

Designing a synchrotron - A real life example

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- Review several aspects of beam dynamics (mostly) presented in the introductory CAS lectures, applied to the design and operation of a real synchrotron
 - Choice of energy, bending field and circumference
 - Optics design
 - Cell optics, insertions, transition energy
 - Collective effects
 - Instabilities, Space-charge
 - Electron / Positron beam dynamics
 - Equilibrium beam properties, energy loss/turn, damping time

- Our choice is the **CERN Super Proton Synchrotron (SPS)**
- From its design and operation, it has shown enormous versatility used for several purposes and serving various applications
 - High energy synchrotron serving fixed target experiments (West Area, North Area, CNGS, HIRADMAT)
 - Collider of protons and anti-protons (W and Z bosons discovery in 1983)
 - Accelerating electrons and positrons and injecting them to the Large Electron-Positron (LEP) Collider
 - Accelerating protons for the Large Hadron Collider (LHC)
 - Accelerating ions for fixed target physics and the LHC
 - Extracting protons for exciting plasma for a plasma wakefield acceleration experiment (AWAKE)

Basic parameters energy, bending field and circumference

- Consider accelerator ring for particles with energy E with N dipoles of length L or effective length l , i.e. measured on beam path

- **Bending angle** $\theta = \frac{2\pi}{N}$

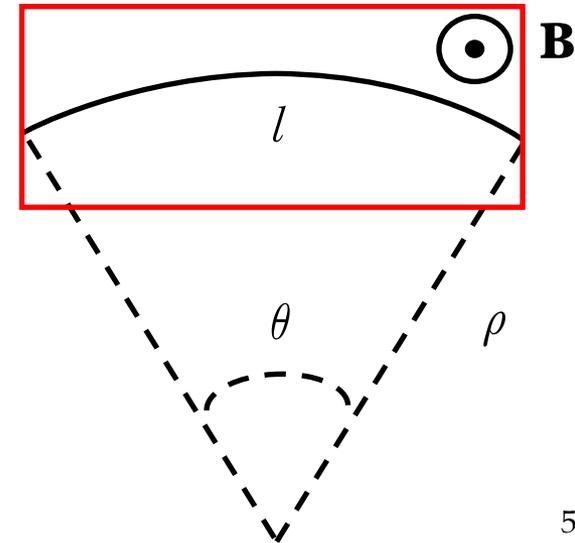
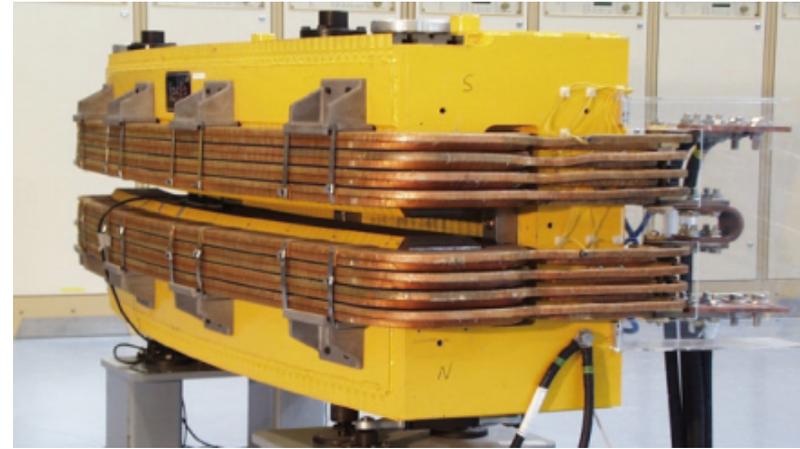
- **Bending radius** $\rho = \frac{l}{\theta}$

- The magnetic rigidity is $B\rho = \frac{\beta E}{q}$

- The **integrated dipole strength** is

$$Bl = \frac{2\pi}{N} \frac{\beta E}{q}$$

- By imposing a dipole field, the dipole length is fixed and vice versa
- The higher the field, the shorter or smaller number of dipoles can be used



- The **filling factor**, is defined as the ratio of the total length of the bending path, with respect to the circumference

$$k_f = \frac{Nl}{C}$$

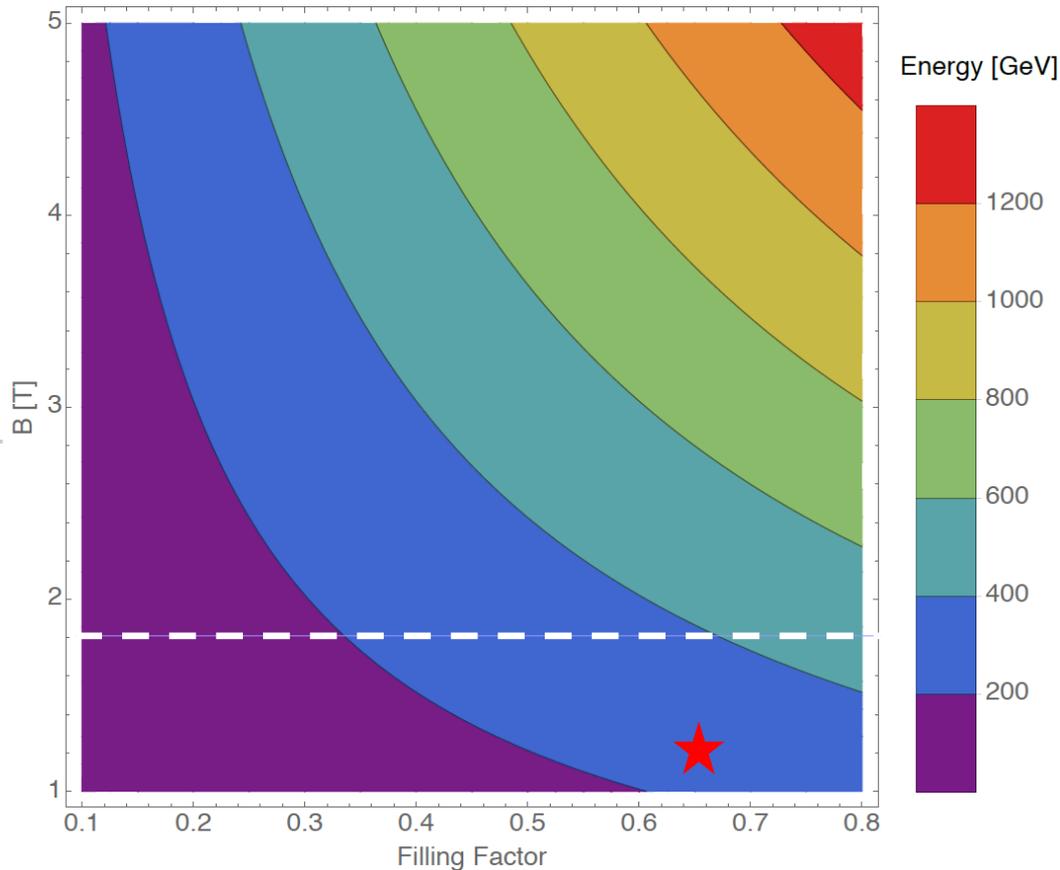
- The ring circumference becomes

$$C = \frac{2\pi}{k_f B} \frac{\beta E}{q}$$

- The ring circumference (cost) is driven by the bending field choice (technology), the energy reach (physics case, applications) and the design of the lattice cells (optics)

- The maximum possible circumference between the CERN I (Meyrin) and CERN II (Prevessin) site was

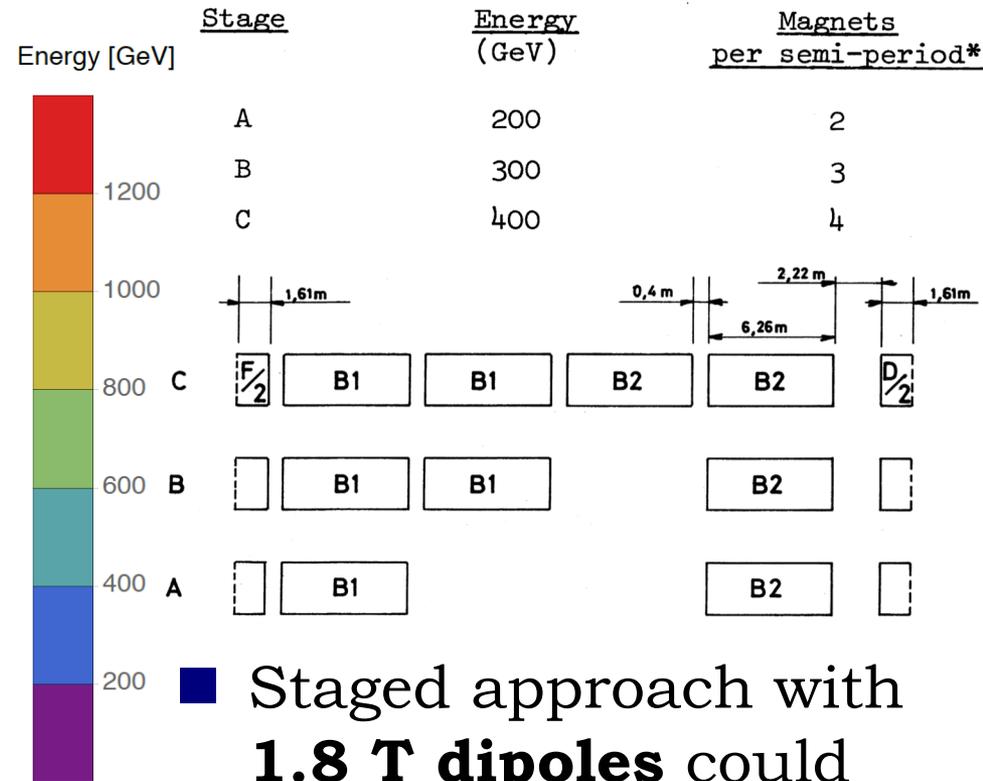
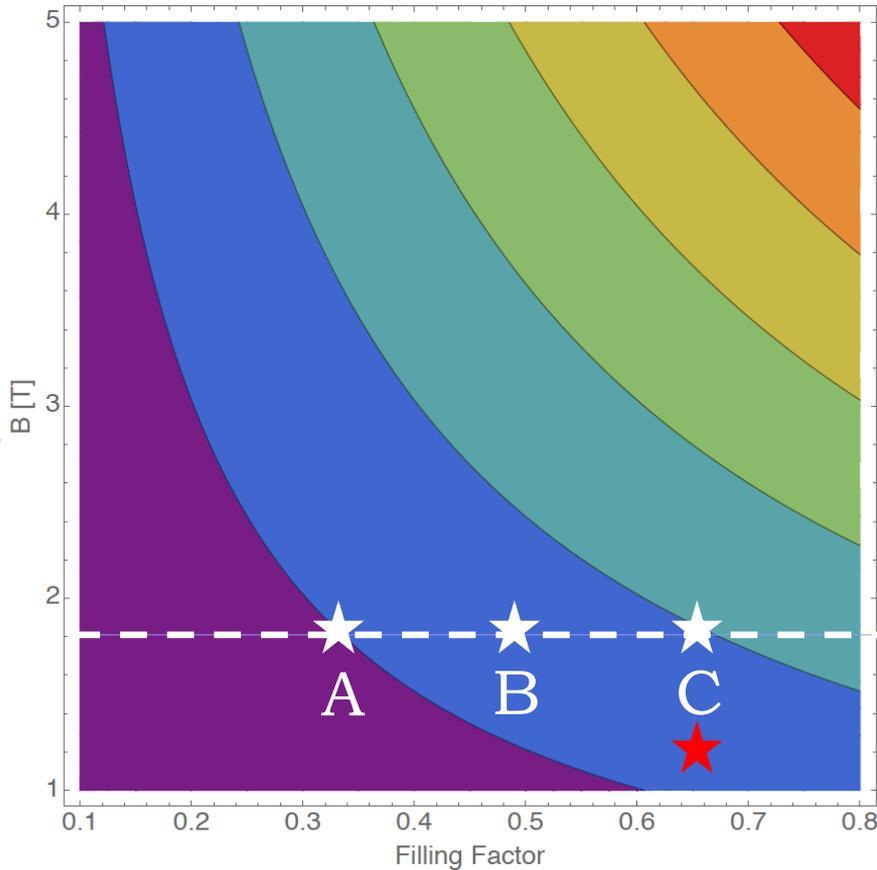
$$C_{SPS} = 11C_{PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$$



- Combined function magnets with **1.2 T** field (PS-like) would give an energy of no more than ~260 GeV for a highly packed lattice

- The maximum possible circumference between the CERN I (Meyrin) and CERN II (Prevessin) site was

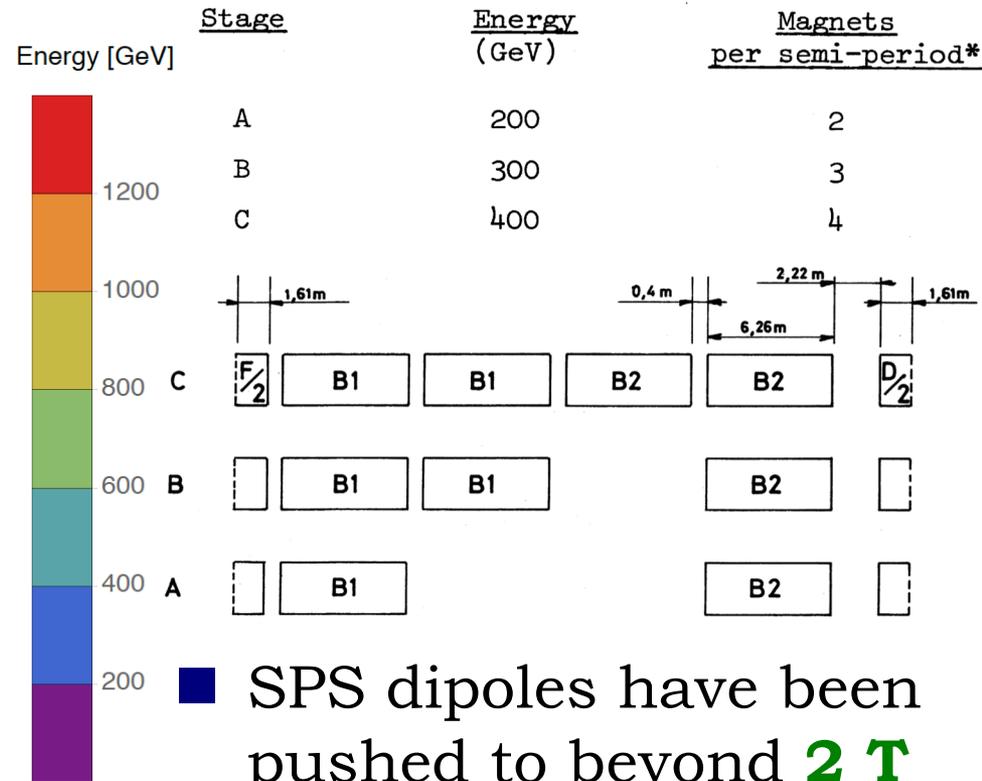
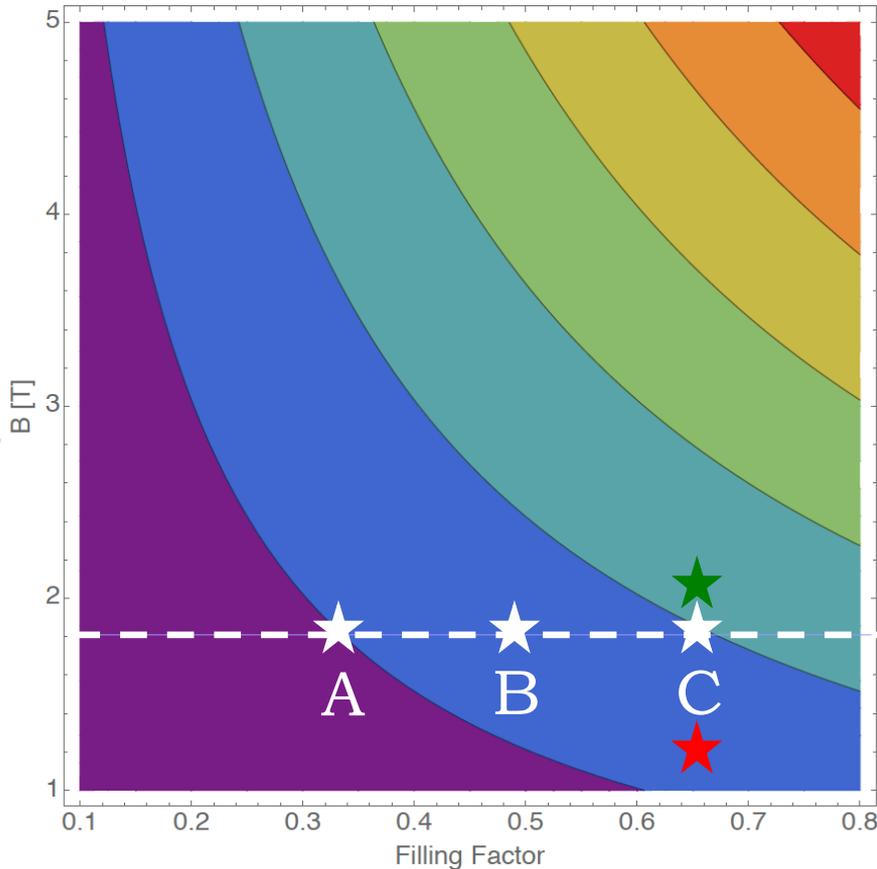
$$C_{SPS} = 11C_{PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$$



- Staged approach with **1.8 T dipoles** could reach from 200 to 400 GeV

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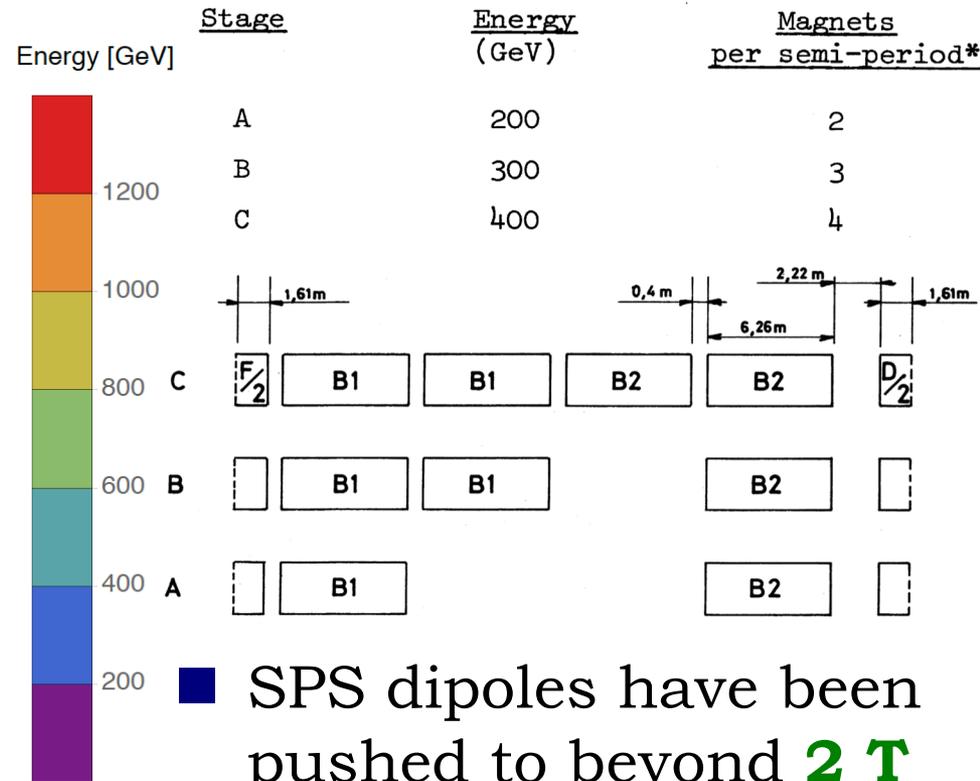
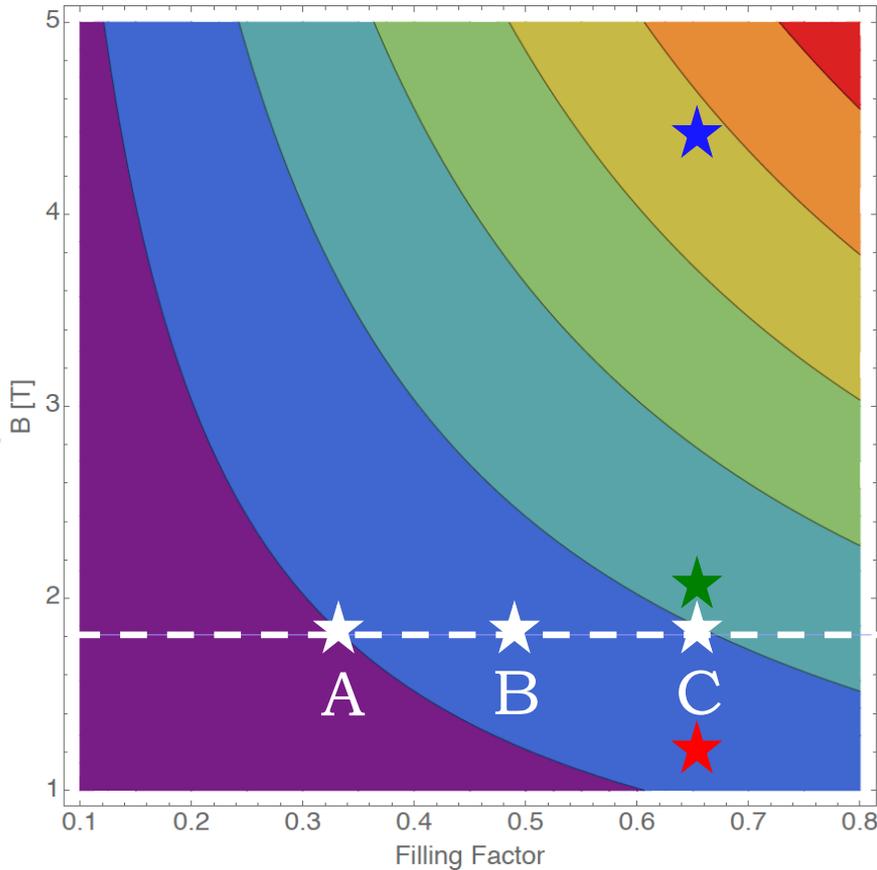
$$C_{SPS} = 11C_{PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$$



- SPS dipoles have been pushed to beyond **2 T** to reach from **450 GeV**

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$$C_{SPS} = 11C_{PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$$



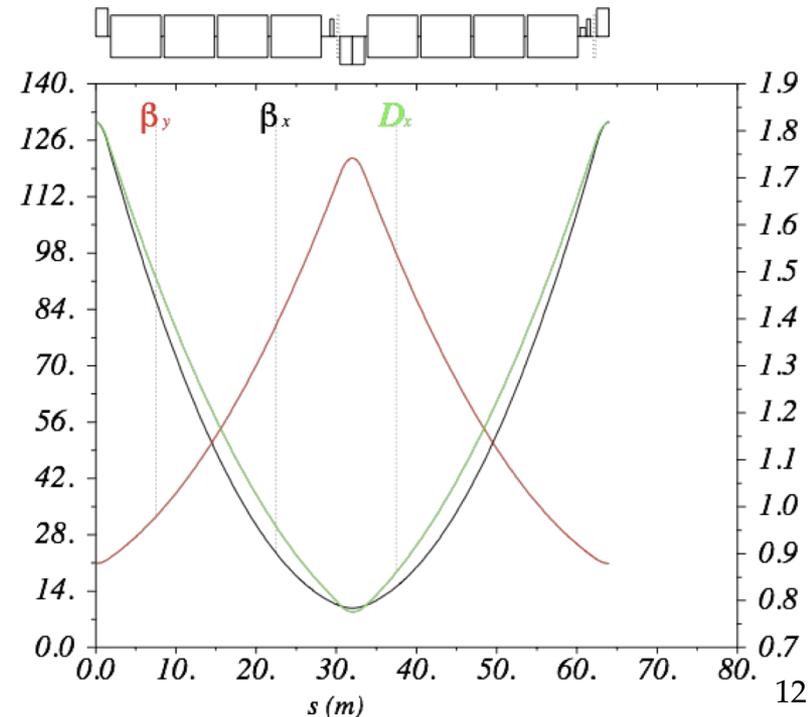
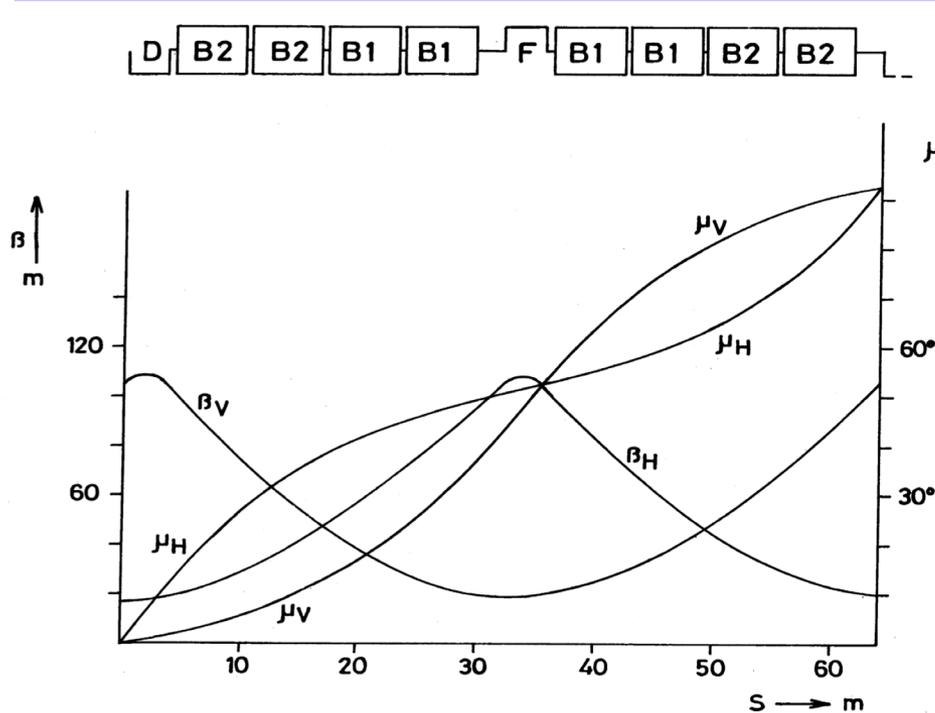
- SPS dipoles have been pushed to beyond **2 T** to reach from **450 GeV**

- Super-conducting option could raise the energy to **1 TeV** ¹⁰

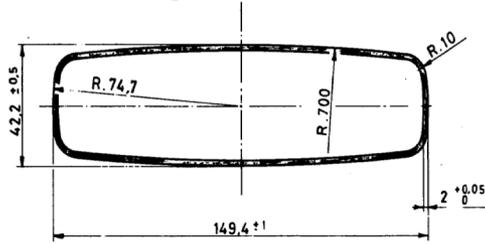
Optics design

- FODO cell of around 65 m long with phase advances of $\pi/2$
- Beta function maxima slightly above 100 m

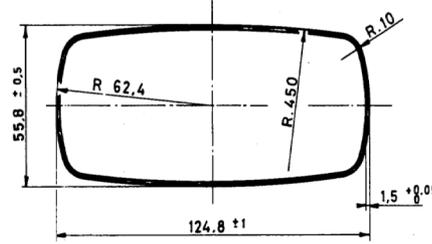
The 300 GeV Program, CERN/1050, 14/01/1972



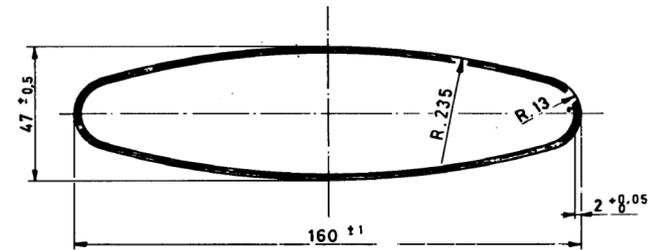
- Magnet apertures follow beta function and dispersion evolution



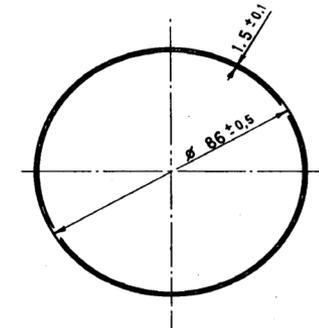
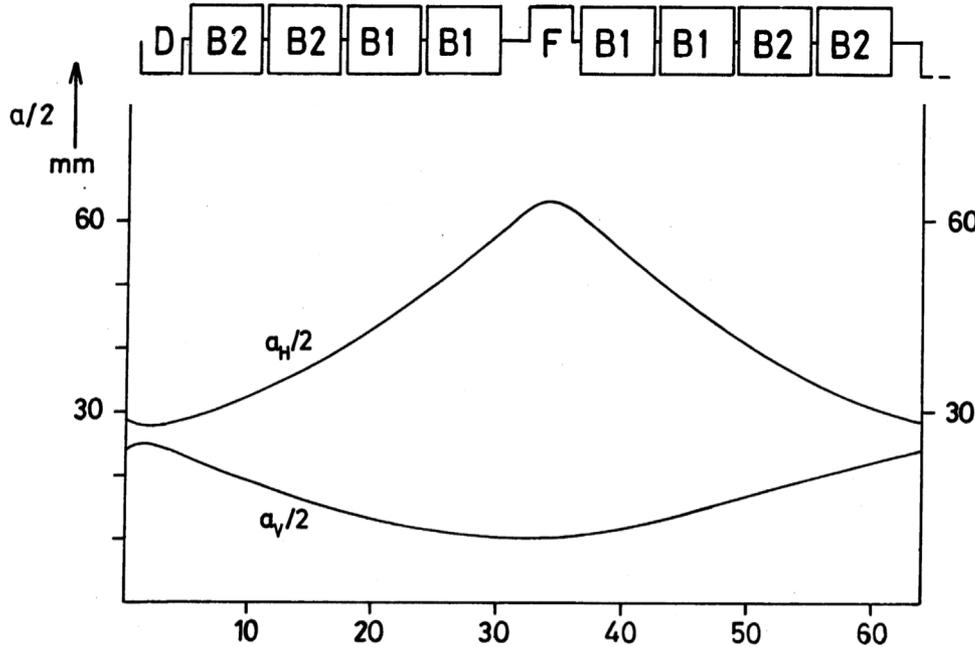
Dipole B1



Dipole B2



Quadrupole F



Quadrupole D

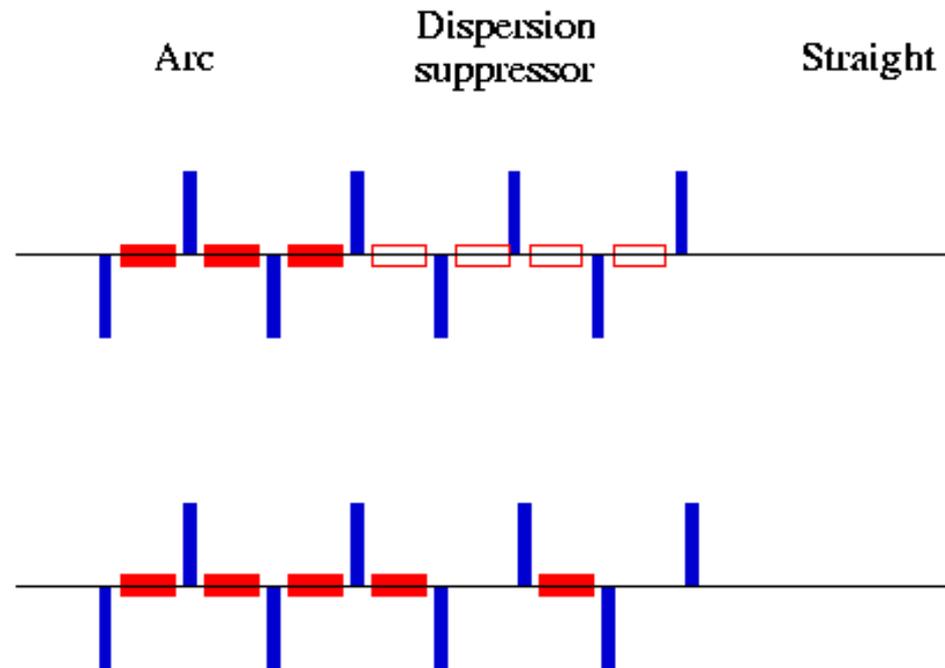
- Dispersion has to be eliminated in special areas like injection, extraction or interaction points (orbit independent to momentum spread)
- Use **dispersion suppressors**
- Two methods for suppressing dispersion
 - Eliminate two dipoles in a FODO cell (missing dipole)

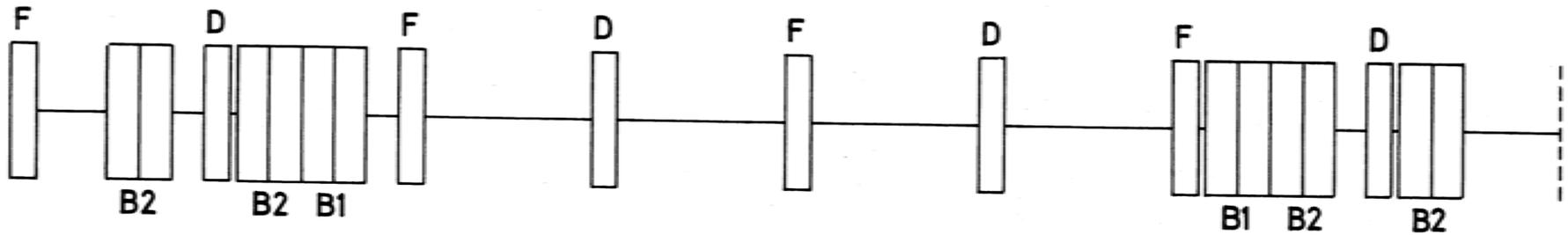
- Set last dipoles with different bending angles

$$\theta_1 = \theta \left(1 - \frac{1}{4 \sin^2 \mu_{\text{HFODO}}} \right)$$

$$\theta_2 = \frac{\theta}{4 \sin^2 \mu_{\text{HFODO}}}$$

- For equal bending angle dipoles, the FODO phase advance should be equal to $\pi/2$

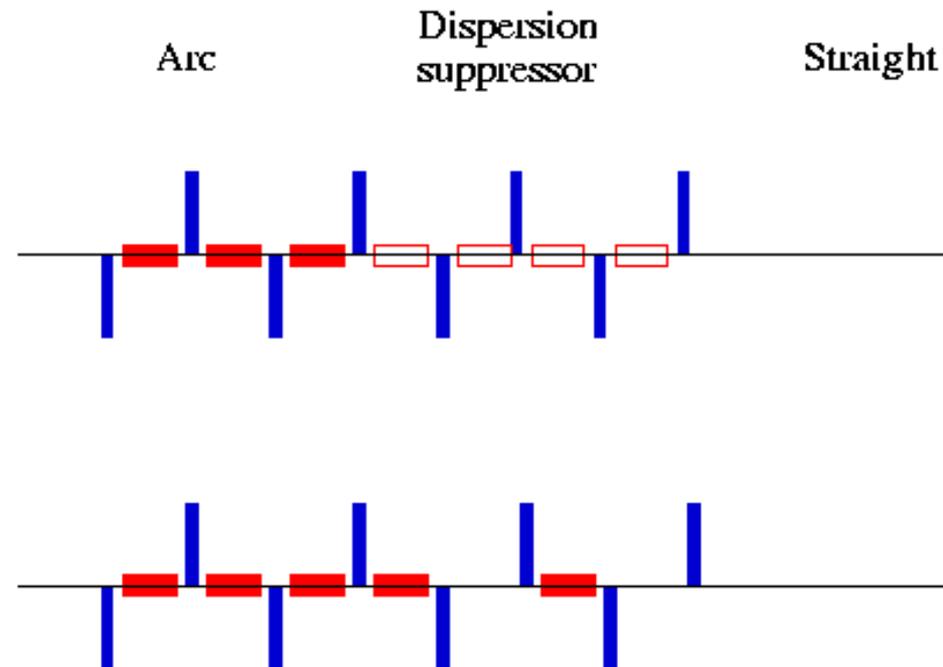




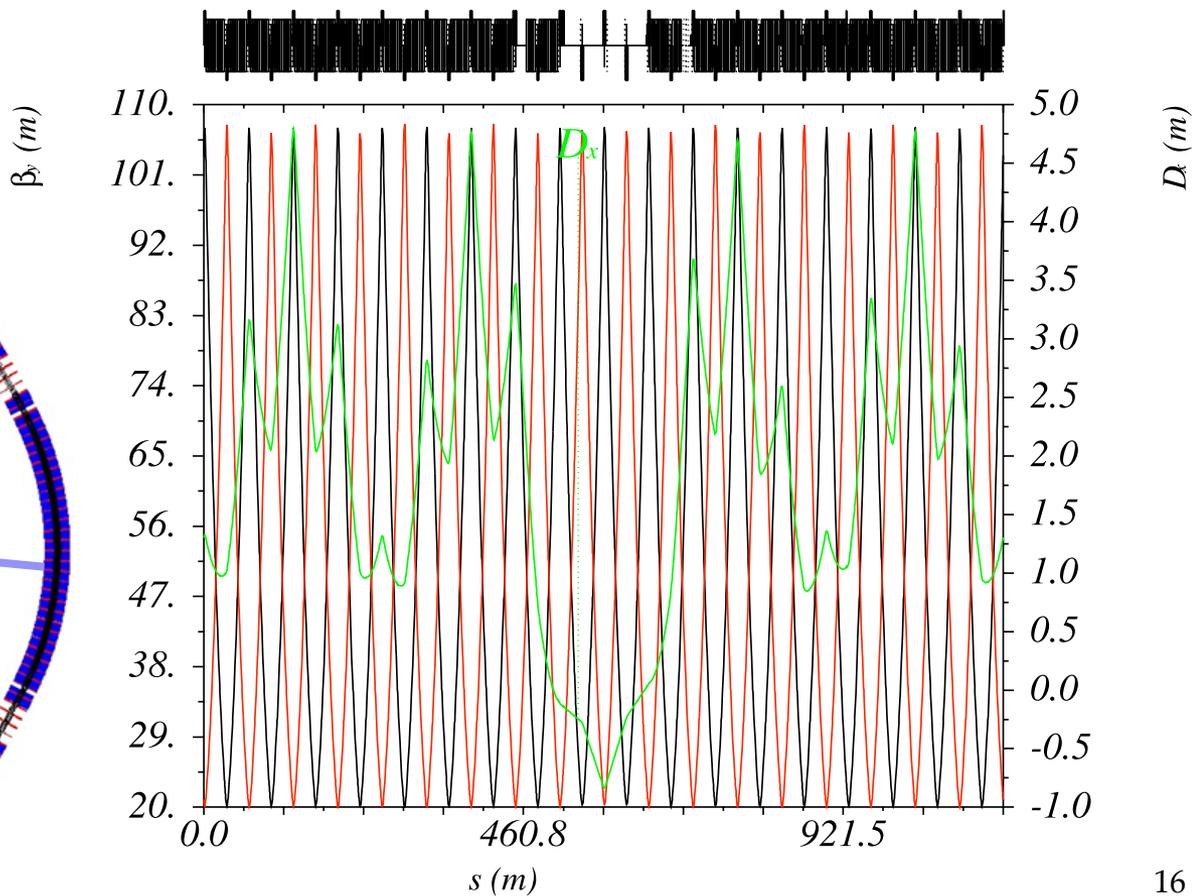
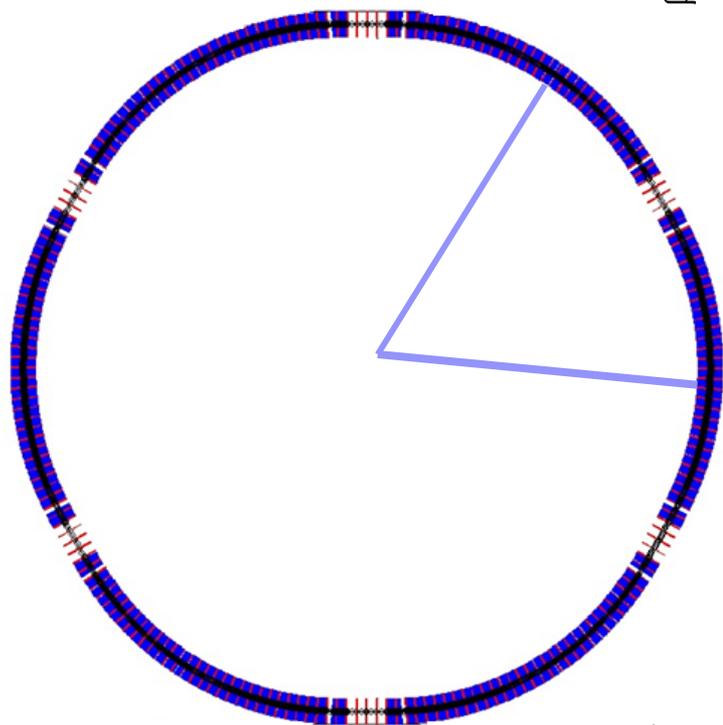
- In the SPS, all dipoles are powered in series, i.e. dispersion suppressor cells looks like a missing dipole, but they are not!

- Dispersion suppression is achieved by tuning the phase advance of the arc, to a multiple of 2π

- Dispersion oscillates through the arc and vanishes at the edges

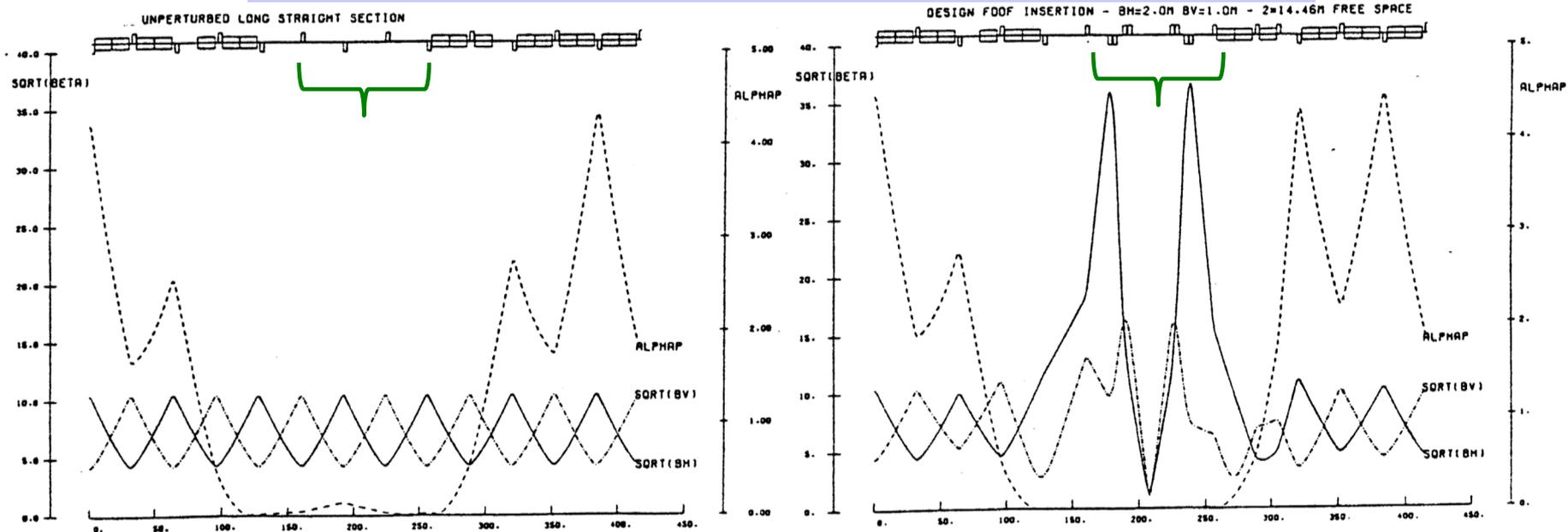


- Ring is composed by 6 identical sectors (“sextants”) with 16 cells in the arc and 2 cells in the straight section
- The cell phase advance of $\pi/2$ brings the tunes to 26-27 (**Q26**)



- Replace two straight section quadrupoles with **2 doublets** (4 quadrupoles)
- Equip adjacent left/right quadrupoles with individual bipolar power convertors
- Achieved low β^* of 1.3/0.65 m

P. Faugeras et al., CERN-SPS-80/11, CERN-SPS-83/29



Magnet system

- 744 dipoles (MBAs and MBBs) with 6.26 m length and different gaps

Number of magnets	744
Year of 1 st operation	1976
Maximum field on beam axis [T]	2.02
Physical vertical aperture [mm] MBA/MBB	38.5/51.5
Yoke assembly [Solid,Laminated,Welded,Glued]	L,W
Coil technology [Copper,Aluminium,Glass-epoxy,Mica,Other]	C,G
Maximum voltage to ground [V] (worst case 2 spare converters)	4150
Operation	Cycled
Maximum cooling water velocity [m/s]	9
Operational temperature [C°]	40



D. Tommasini CERN/TE-Note-2010-003

- Maximum field of **2.02 T**, for reaching 450 GeV
- High mechanical stress on coils

- 216 quadrupoles (102 QF, 100 QD, 6 QFA and 8 QDA)
- Maximum gradient of 22 T/m, corresponding to a pole-tip field of around **1 T**
- Normal operation necessitates almost the full gradient @ 450 GeV

D. Tommasini CERN/TE-Note-2010-003

Number of magnets	216
Year of 1 st operation	1976
Maximum gradient [T/m]	22
Physical vertical aperture [mm]	88
Yoke assembly [Solid,Laminated,Welded,Glued]	L,W
Coil technology [Copper,Aluminium,Glass-epoxy,Mica,Other]	C,G
Maximum voltage to ground [V]	3450
Operation	Cycled
Maximum cooling water velocity [m/s]	3.6
Operational temperature [C°]	40

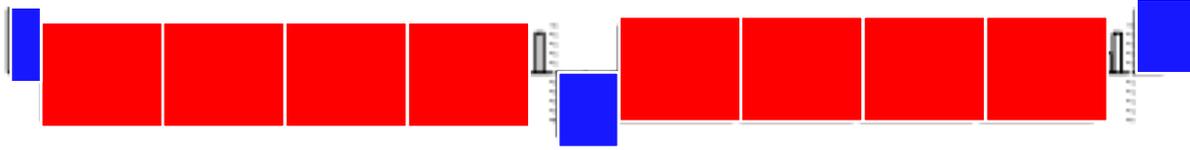


M. Giesch, CERN/SPS/80-3/AMS, 1980

MAIN PARAMETERS OF SEXTUPOL		LSFN	LSDN
Basic	: Nominal rms current [A]	350	350
	Peak Current [A]	500	450
* Strength at peak current			
	1) Sextup. $\int a_3 dl$ ($a_3 = B/r_2 = B''/2$) [T/m]	85.8	176.6
	n^2		
	* Magnetic length [m]	0.435	0.426
	Aperture, radius of inscr.circle [mm]	60.7	44.0
Core	: Length [m]	0.4	0.4



- 54 “focusing” and 54 “defocusing” 0.4 m long sextupoles in two (three for F) families (24 and 30), with different apertures
- Maximum pole-tip field of around 0.8 T
- Around 80% and 60% in operational conditions



Transition energy and slippage factor

- Transition energy (or momentum compaction factor) is defined as

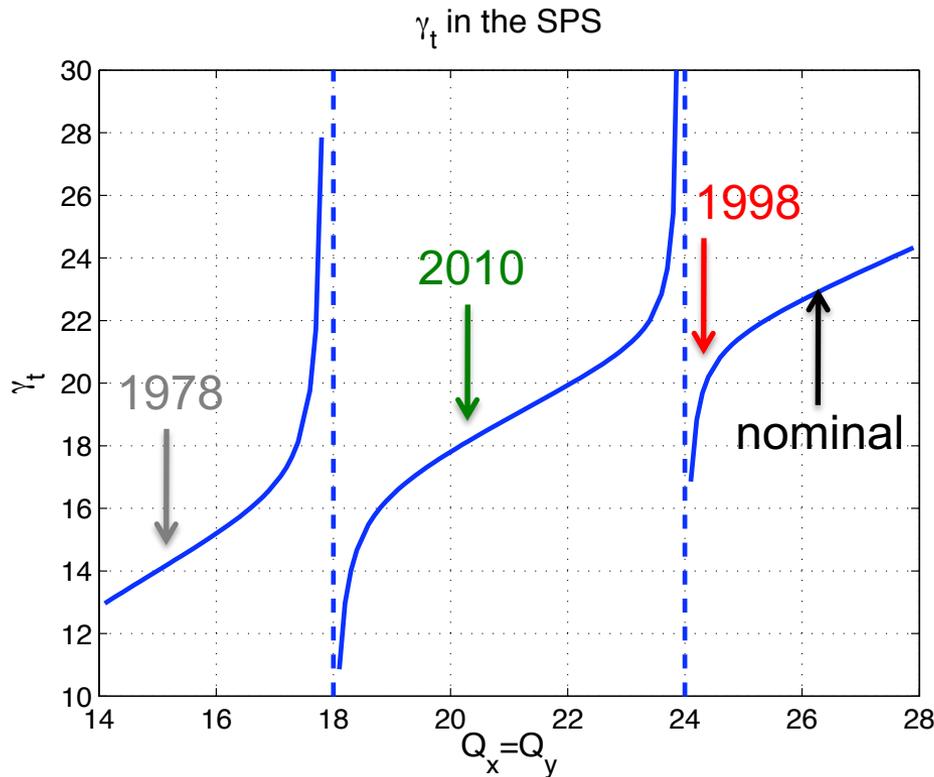
$$\frac{1}{\gamma_t^2} = \alpha_p = \frac{1}{C} \oint \frac{D(s)}{\rho(s)} ds$$

- The higher the dispersion oscillation in the bends, the lower the transition energy



Quadrupoles

- Note also that, for FODO cells (SPS lattice), $\gamma_t \approx Q_x$, meaning that lowering the transition energy implies **lowering** the **horizontal tune**
- High intensity beams can be injected in the SPS above transition avoiding losses and operational complexity of transition jump scheme

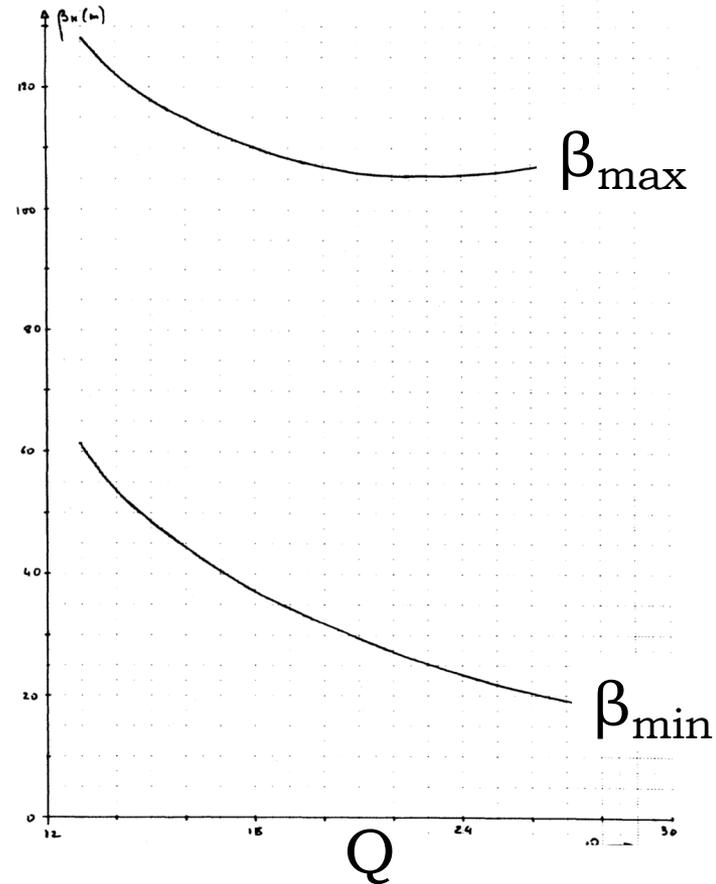
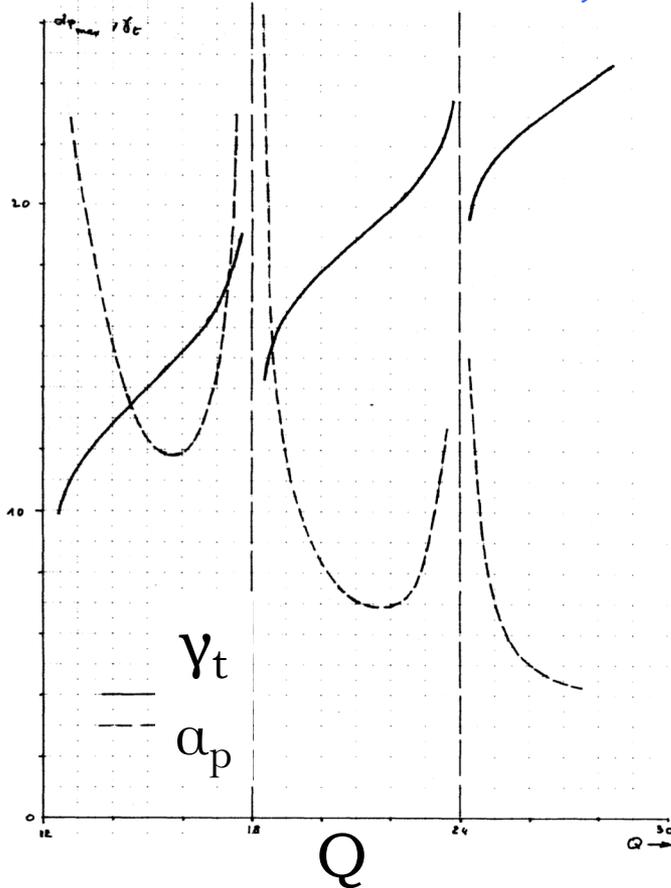


- Resonant oscillation of dispersion function close to the “Resonant integer tunes” (multiples of **super-periodicity 6**) \rightarrow asymptotic behavior of γ_t , (difficult for routine operation)
- γ_t is a linear function of horizontal tune Q_x elsewhere

- Nominal SPS working point for LHC proton beams ($\gamma_t \sim 23$)
- D. Boussard et al., SPS improvement note No 147, 1978; Injection above transition as TT10 was not ready for 26 GeV/c ($\gamma_t \sim 14$)
- G. Arduini et al., CERN/SL-Note 98-001, 1998; “Resonant tune” ($\gamma_t \sim 20$)
- Low γ_t , 2010 - “Resonant arc” with small dispersion in long straight sections ($\gamma_t \sim 18$)

- Injection beam line TT10 has not been upgraded to 26 GeV in 1978 and limited to 16 GeV
- Injection above transition is possible if SPS integer part of the tune is lowered to 15 ($\gamma_t \sim 14$)

D. Boussard et al., SPS improvement note No 147, 1978



Manipulating optics for curing instabilities

□ Transverse instabilities

- TMCI at injection - single bunch instability in vertical plane
 - Threshold at 1.6×10^{11} p/b ($\epsilon_l = 0.35$ eVs, $\tau = 3.8$ ns) with low vertical chromaticity

$$N_{th} \propto \frac{\epsilon_l}{\beta_y} \eta$$

- E-cloud vertical instability for 25 ns beam
 - Threshold higher than 1.2×10^{11} p/b

$$N_{th} \propto Q_s \propto \sqrt{\eta}$$

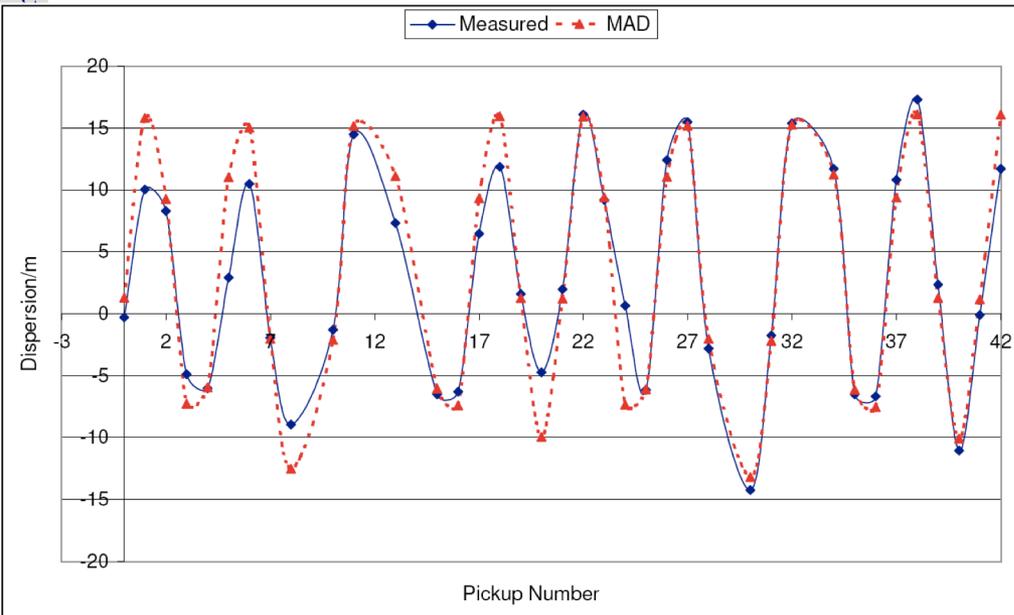
□ Longitudinal instabilities

- Single bunch and coupled bunch
 - Threshold at 2×10^{10} p/b for single harmonic RF (800 MHz cavity use is mandatory)

$$N_{th} \propto \epsilon_l^{5/2} \eta$$

- By setting the SPS integer tune to a multiple of 6, large dispersion wave can be introduced (dispersion becomes even negative) by overall reducing transition energy
- Successfully establishing cycle in the SPS and measuring dispersion very close to the one of MAD
- 3-fold increase of the slippage factor can be achieved (model)
- “Difficult” beam conditions (especially for injection)
- Need optics where dispersion is suppressed in straight section

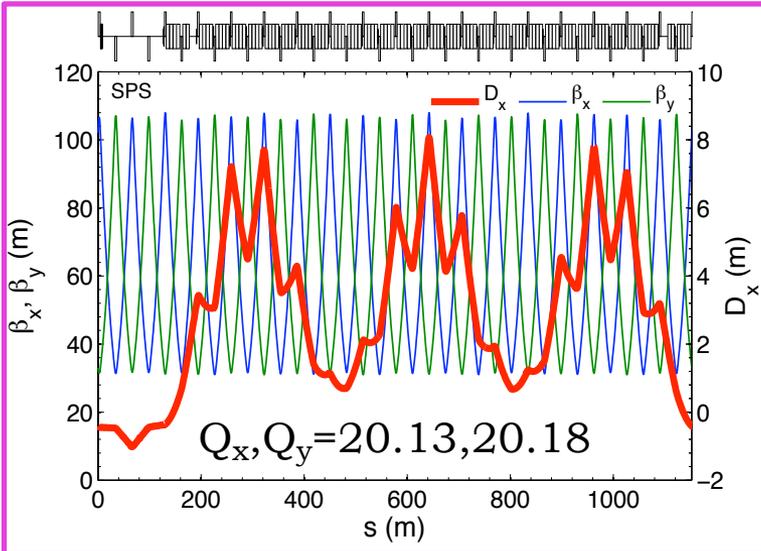
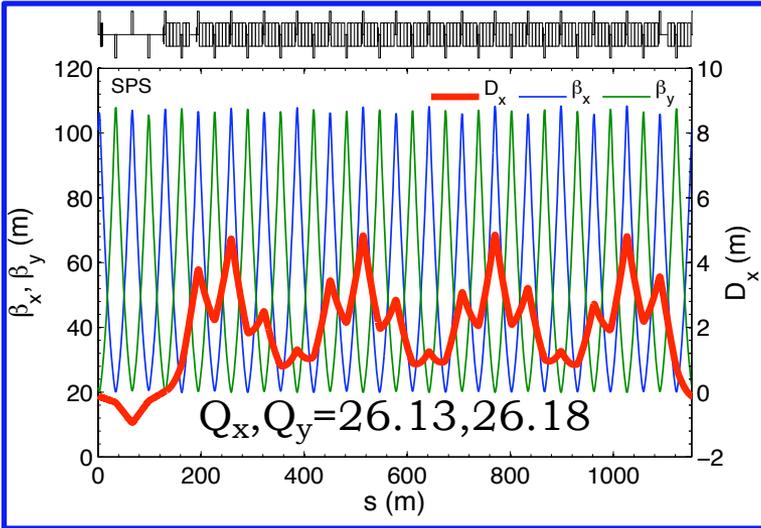
CERN Accelerator School, September 2018



G. Arduini et al., CERN/SL-Note 98-001 (MD), 1998

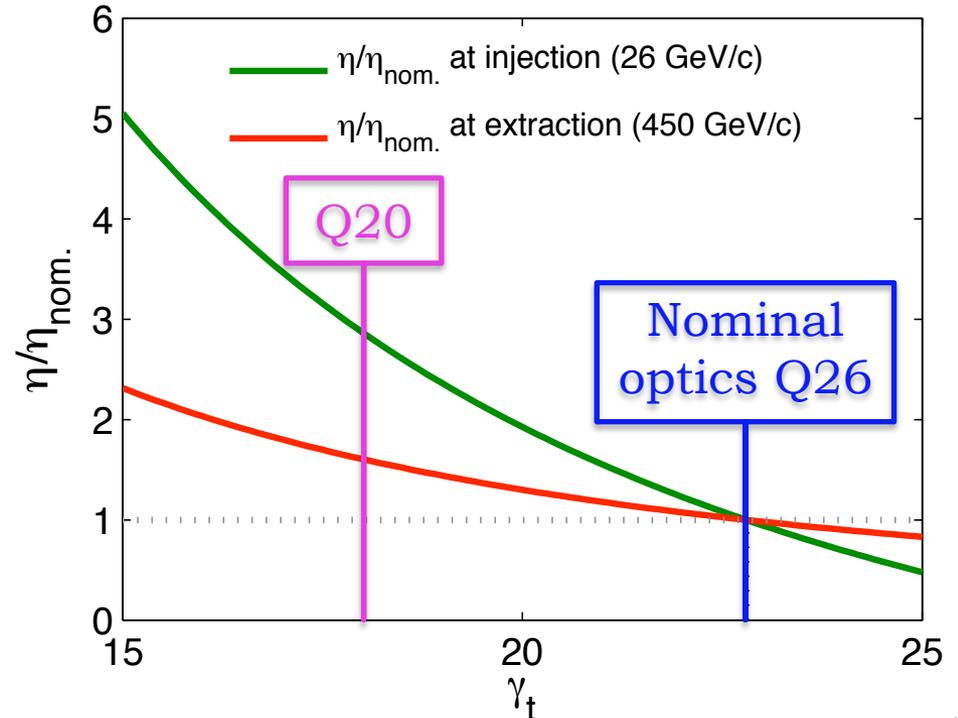
Q_h	Q_v	γ_{tr}	η (10^{-3})
24.18	24.22	18.54	1.61
24.29	24.32	19.59	1.30
26.62	26.58	23.23	0.551

Moving FODO phase advance from $4/16 \cdot 2\pi$ ($\pi/2$) to $3/16 \cdot 2\pi$ ($3\pi/8$)

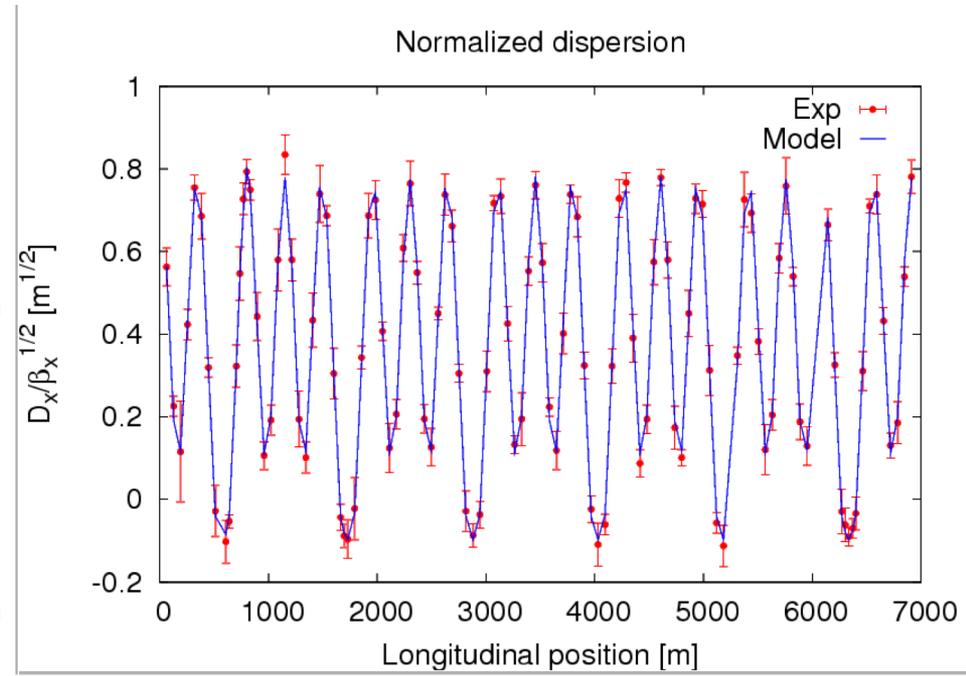
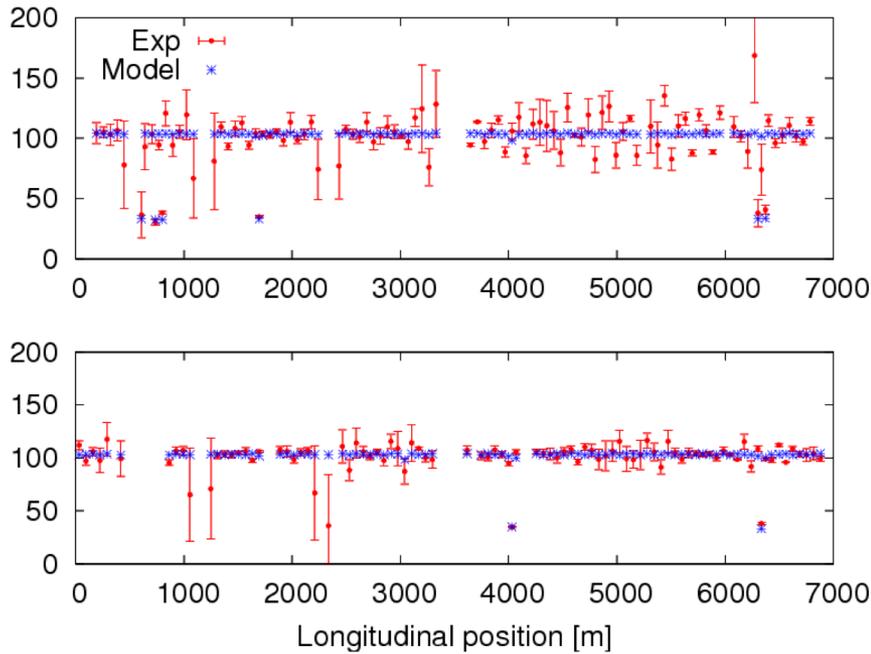


$$\eta = \frac{1}{\gamma_t^2} - \frac{1}{\gamma^2} \quad \longleftrightarrow \quad \gamma_{tFODO} \approx Q_x$$

Slippage factor increased by a factor of **2.8** at **injection** and **1.6** at **flat top**



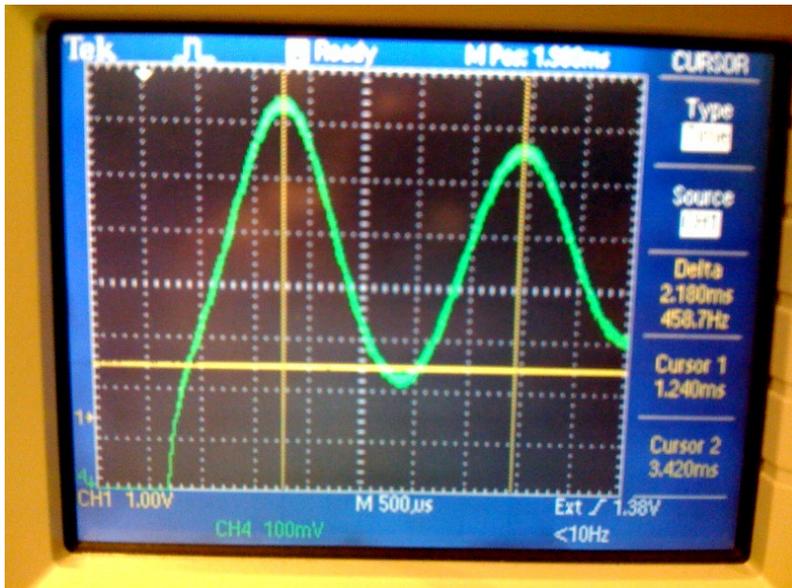
Slip factor relative to nominal SPS optics ^ω



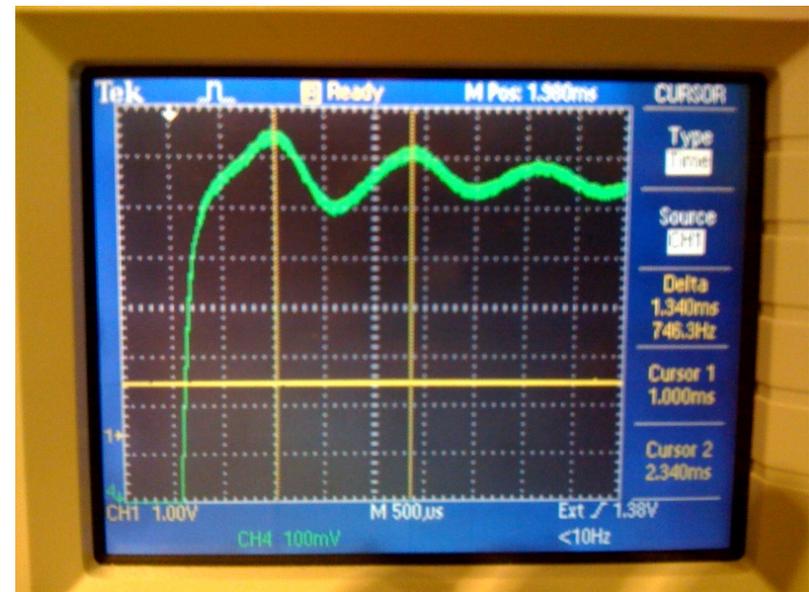
Measurement of the optics functions of the new lattice

- Beta beating around 20% in horizontal and 10% in vertical plane
- Normalized dispersion in striking agreement with the model

- ❑ Measured synchrotron frequency from “quadrupole” oscillations at injection
 - ❑ Same RF-voltage for both optics
- ❑ Ratio of Synchrotron frequencies ~ 1.63 corresponds to an **increase in slippage factor η by factor 2.65** (MADX prediction: 2.86)



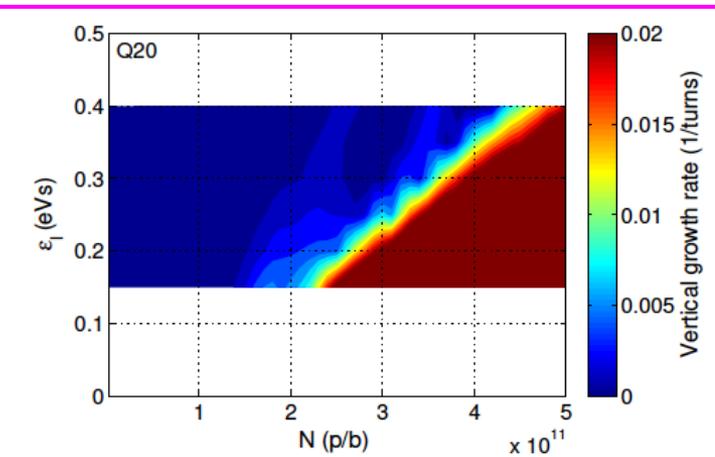
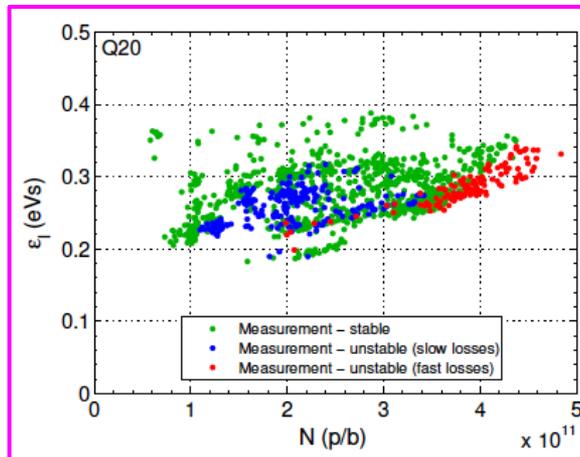
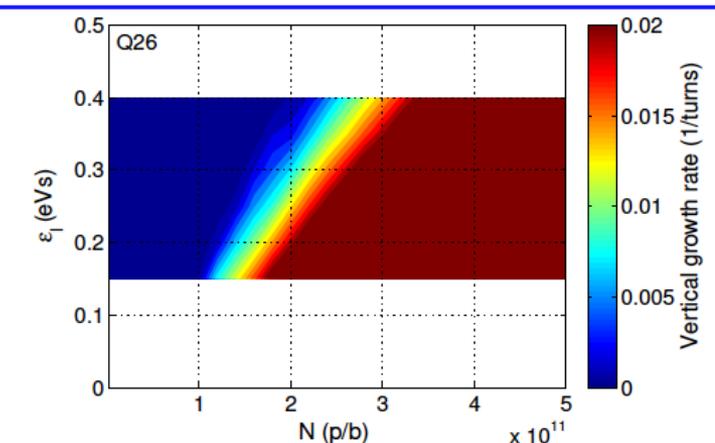
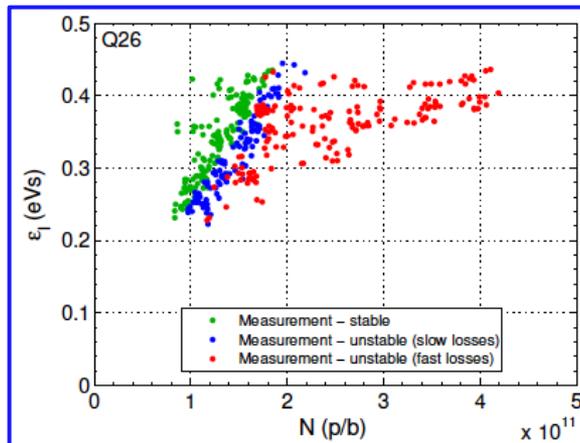
Q26: $F_s = 458/2 = 229\text{Hz}$,
 $Q_s = 0.0106/2 = 0.0053$



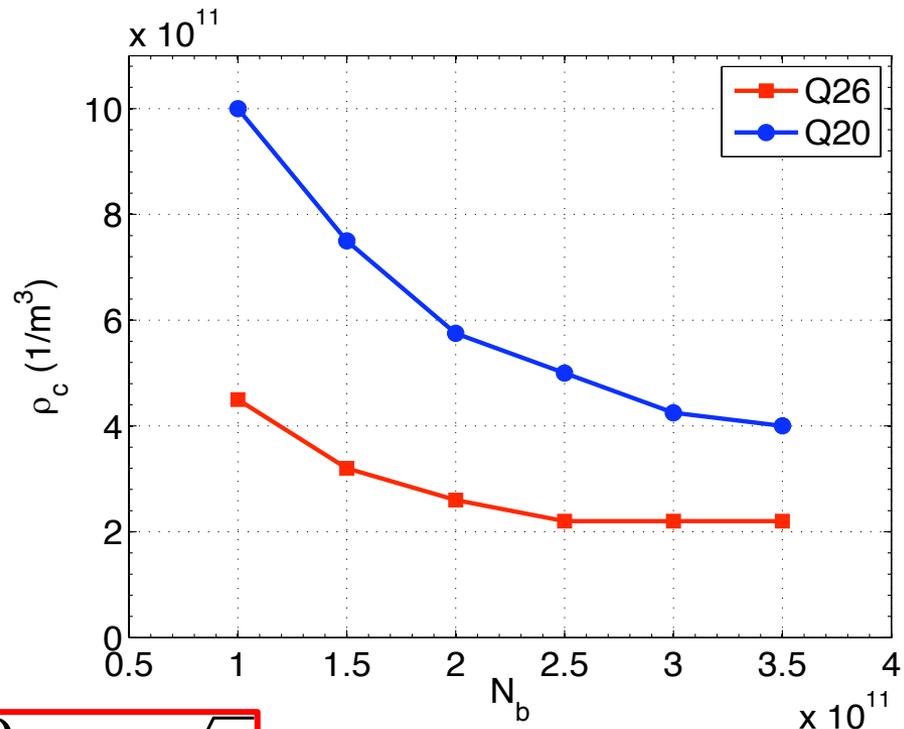
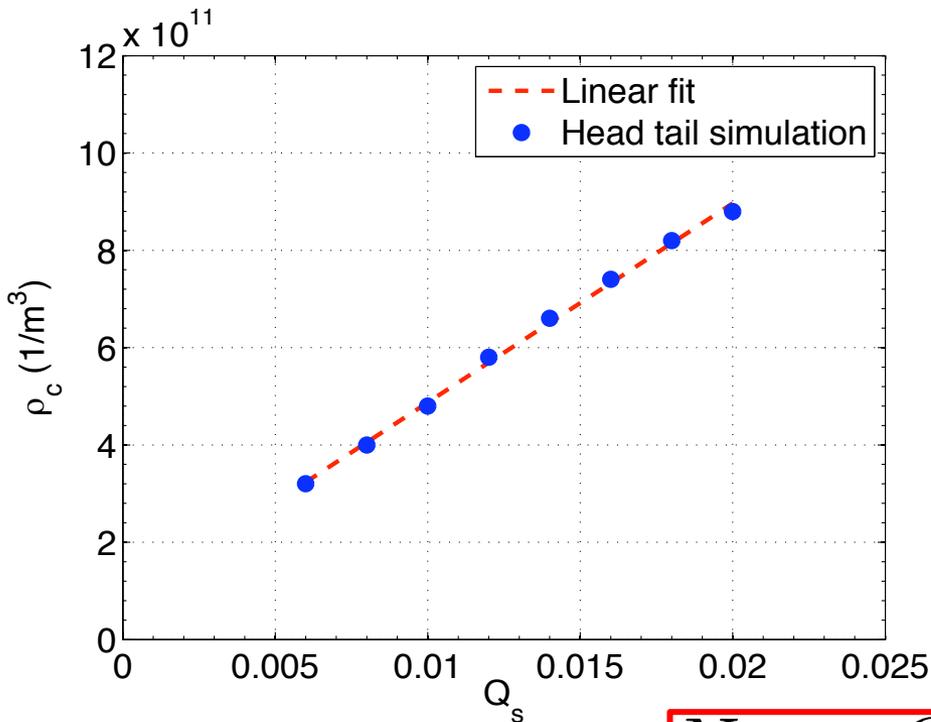
Q20: $F_s = 746/2 = 373\text{Hz}$,
 $Q_s = 0.0172/2 = 0.0086$

- ❑ In **nominal optics**, measured/simulated threshold at 1.6×10^{11} p/b for low chromaticity
 - ❑ High-chromaticity helps increasing threshold, but also losses along the cycle become excessive
- ❑ Measured/simulated threshold in **Q20** $> 4 \times 10^{11}$ p/b!!!

$$N_{th} \propto \frac{\epsilon_l}{\beta_y} \eta$$



H. Bartosik et al,
IPAC 2014



$$N_{th} \propto Q_s \propto \sqrt{\eta}$$

- ❑ Simulations with HEADTAIL code
 - ❑ Injection energy, uniform cloud distribution, located in dipole regions
- ❑ Linear scaling with Synchrotron tune demonstrated
 - ❑ Clearly higher thresholds predicted for **Q20**

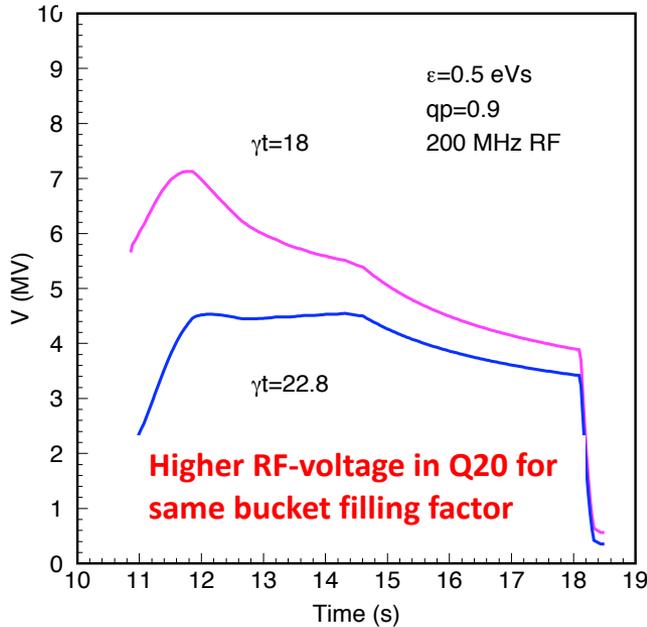
More margin with Q20 if e-cloud becomes issue for high intensity

H. Bartosik et al, IPAC2011

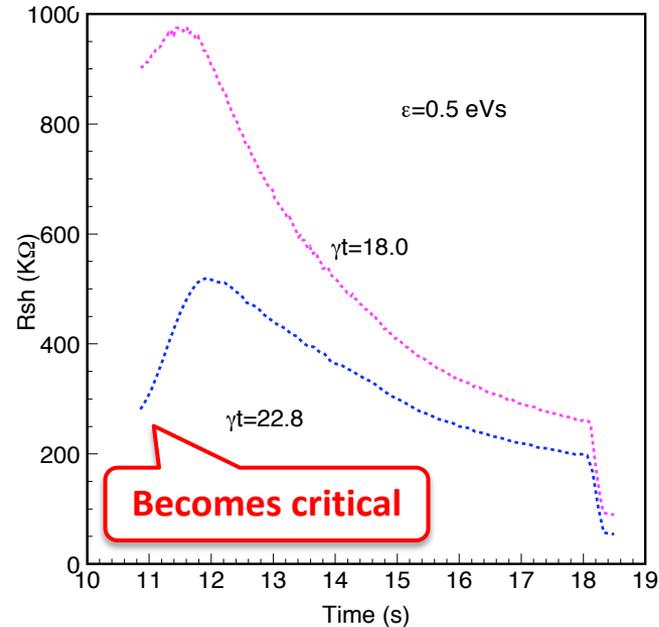
$$N_{th} \propto \epsilon_l^{5/2} \eta$$

E. Shaposhnikova

Voltage program 200 MHz RF system



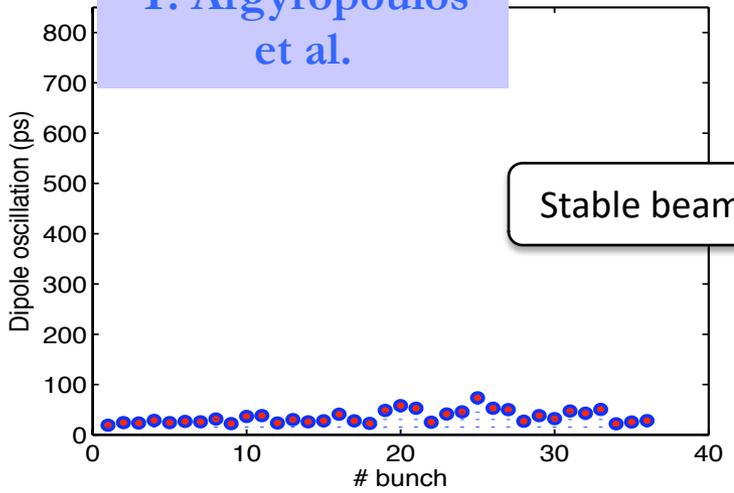
Narrow band impedance threshold



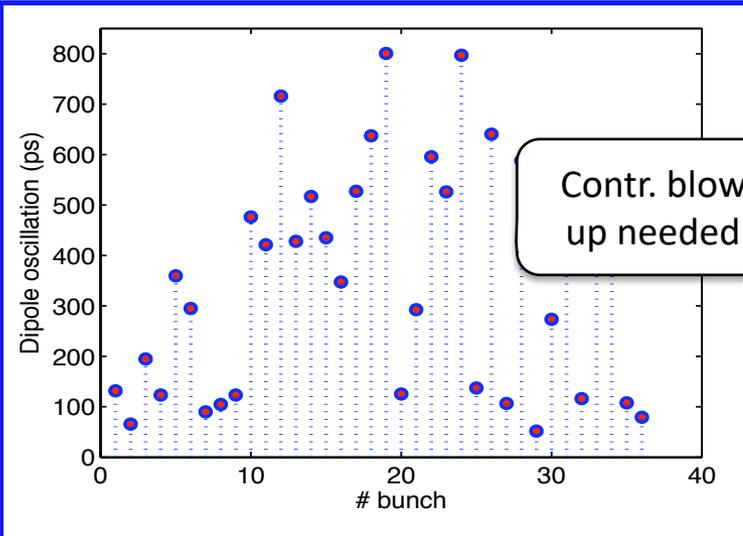
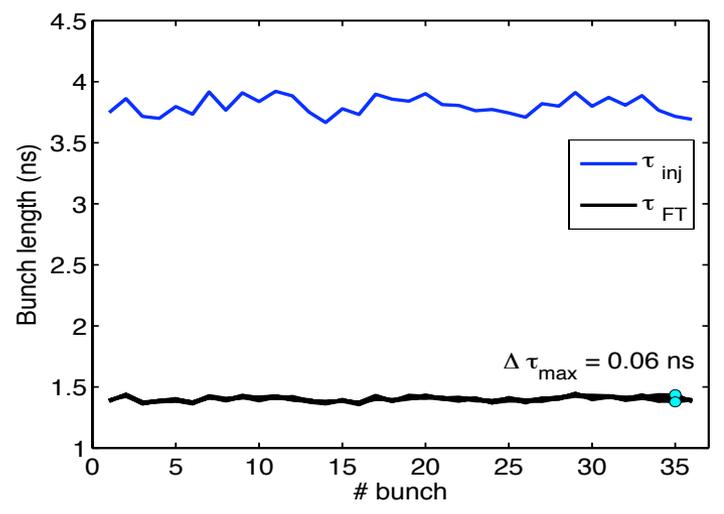
- ❑ Impedance threshold has minimum at flat top
 - ❑ Controlled longitudinal emittance blow-up during ramp for **Q26**
 - ❑ Less (or no) longitudinal emittance blow-up needed in **Q20**
- ❑ Instability limit at flat bottom
 - ❑ Critical with **Q26** when pushing intensity
 - ❑ Big margin with **Q20** (factor of 3)

T. Argyropoulos et al.

Stable beam

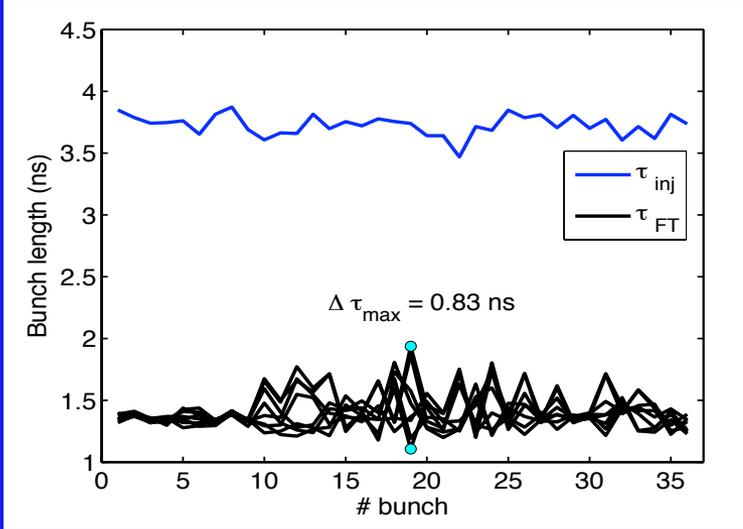


SPS-Q20 (1.6×10^{11} p/b)
double harmonic RF

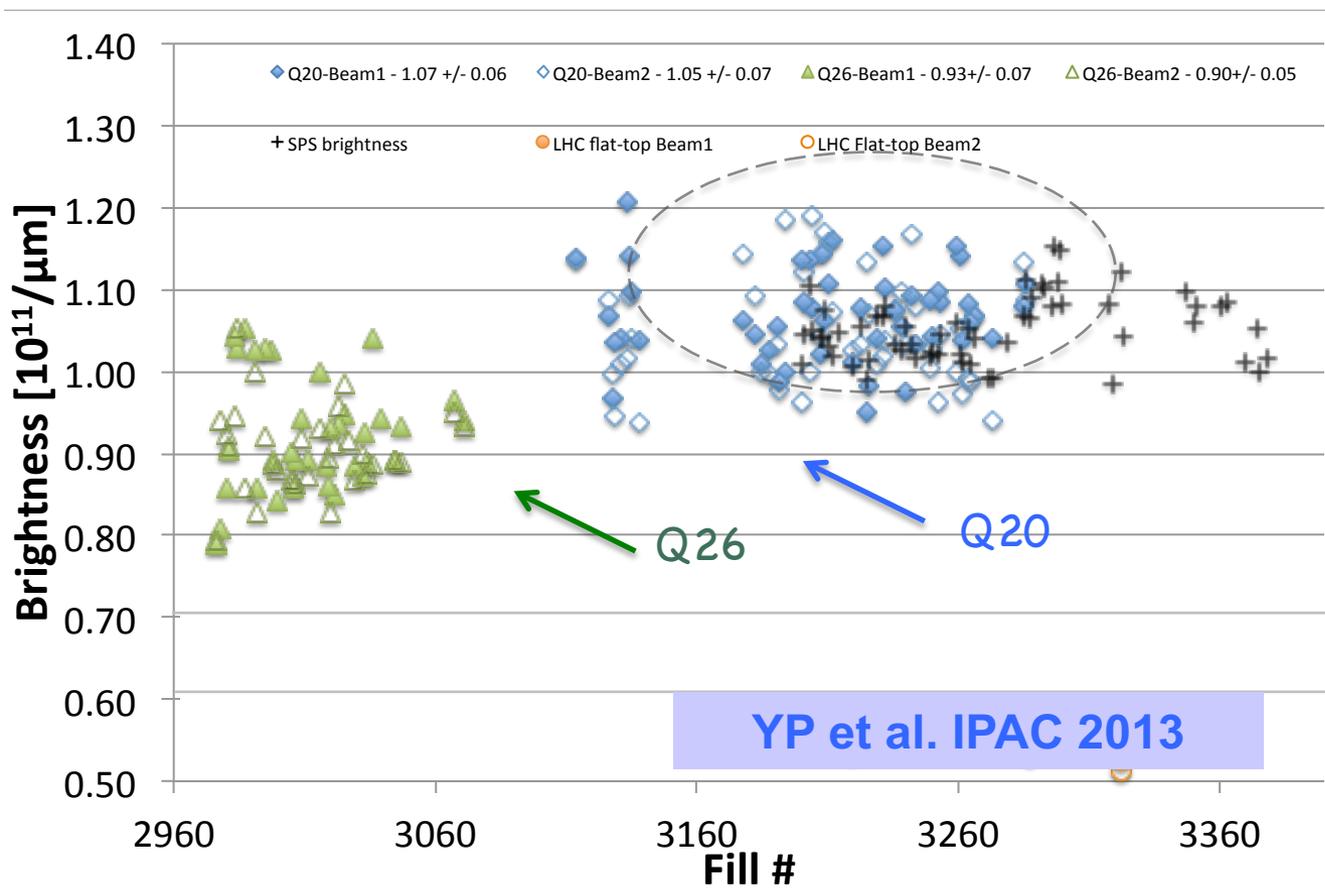


Contr. blow-up needed!

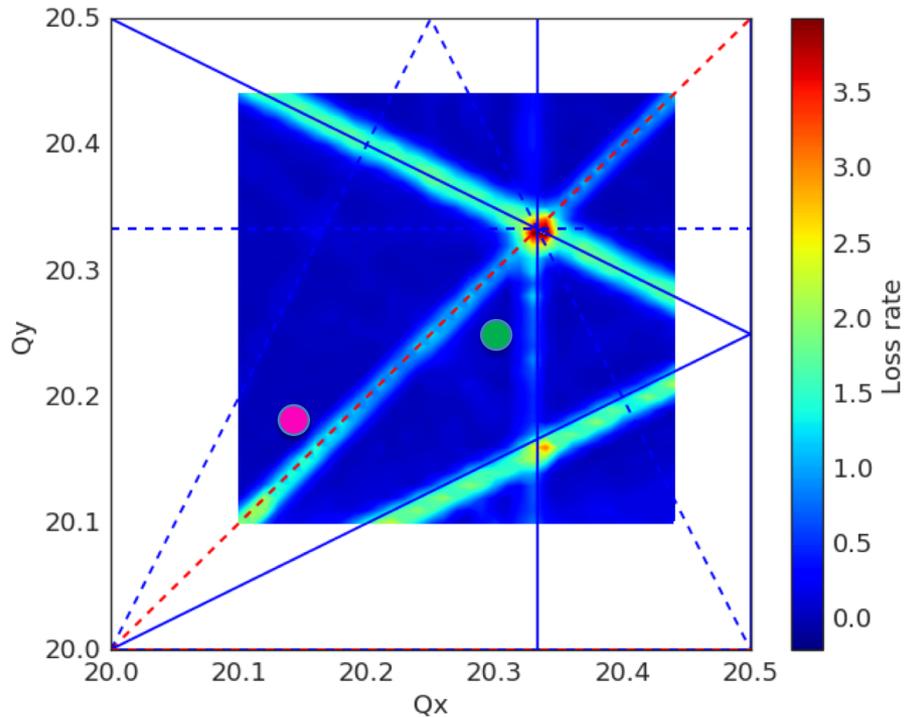
SPS-Q26 (1.6×10^{11} p/b)
double harmonic RF



- Operational deployment of Q20 optics for LHC beams since 2012 allowing around **20% brighter beams on LHC flat bottom**
- Opened way for **ultra-high brightness beams** of HL-LHC era



Non-linear dynamics



- proton working point
- ion working point

resonances:

red: systematic

blue: non-systematic

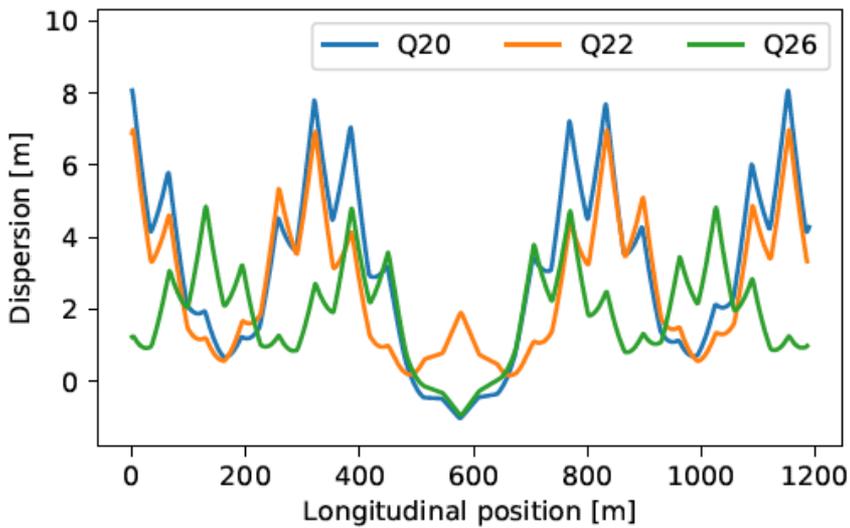
— upright

- - skew

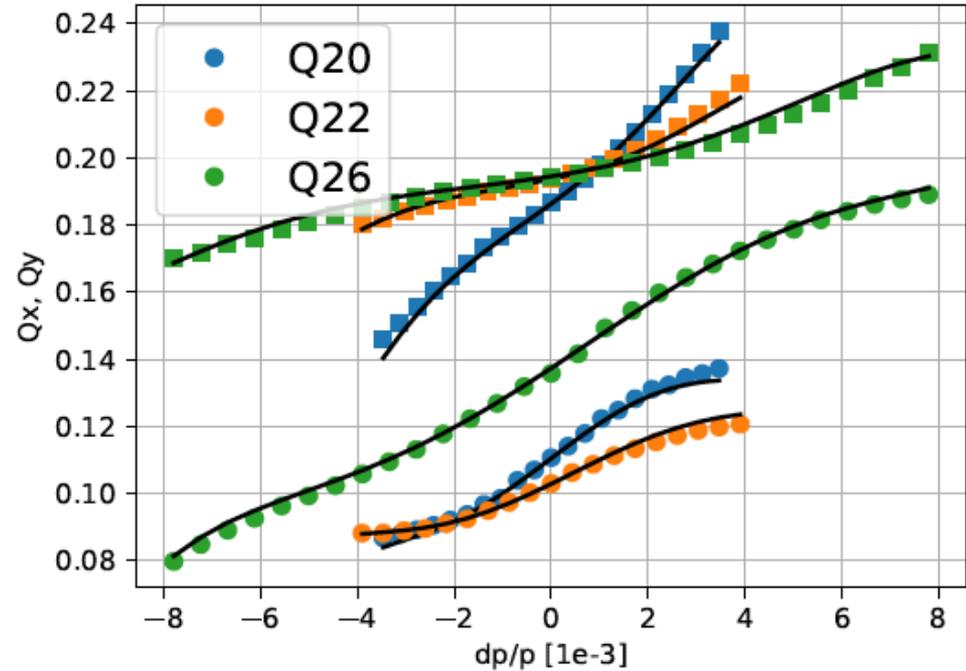
H. Bartosik et al. HB2018

■ Dynamic tune scan for identification of resonances

- Losses around 3rd order (normal) resonances and the diagonal clearly observed
- Faint traces of 4th order resonances
- Operational working point for protons 20.13/20.18 (moved up for high brightness beams)



M. Cartla' et al. IPAC2018



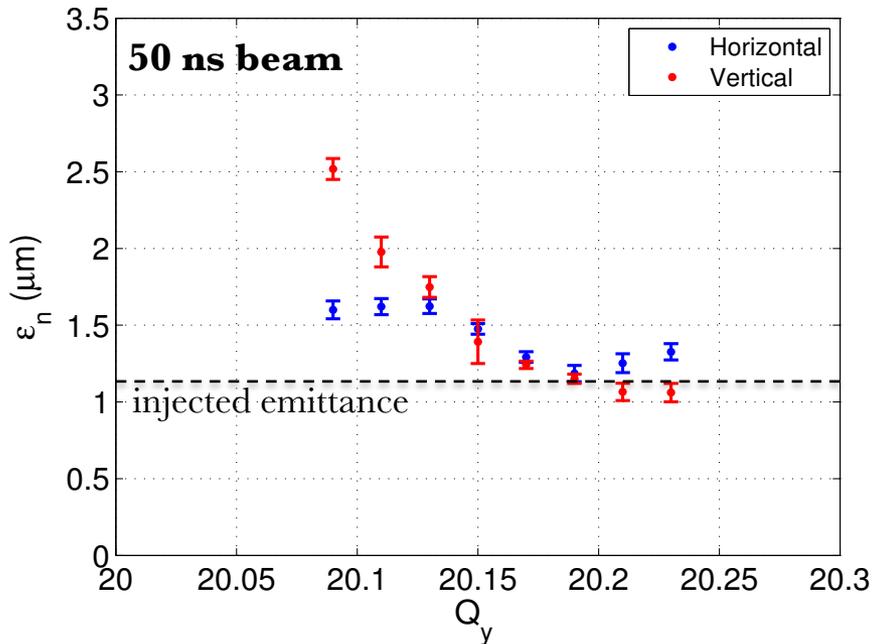
Sextupoles

$$Q = \underbrace{Q_0}_{\text{Quadrupoles}} + \underbrace{Q' \left(\frac{\delta p}{p}\right)}_{\text{Octupoles}} + \underbrace{Q'' \left(\frac{\delta p}{p}\right)^2}_{\text{Decapoles}} + Q''' \left(\frac{\delta p}{p}\right)^3 + \dots$$

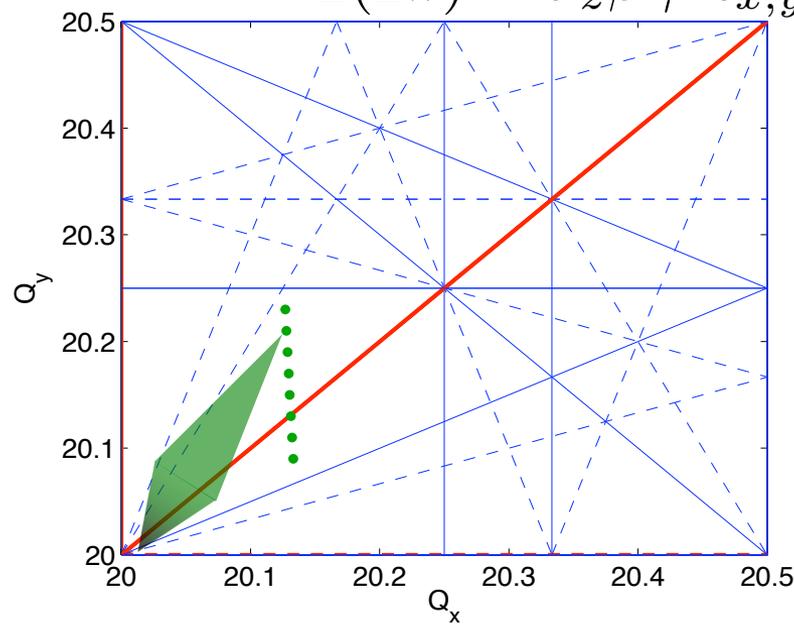
- Estimate “effective” magnet multi-poles that reproduce non-linear chromaticity measurement for three different optics

Space-charge

emittance at end of flat bottom

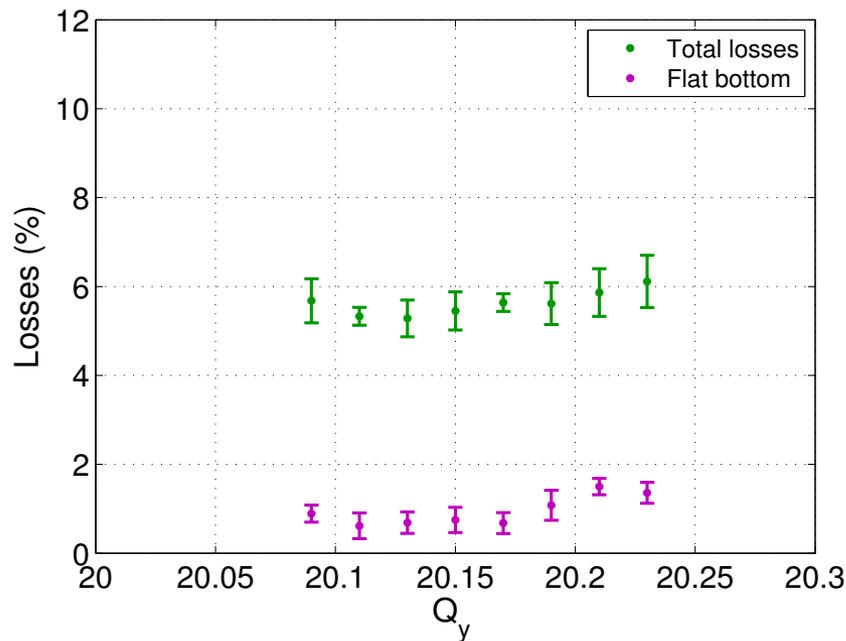
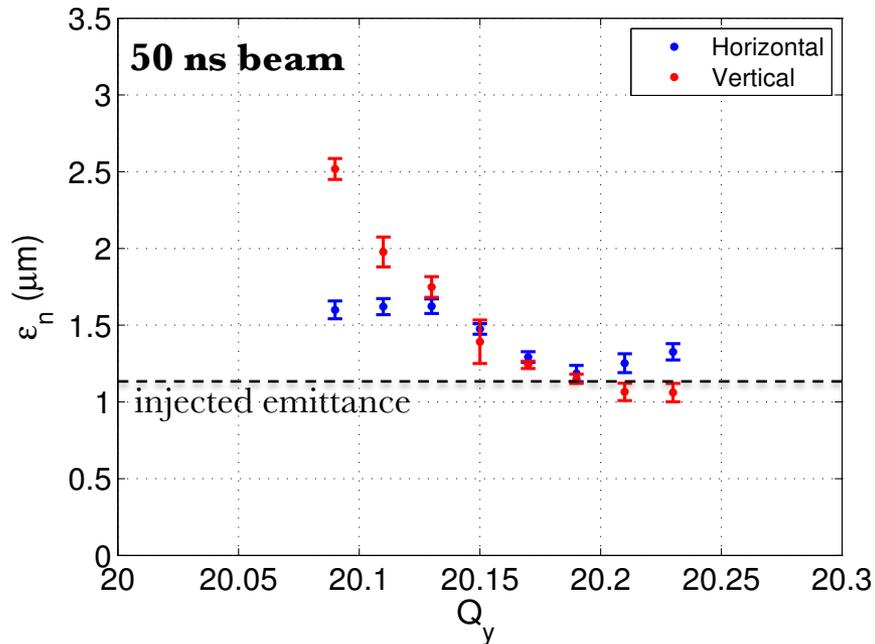


$$\Delta Q_{x,y} = - \frac{r_0 N_p C}{2(2\pi)^{3/2} \sigma_z \beta \gamma^2 \epsilon_{x,y}}$$



- Vertical tune scan with high brightness beam for **10 s storage time**
 - $N = 1.95 \times 10^{11}$ p/b (at injection)
 - $\epsilon \sim 1.1 \mu\text{m}$ (at injection)
 - Transmission to flat top around 94% (very small losses on flat bottom)
 - Budget of 10% losses and 10% blow-up allows for tune spread of $\Delta Q_y = 0.21$
- $\Delta Q_x / \Delta Q_y \sim 0.10 / 0.20$

emittance at end of flat bottom

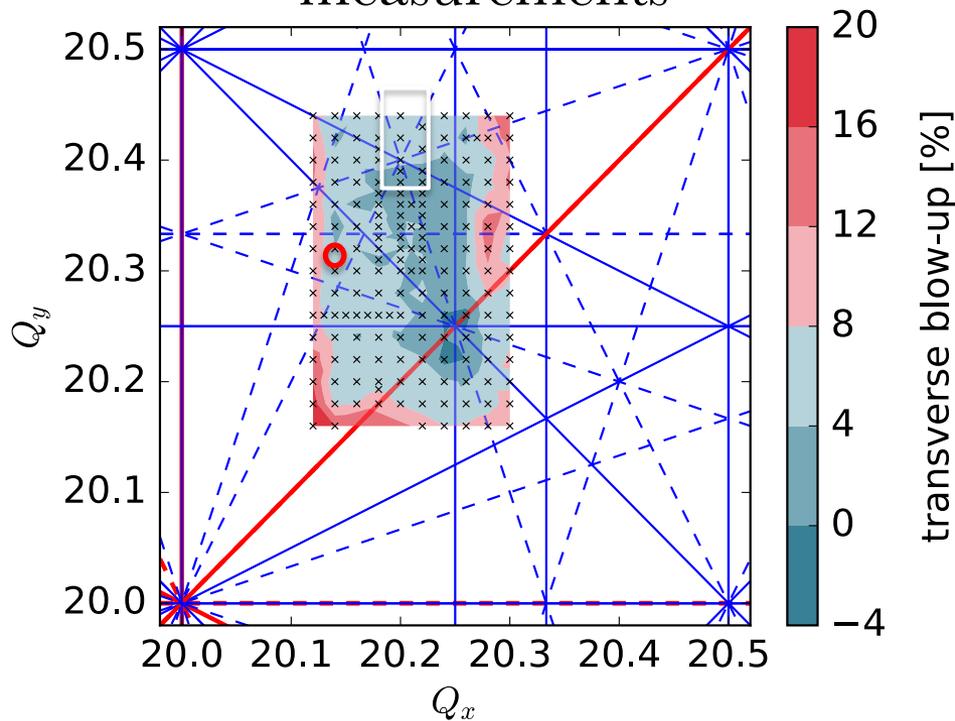


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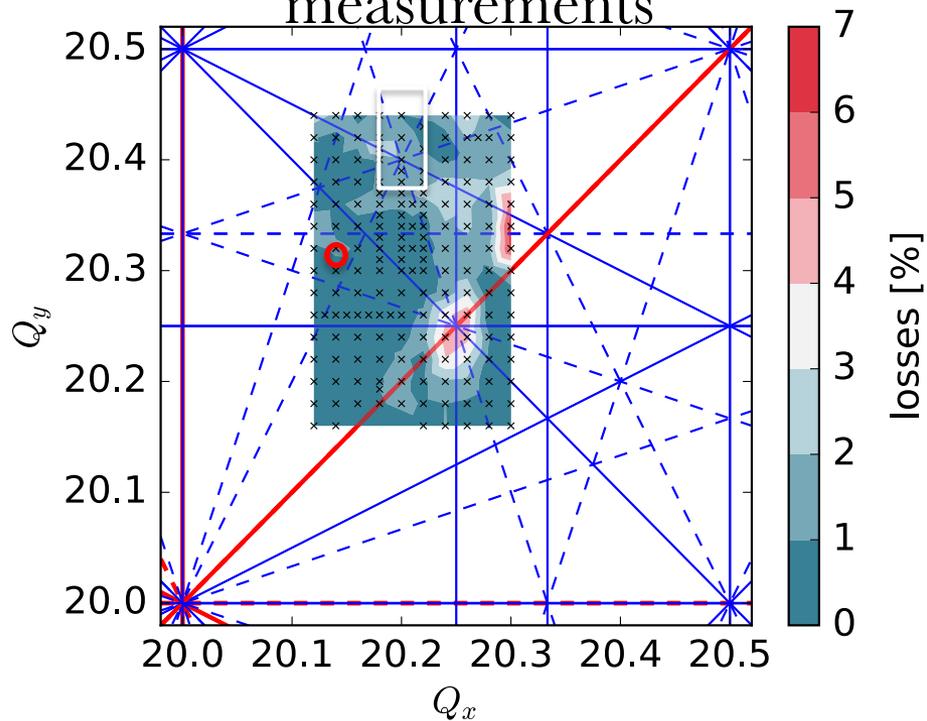
- Tune scan with high brightness single bunch beam for **3 s storage time**
 - Blow-up at integer resonances as expected (tune spread $\Delta Q_x, \Delta Q_y \sim 0.10, 0.19$)
 - **Margin for higher brightness for working points in white box** (enhanced losses only close to $Q_x + 2Q_y = 61$ normal 3rd order resonance and around $4Q_x = 81$ normal 4th order resonance)

September 2018

measurements

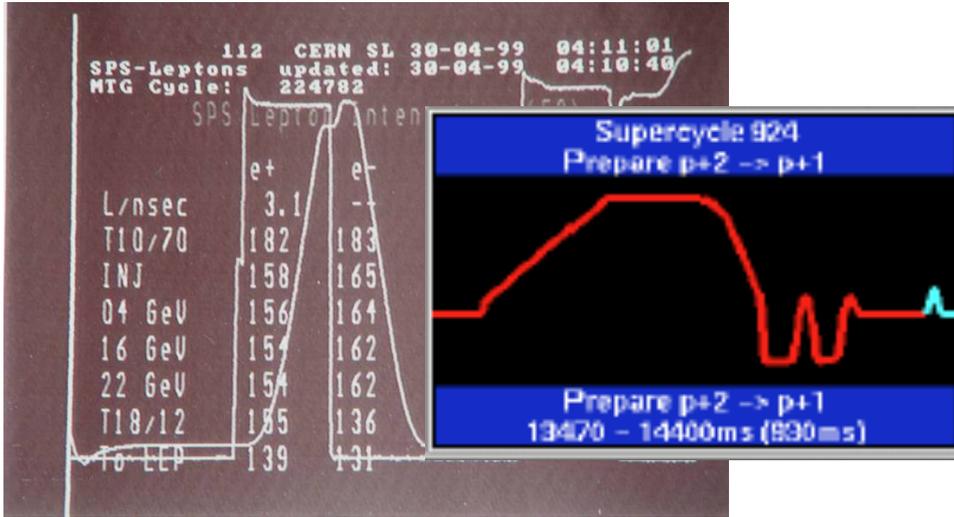


measurements

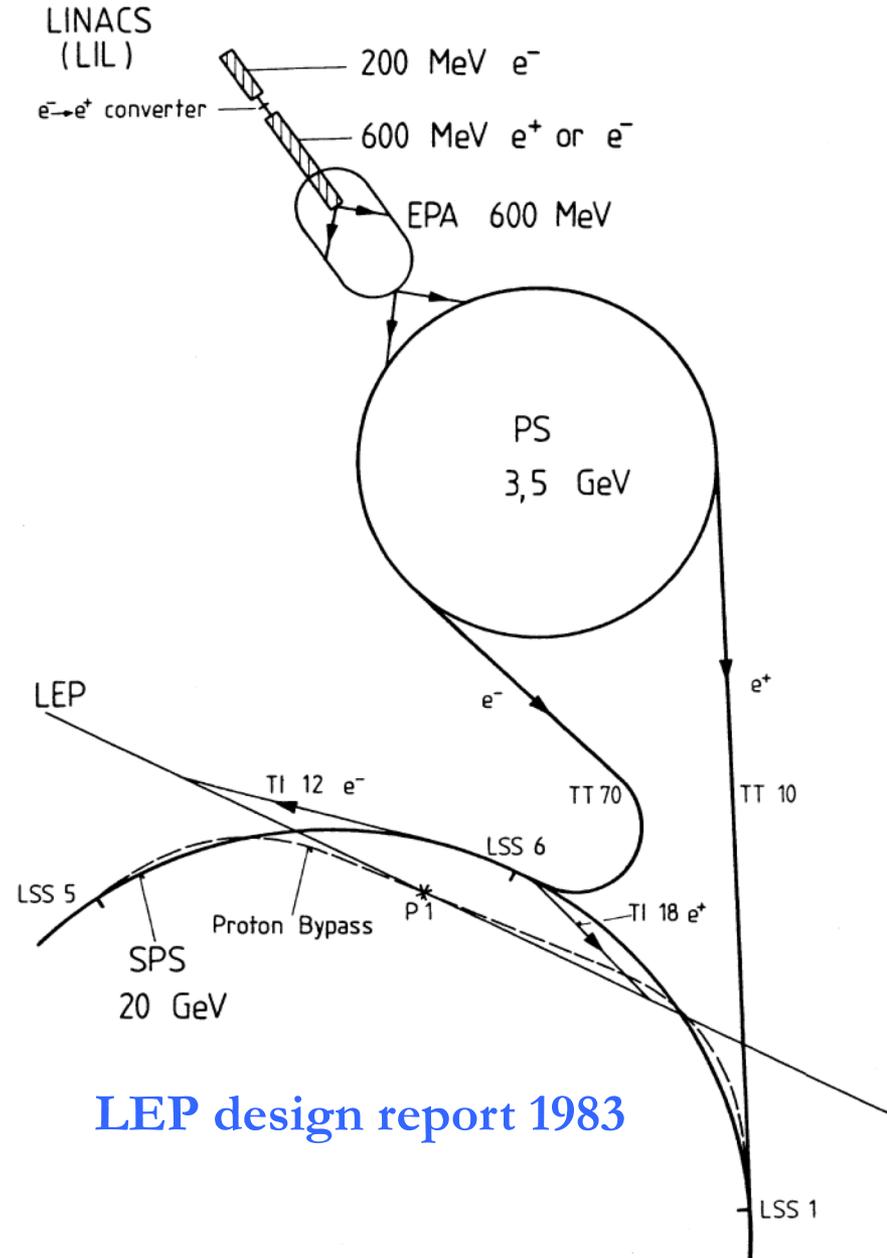


Electron-positron dynamics

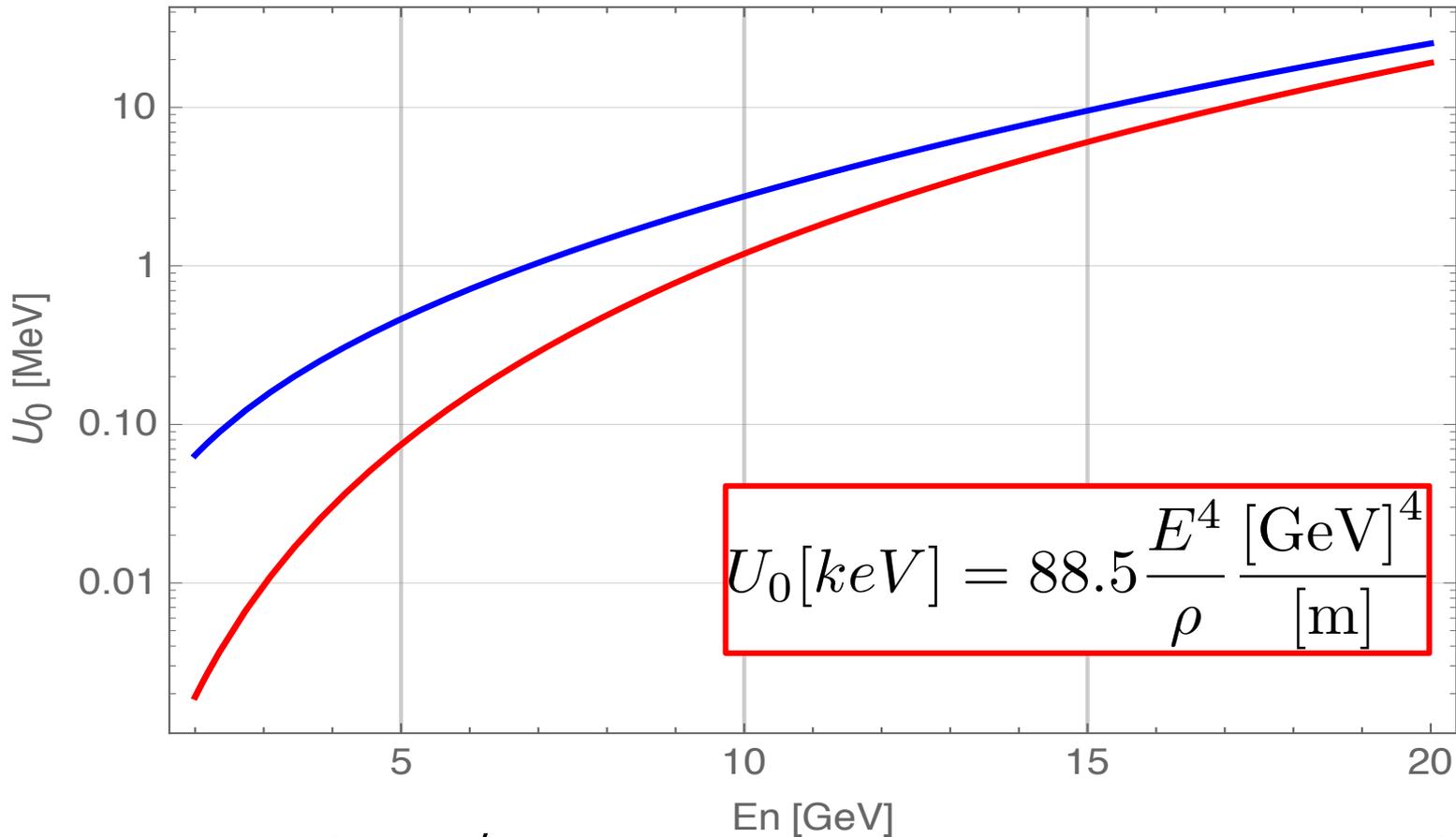
P. Collier – Academic Training 2005



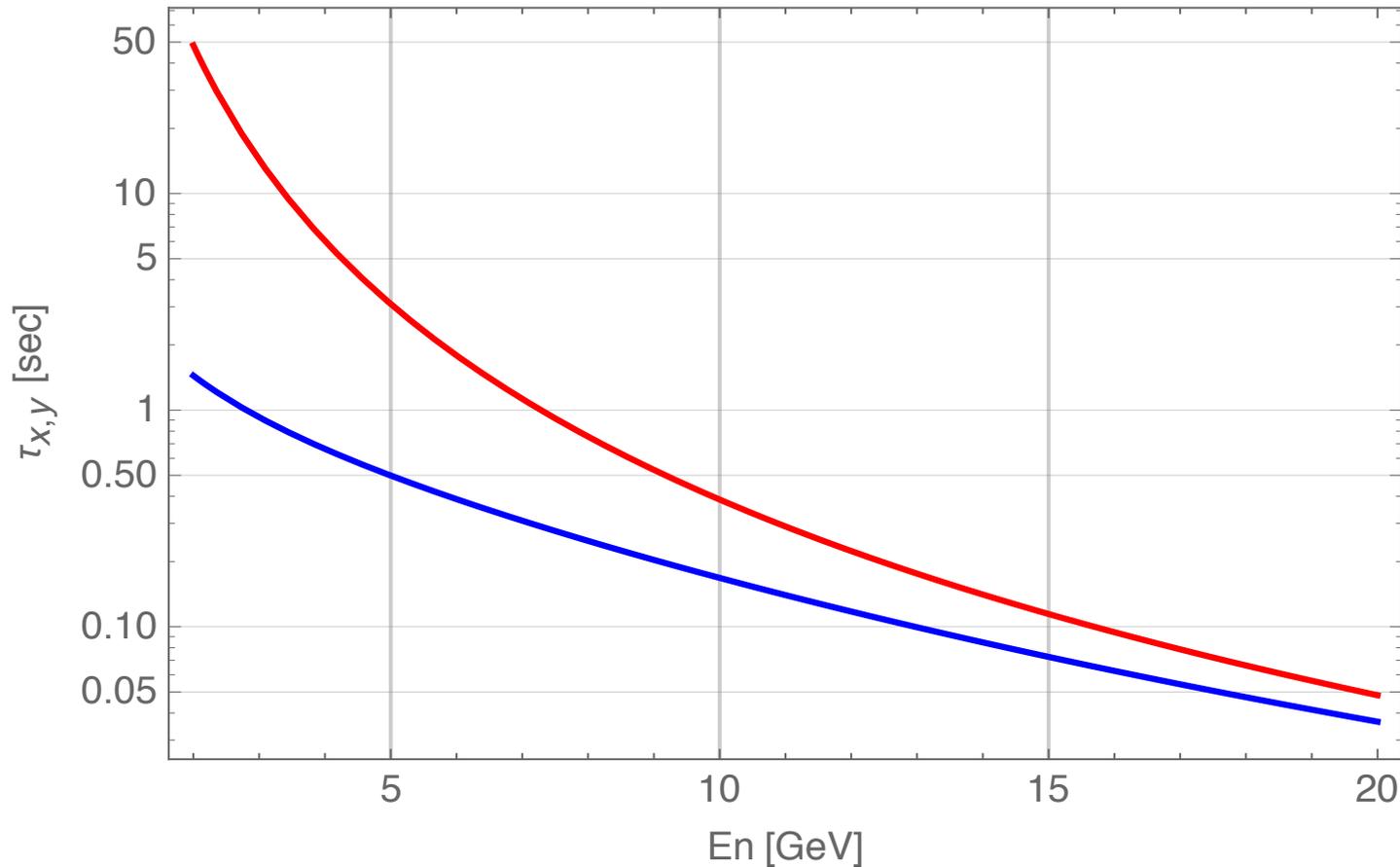
- LEP filling interleaved with proton operation
- 4 cycles with 4 bunches ($2e^+$, $2e^-$) evolved to 2 cycles with 8 bunches ($\sim 2.5 \times 10^{10}$ p/b)
- Energy to LEP: 18 \rightarrow 20 \rightarrow 22 GeV
- Lots of RF for leptons (200MHz SWC, 100MHz SWC, 352MHz SC),
- 2 Extractions in Point 6 towards LEP



LEP design report 1983

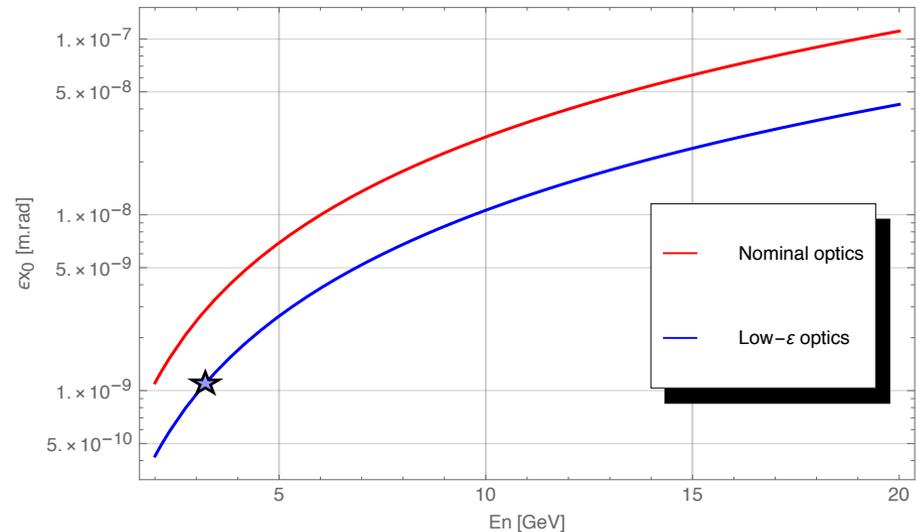
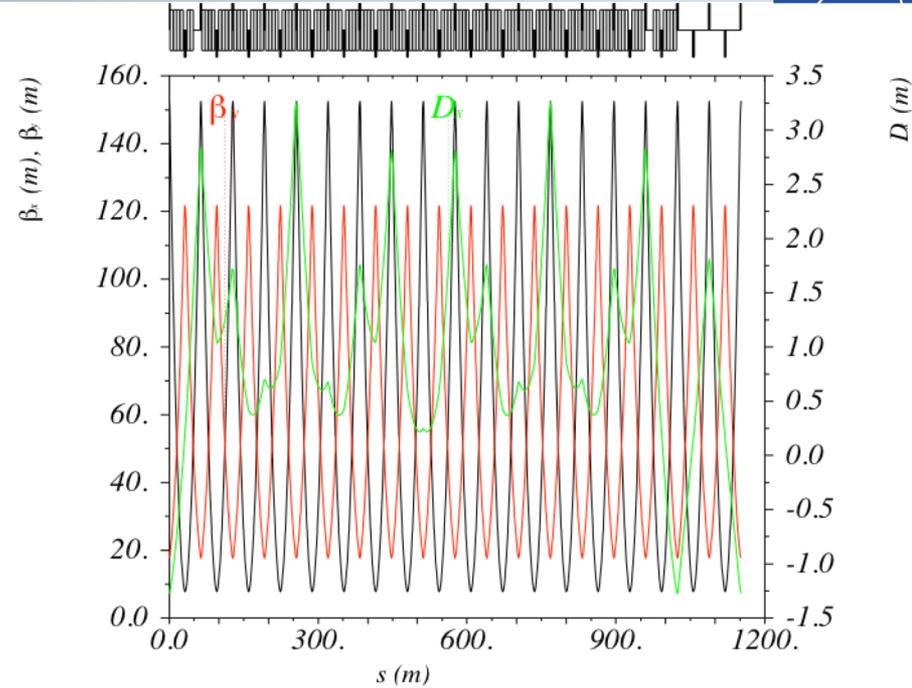


- Energy loss / turn necessitate **large RF voltage** (30 MV) at high energy
- Impact of a 2-m 3.5T **damping wiggler** is mild at high energies



- Damping time at injection (3.5 GeV) very large (9 s)
- A 2-m 3.5T **damping wiggler** could enhance damping for low energies to below 1 s (good for instabilities)

- Move horizontal phase advance to 135 deg. i.e. $3\pi/4$ (**Q40 optics**) which is optimal for low emittance in a FODO cell
- **Emittance with nominal optics @ 3.5 GeV of 3.4 nm drops to 1.3nm**
- Further reduction can be achieved with damping wiggler



- Using the 40+ years experience since the design and operation of the Super Proton Synchrotron (SPS), reviewed several beam dynamics concepts
 - Choice of energy, bending field and circumference
 - Optics design
 - Cell optics, insertions, transition energy
 - Collective effects
 - Instabilities, Space-charge
 - Electron / Positron beam dynamics
 - Equilibrium beam properties, energy loss / turn, damping time