Introduction to Non-linear Longitudinal Beam Dynamics

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CERN

Introduction to Accelerator Physics

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Outline

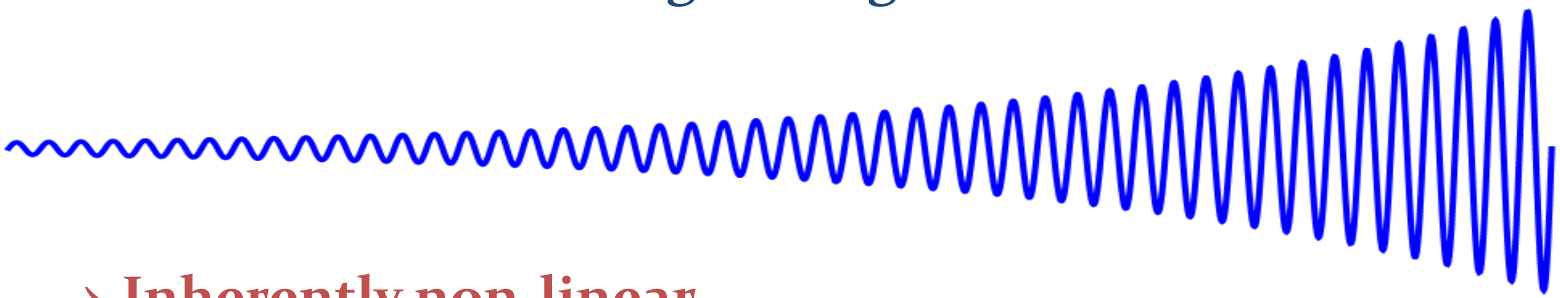
- Introduction
- Linear and non-linear longitudinal dynamics
  - Equations of motion, Hamiltonian, RF potential
- Longitudinal manipulations
  - Bunch length and distance control by multiple RF systems
  - Bunch rotation
- Synchrotron frequency distribution
  - Effect on longitudinal beam stability
- Summary

Introduction

Introduction

- **Signals generated by radio-frequency systems in particle accelerators are of the form** $V\sin(h\omega_{\mathrm{rev}}t)$

  → **Resonance effect:** **large voltage with little effort**



→ **Inherently non-linear**

→ **Linear longitudinal beam dynamics only an approximation**

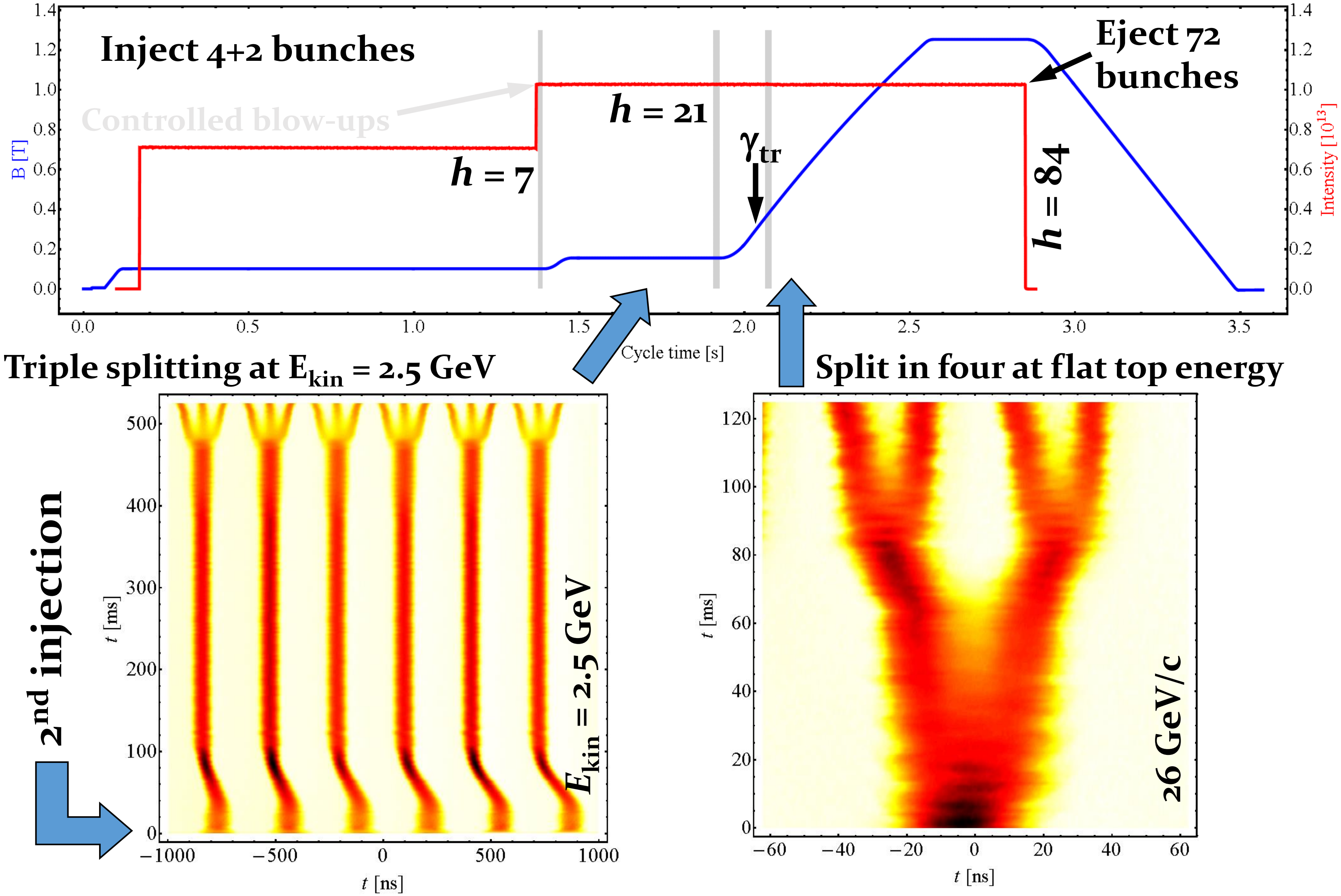
→ **Effect of non-linearity on beam?**

→ **Tools to describe and analyse non-linearity**

→ **Use non-linearity to improve beam conditions**

Non-linear longitudinal dynamics

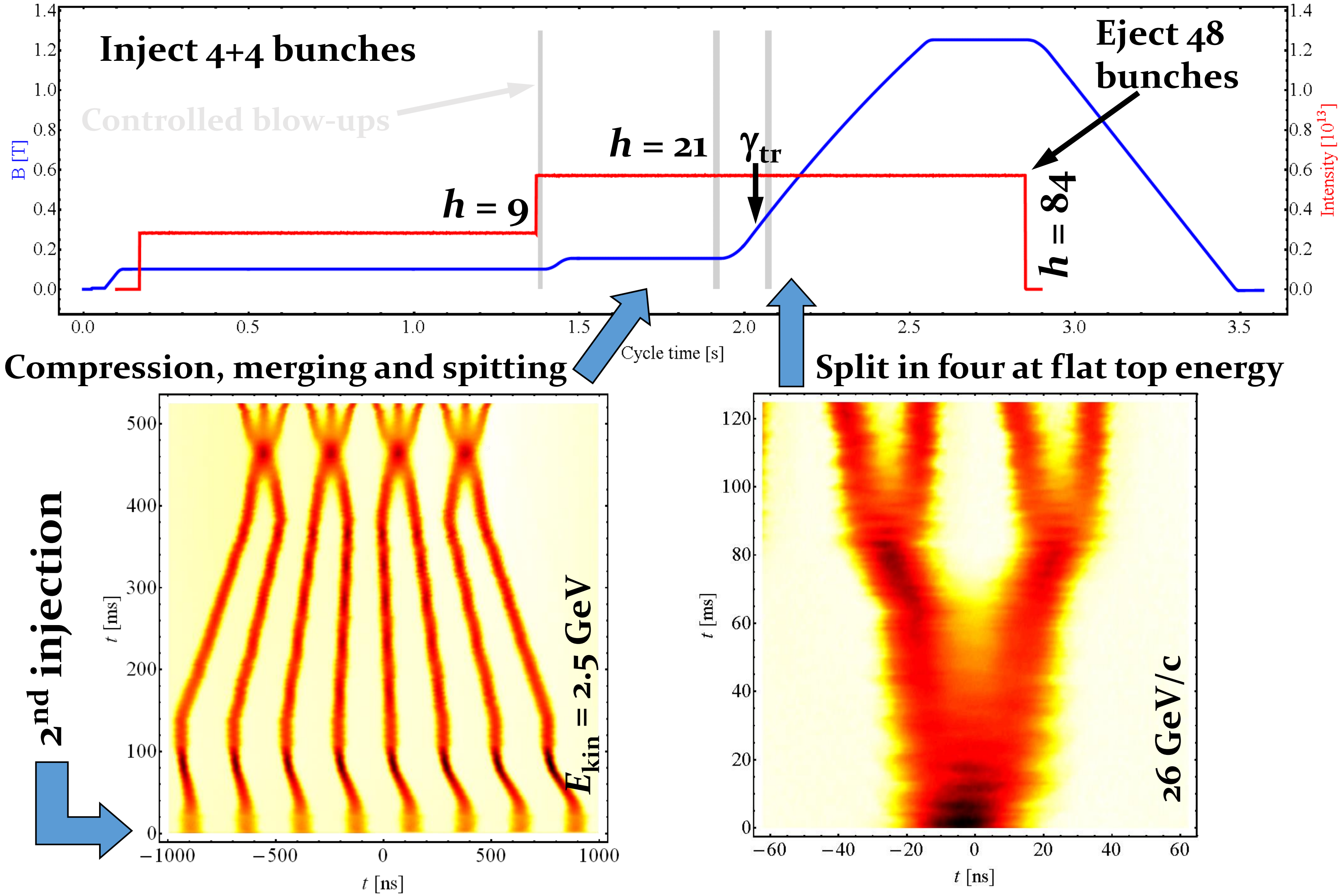
Example: LHC-type beam in the CERN PS



- **Non-linear RF allows to control all longitudinal parameters**

→ Number of bunches, bunch length and emittance, longitudinal stability

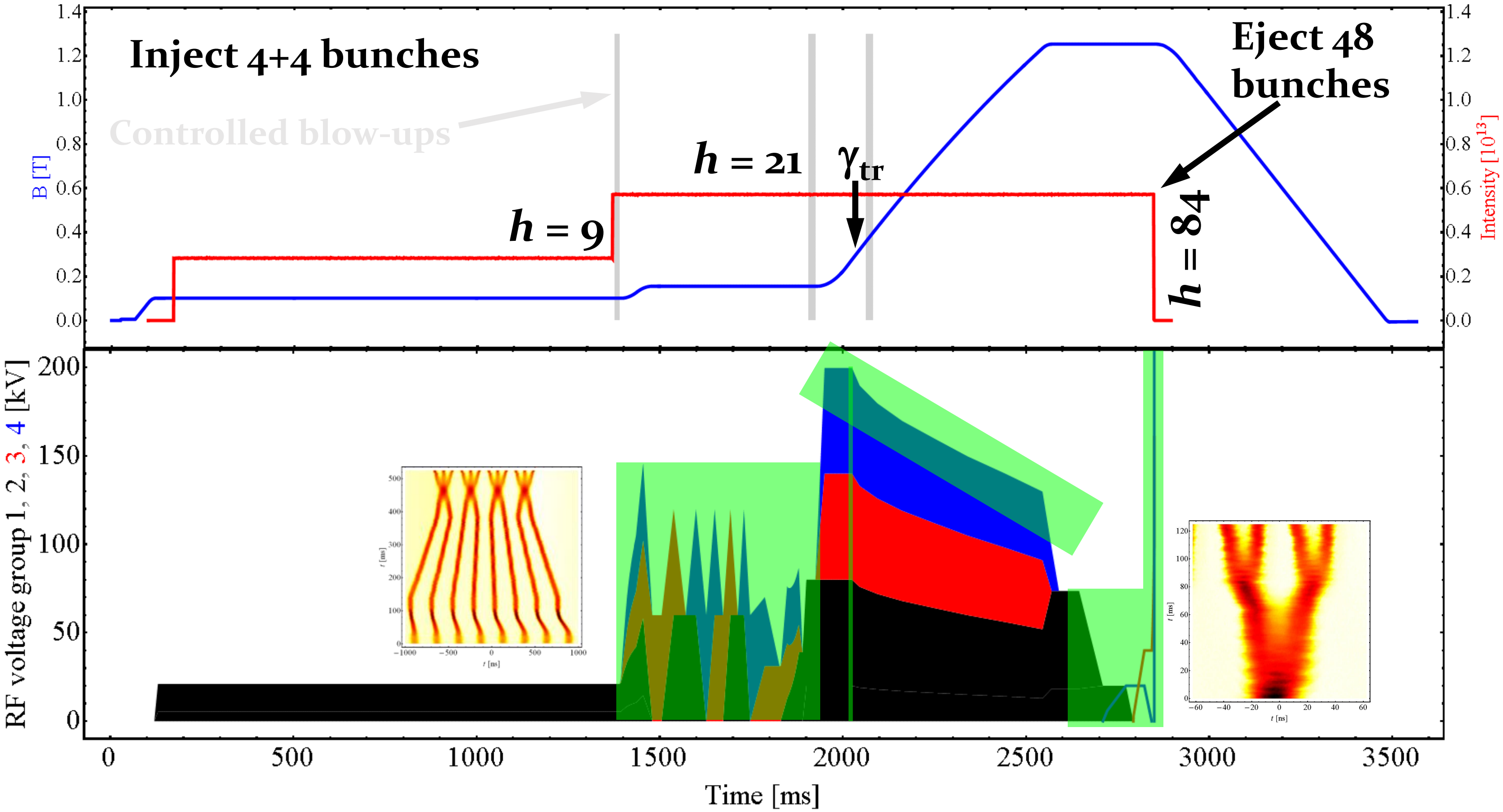
Example: LHC-type beam in the CERN PS



- **Non-linear RF allows to control all longitudinal parameters**

→ **Number of bunches, bunch length and emittance, longitudinal stability**

Where profit from non-linear RF?



→ RF manipulation from 8 bunches in *h* = 9 to 12 in *h* = 21

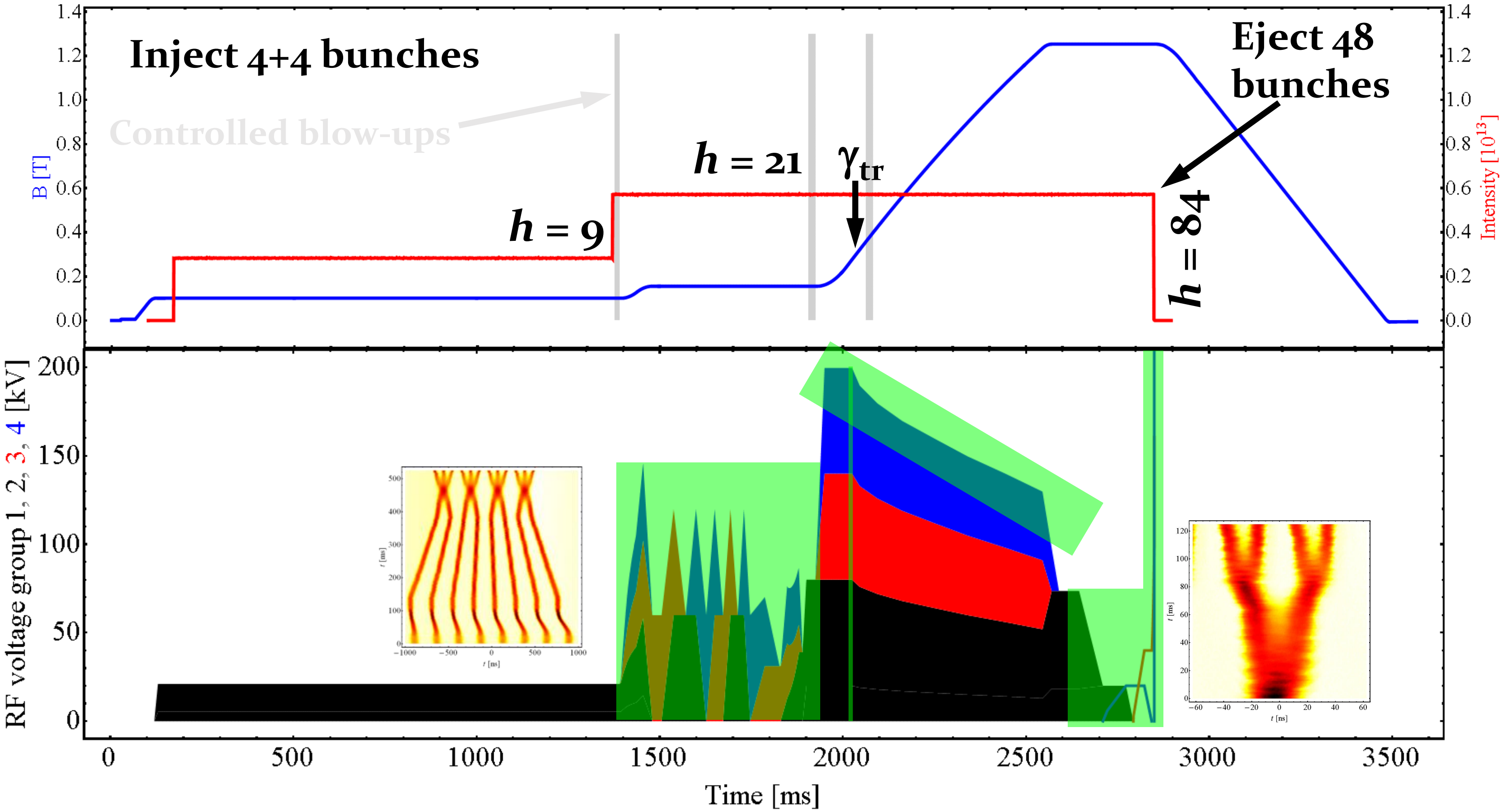
→ Transition crossing

→ RF voltage reduction during acceleration

→ Splitting at the flat-top

→ Bunch shortening (rotation) before extraction

Where profit from non-linear RF?



→ RF manipulation from 8 bunches in $h = 9$ to 12 in $h = 21$

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Applications

- **Introduce extra non-linearity**
  - **Bunch lengthening in double-harmonic RF system to reduce peak current (space charge)**

$$V_1 \sin(h_1 \omega_{\mathrm{rev}} t + \phi_1) + V_2 \sin(h_2 \omega_{\mathrm{rev}} t + \phi_2)$$

  - **Short and long bunches with multi-harmonic RF systems**

$$\sum_n V_n \sin(h_n \omega_{\mathrm{rev}} t + \phi_n)$$

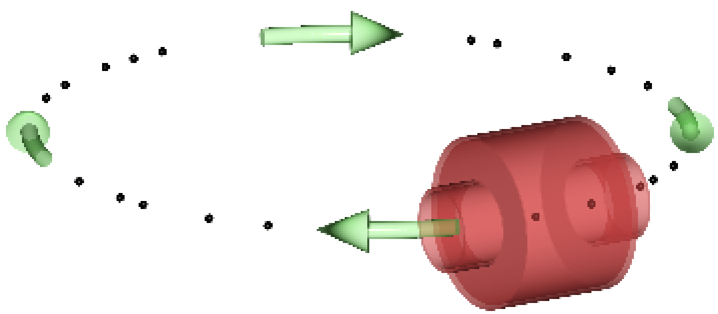
  - **Adapt bunch-to-bunch distance**

- **Profit from non-linearity for beam stabilization**
  - **Stabilize beam using higher-harmonic RF**
  - **Controlled longitudinal emittance blow-up**

Linear longitudinal beam dynamics

Interaction between particles and RF

Simple accelerator model:

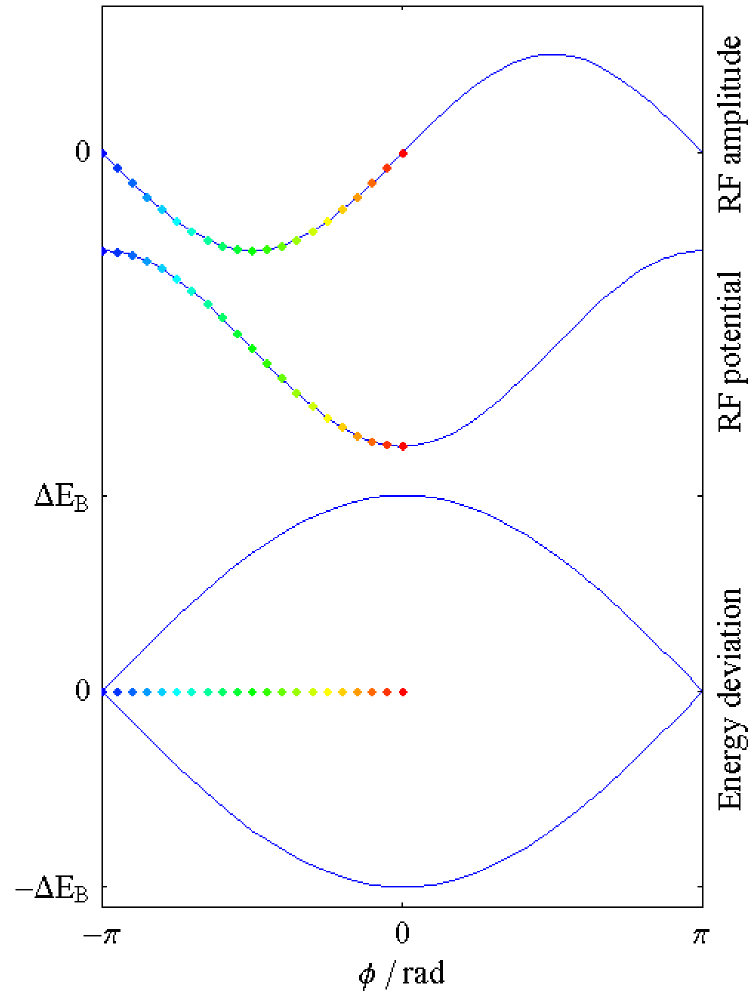


Energy dependent phase advance, ϕ:

$$\phi_{n+1} = \phi_n - 2\pi h\eta/\beta^2 \frac{\Delta E_n}{E_0}, \; \eta = \frac{1}{\gamma^2} - \frac{1}{\gamma_{\mathrm{tr}}^2}$$

Phase dependent energy gain, Δ*E*:

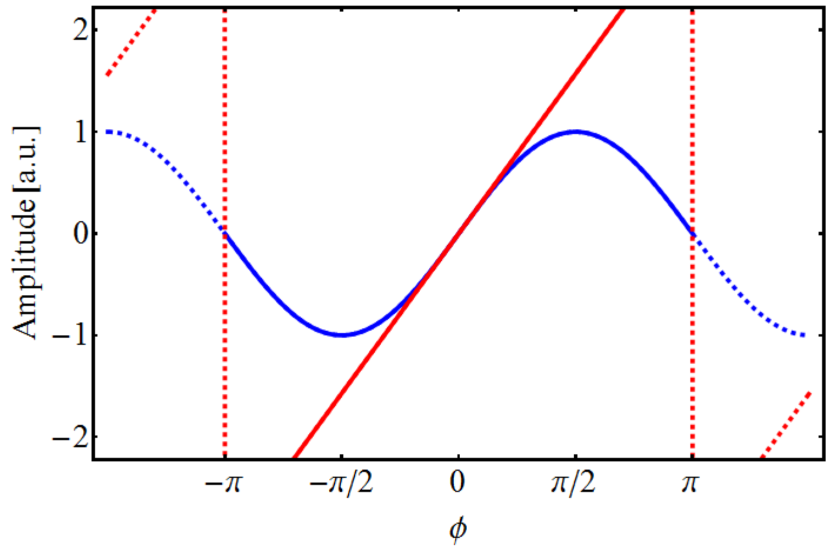
$$\Delta E_{n+1} = \Delta E_n + qVg(\phi_{n+1})$$



Works for arbitrary shape of acceleration amplitude *g*(ϕ)

Linear longitudinal beam dynamics

- **Usual longitudinal beam dynamics already non-linear, since RF system usually provides sinusoidal amplitude**
- **Linear longitudinal beam dynamics?**



$$\frac{d}{dt}\phi = -\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)$$

$$\frac{d}{dt}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = \frac{qV}{2\pi}\phi$$

same structure

$$\frac{dq}{dt} = \frac{\partial H}{\partial p}$$

$$\frac{dp}{dt} = -\frac{\partial H}{\partial q}$$

Linear longitudinal beam dynamics

$$\frac{d}{dt}\phi = -\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)$$

$$\frac{d}{dt}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = \frac{qV}{2\pi}\phi$$

The Hamiltonian from the equations can be written as

$$H\left(\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 - \frac{1}{2}\frac{qV}{2\pi}\phi^2$$

$$= -\frac{1}{2}\frac{pR}{h\eta\omega_{\mathrm{rev}}}\dot{\phi}^2 - \frac{1}{2}\frac{qV}{2\pi}\phi^2$$

$$\eta = \frac{1}{\gamma^2} - \frac{1}{\gamma_{\mathrm{tr}}^2}$$

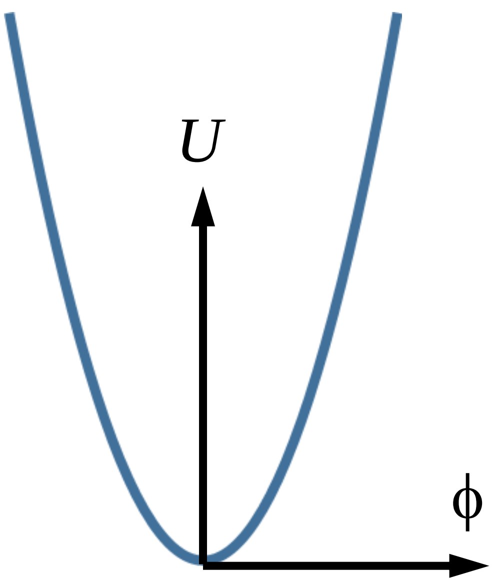
Linear longitudinal beam dynamics

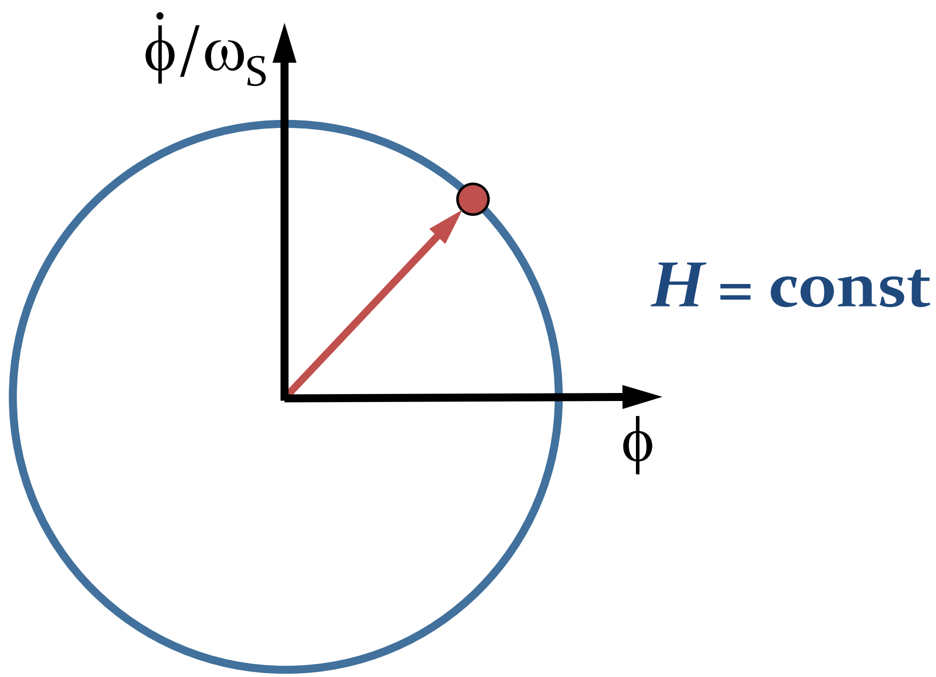
$$H\left(\phi, \frac{\dot{\phi}}{\omega_S}\right) = \left(\frac{\dot{\phi}}{\omega_S}\right)^2 + \phi^2 = T + U$$

→ **Particles move on circular trajectories in ϕ-$\dot{\phi}/\omega_S$ phase space**

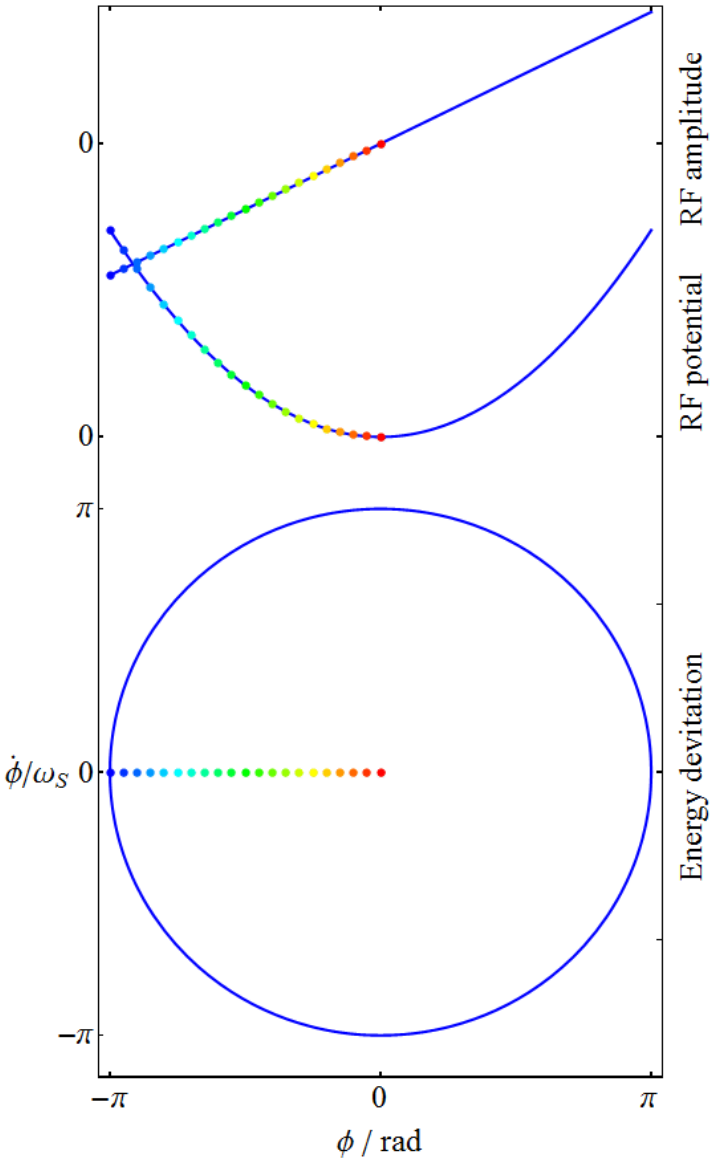
→ **RF potential is parabolic, $U = W(\phi)$**

→ **Hamiltonian is constant on these trajectories**





Linear longitudinal phase space



- **Simple model**
- **Circular trajectories**
- **All particles have same synchrotron frequency**
- **Normalized bucket area:** $A_b = \pi r^2 = \pi^3$

→ **Harmonic oscillator**

Non-linear longitudinal beam dynamics

Introduce most simple non-linearity

RF amplitude function $V\phi \rightarrow V\sin\phi$

$$\frac{d}{dt}\Delta\phi = -\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)$$

$$\frac{d}{dt}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = \frac{qV}{2\pi}(\sin\phi - \sin\phi_0)$$

$$H\left(\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 + \frac{qV}{2\pi}\left[\cos\phi - \cos\phi_0 + (\phi - \phi_0)\sin\phi_0\right]$$

with $\phi = \phi_0 + \Delta\phi$ **this becomes**

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 + \frac{qV}{2\pi}\left[\cos(\phi_0 + \Delta\phi) - \cos\phi_0 + \Delta\phi\sin\phi_0\right]$$

→ Standard longitudinal beam dynamics → Lectures F. Tecker

Introduce most simple non-linearity

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 + \frac{qV}{2\pi}\left[\cos(\phi_0 + \Delta\phi) - \cos\phi_0 + \Delta\phi\sin\phi_0\right]$$

using

$$\begin{aligned}\cos(\phi_0 + \Delta\phi) &= \cos\phi_0\cos\Delta\phi - \sin\phi_0\sin\Delta\phi \\ &\simeq \cos\phi_0\left(1 - \frac{1}{2}\Delta\phi^2\right) - \sin\phi_0\Delta\phi\end{aligned}$$

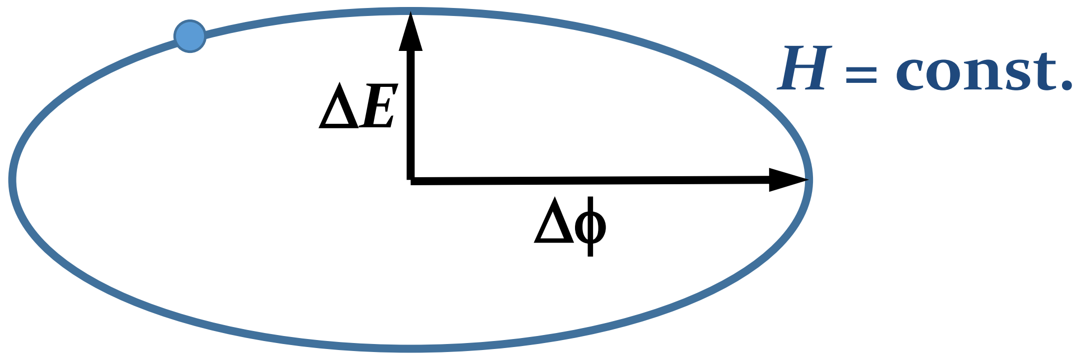
this Hamiltonian simplifies to

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) \simeq -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 - \frac{1}{2}\frac{qV}{2\pi}\Delta\phi^2\cos\phi_0$$

Linear part of non-linear bucket

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) \simeq -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 - \frac{1}{2}\frac{qV}{2\pi}\Delta\phi^2\cos\phi_0$$

- **In the centre of the bucket, particles move on elliptical trajectories in Δϕ-Δ*E* phase space**
- **Hamiltonian is constant on these trajectories**



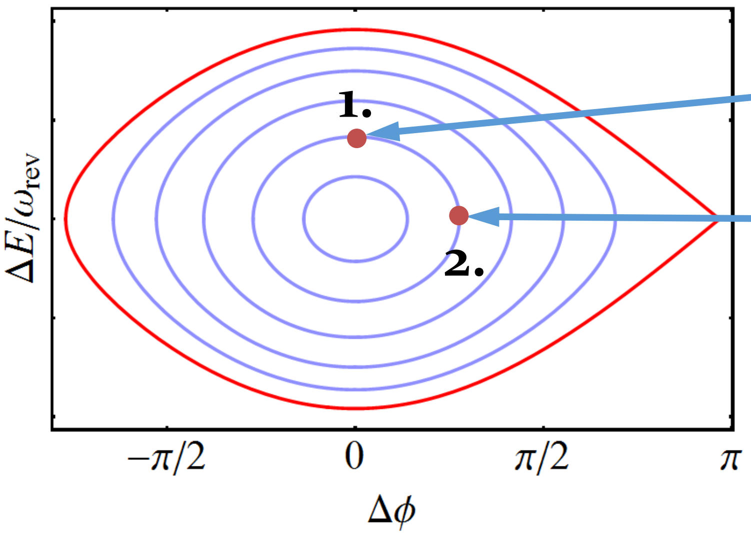
- **In the bucket centre, particles oscillate with the synchrotron frequency, $\omega_S = 2\pi f_S$**

$$\omega_S^2 = \frac{h\eta\omega_{\mathrm{rev}} qV\cos\phi_0}{2\pi pR}$$

$$\eta = \frac{1}{\gamma^2} - \frac{1}{\gamma_{\mathrm{tr}}^2}$$

Longitudinal emittance

- **Compare two particles on the same trajectory**
  1. **No phase deviation**
  2. **No energy deviation**



$$H\left(\Delta\phi = 0, \frac{\Delta E}{\omega_{\rm rev}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\rm rev}}{pR}\left(\frac{\Delta E}{\omega_{\rm rev}}\right)^2$$

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\rm rev}} = 0\right) = -\frac{1}{2}\frac{qV}{2\pi}\Delta\phi^2\cos\phi_0$$

Longitudinal emittance

- **Compare two particles on the same trajectory**

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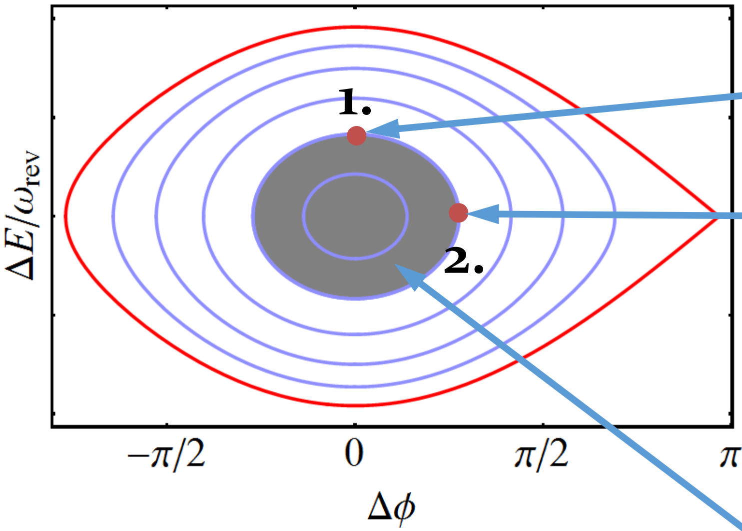


$$H\left(\Delta\phi = 0, \frac{\Delta E}{\omega_{\text{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\text{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\text{rev}}}\right)^2$$

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\text{rev}}} = 0\right) = -\frac{1}{2}\frac{qV}{2\pi}\Delta\phi^2 \cos\phi_0$$

$$\varepsilon_l = \frac{2}{h\omega_{\text{rev}}}\int_{\Delta\phi_i}^{\Delta\phi_f} \Delta E(\Delta\phi)\, d(\Delta\phi)$$

Longitudinal emittance, ε_l

~ **Surface occupied by particles in longitudinal phase space**

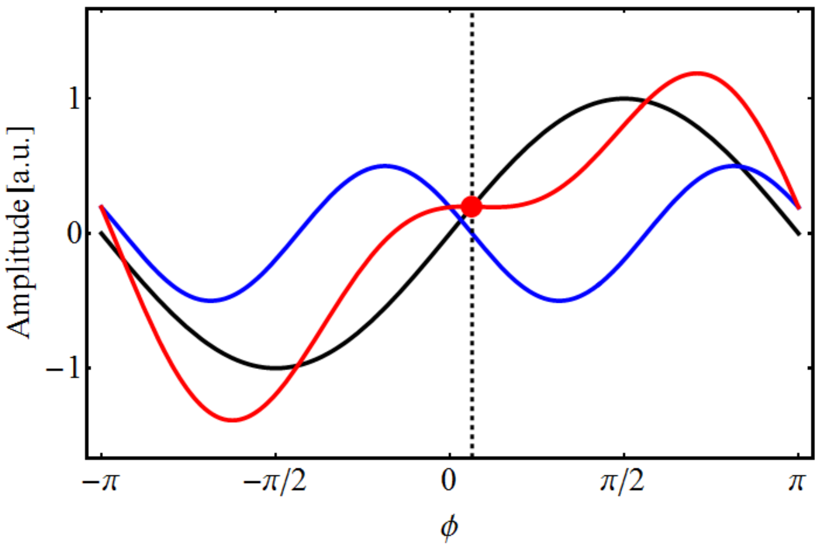
→ **Preserved in physical $[\pi\Delta\tau\,\Delta E]$ = eVs**

More non-linearity: multi-harmonic RF

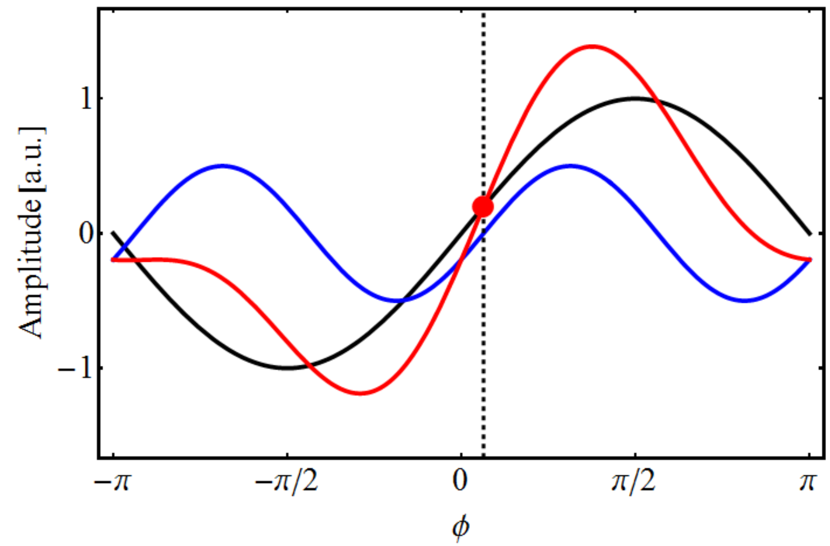
RF amplitude $V \sin\phi \rightarrow V\left[\sin\phi + r\sin(n\phi + \phi_1)\right]$

- **Example case $n = 2$ and $r = 0.5$**

Both RF systems in phase at bunch



Both RF systems in counter-phase



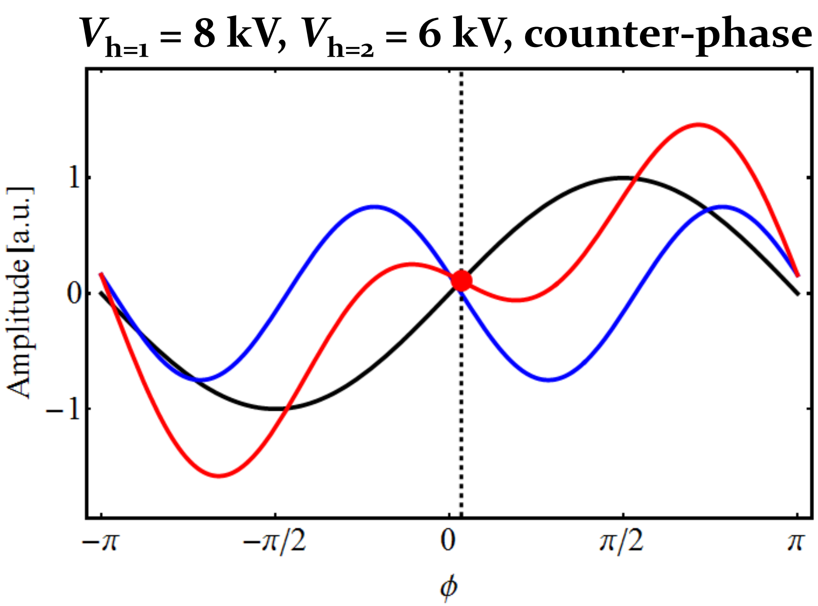
→ **Local voltage gradient decreased**
→ **Bunch is stretched**
→ **Lower peak current**

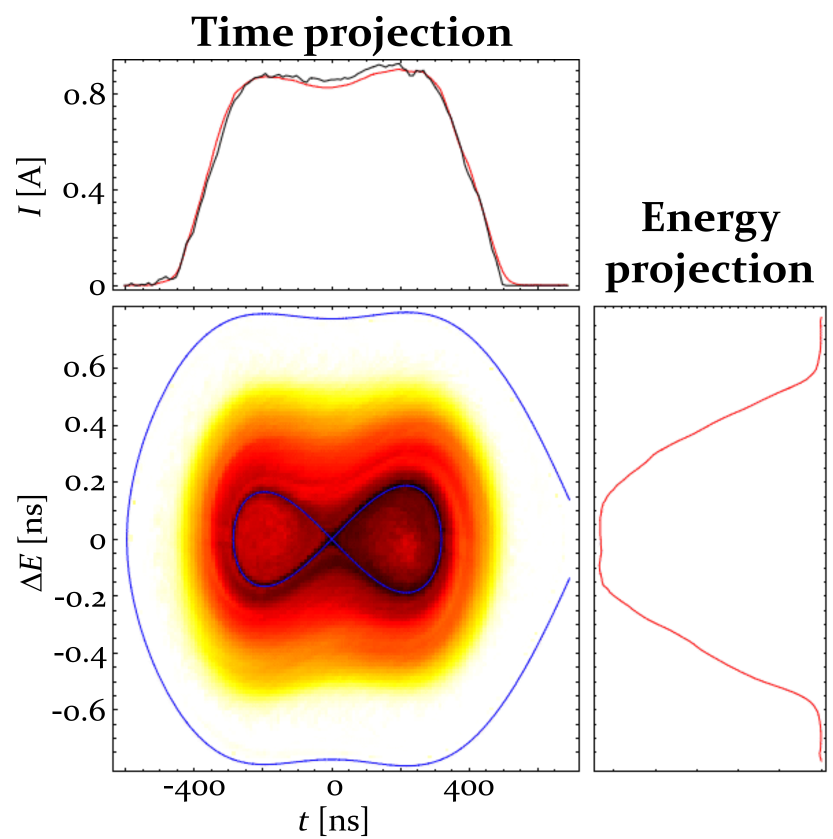
→ **Local voltage gradient increased**
→ **Bunch is compressed**
→ **Higher peak current**

Example application: space charge in PSB

RF amplitude $V \sin\phi \rightarrow V\left[\sin\phi + r\sin(n\phi + \phi_1)\right]$

→ Space charge ∝ instantaneous current





- **Inverted gradient at bucket centre**
- **Flattened bunch with reduced peak current** → Space charge reduction at low energy

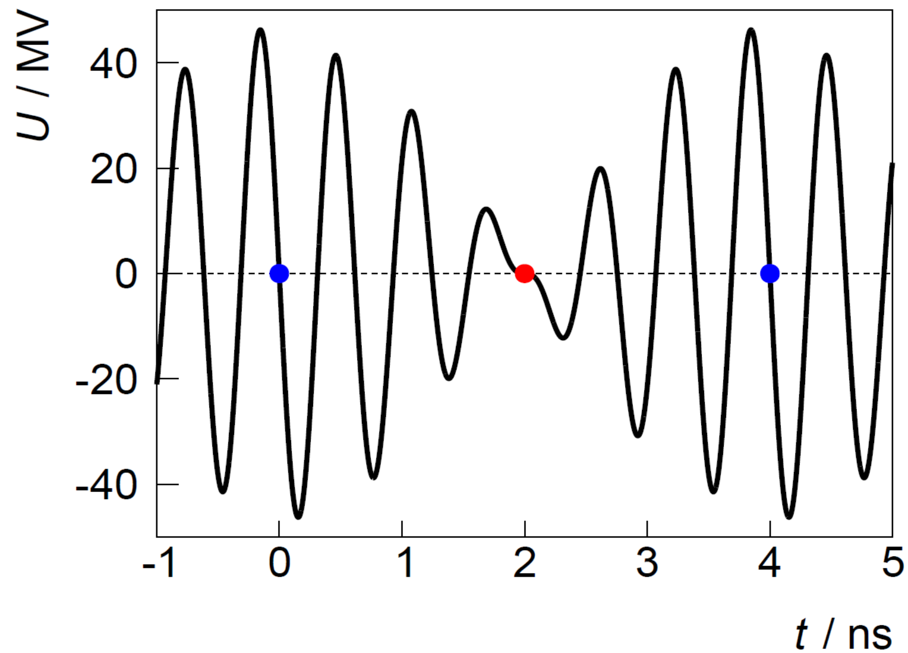
Long and short bunches simultaneously

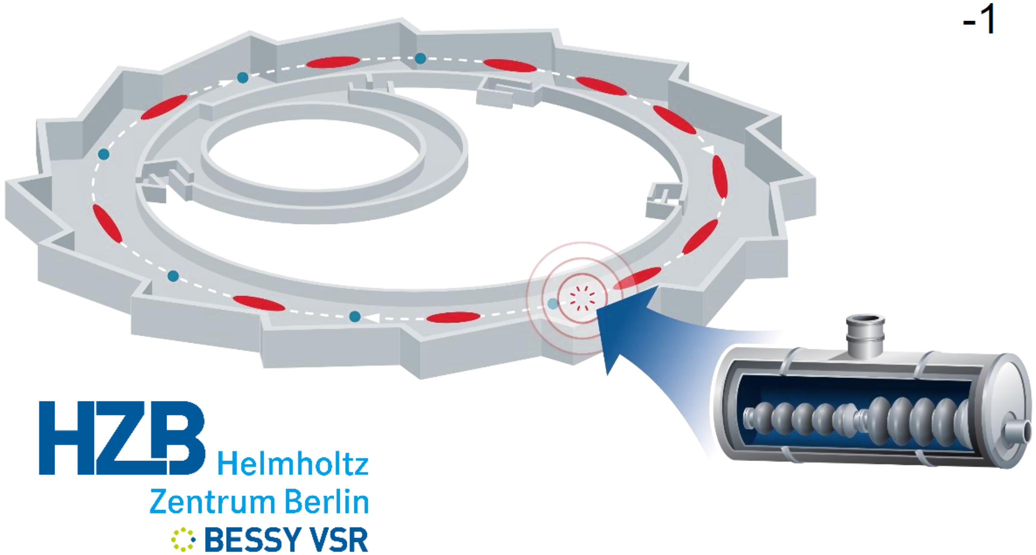
Markus Ries et al.

- Example BESSY VSR

→ Depending on user of synchrotron radiation: need long or short bunches

Do long and short bunches simultaneously!





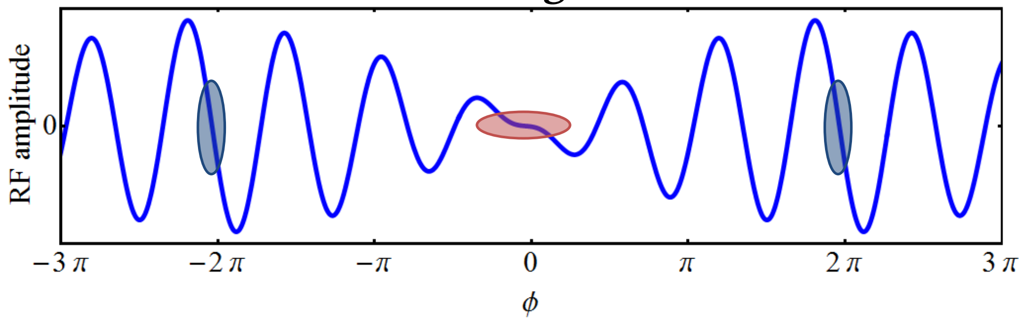
- 4 × 0.5 GHz NC (existing)
- 4 × 1.5 GHz supercond.
- 4 × 1.75 GHz supercond.

Bunch length modulation

Markus Ries et al.

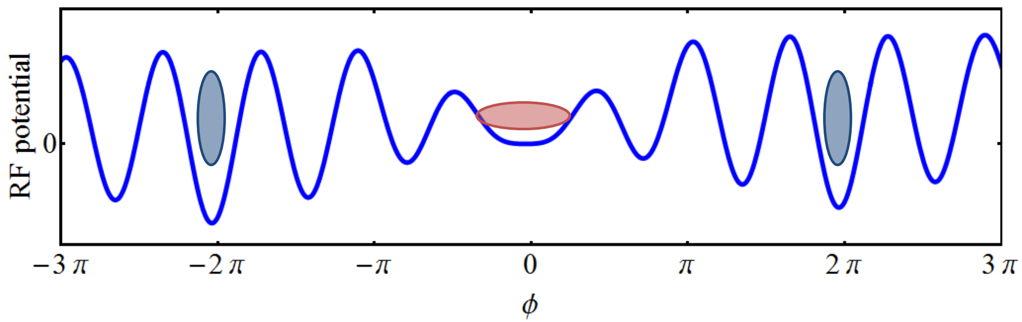
- **Future 3-harmonic RF system for BESSY VSR**

RF voltage sum



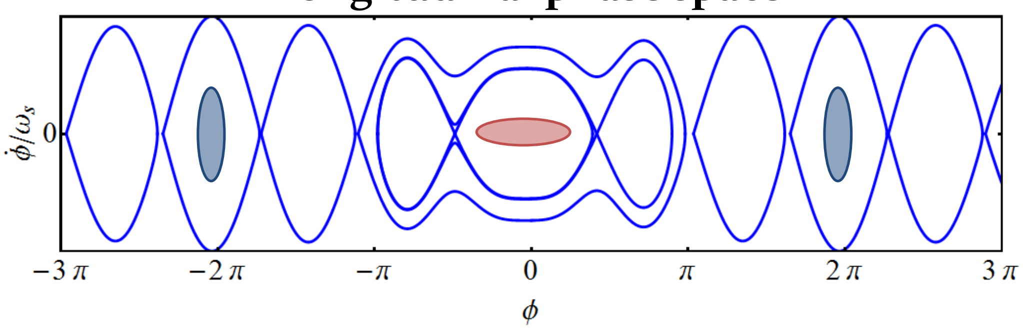
RF voltages sum up or cancel

RF Potential



Short deep or bathtub potential well

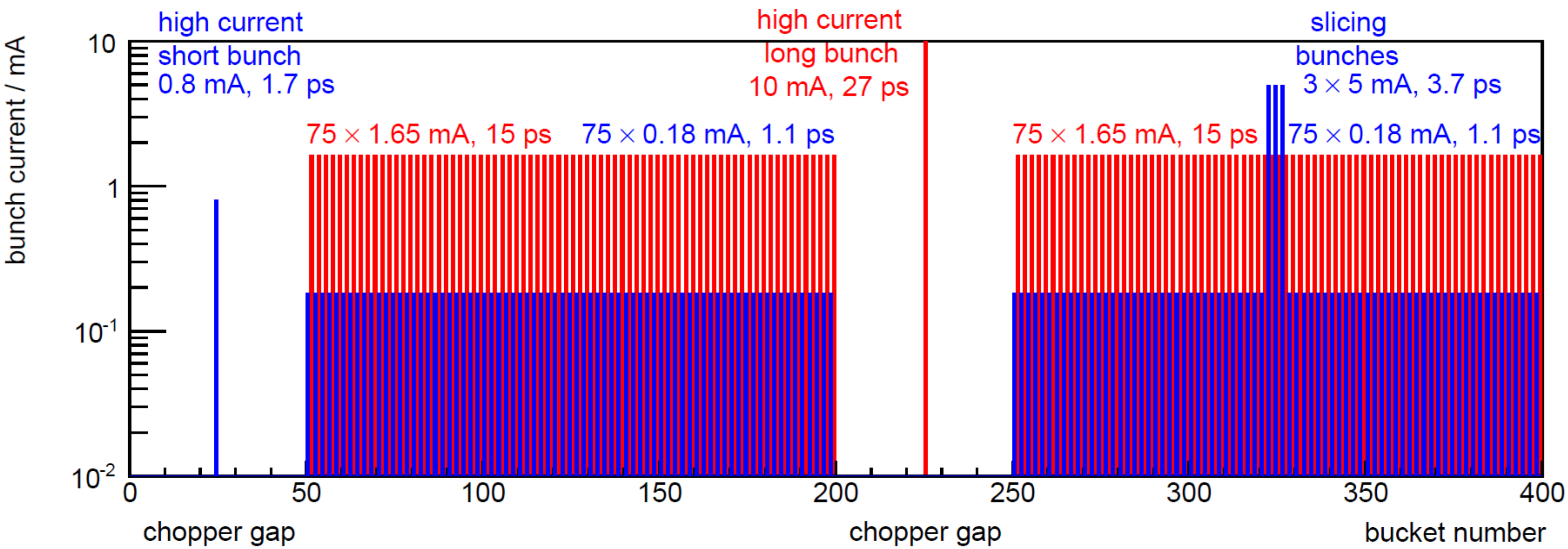
Longitudinal phase space



Short or stretched bucket

Filling pattern

Markus Ries et al.



- 300 mA average current

→ High-current single bunches

→ short (0.8 mA) & long (10 mA)

→ Special high-current density bunches

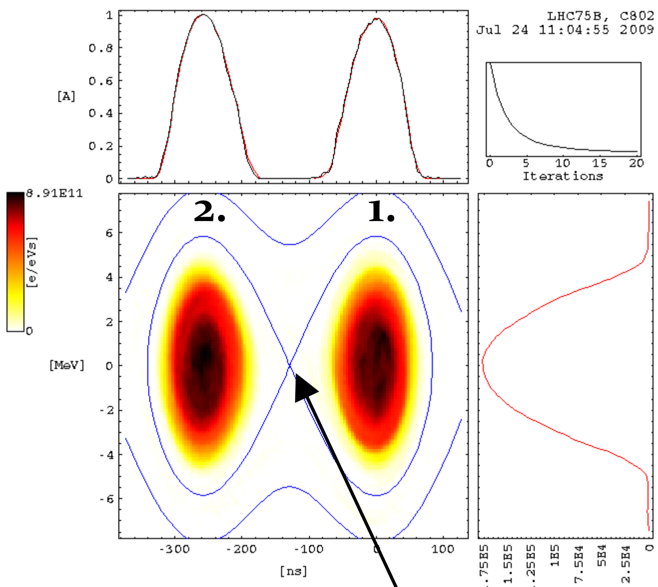
☝ Two electron storage ring in one

☝ Thanks to longitudinal beam dynamics trick

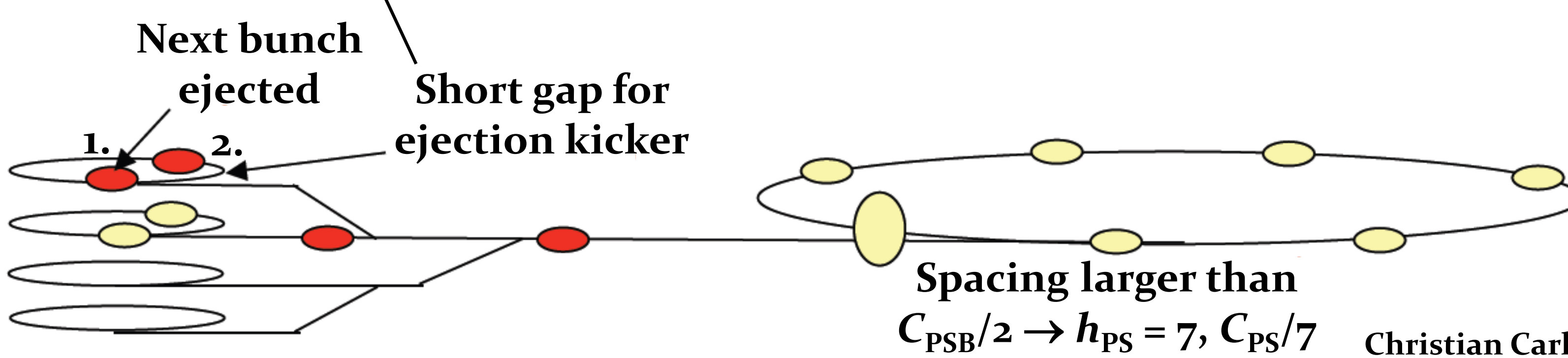
Example: adjust bunch spacing

- Was used at CERN PSB-to-PS to transfer 2 bunches at once
- Circumference ratio $C_{PS}/C_{PSB} = 4$

→ Ratio virtually moved to 2/7: use $h_{RF} = 2 + 1$



1. Add $h1$ component such that bunches approach to 245 ns (small spacing) → big spacing becomes 327 ns
2. Synchronize on $h1$ to the PS
3. Trigger extraction kicker in-between the small spacing
4. Eject two bunches per ring at a distance of 327 ns



Christian Carli

Introduce general non-linearity

Replace $V\sin\phi \rightarrow V g(\phi)$ $\rightarrow$ **arbitrary amplitude**

Equations of motion

$$\frac{d\phi}{dt} = -\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)$$

$$\frac{d}{dt}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = \frac{qV}{2\pi}[g(\phi) - g(\phi_0)]$$

same structure

$$\frac{dq}{dt} = \frac{\partial H}{\partial p}$$

$$\frac{dp}{dt} = -\frac{\partial H}{\partial q}$$

The Hamiltonian describing the system becomes

$$H\left(\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2 + \frac{qV}{2\pi}\left[g(\phi_0)\phi - \int g(\phi)\,d\phi\right]$$

$$= \quad \text{kinetic} \quad + \quad \text{potential terms}$$

$$\eta = \frac{1}{\gamma^2} - \frac{1}{\gamma_{\mathrm{tr}}^2}$$

Arbitrary RF waveform

$$H\left(\phi, \frac{\Delta E}{\omega_{\text{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\text{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\text{rev}}}\right)^2 + \frac{qV}{2\pi}\left[g(\phi_0)\phi - \int g(\phi)\,d\phi\right]$$

Using $$\dot{\phi} = -\frac{h\eta\omega_{\text{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\text{rev}}}\right)$$

The Hamiltonian can be rewritten, with RF potential *W*(ϕ)

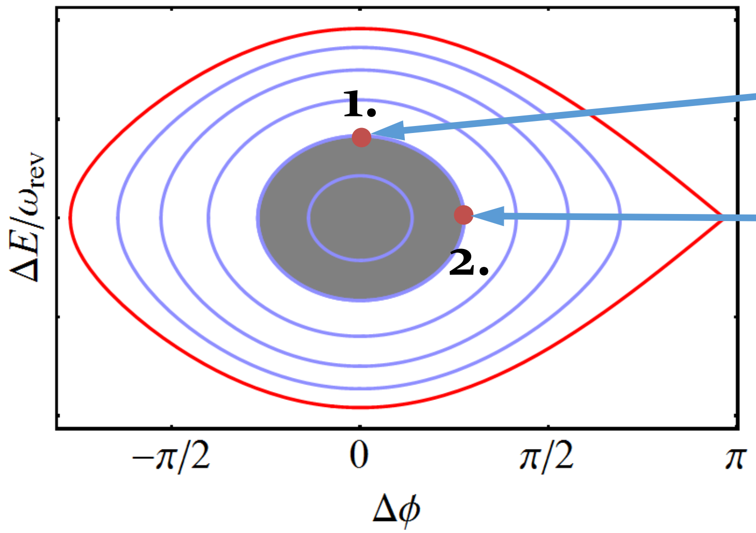
$$H(\phi, \dot{\phi}) = \frac{1}{2}\left(\frac{\dot{\phi}}{\omega_S}\right)^2 + W(\phi)$$

$$W(\phi) = \frac{1}{\cos\phi_0}\left[\int g(\phi)\,d\phi - g(\phi_0)\phi\right]$$

Longitudinal beam manipulations using non-linearity

Change RF voltage to change bunch length?

→ **Calculate aspect ratio of bucket trajectories**



$$H\left(\Delta\phi = 0, \frac{\Delta E}{\omega_{\mathrm{rev}}}\right) = -\frac{1}{2}\frac{h\eta\omega_{\mathrm{rev}}}{pR}\left(\frac{\Delta E}{\omega_{\mathrm{rev}}}\right)^2$$

$$H\left(\Delta\phi, \frac{\Delta E}{\omega_{\mathrm{rev}}} = 0\right) = -\frac{1}{2}\frac{qV}{2\pi}\Delta\phi^2 \cos\phi_0$$

Equating both sides gives

$$\left(\frac{\Delta E}{\Delta\tau}\right)^2 = \frac{qV}{2\pi}E\beta^2 h\omega_{\mathrm{rev}}^2 \frac{\cos\phi_0}{\eta}$$

with emittance as $\varepsilon_l = \pi\Delta\tau\Delta E = \mathrm{const.}$ $\rightarrow$ $\Delta\tau \propto \frac{1}{\sqrt[4]{V}}$

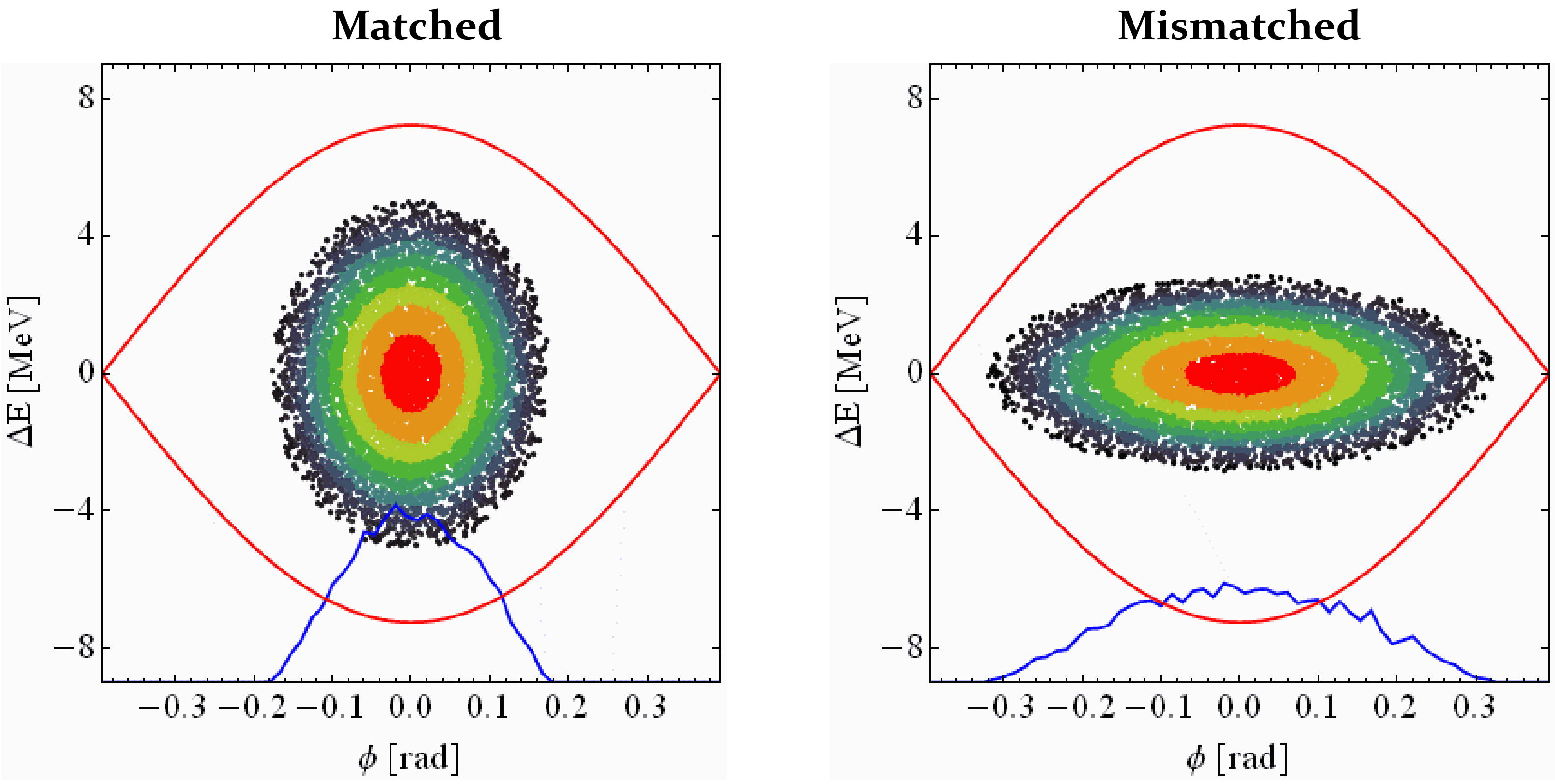
→ **Not efficient at all**

→ **16 times more RF voltage needed to cut bunch length in half**

Abrupt change of RF voltage

→ Individual particles in matched bunch oscillate but no macroscopic motion

→ Abruptly changing the RF voltage flips particles to new trajectories



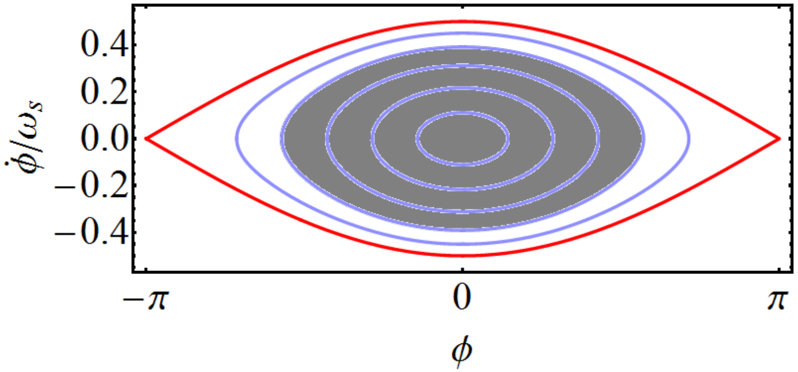
→ The bunch distribution seems to rotate

→ Exchange of bunch length and momentum spread

Introduce sudden change: bunch rotation

→ **Quickly exchange longitudinal phase space behind bunch**

→ **Increase RF voltage much faster than period of f_S**



Introduce sudden change: bunch rotation

→ **Quickly exchange longitudinal phase space behind bunch**
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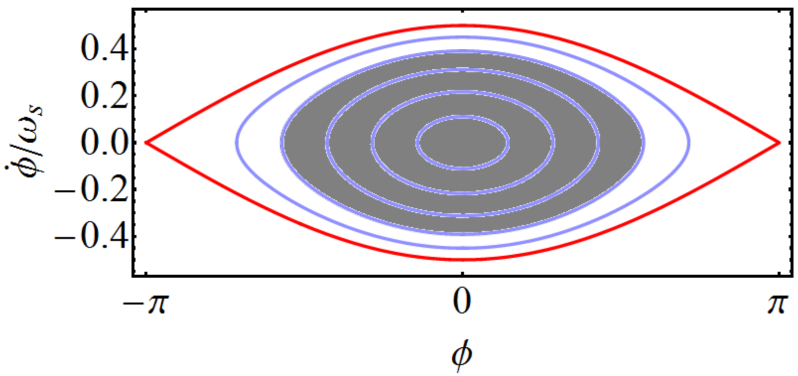
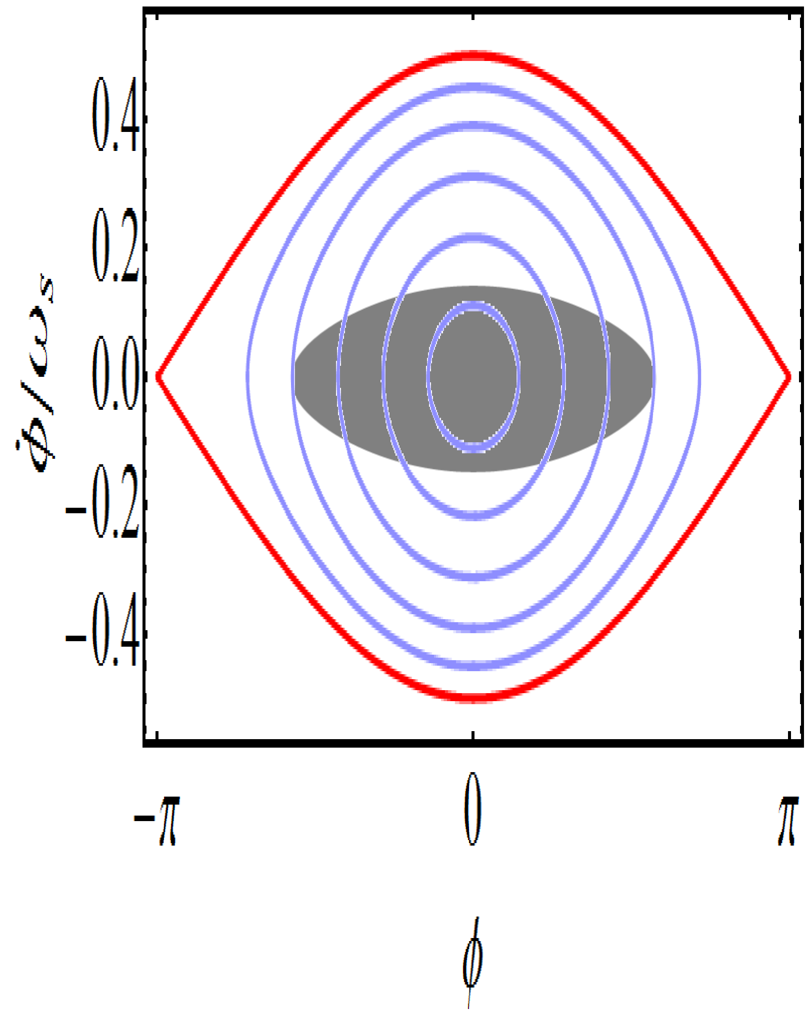
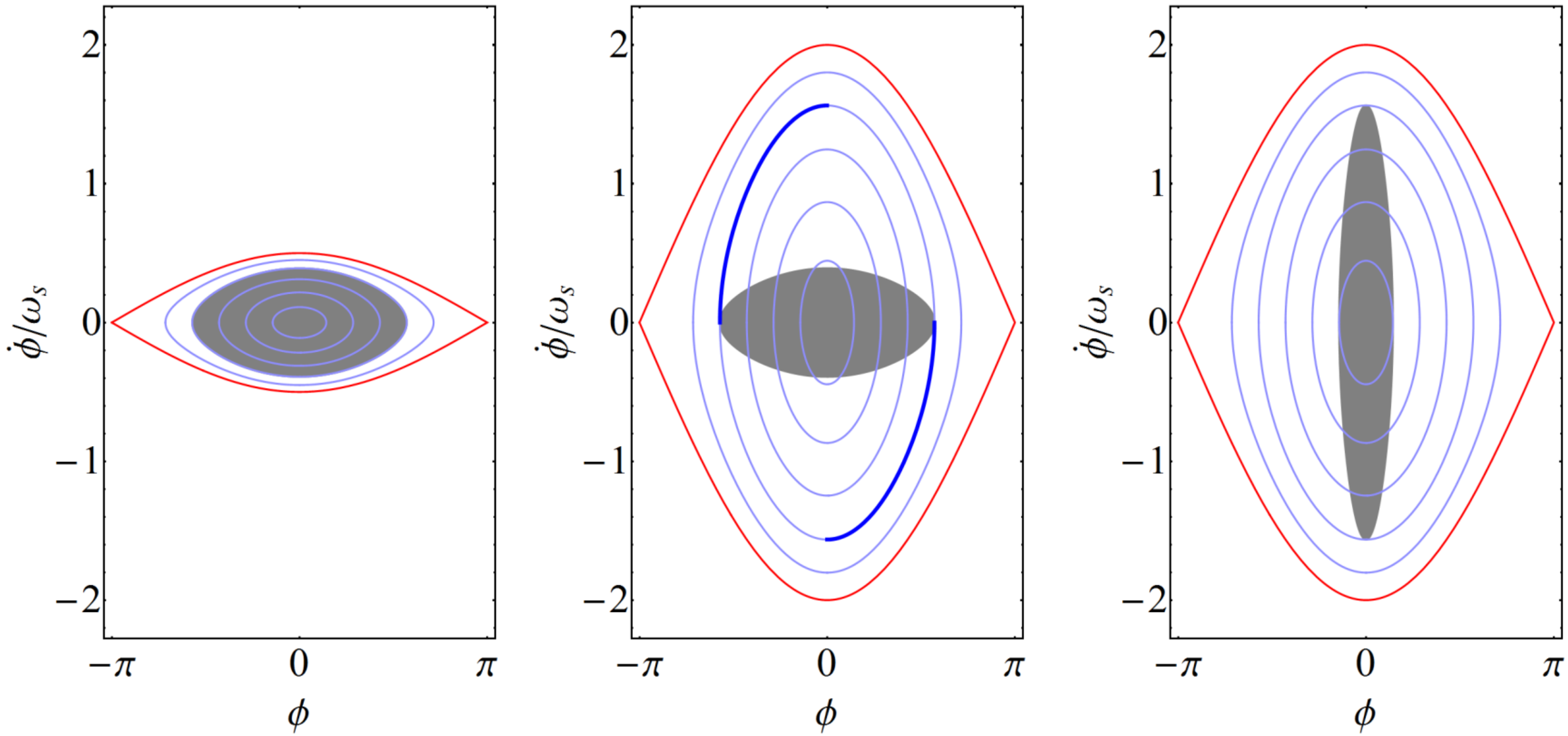




Introduce sudden change: bunch rotation

→ **Switch RF voltage much faster than period of f_s**

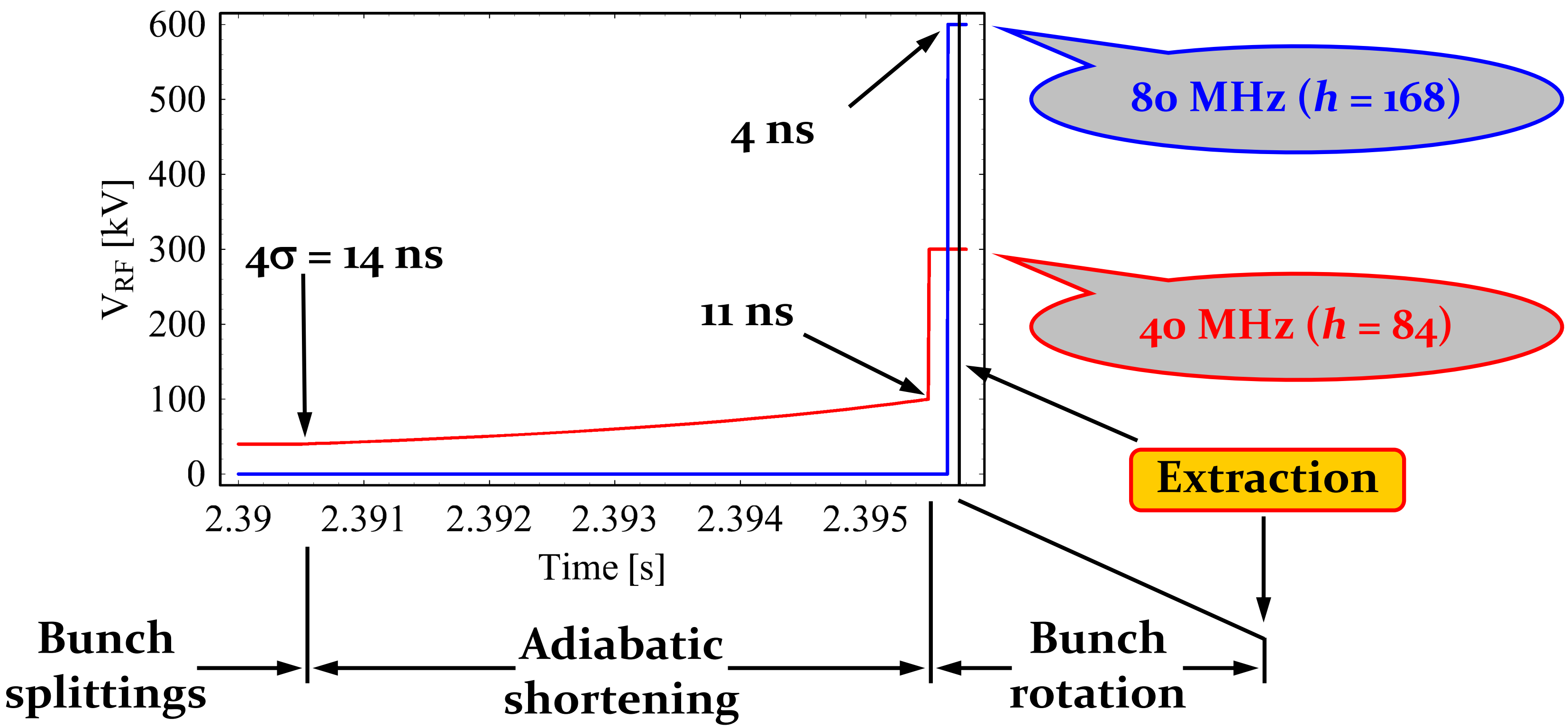


$$V_i \propto \left(\frac{\Delta E_i}{\Delta \tau_i}\right)^2 \qquad V_f \propto \left(\frac{\Delta E_f}{\Delta \tau_i}\right)^2 \qquad \frac{\Delta \tau_f}{\Delta \tau_i} = \frac{\Delta E_i}{\Delta E_f} = \sqrt{\frac{V_i}{V_f}}$$

Example: PS to SPS transfer at CERN

- **Fit 14 ns long bunches into 5 ns long buckets in the SPS**

→ Double-step bunch rotation



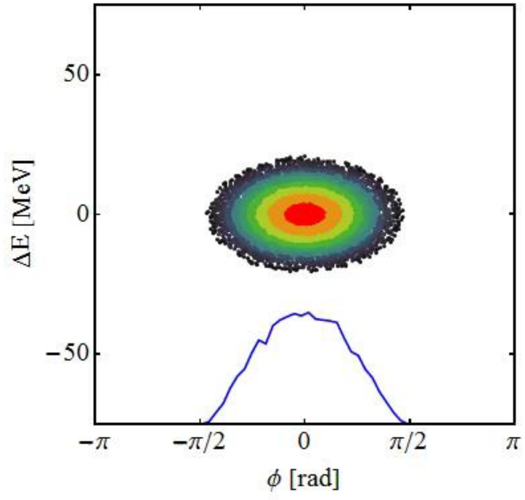
Example: rotation at PS-SPS transfer

→ **Bunch length now proportional to $\sqrt{V}$ and not $\sqrt[4]{V}$**
→ **Can save enormous RF voltage**

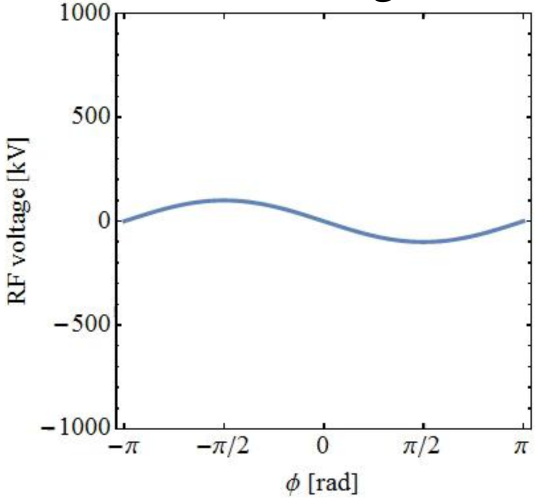
→ **Bunch shortening from 14 ns to 4 ns (ratio ~3.5)**
→ **Starting from 100 kV at 40 MHz**

→ **Slow shortening would require $100\ \mathrm{kV} \cdot 3.5^4 \sim 15\ \mathrm{MV}$**
→ **Installed RF voltage is only about 1.2 MV**

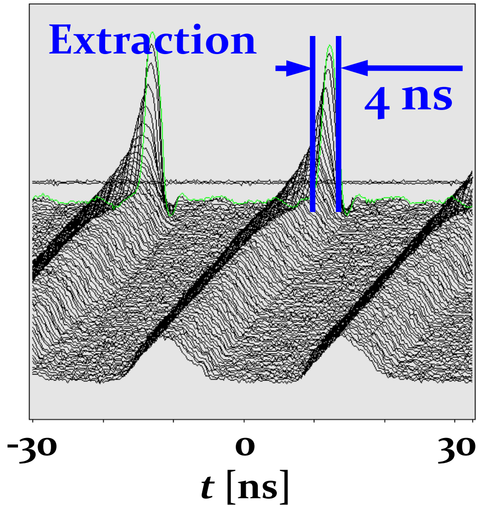
Simulation



RF voltage



Measurement

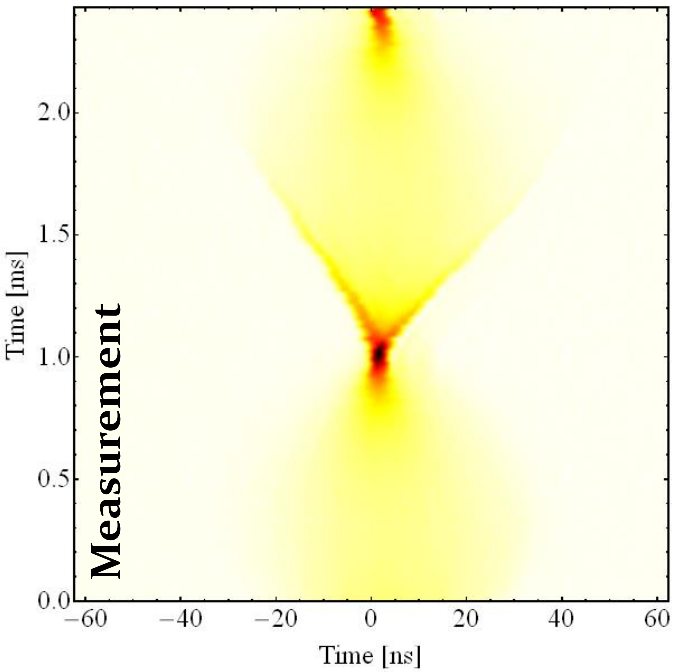


Profiting from the non-linear rotation

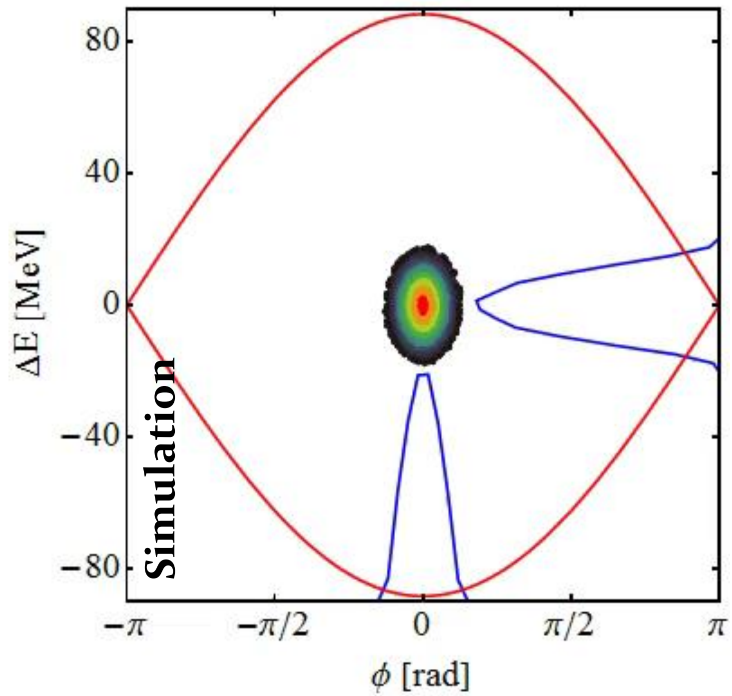
Need large momentum spread for slow extraction

1. **Jump RF phase such that bunch at unstable fixed point**
2. **Jump back**
3. **Let bunch rotate, switch RF off at large momentum spread**

Bunch profile evolution



Longitudinal phase space

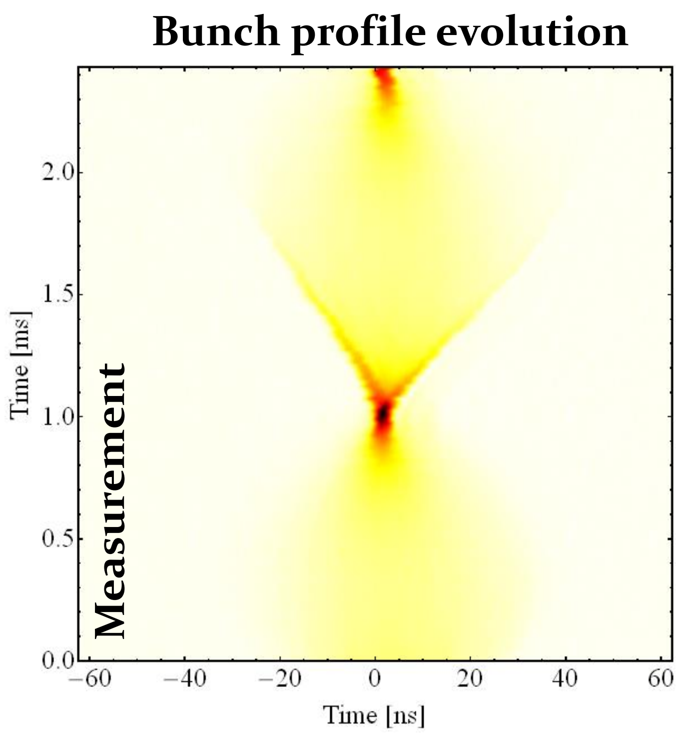


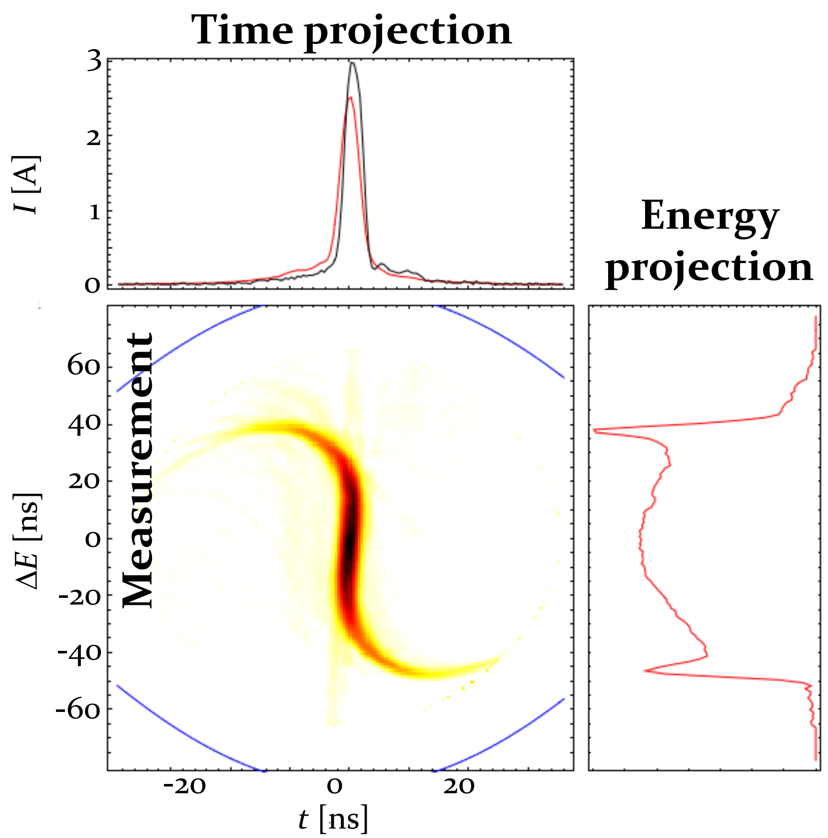
→ **Non-linearly of bunch rotation helps**

Example: using the non-linearity

Need large momentum spread for slow extraction

1. **Jump RF phase such that bunch at unstable fixed point**
2. **Jump back**
3. **Let bunch rotate, switch RF off at large momentum spread**





→ **Almost constant momentum distribution after rotation**

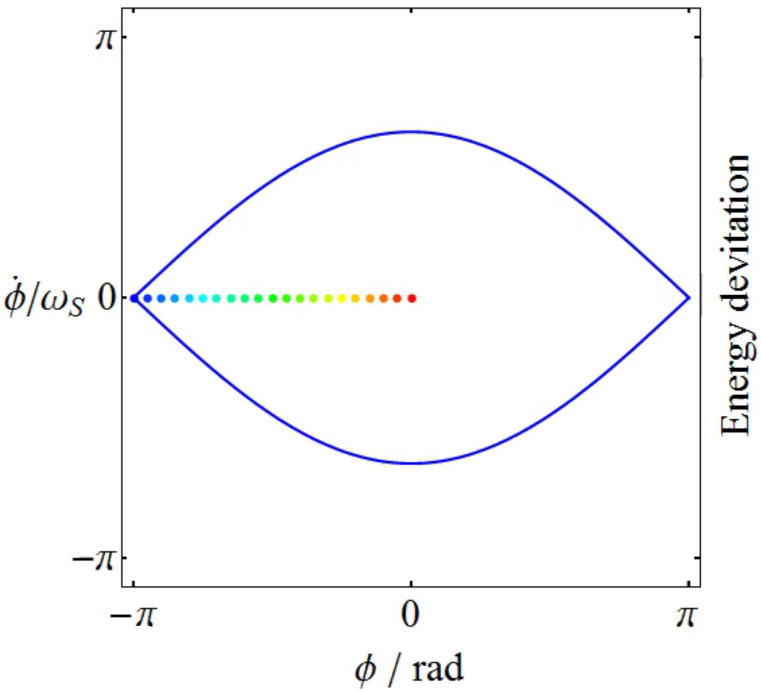
Synchrotron frequency distribution

General synchrotron frequency

- **Synchrotron frequency depends on trajectory**

→ Calculate average velocity for given trajectories in longitudinal phase space → Action angle, *J*

$$J(H) = \frac{1}{2\pi\omega_S} \oint \dot{\phi}(\phi)\, d\phi$$



The angular frequency becomes

$$\omega(H) = \frac{d}{dJ} H$$

General expression for ω_S

$$\frac{\omega(H)}{\omega_S} = \frac{\sqrt{2}\pi}{\int_{\phi_l}^{\phi_u} \frac{1}{\sqrt{H/\omega_S^2 - W(\phi)}}\, d\phi}$$

(for bucket boundaries $\phi_l \rightarrow \phi_u$)

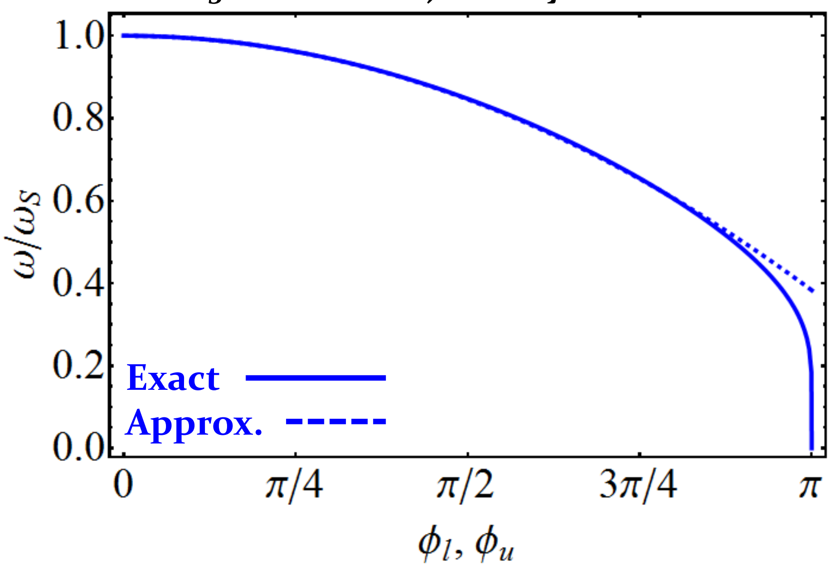
Distribution for stationary bucket

- **Single-harmonic RF in stationary bucket**

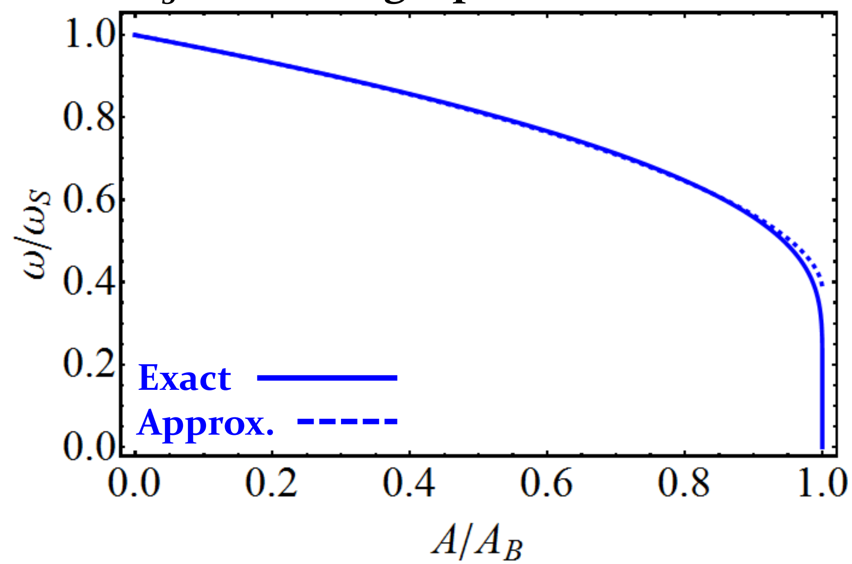
$$\frac{\omega(\Delta\phi_u)}{\omega_S} = \frac{\pi}{2K[\sin(\phi_u/2)]} \simeq 1 - \frac{\phi_u^2}{16}$$

***K*(*x*): 1st kind elliptical integral function**

ω_S versus trajectory duration



ω_S versus single particle emittance



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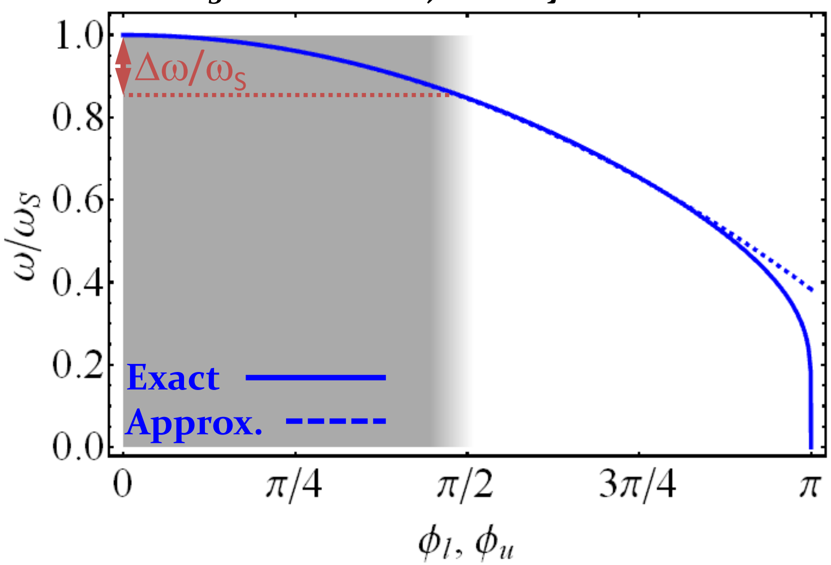
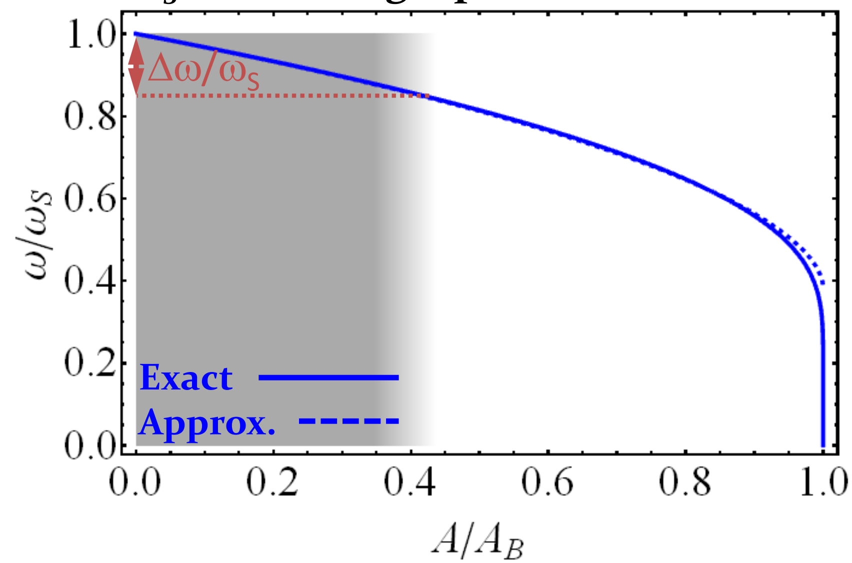


ω_S versus single particle emittance



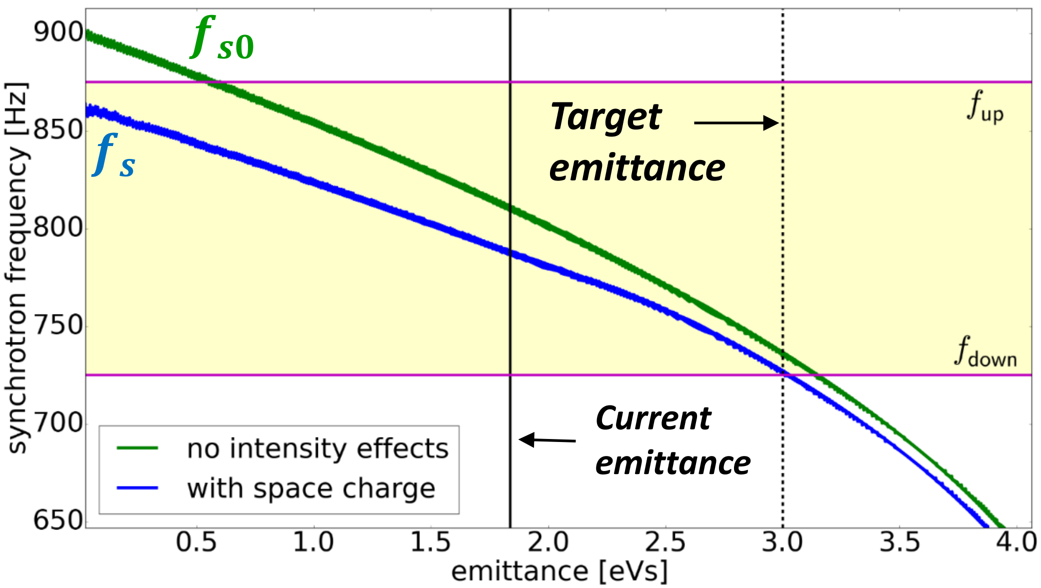
→ Different synchrotron frequencies of particles in bunch

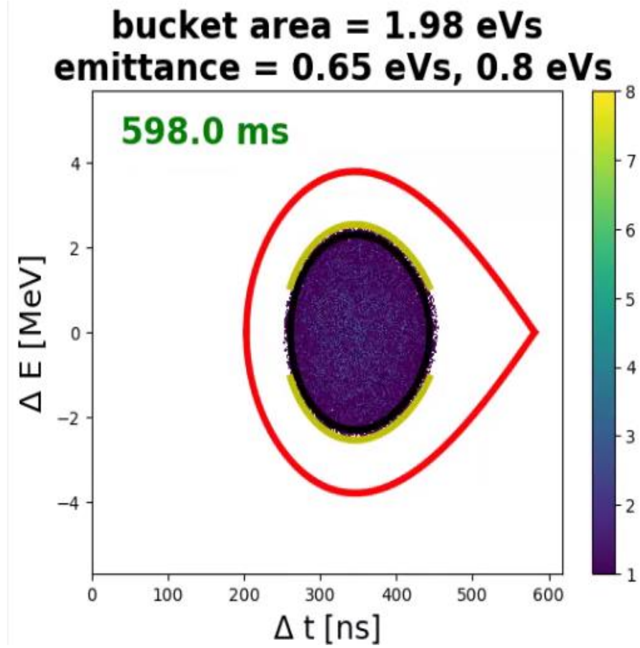
→ Total spread $\Delta\omega/\omega_S$ depends on filling factor of bucket

Example: Emittance control with noise

- **Noise excitation of bunch by band-width limited noise**

→ Controlled longitudinal blow-up in the PSB



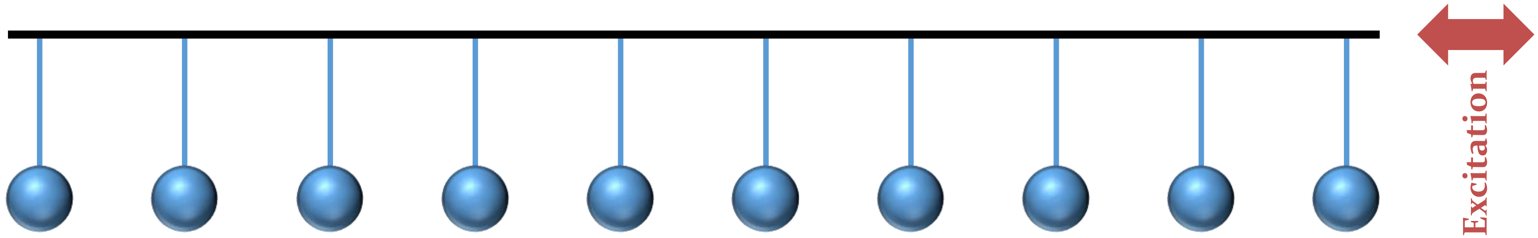


1. **Choose upper frequency to cover synchrotron frequency at bunch centre**
2. **Choose lower frequency to match target emittance**
3. **Excite**

D. Quartullo

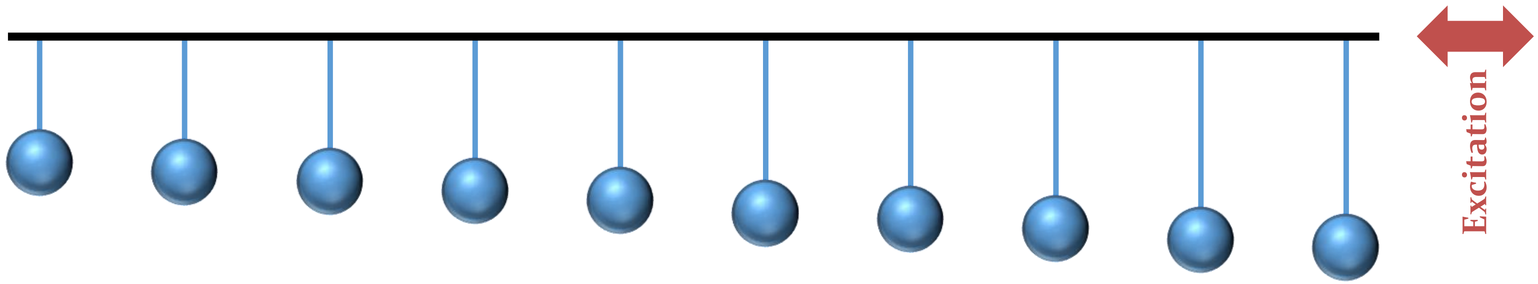
Analogy: pendulums mounted on a bar

- All particles have the same resonance frequency



→ Easy to excite macroscopic oscillation

- Resonance frequencies of individual particles varies

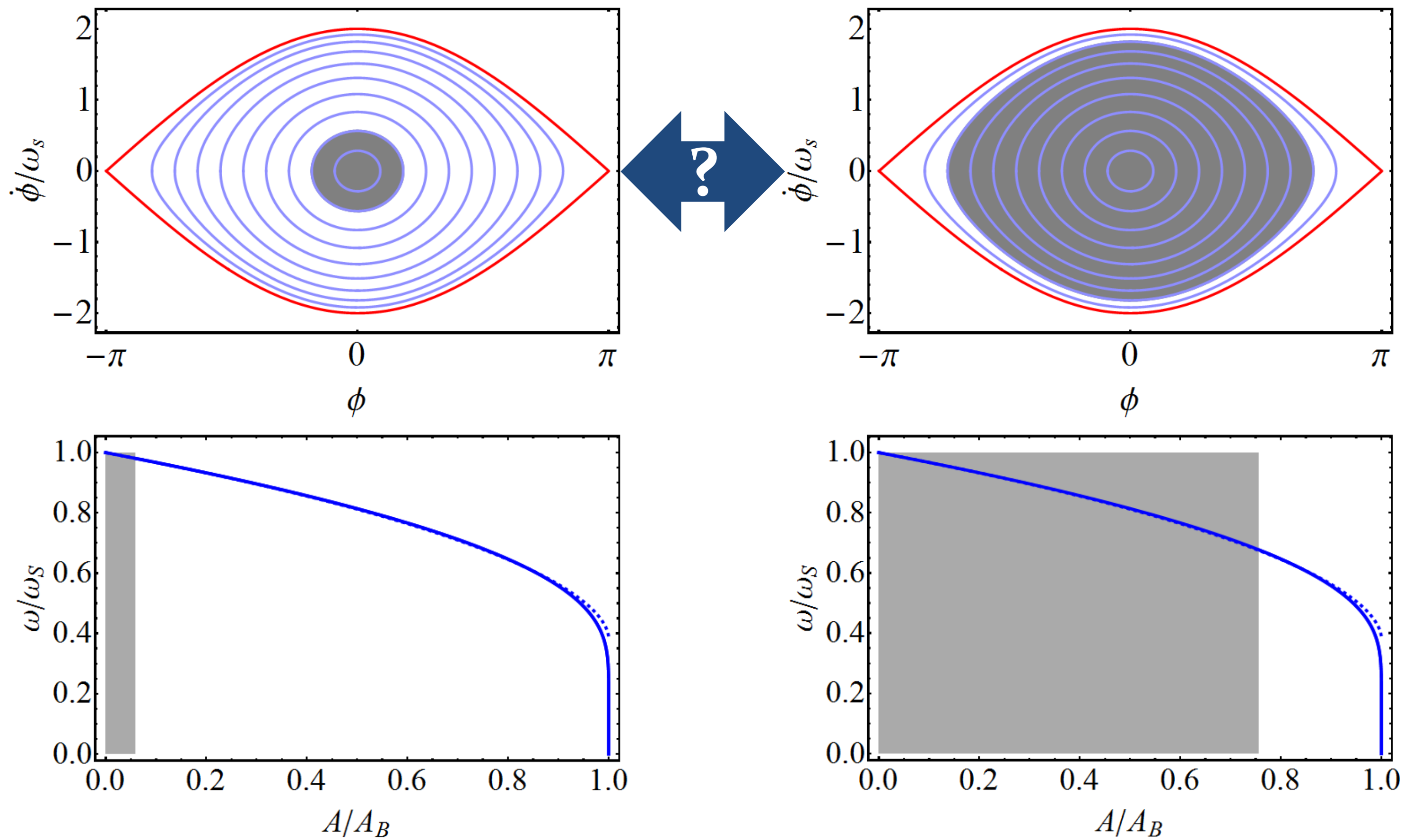


→ Difficult to excite macroscopic oscillation

→ Large synchrotron frequency spread increases stability

Bucket filling ratio

Smaller or larger bunch or bucket? What is more stable?



- Easy to excite

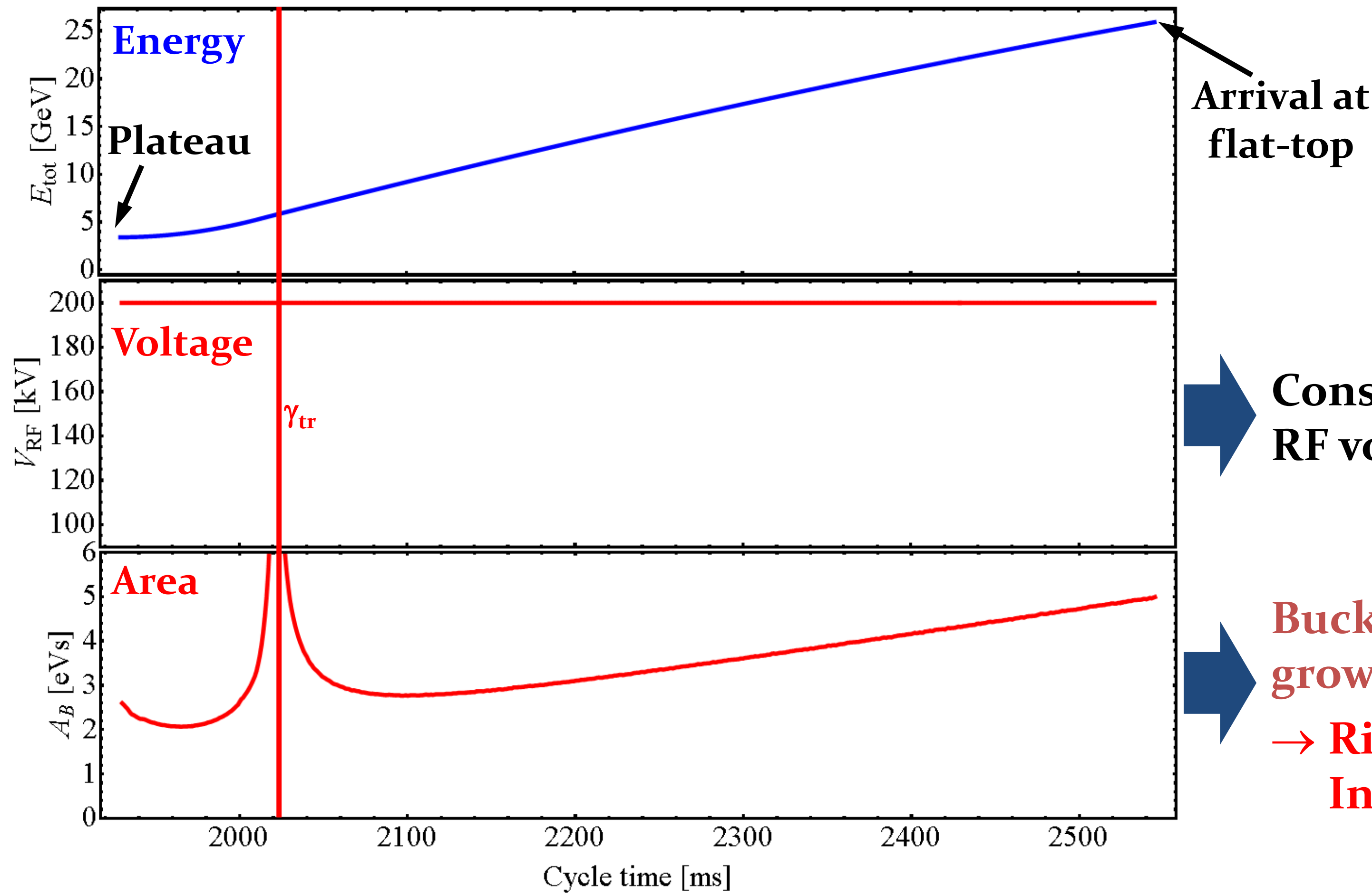
→ **Prone to instability**

- Large f_s spread

→ More stable

Example: stabilization with lower voltage

→ **Acceleration of protons in the CERN PS (3.4 → 26 GeV total)**



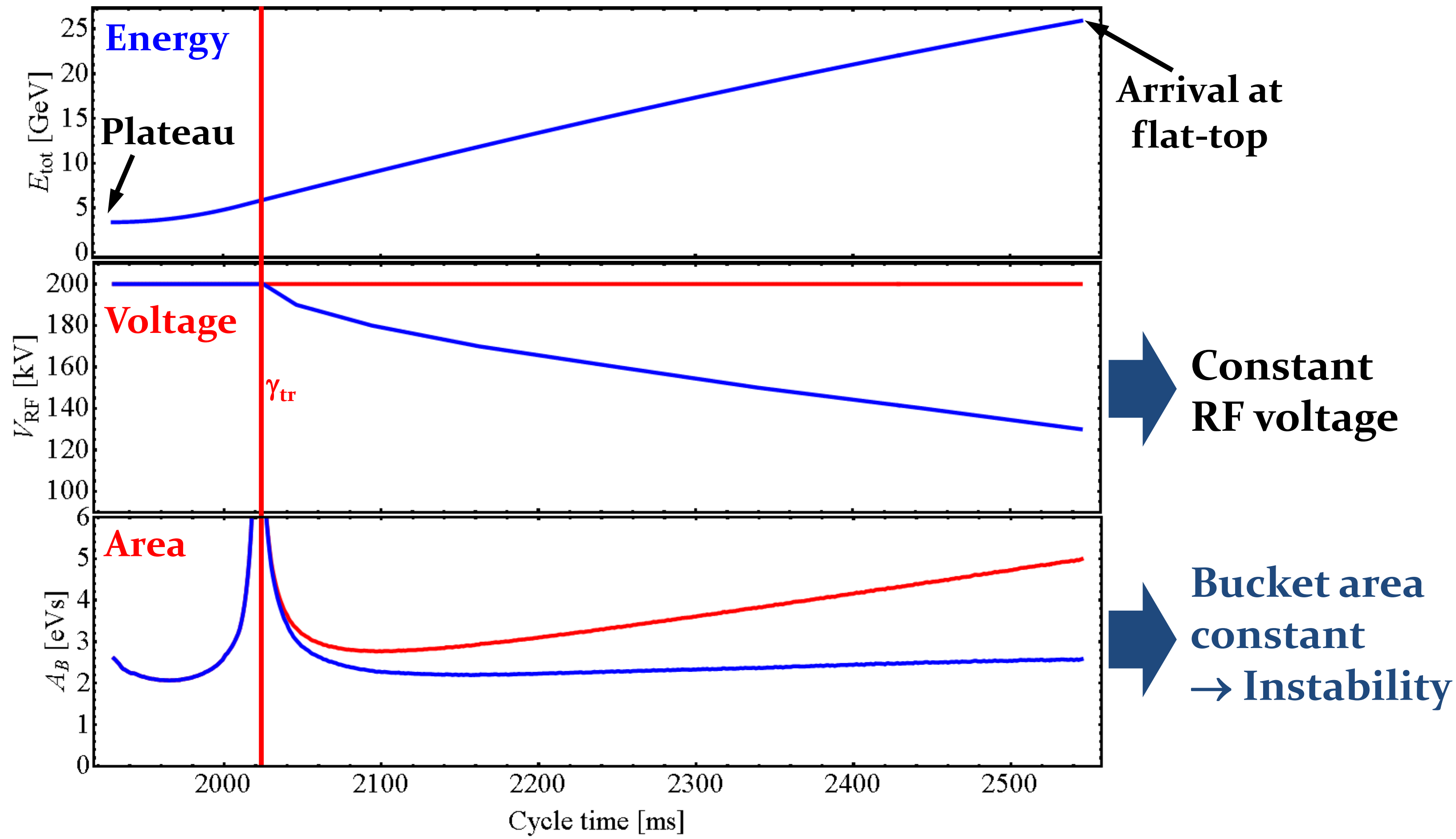
Constant RF voltage

Bucket area grows

→ Risk of Instability

Example: stabilization with lower voltage

→ **Acceleration of protons in the CERN PS (3.4 → 26 GeV total)**



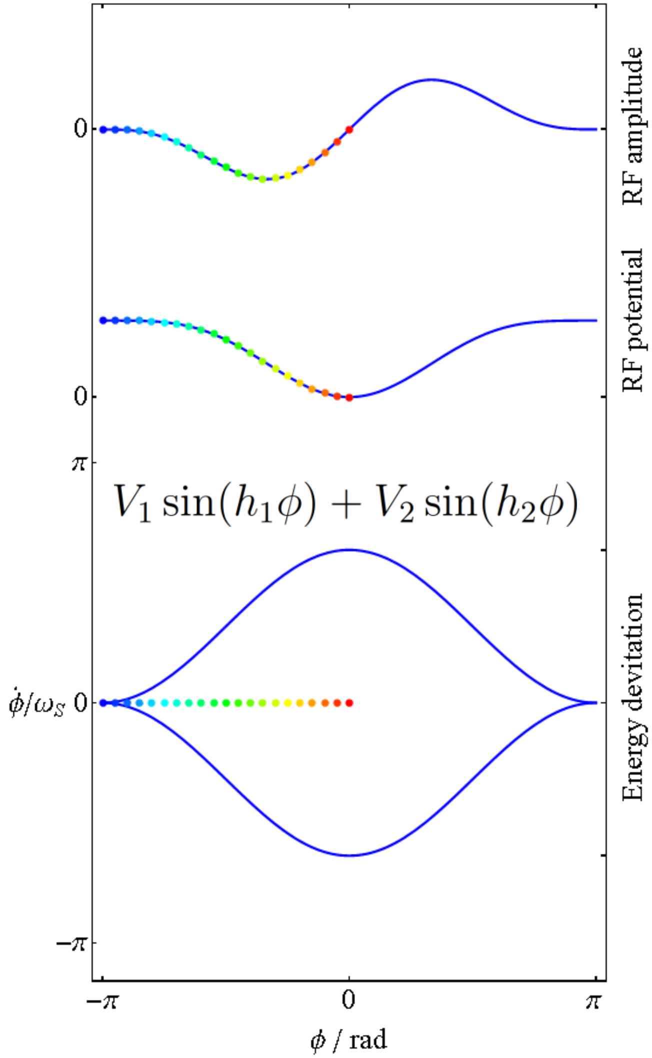
- **Same principle also applied in SPS and LHC**

→ **Prevent bucket filling to decrease**

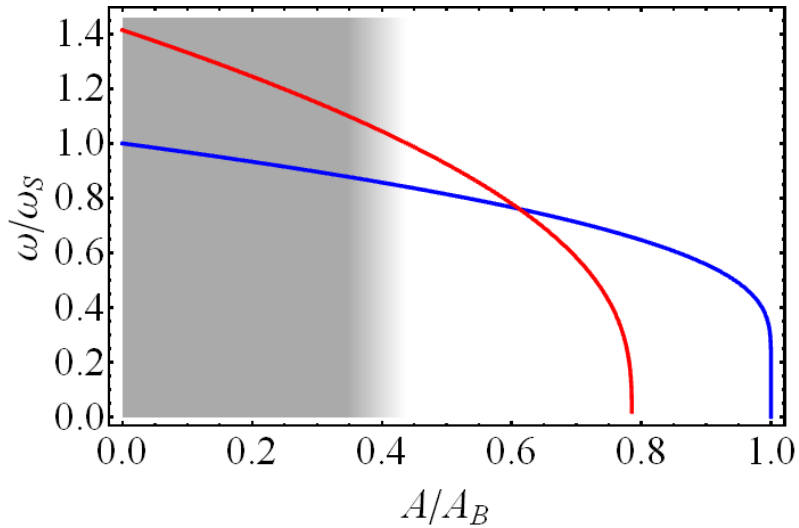
Additional non-linearity by double RF

→ **RF system at twice the main frequency and at half amplitude**

$V_2/V_1 = 1/2$, $h_2/h_1 = 2$



Synchrotron frequency



- **Both RF systems in phase**

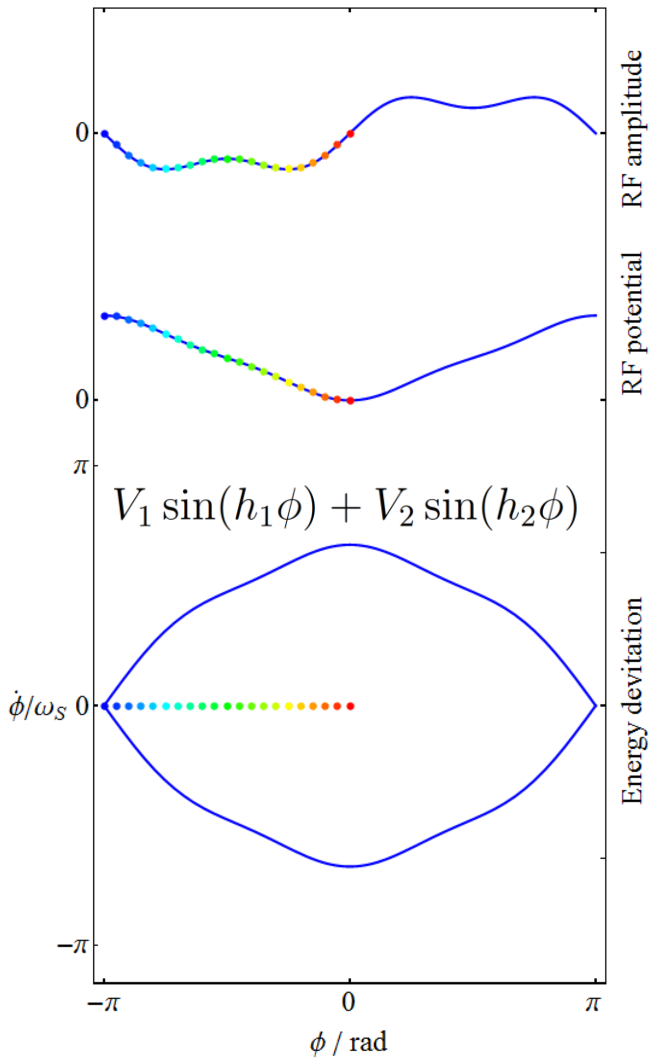
→ **Important increase in synchrotron frequency spread**

→ **Improves stability**

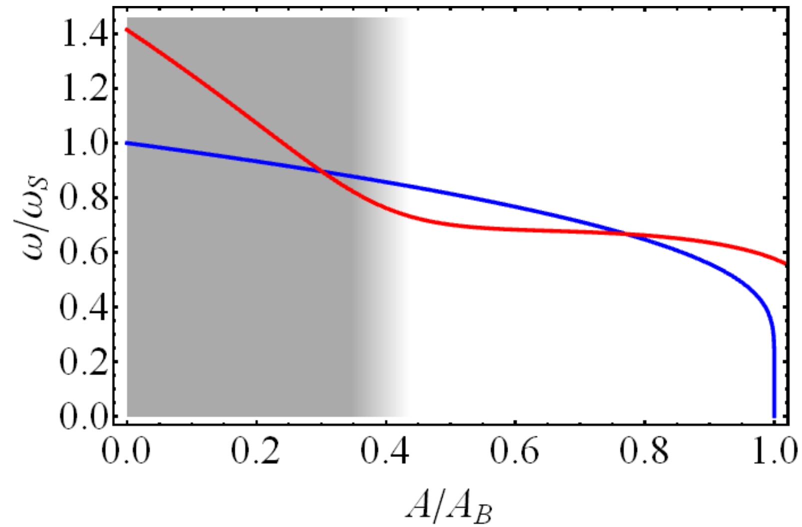
Additional non-linearity by double RF

→ RF system at twice the main frequency and at half amplitude

$V_2/V_1 = 1/3$, $h_2/h_1 = 3$



Synchrotron frequency



- **Both RF systems in phase**

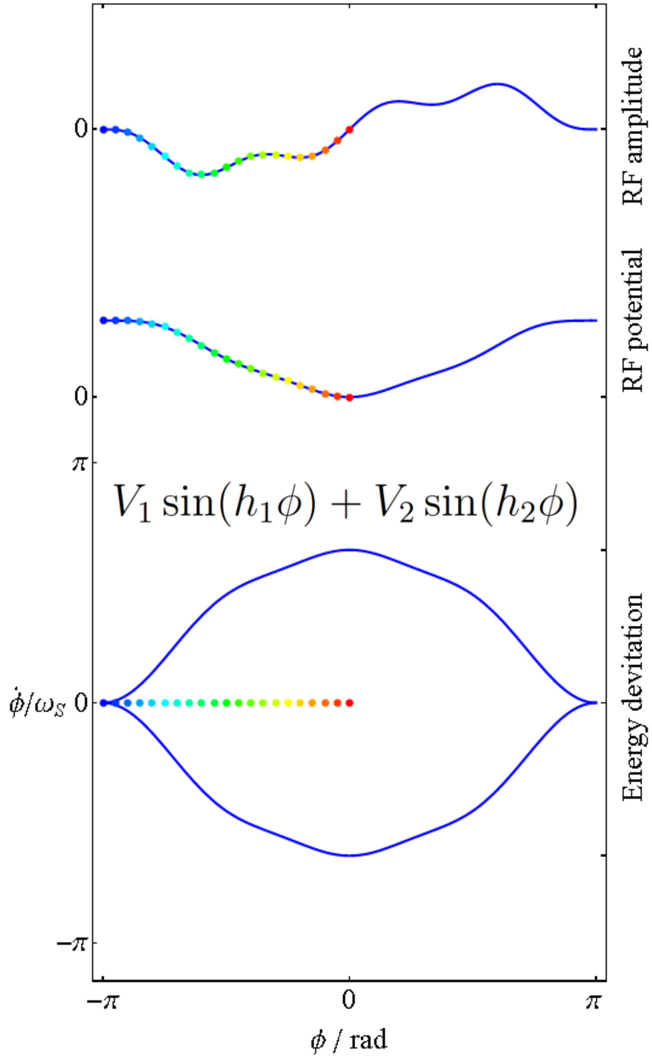
→ **Important increase in synchrotron frequency spread**

→ **Improves stability**

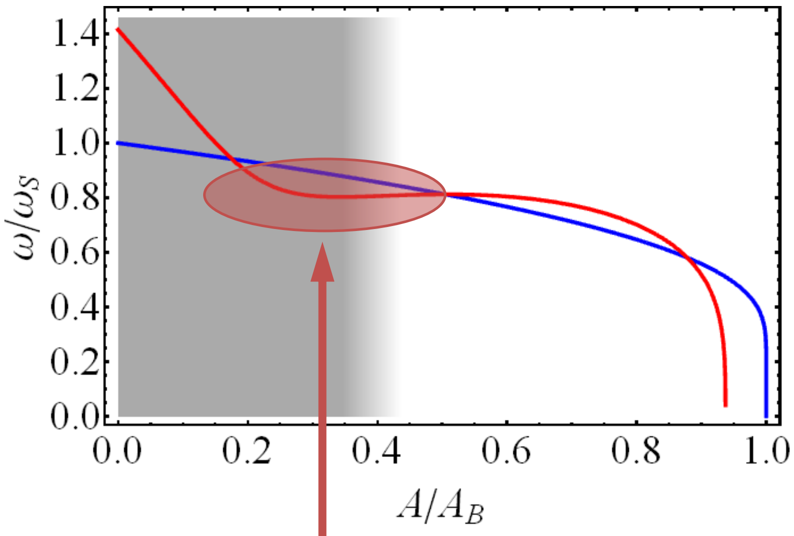
Additional non-linearity by double RF

→ RF system at twice the main frequency and at half amplitude

$V_2/V_1 = 1/4$, $h_2/h_1 = 4$



Synchrotron frequency



- **Local regions of bunch with no f_S gradient**

→ Again prone to instability

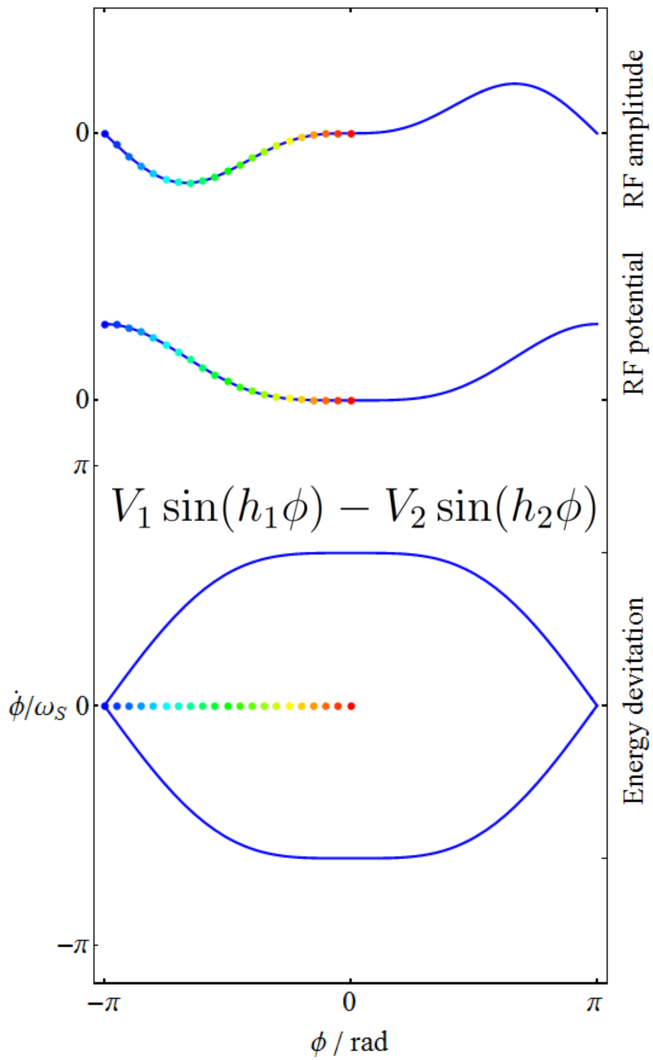
→ Reduce voltage of 2nd harmonic RF system

→ Improving stability depends on appropriate voltage ratio

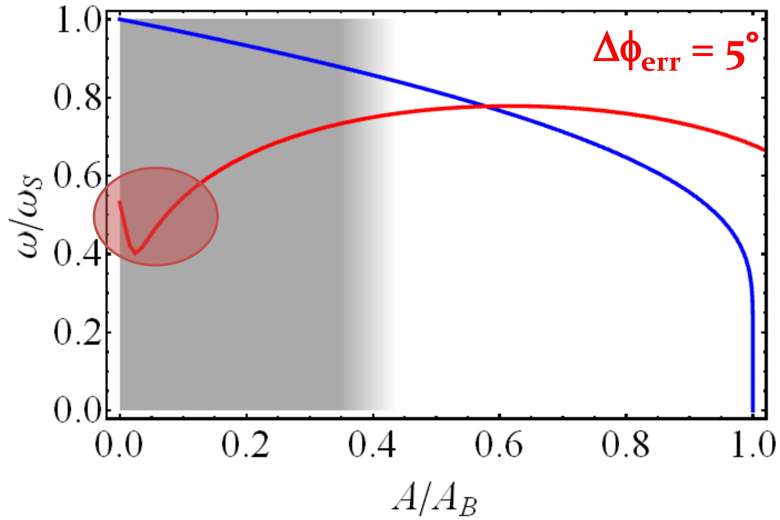
Two RF systems in counter-phase?

→ 2^{nd} RF twice frequency, half amplitude in counter-phase

$V_2/V_1 = 1/2$, $h_2/h_1 = 2$



Synchrotron frequency



- **Large frequency spread at bunch centre with perfectly adjusted phases**

→ **Minor phase offset causes locally unstable regions**

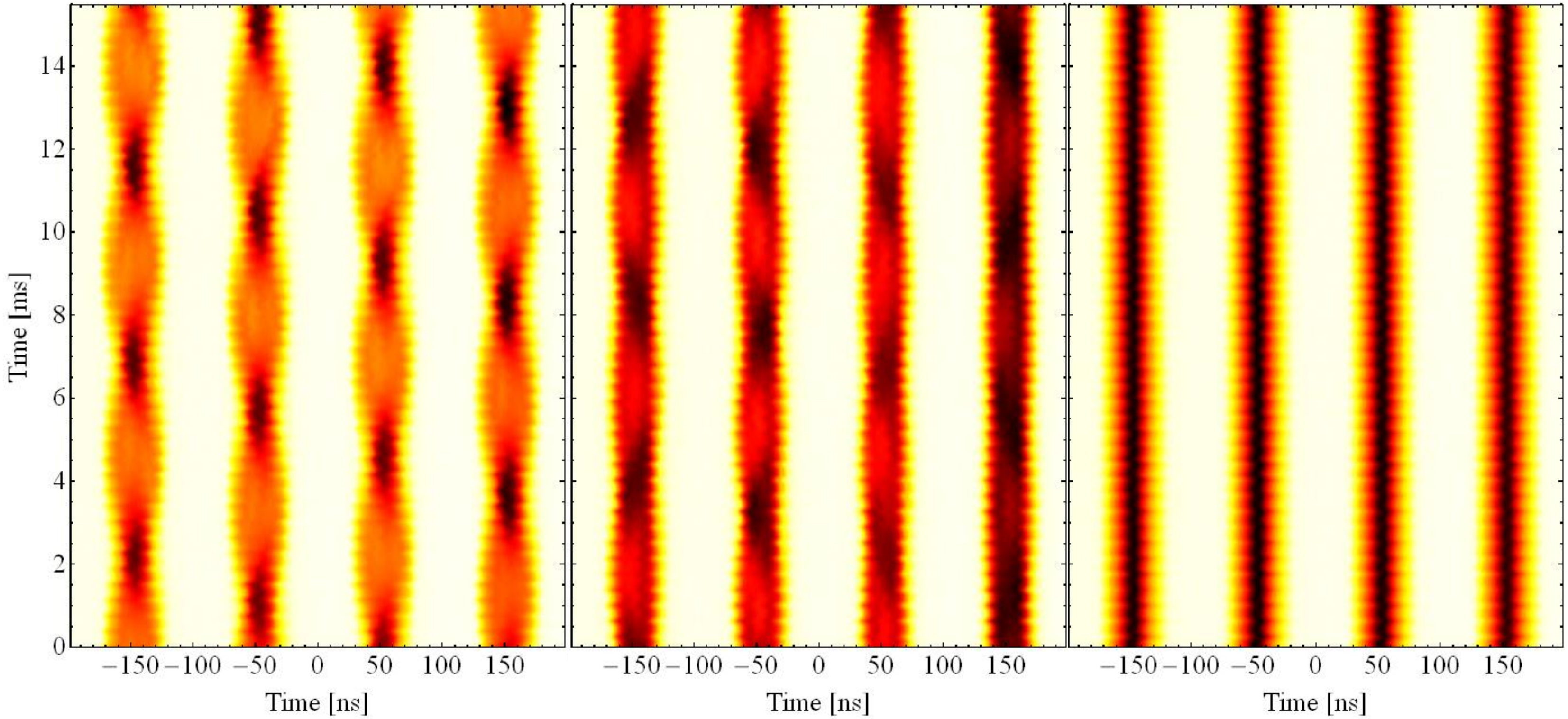
→ **Works only for very short bunches**

→ **Electron accelerators**

Example: damping observations in the PS

- **Quadrupolar coupled-bunch oscillations at flat-top**
- **Main RF system:** $h_1 = 21$, 10 MHz, 4 out of 18 bunches
- **Higher-harmonic RF system:** $h_2 = 84$, 40 MHz

Single-harmonic RF | **(Imperfect) counter-phase** | **Double RF, in phase**



Both RF systems in phase:

→ Highest peak current, but most stable

Summary

- **Longitudinal beam dynamics**
  $\rightarrow$ **Everything non-linear**

- **Longitudinal manipulations**
  $\rightarrow$ **Tricks to adjust length and distance of bunches**
  $\rightarrow$ **Do more with less RF**

- **Synchrotron frequency spread**
  $\rightarrow$ **More RF voltage may be less stability**
  $\rightarrow$ **Higher peak density may be more stable**
  $\rightarrow$ **Improve stability and control emittance**

A big Thank You

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Danilo Quartullo, Markus Ries, Elena Shaposhnikova,
Frank Tecker**

Thank you very much for your attention!

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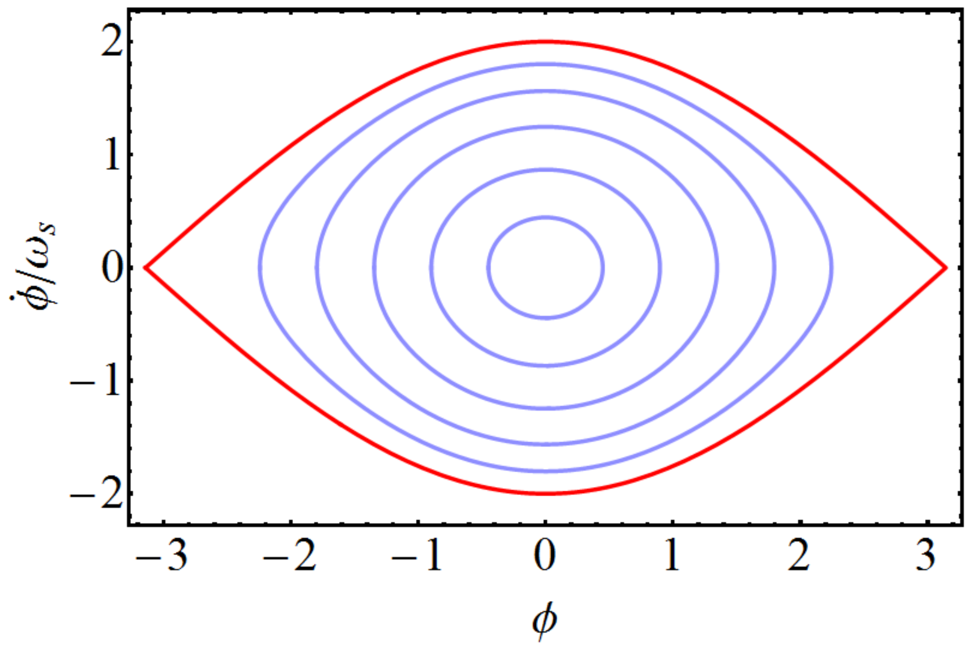
Spare slides

Stationary bucket in normalized coordinates

→ **RF bucket properties become independent from accelerator parameters**

→ **Significant simplification of equations, easy to use**

Example of stationary bucket



→ **Bucket height**

$$\frac{\dot{\phi}_B}{\omega_S} = 2\,\mathrm{rad}$$

→ **Bucket area**

$$\frac{A_B}{\omega_S} = 16\,\mathrm{rad}^2$$

→ **Exception: conservation of longitudinal phase space**