

Beam-beam effects

(an introduction)

Slide 1

Werner Herr
CERN

http://cern.ch/Werner.Herr/CAS2011/lectures/Chios_beambeam.pdf
http://cern.ch/Werner.Herr/CAS2011/proceedings/bb_proc.pdf

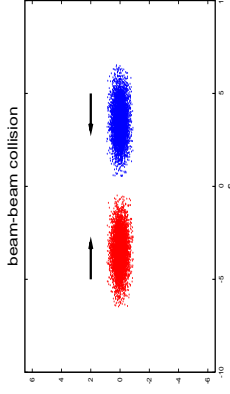
Werner Herr, beam-beam effects, CAS 2011, Chios

What are beam-beam effects ?

- They occur when two beams collide
- Two types of beam-beam effects:
 - High energy collisions between two particles (wanted)
 - Distortions of beams by electromagnetic forces (unwanted)
- Unfortunately: usually both go together ...

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Beam-beam collision



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Typically:

- 🤔 0.001% (or less) of particles collide
 - 😬 99.999% (or more) of particles are distorted
-

Beam-beam effects

- In circular colliders: interactions happen (at least) once per turn !
 - Many different effects and problems
 - Try to understand some of them
 - In linear collider: **VERY** different problems
 - Two main questions:
 - What happens to a single particle ?
 - What happens to the whole beam?
-

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BEAMS: moving charges

- Beam is a collection of charges
- Represent electromagnetic potential for other charges
 - Forces on itself (**space charge**) and opposing beam (**beam-beam effects**)
 - Main limit in past, present and future colliders
 - Important for high density beams, i.e. high intensity and/or small beams: for **high luminosity** !

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Beam-beam effects

Remember:

$$\mathcal{L} = \frac{N_1 N_2 f n_B}{4\pi \sigma_x \sigma_y} = \frac{N_1 N_2 f n_B}{4\pi \cdot \sigma_x \sigma_y}$$

- Slide 6
- Overview: which effects are important for present and future machines (LEP, PEP, Tevatron, RHIC, LHC, ...)
 - Qualitative and physical picture of the effects
 - Mathematical derivations in:

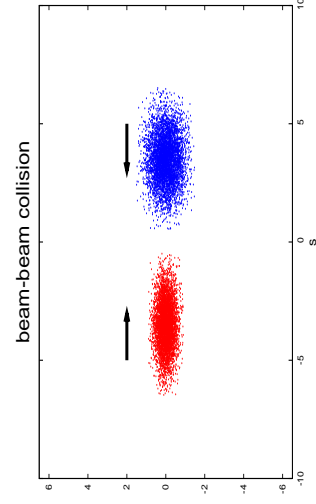
http://cern.ch/Werner.Herr/CAS2011/proceedings/bb_proc.pdf

Beam-beam effects

- A beam acts on particles like an electromagnetic lens, but:
 - Does not represent simple form, i.e. well defined multipoles
 - Very non-linear form of the forces, depending on distribution
 - Can change distribution as result of interaction (time dependent forces ..)

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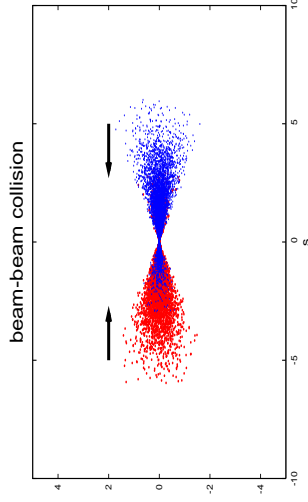
Beam-beam collision



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- Two beams can have different parameters ($I, \sigma \dots$)
- Very detrimental effects on the beams

Beam-beam collision



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- They can change as a result of the beam-beam interaction
- Very detrimental effects on the beams

Studying beam-beam effects

- Need knowledge of the forces
- Apply concepts of non-linear dynamics
- Apply concepts of multi-particle dynamics
- Analytical models and simulation techniques well developed in last 10 years

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Fields and Forces (I)

- Need fields \vec{E} and \vec{B} of opposing beam with a particle distribution $\rho(x, y, z)$
- In rest frame only electrostatic field: \vec{E}' , $\vec{B}' \equiv 0$
- Derive potential $U(x, y, z)$ from Poisson equation:

$$\Delta U(x, y, z) = -\frac{1}{\epsilon_0} \rho(x, y, z)$$

- The electrostatic fields become:

$$\vec{E}' = -\nabla U(x, y, z)$$

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Fields and Forces (II)

- Transform into moving frame and calculate Lorentz force \vec{F} on particle with charge $q = Z_2 e$

$$E_{\parallel} = E'_{\parallel}, \quad E_{\perp} = \gamma \cdot E'_{\perp} \quad \text{with:} \quad \vec{B} = \vec{\beta} \times \vec{E}/c$$

$$\vec{F} = q(\vec{E} + \vec{\beta} \times \vec{B})$$

- Example Gaussian distribution:

$$\rho(x, y, z) = \frac{NZ_1 e}{\sigma_x \sigma_y \sigma_z \sqrt{2\pi}^3} \exp\left(-\frac{x^2}{2\sigma_x^2} - \frac{y^2}{2\sigma_y^2} - \frac{z^2}{2\sigma_z^2}\right)$$

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Simple example: Gaussian

- For 2D case the potential becomes (see proceedings):

$$U(x, y, \sigma_x, \sigma_y) = \frac{NZ_1 e}{4\pi\epsilon_0} \int_0^\infty \frac{\exp\left(-\frac{x^2}{2\sigma_x^2+q} - \frac{y^2}{2\sigma_y^2+q}\right)}{\sqrt{(2\sigma_x^2+q)(2\sigma_y^2+q)}} dq$$

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- Can derive \vec{E} and \vec{B} fields and therefore forces
- For arbitrary distribution (non-Gaussian): difficult (or impossible, numerical solution required)

Simple example: round Gaussian beams

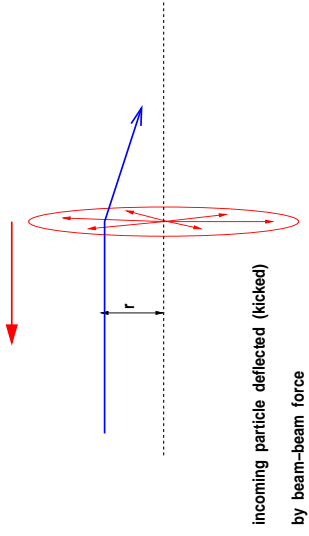
- Assumption 1: $\sigma_x = \sigma_y = \sigma$, $Z_1 = -Z_2 = 1$
- Assumption 2: very relativistic $\rightarrow \beta \approx 1$
- Only components E_r and B_ϕ are non-zero
- Force has only radial component, i.e. depends only on distance \mathbf{r} from bunch centre (where: $r^2 = x^2 + y^2$) (see proceedings)

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$$F_r(\mathbf{r}) = -\frac{Ne^2(1+\beta^2)}{2\pi\epsilon_0 \cdot r} \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

Beam-beam kick:

- We use (x, x', y, y') as coordinates
- We need the deflections (kicks $\Delta x', \Delta y'$) of the particles:



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Beam-beam kick:

- Kick ($\Delta r'$): angle by which the particle is deflected during the passage
- Integration of force over the collision, i.e. time of passage Δt (assuming: $m_1 = m_2$ and $Z_1 = -Z_2 = 1$):

$$\text{Newton's law: } \Delta r' = \frac{1}{mc\beta\gamma} \int_{-\frac{\Delta t}{2}}^{+\frac{\Delta t}{2}} F_r(r, s, t) dt$$

with:

$$F_r(r, s, t) = -\frac{Ne^2(1 + \beta^2)}{\sqrt{(2\pi)^3} \epsilon_0 r \sigma_s} \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right] \cdot \left[\exp\left(-\frac{(s + vt)^2}{2\sigma_s^2}\right) \right]$$

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Beam-beam kick:

- Using the classical particle radius (implies $Z_1 = \pm Z_2$):

$$r_0 = e^2 / 4\pi\epsilon_0 m c^2$$

we have (radial kick and in Cartesian coordinates):

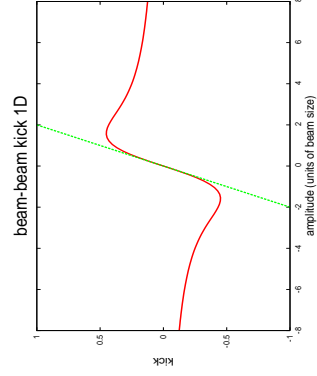
$$\Delta r' = -\frac{2Nr_0}{\gamma} \cdot \frac{r}{r^2} \cdot \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

$$\Delta x' = -\frac{2Nr_0}{\gamma} \cdot \frac{x}{r^2} \cdot \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

$$\Delta y' = -\frac{2Nr_0}{\gamma} \cdot \frac{y}{r^2} \cdot \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

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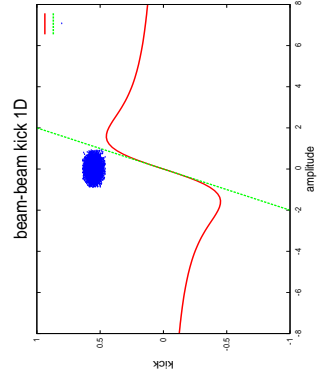
Beam-beam force/kick



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- For small amplitude: linear force (like quadrupole)
- For large amplitude: very non-linear force

Beam-beam force/kick

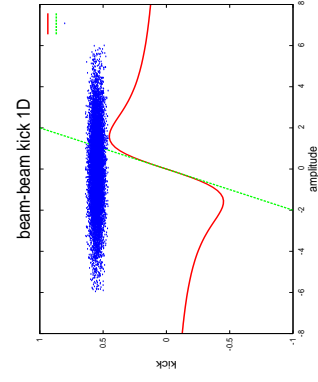


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➤ For small amplitude: tune shift



Beam-beam force/kick



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➤ For small amplitude: tune shift
➤ For large amplitude: amplitude dependent tune shift



Can we quantify the beam-beam strength ?

- Try the slope of force (kick $\Delta r'$) at zero amplitude
- This defines: beam-beam parameter ξ
- For head-on interactions and round beams ($\beta_x^* = \beta_y^* = \beta_y^*$) we get:

$$\xi = \frac{\beta^*}{4\pi} \cdot \frac{\delta(\Delta r')}{\delta r} = \frac{N \cdot r_0 \cdot \beta^*}{4\pi \gamma \sigma^2}$$

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LEP - LHC

	LEP (e ⁺ e ⁻)	LHC (pp)
Beam sizes	160 - 200 μm · 2 - 4 μm	16.6 μm · 16.6 μm
Intensity N	4.0 · 10 ¹¹ /bunch	1.15 · 10 ¹¹ /bunch
Energy	100 GeV	7000 GeV
$\epsilon_x \cdot \epsilon_y$	(\approx) 20 nm · 0.2 nm	0.5 nm · 0.5 nm
$\beta_x^* \cdot \beta_y^*$	(\approx) 1.25 m · 0.05 m	0.55 m · 0.55 m
Crossing angle	0.0	285 μrad
Beam-beam parameter (ξ)	0.0700	0.0037

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Can we quantify the beam-beam strength ?

- In general for non-round beams ($\beta_x^* \neq \beta_y^*$):

$$\xi_{x,y} = \frac{N r_0 \beta_{x,y}^*}{2\pi\gamma\sigma_{x,y}(\sigma_x + \sigma_y)}$$

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- Proportional to (linear) tune shift ΔQ_{bb} from beam-beam interaction: $\Delta Q_{bb} \propto \pm \xi$
- Good measure for strength of beam-beam interaction
- BUT: does not describe
 - changes to optical functions
 - non-linear part of beam-beam force

Linear beam-beam tune shift

- For small amplitudes linear force like a quadrupole with focal length f

$$\frac{1}{f} = \frac{\Delta x'}{x} = \frac{Nr_0}{\gamma\sigma^2} = \left[\frac{\xi \cdot 4\pi}{\beta^*} \right]$$

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- Transformation matrix over the interaction becomes:

$$\begin{pmatrix} 1 & 0 \\ \frac{1}{-f} & 1 \end{pmatrix}$$

Linear beam-beam tune shift

- Full turn matrix including the tune shift ΔQ computed from unperturbed full turn matrix plus interaction

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$$\begin{aligned} & \begin{pmatrix} \cos(2\pi(Q+\Delta Q)) & \beta^* \sin(2\pi(Q+\Delta Q)) \\ -\frac{1}{\beta^*} \sin(2\pi(Q+\Delta Q)) & \cos(2\pi(Q+\Delta Q)) \end{pmatrix} \\ = & \begin{pmatrix} 1 & 0 \\ \frac{1}{-2J} & 1 \end{pmatrix} \circ \begin{pmatrix} \cos(2\pi Q) & \beta_0^* \sin(2\pi Q) \\ -\frac{1}{\beta_0^*} \sin(2\pi Q) & \cos(2\pi Q) \end{pmatrix} \circ \begin{pmatrix} 1 & 0 \\ \frac{1}{-2J} & 1 \end{pmatrix} \end{aligned}$$

Linear beam-beam tune shift

- Solving this equation gives us:

$$\cos(2\pi(Q + \Delta Q)) = \cos(2\pi Q) - \frac{\beta_0^*}{2f} \sin(2\pi Q)$$

and

$$\frac{\beta^*}{\beta_0^*} = \sin(2\pi Q) / \sin(2\pi(Q + \Delta Q))$$

- Tune is changed by ΔQ
- β -function is changed (β -beating)

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Linear beam-beam tune shift

- For small ξ and Q not too close to 0.0 and 0.5 we have:

$$\Delta Q \approx \xi$$

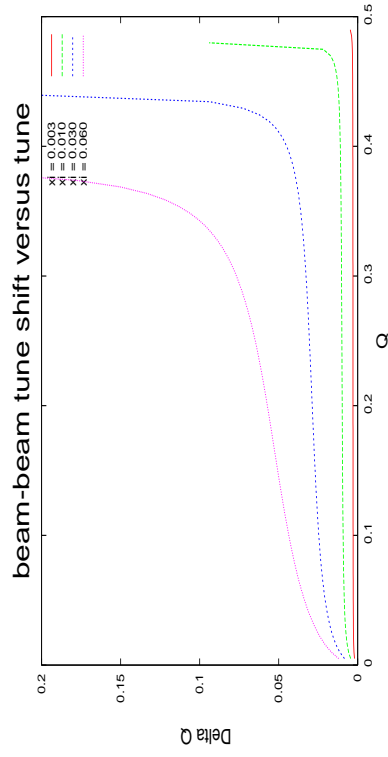
and

$$\frac{\beta^*}{\beta_0^*} = \frac{\sin(2\pi Q)}{\sin(2\pi(Q + \Delta Q))} = \frac{\beta_0}{\sqrt{1 + 4\pi\xi \cot(2\pi Q) - 4\pi^2\xi^2}}$$

- β can become smaller or larger at interaction point (dynamic β)

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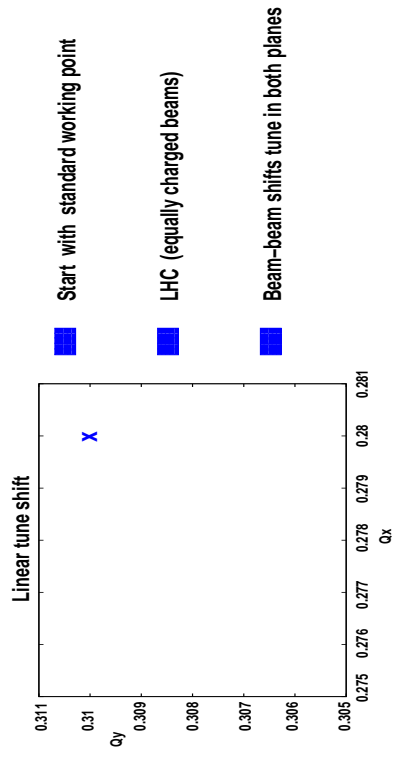
Tune dependence of tune shift



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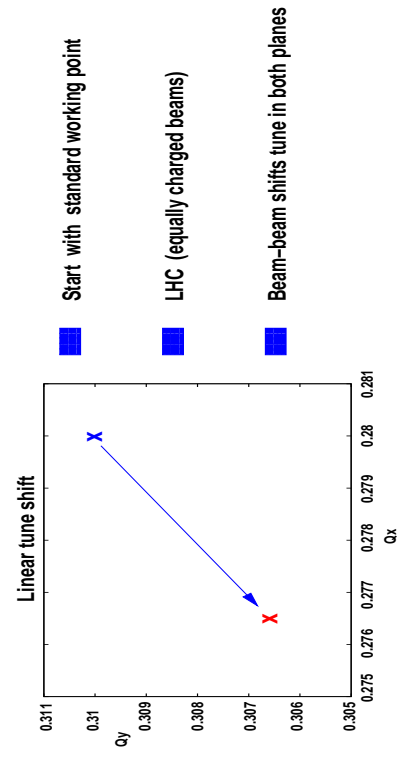
- Strong dependence on Q for larger ξ (dynamic β)

Linear tune shift - two dimensions



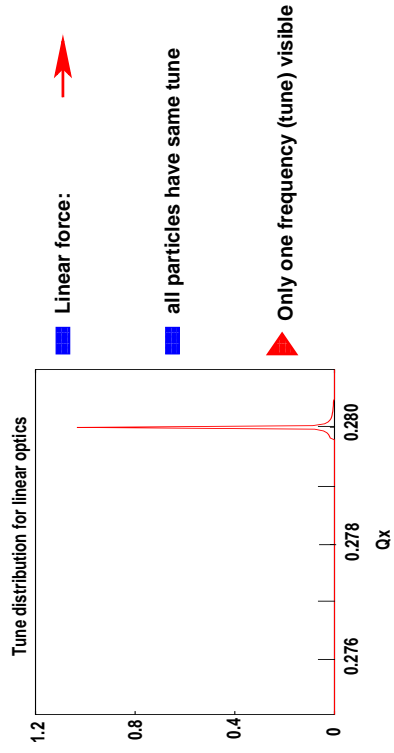
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Linear tune shift - two dimensions



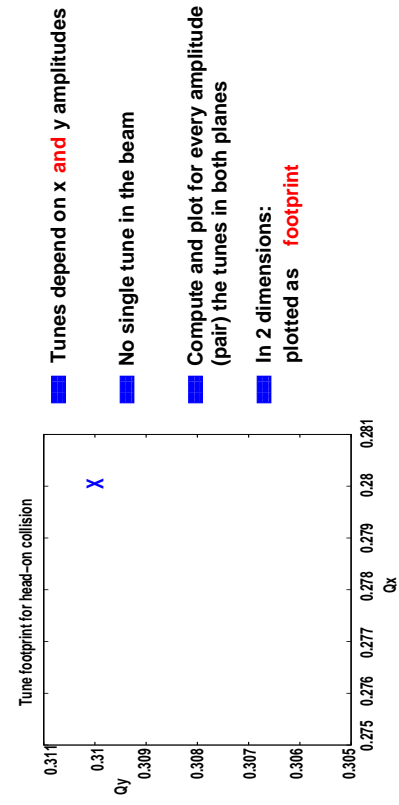
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Tune measurement: linear optics



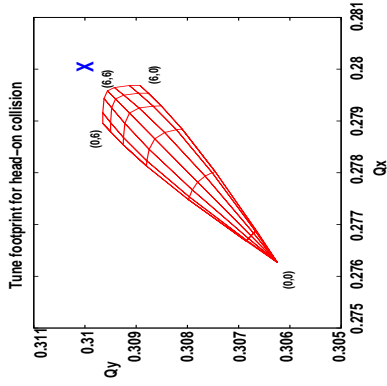
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Non-linear tune shift - two dimensions



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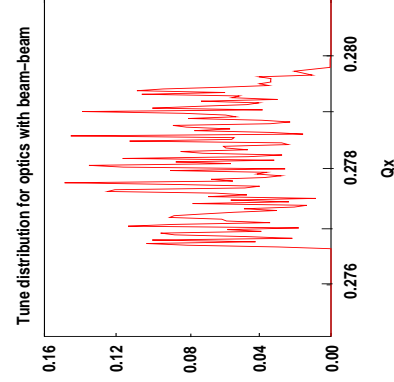
Non-linear tune shift - two dimensions



- Tunes depend on x and y amplitudes
- No single tune in the beam
- Compute and plot for every amplitude (pair) the tunes in both planes
- In 2 dimensions: plotted as footprint

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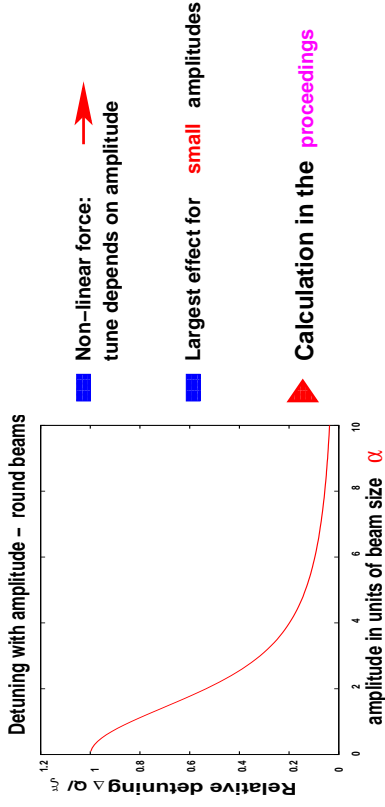
Tune measurement: with beam-beam



- Non-linear force: \rightarrow
- particles with different amplitudes have different frequencies (tunes)
- We get frequency (tune) spectra
- Width of the spectra: about ξ

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Amplitude detuning



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→ with $\alpha \equiv \frac{a}{\sigma_*}$ we get: $\Delta Q/\xi = \frac{4}{\alpha^2} \left[1 - I_0\left(\frac{\alpha^2}{4}\right) \cdot e^{-\frac{\alpha^2}{4}} \right]$

Weak-strong and strong-strong

- Both beams are very strong (strong-strong):
- Both beam are affected and change due to beam-beam interaction
- Examples: LHC, LEP, RHIC, ...
- Evaluation of effects challenging
- One beam much stronger (weak-strong):
- Only the weak beam is affected and changed due to beam-beam interaction
- Examples: SPS collider, Tevatron, ...

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Incoherent effects

(single particle effects)

- Single particle dynamics: treat as a particle through a static electromagnetic lens
- Basically non-linear dynamics
- All single particle effects observed:
 - Unstable and/or irregular motion
 - Beam blow up
 - Bad lifetime, particle loss

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Observations hadrons

- Non-linear motion can become chaotic
 - reduction of "dynamic aperture"
 - particle loss and bad lifetime
- Strong effects in the presence of noise or ripple
- Very bad: unequal beam sizes (studied at SPS, HERA)
- Evaluation is done by simulation

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Observations leptons

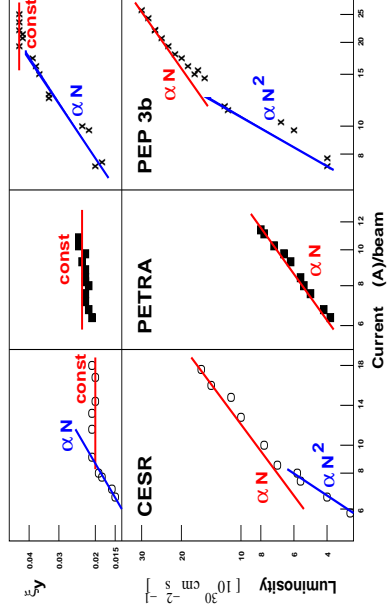
Remember:

$$\Rightarrow \mathcal{L} = \frac{N_1 N_2 f n_B}{4\pi\sigma_x\sigma_y}$$

- Luminosity should increase $\propto N_1 N_2$
- ➡ for: $N_1 = N_2 = N \rightarrow \propto N^2$
- Beam-beam parameter should increase $\propto N$
- But:

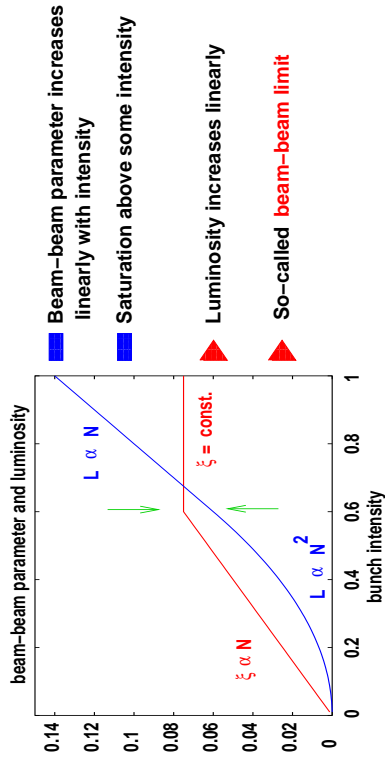
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Examples: beam-beam limit



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Beam-beam limit (schematic)



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What is happening ?

we have
$$\xi_y = \frac{N r_0 \beta_y}{2\pi\gamma\sigma_y(\sigma_x + \sigma_y)} \quad (\sigma_x \gg \sigma_y) \approx \frac{r_0 \beta_y}{2\pi\gamma(\sigma_x)} \cdot \frac{N}{\sigma_y}$$

and
$$\mathcal{L} = \frac{N^2 f n_B}{4\pi\sigma_x\sigma_y} = \frac{N f n_B}{4\pi\sigma_x} \cdot \frac{N}{\sigma_y}$$

- Above beam-beam limit: σ_y increases when N increases to keep ξ constant \rightarrow **equilibrium emittance !**
- Therefore: $\mathcal{L} \propto N$ and $\xi \approx$ constant
- ξ_{limit} is **NOT** a universal constant !
- Difficult to predict

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What is happening ?

- Where does it come from ?
- From synchrotron radiation: vertical plane damped, horizontal plane excited
- Horizontal beam size usually (much) larger
- Vertical beam-beam effect depends on horizontal (large) amplitude
- Coupling from horizontal to vertical plane
- Equilibrium between this excitation and damping determines ξ_{limit}

Lesson: Keep the coupling small !

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The next problem

Remember:

$$\Rightarrow \mathcal{L} = \frac{N_1 N_2 f \cdot \eta_B}{4\pi\sigma_x\sigma_y}$$

- How to collide many bunches (for high \mathcal{L}) ??
- Must avoid unwanted collisions !!
- Separation of the beams:
 - Pretzel scheme (SPS, LEP, Tevatron)
 - Bunch trains (LEP, PEP)
 - Crossing angle (LHC)

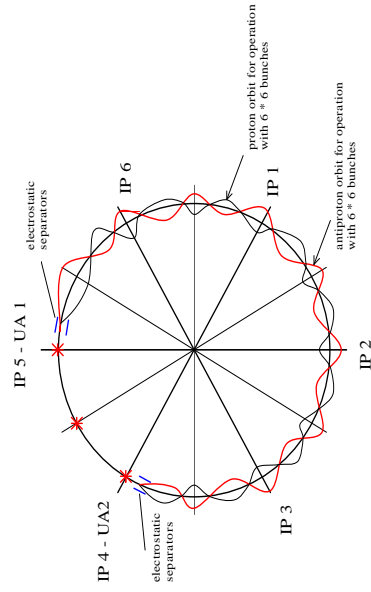
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Separation: SPS

- Few equidistant bunches
(6 against 6)
- Beams travel in same beam pipe
(12 collision points !)
- Two experimental areas
- Need **global** separation
- Horizontal pretzel around most of the circumference

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Separation: SPS



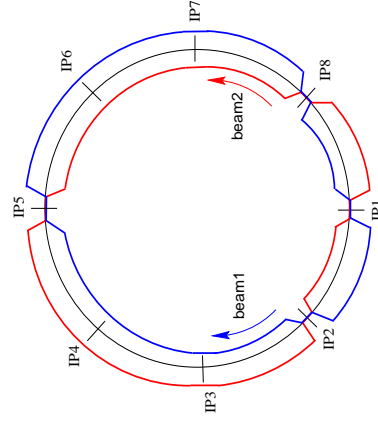
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Separation: LHC

- Many equidistant bunches (2808 per beam)
- Two beams already separated in two separate beam pipes except:
 - Four experimental areas
 - Need **local** separation
- Two horizontal and two vertical crossing angles

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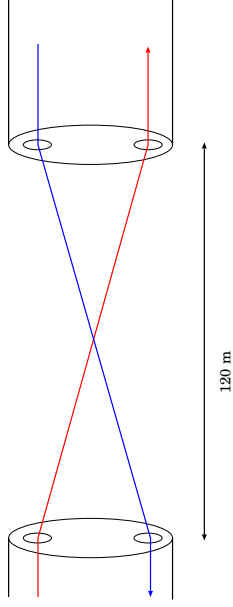
Layout of LHC



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Example: LHC

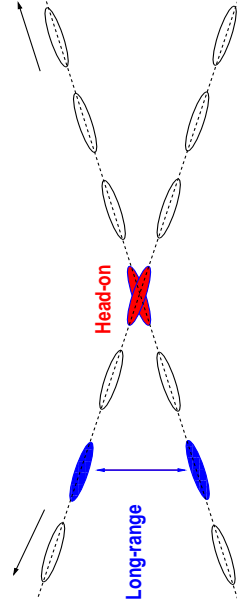
- Two beams, 2808 bunches each, every 25 ns
- In common chamber around experiments



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- Over 120 m: about 30 parasitic interactions

Crossing angles (example LHC)



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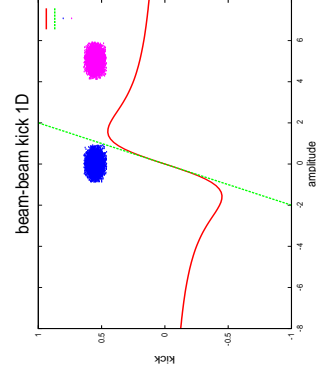
- Particles experience distant (weak) forces
- Separation typically 6 - 12 σ
- We get so-called long range interactions

What is special about them ?

- Break symmetry between planes, stronger resonance excitation
- Mostly affect particles at **large** amplitudes
- Cause effects on closed orbit
- PACMAN effects
- Tune shift has **opposite** sign in plane of separation

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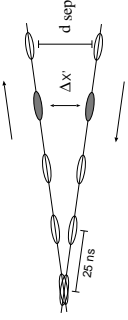
Why opposite tuneshift ???



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- Local slope has opposite sign for large separation
- Opposite sign for focusing !

Long range interactions (LHC)



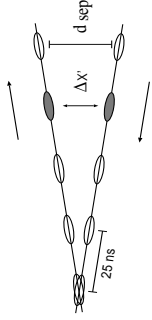
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→ For horizontal separation d :

$$\Delta x'(x+d, y, r) = -\frac{2Nr_0}{\gamma} \cdot \frac{(x+d)}{r^2} \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

(with: $r^2 = (x+d)^2 + y^2$)

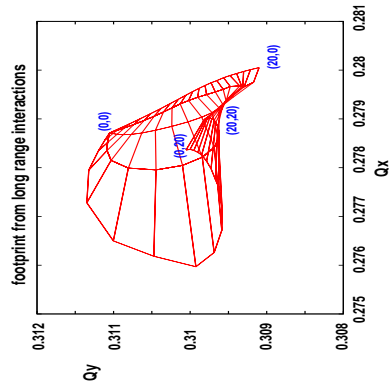
Long range interactions (LHC)



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- Number of long range interactions depends on spacing and length of common part
- In LHC 15 collisions on each side, 120 in total !
- Effects depend on separation: $\Delta Q \propto -\frac{N}{d^2}$ (for large enough d !) footprints ??

Footprints

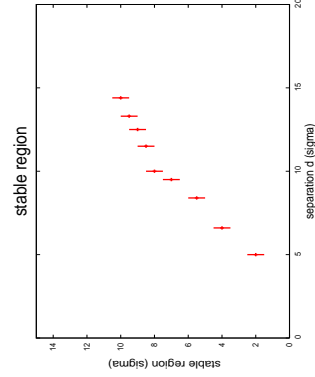


- ▲ Large for largest amplitudes where non-linearities are strong
- ▲ Size proportional to $\frac{1}{d^2}$
- ▲ Must expect problems at small separation
- ▲ Footprint very asymmetric

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Particle losses

- Small crossing angle \iff small separation
- Small separation: particles become unstable and get lost



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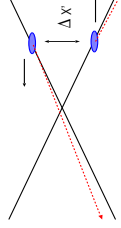
- Minimum separation for LHC: $\approx 10 \sigma$

Closed orbit effects

$$\Delta x'(x + d, y, r) = -\frac{2Nr_0}{\gamma} \cdot \frac{(x + d)}{r^2} \left[1 - \exp\left(-\frac{r^2}{2\sigma^2}\right) \right]$$

For well separated beams ($d \gg \sigma$) the force (kick) has an amplitude independent contribution: \rightarrow orbit kick

$$\Delta x' = \underbrace{\frac{const.}{d}} \cdot \left[1 - \frac{x}{d} + O\left(\frac{x^2}{d^2}\right) + \dots \right]$$



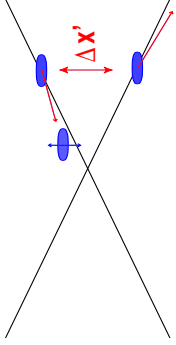
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Closed orbit effects

- Beam-beam kick from long range interactions changes the orbit
- Has been observed in LEP with bunch trains
- Self-consistent calculation necessary
- Effects can add up and become important
- The two beams separate, more than 1σ not unusual !

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Coherent beam-beam effect

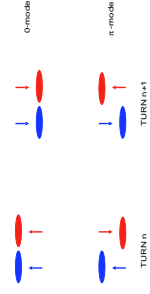


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- Whole bunch sees a kick as an entity (coherent kick)
- The coherent kick of separated beams can excite coherent dipole oscillations
- All bunches couple because each bunch "sees" many opposing bunches: many coherent modes possible !

Coherent beam-beam effect

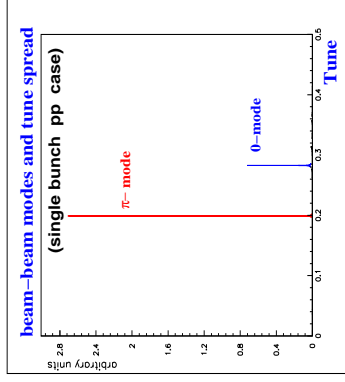
Simplest case: one bunch per beam:



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- Coherent mode: two bunches are "locked" in a coherent oscillation
- 0-mode is stable (Mode with **NO** tune shift)
- π -mode can become unstable (Mode with **LARGEST** tune shift)

Coherent beam-beam frequencies (schematic)



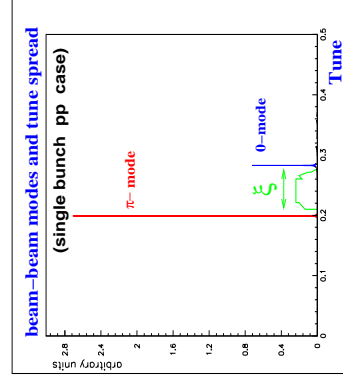
▲ 0-mode is at unperturbed tune

▲ π -mode is shifted by $1.1 - 1.3 \cdot \xi$

- Two separate modes visible
- But we have many particles and tune spread ... !

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Coherent beam-beam frequencies (schematic)



▲ 0-mode is at unperturbed tune

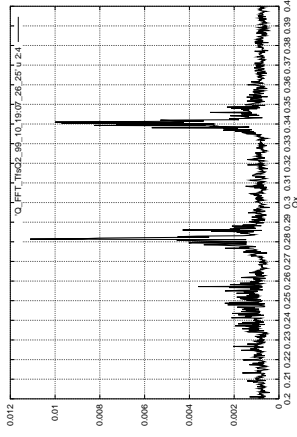
▲ π -mode is shifted by $1.1 - 1.3 \cdot \xi$

▲ Incoherent spread between $[0.0, 1.0] \cdot \xi$

- Strong-strong case: π -mode shifted outside tune spread
- No Landau damping possible

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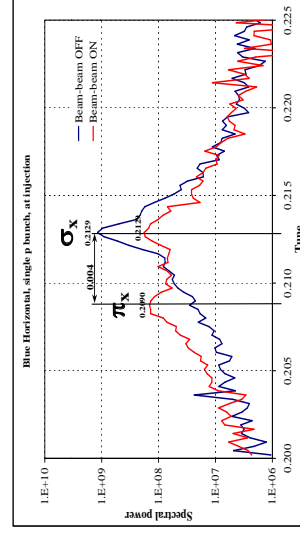
What we measure: LEP



- Two modes clearly visible
- Can be distinguished by phase relation, i.e. sum and difference signals

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What we measure: RHIC

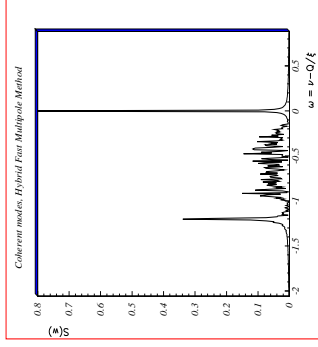


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Courtesy W. Fischer (BNL)

- Compare spectra with and without beams : two modes visible with beams

Simulation of coherent spectra



- ▲ Full simulation of both beams required
- ▲ Use up to 10^8 particles in simulations
- ▲ Must take into account changing fields
- ▲ Requires computation of arbitrary fields

➤ Time consuming for many particles ..

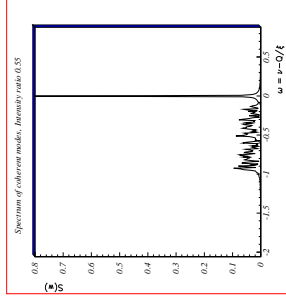
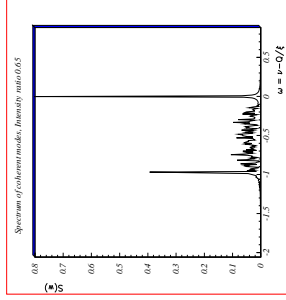
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What can be done to avoid problems ?

- Coherent motion requires 'organized' motion of many particles
- Therefore high degree of symmetry required
- Possible countermeasure: (**symmetry breaking**)
 - Different bunch intensity
 - Different tunes in the two beams

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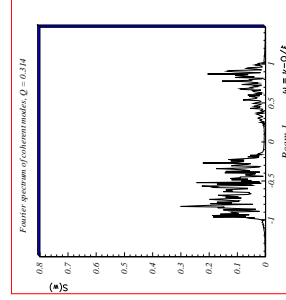
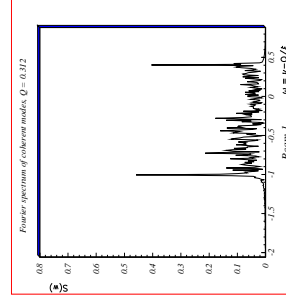
Beams with different intensity



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- Bunches with different intensities cannot maintain coherent motion
- Landau damping restored

Beams with different tunes



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- Bunches with different tunes cannot maintain coherent motion
- Landau damping restored

Can we suppress beam-beam effects ?

- Find 'lenses' to correct beam-beam effects
 - Head on effects:
 - Linear "electron lens" to shift tunes
 - Non-linear "electron lens" to reduce spread
 - Tests in progress at Tevatron and RHIC
 - Long range effects:
 - At very large distance: force is $1/r$
 - Same force as a wire !
- So far: mixed success with **active** compensation

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Others: Möbius lattice

- Principle:
 - Interchange horizontal and vertical plane each turn
- Effects:
 - Round beams (even for leptons)
 - Some compensation effects for beam-beam interaction
 - First test at CESR at Cornell

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Not mentioned:

- Effects in linear colliders
- Asymmetric beams
- Coasting beams
- Beamstrahlung
- Synchrotron coupling
- Monochromatization
- Beam-beam experiments
- ... and many more

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