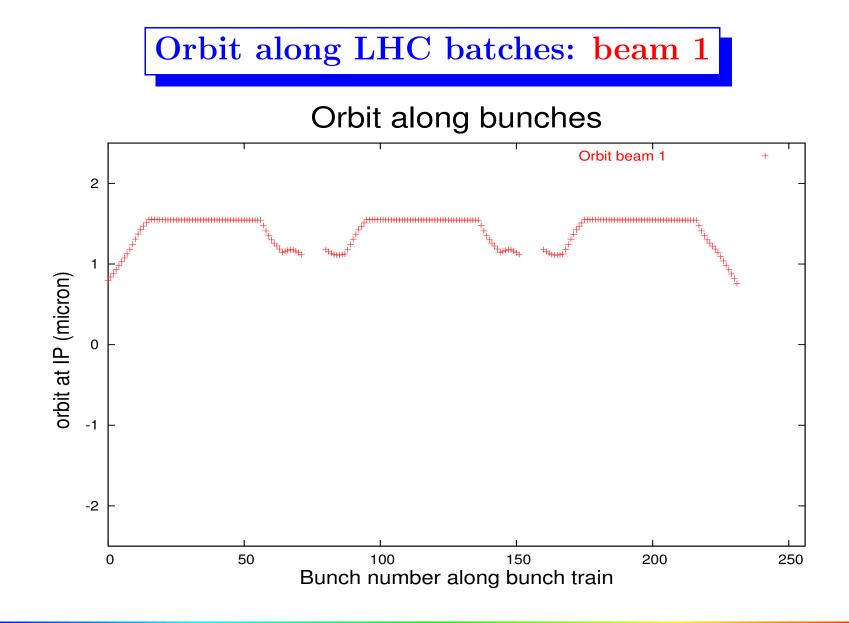


- A bunch can meet a hole (at beginning and end of bunch train)
 - **Results in left-right asymmetry**
- Example LHC: between 120 (max) and 40 (min) long range collisions for different bunches

PACMAN bunches

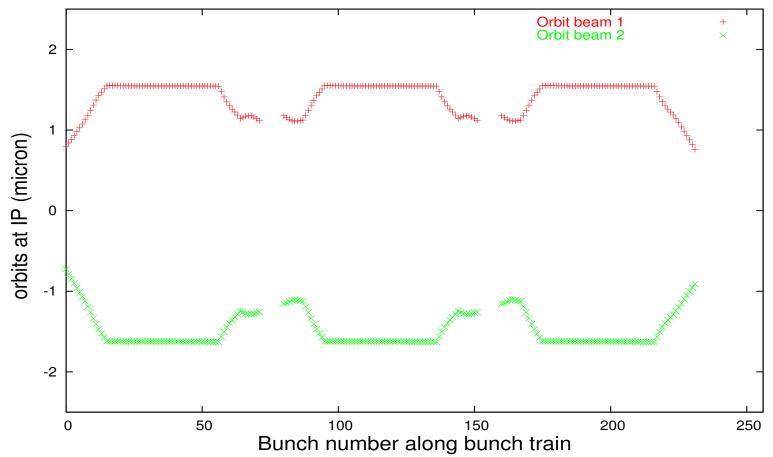
When a bunch meets a "hole":

- → Miss some long range interactions, PACMAN bunches
- → They see fewer unwanted interactions in total
- → Different integrated beam-beam effect
- In general: when different bunches have different beam-beam effects
- **Example:** orbit and tune effects



Orbit along LHC batches: beam 1 and beam 2

Orbits along bunches



Tune along LHC batches

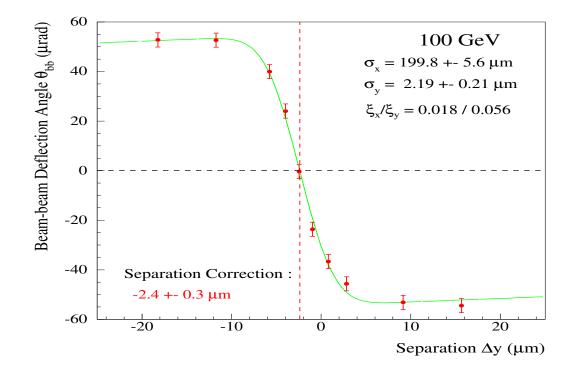
Fune along bunches 0.314 0.312 0.314 0.312 0.314 0.314 0.312 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.314 0.316 0.316 0.316 0.308 0.308 0.306 0.3

Spread is too large for safe operation

Beam-beam deflection scan

- The orbit effect can be useful when one has only a few bunches, i.e. not PACMAN effects
 - Effect can be used to optimize luminosity
 - **Scanning two beams against each other**
- Two beams get a orbit kick, depending on distance

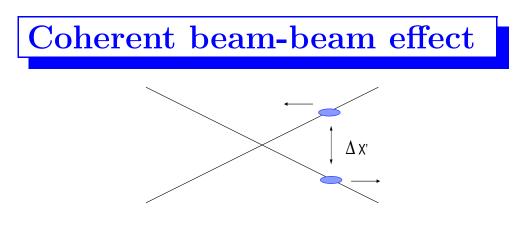
Deflection scan (LEP measurement)



(Courtesy J. Wenniger)

Deflection scan

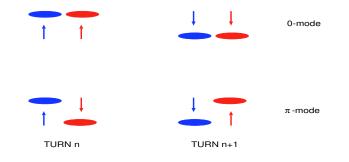
- Calculated kick from orbit follows the force function
- Allows to calculate parameters
- Allows to centre the beam
- **Standard procedure at LEP**



- Whole bunch sees a kick as an entity (coherent kick)
- The coherent kick of separated beams can excite coherent dipole oscillations
- All bunches couple together because each bunch "sees" many opposing bunches: many coherent modes possible !

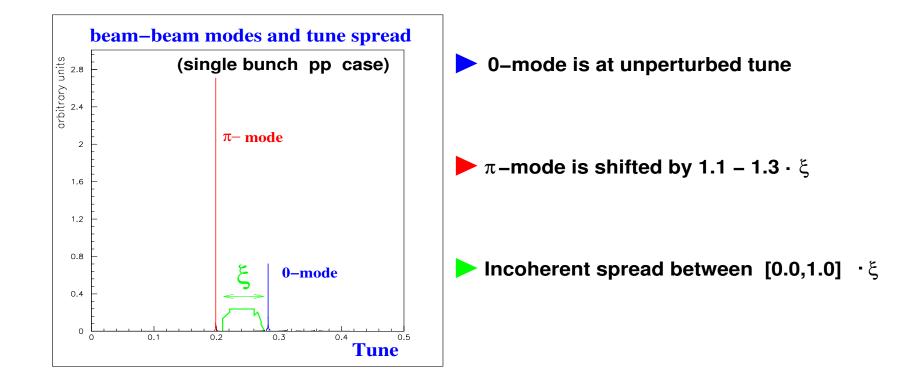


Simplest case: one bunch per beam:



- Coherent mode: two bunches are "locked" in a coherent oscillation
 - **0-mode is stable (Mode with NO tune shift)**
 - π -mode can become unstable (Mode with LARGEST tune shift)

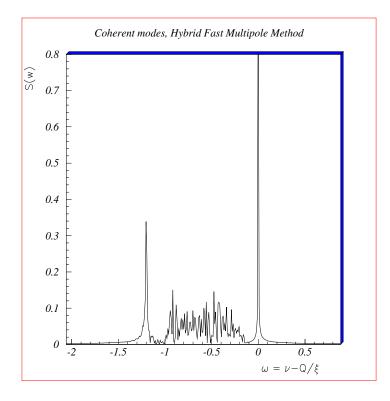
Coherent beam-beam frequencies (schematic)



Strong-strong case: π -mode shifted outside tune spread

No Landau damping possible

Simulation of coherent spectra





Full simulation of both beams required

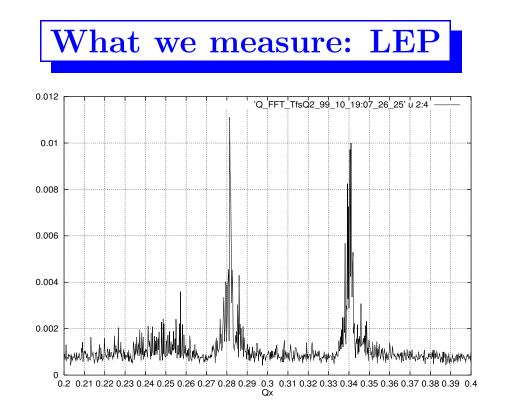




Must take into account changing fields

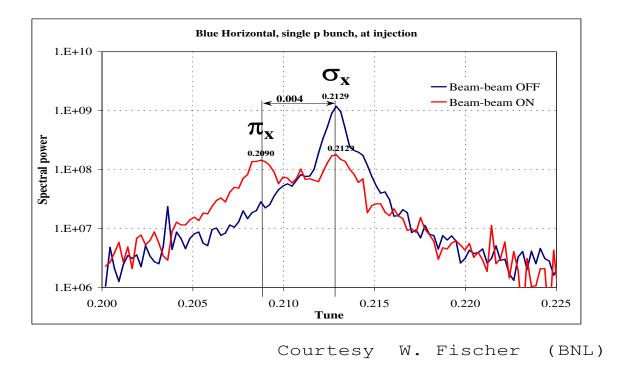


Requires computation of arbitrary fields



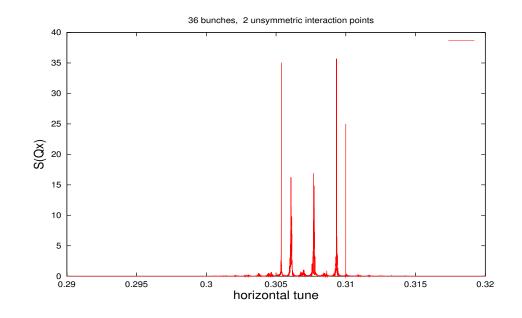
- Two modes clearly visible
- Can be distinguished by phase relation, i.e. sum and difference signals

What we measure: RHIC



Compare spectra with and without beams : two modes visible with beams

Many bunches and more interaction points

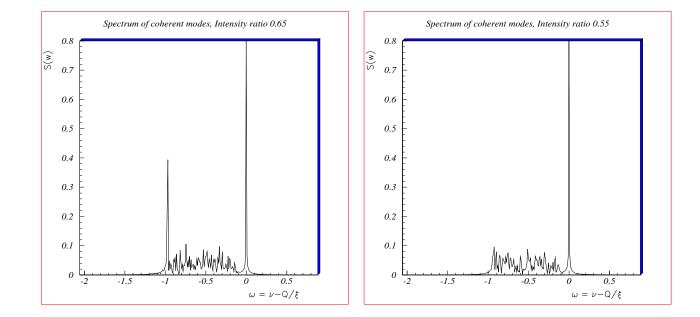


Bunches couple via the beam-beam interaction
Additional coherent modes become visible
Potentially undesirable situation

What can be done to avoid problems ?

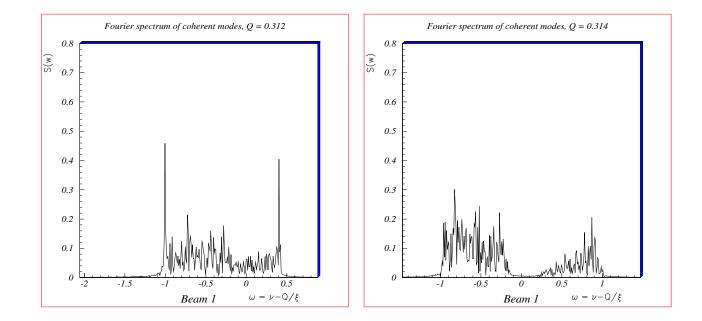
- Coherent motion requires 'organized' motion of many particles
 - **Therefore high degree of symmetry required**
- Possible countermeasure: (symmetry breaking)
 - \rightarrow Different bunch intensity
 - \rightarrow Different tunes in the two beams

Beams with different intensity



→ Bunches with different intensities cannot maintain coherent motion

Beams with different tunes



Bunches with different tunes cannot maintain coherent motion

Can we suppress beam-beam effects ?

Find 'lenses' to correct beam-beam effects

- Head on effects:
 - \rightarrow Linear "electron lens" to shift tunes
 - \rightarrow Non-linear "electron lens" to reduce spread
 - \rightarrow Tests in progress at FNAL
- Long range effects:
 - \rightarrow At very large distance: force is 1/r
 - \rightarrow Same force as a wire !

So far: mixed success with active compensation



Principle:

 \rightarrow Interchange horizontal and vertical plane each turn

Effects:

- \rightarrow Round beams (even for leptons)
- \rightarrow Some compensation effects for beam-beam interaction
- \rightarrow First test at CESR at Cornell

Not mentioned:

- Effects in linear colliders
 Asymmetric beams
 Coasting beams
 Beamstrahlung
 Synchrobetatron coupling
 Monochromatization
 - Beam-beam experiments
- ____ ... and many more

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