



# Designing a synchrotron - A real life example

Yannis PAPAPHILIPPOU

Accelerator and Beam Physics group
Beams Department
CERN

**CERN Accelerator School** 

Introduction to Accelerator Physics 2019 Atrium Hotel, Vysoké Tatry, Slovakia 20 September 2019

# Purpose of the Lectures



Review several **aspects** of **beam dynamics** (mostly) presented in the introductory CAS lectures, applied to the **design** and **operation** of a **real synchrotron** 

#### Purpose of the Lectures



- Review several aspects of beam dynamics (mostly) presented in the introductory CAS lectures, applied to the design and operation of a real synchrotron
  - □ Choice of **basic parameters** 
    - Energy, bending field and circumference
  - Optics design
    - ■Cell optics, insertions, transition energy
  - Collective effects
    - ■Instabilities, Space-charge
  - □ Electron/Positron beam dynamics
    - Equilibrium beam properties, energy loss/turn, damping time



#### Choosing a Synchrotron



- Our choice is the CERN Super Proton Synchrotron (SPS)
- From its design and operation, it has shown enormous versatility used for several purposes and serving various applications

## Choosing a Synchrotron



- Our choice is the CERN Super Proton Synchrotron (SPS)
- From its design and operation, it has shown enormous versatility used for several purposes and serving various applications
  - ☐ High energy **synchrotron** serving **fixed target** experiments (West Area, North Area, CNGS, HIRADMAT)
  - **Collider** of protons and anti-protons (W and Z bosons discovery in 1983)
  - Accelerating **electrons** and **positrons** and injecting them to the Large Electron-Positron (**LEP**) Collider
  - Accelerating protons for the Large Hadron Collider (LHC)
  - Accelerating ions for fixed target physics and the LHC
  - Extracting protons for exciting plasma for a **plasma wakefield acceleration** experiment (AWAKE)





# Basic parameters: energy, bending field and circumference

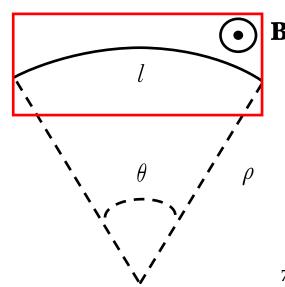


#### Energy and bending field



Consider accelerator ring for particles with energy E with N dipoles of length L or effective length l, i.e. measured on beam path



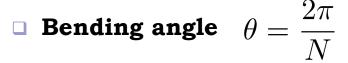




#### Energy and bending field



Consider accelerator ring for particles with energy E with N dipoles of length L or effective length l, i.e. measured on beam path



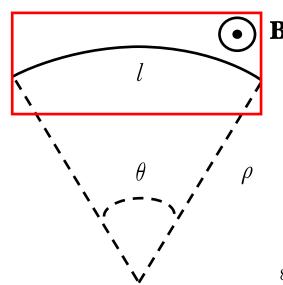
 $oldsymbol{\square}$  Bending radius  $ho = rac{t}{ heta}$ 



□ The integrated dipole strength is

$$Bl = \frac{2\pi}{N} \frac{\beta E}{q}$$





#### Energy and bending field

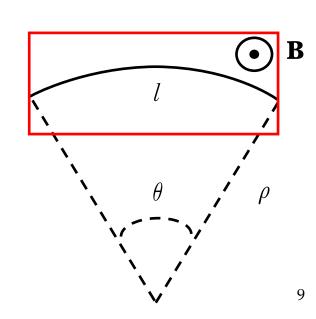


- Consider accelerator ring for particles with energy E with N dipoles of length L or effective length l, i.e. measured on beam path
  - lacksquare Bending angle  $heta=rac{2\pi}{N}$
  - $oldsymbol{\square}$  Bending radius  $ho = rac{t}{ heta}$
  - $lacksymbol{\square}$  The magnetic rigidity is  $B
    ho=rac{eta E}{q}$
  - □ The integrated dipole strength is

$$Bl = \frac{2\pi}{N} \frac{\beta E}{q}$$

- By imposing a **dipole field**, the **dipole length** is **fixed** and vice versa
- The **higher** the **field**, the **shorter** or **less dipoles** can be used







#### Circumference



■ The **filling factor**, is defined as the ratio of the total length of the bending path, with respect to the circumference

$$k_f = \frac{Nt}{C}$$



#### Circumference



■ The **filling factor**, is defined as the ratio of the total length of the bending path, with respect to the circumference

$$k_f = \frac{Nl}{C}$$

The ring circumference becomes

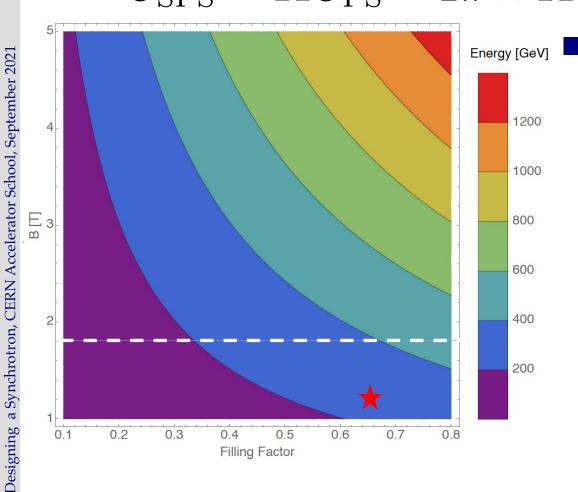
$$C = \frac{2\pi}{k_f B} \frac{\beta E}{q}$$

The ring **circumference** (**cost**) is driven by the bending field choice (technology), the energy reach (physics case, applications) and the design of the lattice cells (optics)





The maximum possible circumference between the CERN I (Meyrin) and CERN II (Prevessin) site was  $C_{\rm SPS} = 11 C_{\rm PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$ 

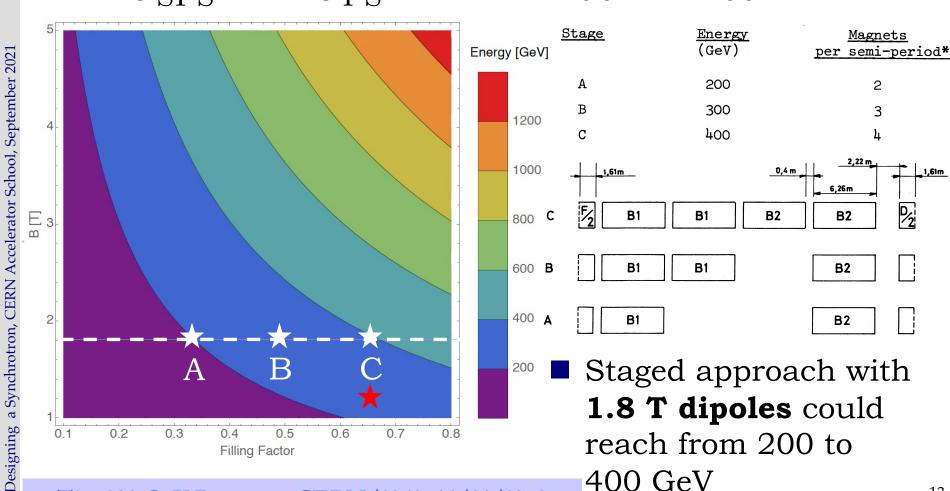


Combined function magnets with 1.2 T field (PS-like) would give an energy of no more then ~260 GeV for a highly packed lattice





The maximum possible circumference between the CERN I (Meyrin) and CERN II (Prevessin) site was  $C_{\rm SPS} = 11C_{\rm PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$ 

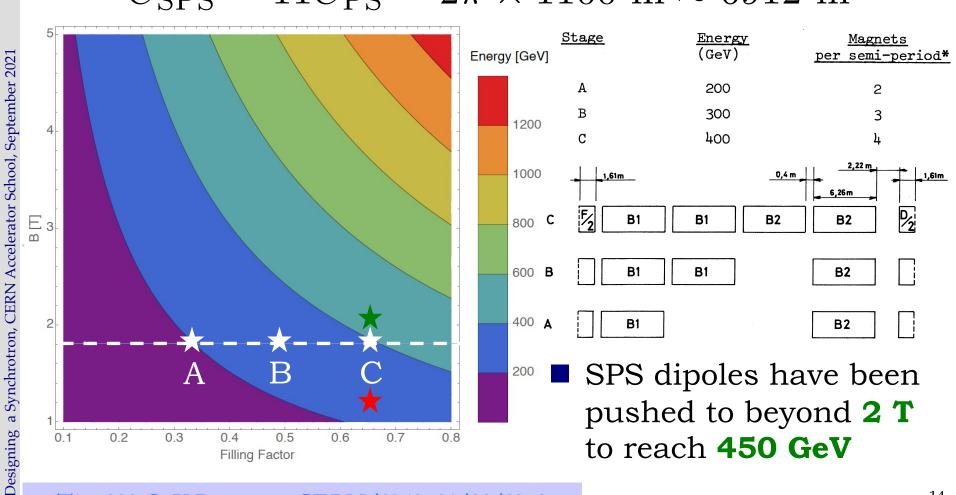


400 GeV





The maximum possible circumference between the CERN I (Meyrin) and CERN II (Prevessin) site was  $C_{\rm SPS} = 11C_{\rm PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$ 



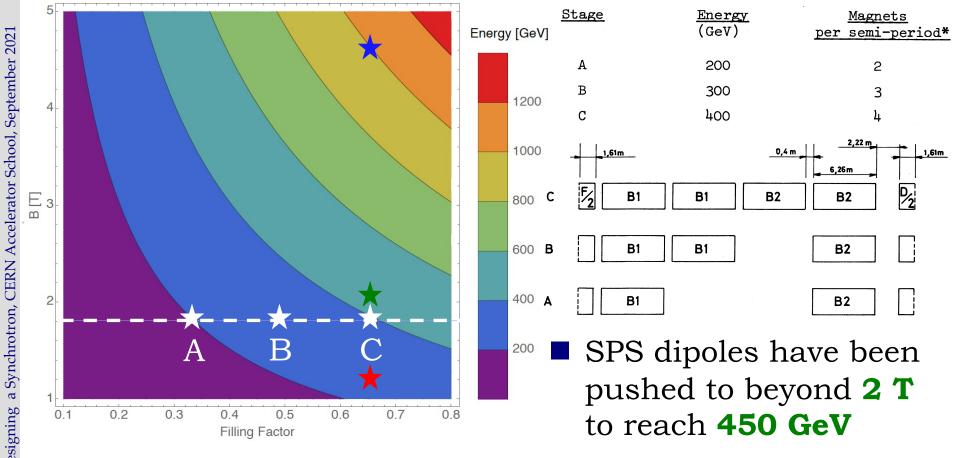
14





The maximum possible circumference between the CERN I (Meyrin) and CERN II (Prevessin) site was

$$C_{\rm SPS} = 11C_{\rm PS} = 2\pi \times 1100 \text{ m} \approx 6912 \text{ m}$$



**Super-conducting** option could raise the energy to 1 TeV<sup>5</sup>





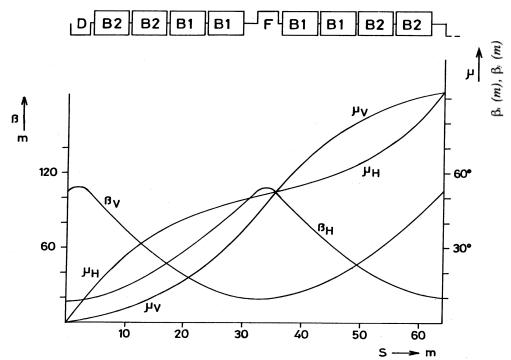
# Optics design

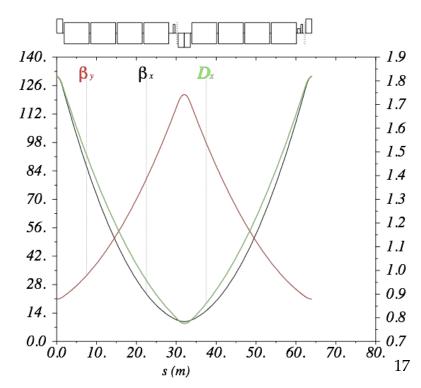
# Basic cell



- **FODO cell** of around **65 m** long with phase advances of  $\pi/2$
- Beta function maxima slightly above 100 m

The 300 GeV Program, CERN/1050, 14/01/1972





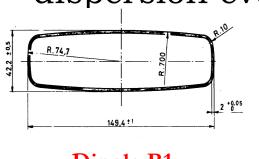


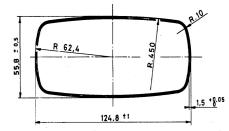


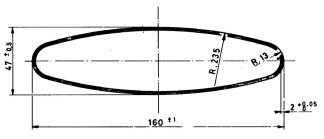
#### Magnet aperture



Magnet apertures follow beta function and dispersion evolution



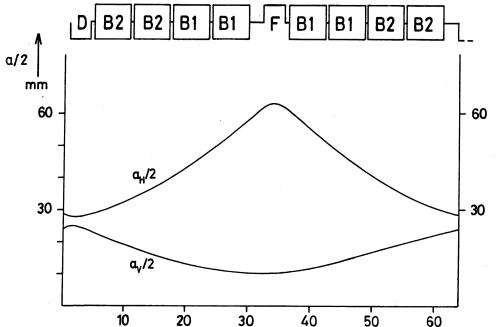


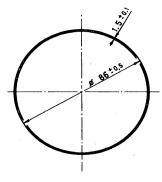




Dipole B2







Quadrupole D

# The CERN Accelerator School

#### Dispersion suppression

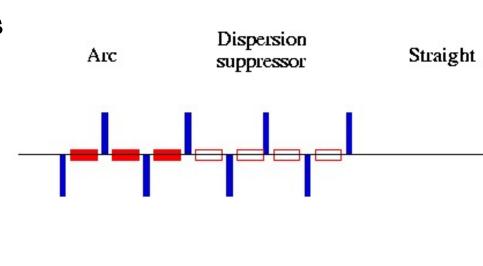


- **Dispersion** has to be **eliminated** in **special areas** like injection, extraction or interaction points (orbit independent to momentum spread)
- Use dispersion suppressors
- Methods for suppressing dispersion
  - □ **Eliminate two dipoles** in a FODO cell (missing dipole)
  - Set last dipoles withdifferent bending angles

$$\theta_1 = \theta (1 - \frac{1}{4\sin^2 \mu_{\text{HFODO}}})$$

$$\theta_2 = \frac{\theta}{4\sin^2 \mu_{\text{HFODO}}}$$

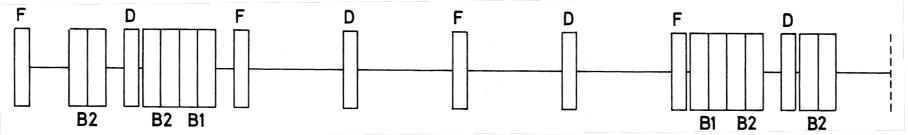
For equal bending angle dipoles, the FODO phase advance should be equal to π/2



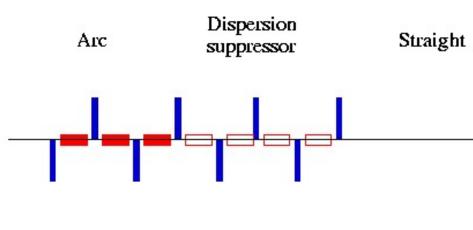


#### Dispersion suppression in the SPS





- In the SPS, all dipoles are powered in series, i.e. dispersion suppressor cells looks like a missing dipole, but they are not!
- Dispersion suppression is achieved by tuning the **phase** advance of the arc, to a multiple of 2 π
- **Dispersion oscillates** through the arc and vanishes at the edges



# a Synchrotron, CERN Accelerator School

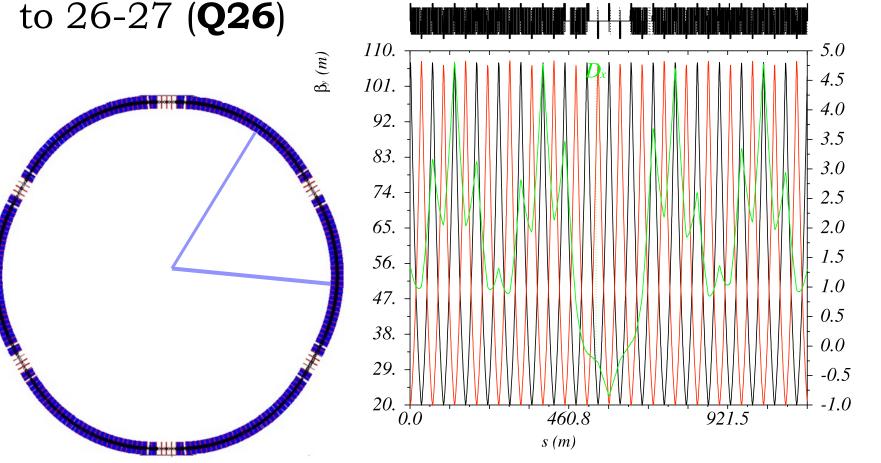
# Ring optics



Ring is composed by 6 identical sectors ("sextants") with 16 arc cells and 2 cells in the straight sections

■ The cell phase advance of  $\pi/2$  brings the tunes

to 26-27 (**Q26**)

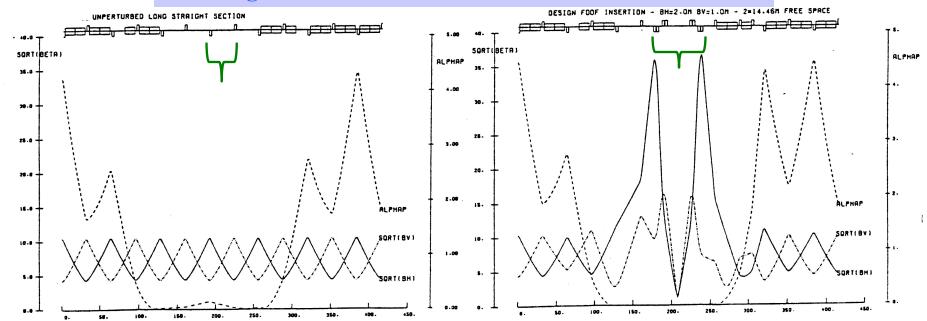


### $Sp\overline{p}S$ collider insertion optics



- Replace two straight section quadrupoles with 2 doublets (4 quadrupoles)
- Equip adjacent left/right quadrupoles with individual bipolar power convertors
- Achieved **low β\*** of 1.3/0.65 m

#### P. Faugeras et al., CERN-SPS-80/11, CERN-SPS-83/29







# Magnet system



## SPS dipole magnets



■ **744 dipoles** (MBAs and MBBs) with 6.26 m length and different gaps



Number of magnets	744
Year of 1 <sup>st</sup> operation	1976
Maximum field on beam axis [T]	2.02
Physical vertical aperture [mm] MBA/MBB	38.5/51.5
Yoke assembly [Solid,Laminated,Welded,Glued]	L,W
Coil technology [Copper,Aluminium,Glass-epoxy,Mica,Other]	C,G
Maximum voltage to ground [V] (worst case 2 spare converters)	4150
Operation	Cycled
Maximum cooling water velocity [m/s]	9
Operational temperature [C°]	40

#### D. Tommasini CERN/TE-Note-2010-003

- Maximum field of2.02 T, for reaching450 GeV
- High mechanical stress on coils



#### SPS quadrupoles



- 216 quadrupoles (102 QF, 100 QD, 6 QFA and 8 QDA)
- Maximum **gradient** of **22 T/m**, corresponding to a pole-tip field of around 1 T
- Normal operation necessitates almost the **full** gradient @ 450 GeV

Tommasini CERN/TE-Note-2010-003

b. 10111111ashii CEITT / 1E-110tt-2010-003			
Number of magnets	216		
Year of 1 <sup>st</sup> operation	1976		
Maximum gradient [T/m]	22		
Physical vertical aperture [mm]	88		
Yoke assembly [Solid,Laminated,Welded,Glued]	L,W		
Coil technology [Copper,Aluminium,Glass-epoxy,Mica,Other]	C,G		
Maximum voltage to ground [V]	3450		
Operation	Cycled		
Maximum cooling water velocity [m/s]	3.6		
Operational temperature [C°]	40		





#### SPS sextupoles



#### M. Giesch, CERN/SPS/80-3/AMS, 1980

MAIN PARAMETERS OF SEXTUP(				LSDN
Basic	: Nominal rms current Peak Current	[A] [A]	350 500	350 450
	* Strength at peak current			
	1) Sextup. $\int a_3  d\ell  (a_3 = B/_{r^2} = B''/2)$	[T/m]	85.8	176.6
		n <sup>2</sup> ]		
	* Magnetic length	[m]	0.435	0.426 44.0
	Aperture, radius of inscr.circle	[mm]	60.7	44.0
Core	: Length	[m]	0.4	0.4



- 54 "focusing" and 54 "defocusing" 0.4 m long sextupoles in two (three for F) families (24 and 30), with different apertures
- Maximum pole-tip field of around 0.8 T
- Around 80% and 60% in operational conditions



# The SPS arc cell











# Transition energy and slippage factor





#### Transition energy



Transition "energy" (or momentum compaction factor) is defined as

$$\frac{1}{\gamma_t^2} = \alpha_p = \frac{1}{C} \oint \frac{D(s)}{\rho(s)} ds$$

The **higher** the **dispersion oscillation** in the bends, the lower the transition energy





#### Transition energy



30

Transition "energy" (or momentum compaction factor) is defined as

$$\frac{1}{\gamma_t^2} = \alpha_p = \frac{1}{C} \oint \frac{D(s)}{\rho(s)} ds$$

The **higher** the **dispersion oscillation** in the bends, the lower the transition energy



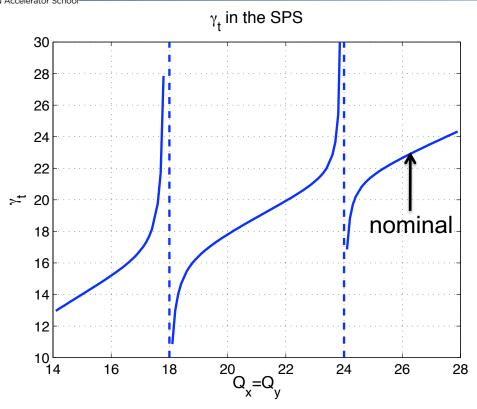
#### Quadrupoles

- Note also that, for FODO cells (SPS lattice),  $\gamma_t pprox Q_x$ , meaning that lowering the transition energy implies lowering the horizontal tune
- High intensity beams can be injected in the SPS above transition avoiding losses and operational complexity of transition jump scheme



#### Transition energy vs SPS working point



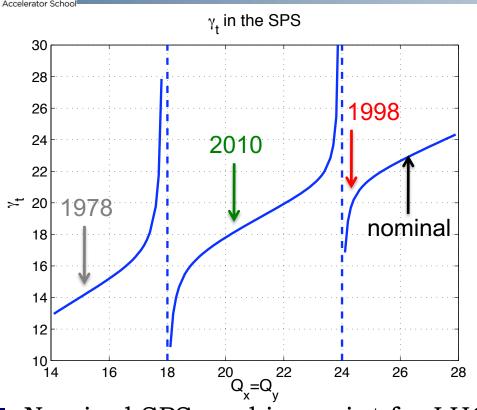


- Resonant oscillation of dispersion function close to the "Resonant integer tunes" (multiples of **super-periodicity 6**)  $\rightarrow$  asymptotic behavior of  $\gamma_{t,}$  (difficult for routine operation)
- lacksquare  $\gamma_t$  is a linear function of horizontal tune  $Q_x$  elsewhere



#### Transition energy vs SPS working point





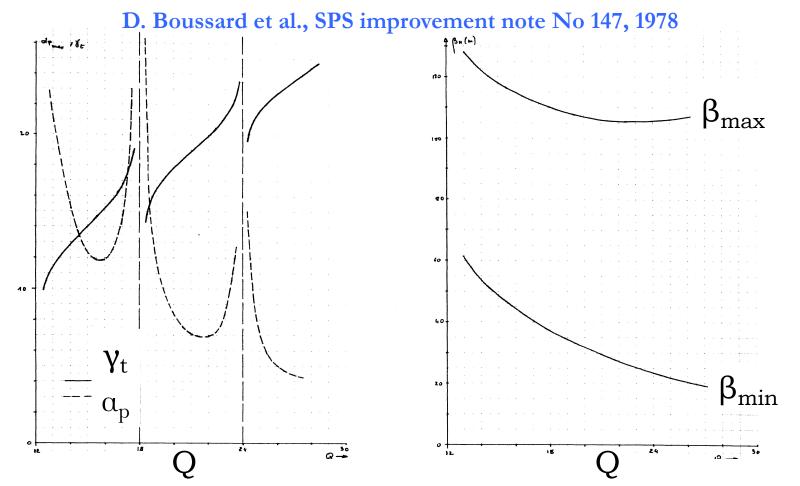
- Resonant oscillation of dispersion function close to the "Resonant integer tunes" (multiples of **super-periodicity 6**)  $\rightarrow$  asymptotic behavior of  $\gamma_{t,}$  (difficult for routine operation)
- lacksquare  $\gamma_t$  is a linear function of horizontal tune  $Q_x$  elsewhere
- Nominal SPS working point for LHC proton beams ( $\gamma_t$ ~23)
- D. Boussard et al., SPS improvement note No 147, 1978; Injection above transition as TT10 was not ready for 26 GeV/c (γ<sub>t</sub>~14)
- G. Arduini et al., CERN/SL-Note 98-001, 1998; "Resonant tune" (γ<sub>t</sub>~20)
- Low  $\gamma_t$ , 2010 "Resonant arc" with small dispersion in long straight sections ( $\gamma_t$ ~18)



#### Avoiding transition energy with Q15



- Injection beam line **TT10** has not been upgraded to 26 GeV in 1978 and limited to **16 GeV**
- Injection above transition is possible if SPS integer part of the tune is lowered to 15 ( $\gamma_t$ ~14)







# Manipulating optics for curing instabilities



#### Instability thresholds and slippage factor



YP et al, IPAC 2013

#### Transverse instabilities

- ☐ **TMCI** at injection single bunch instability in vertical plane
  - Threshold at 1.6x10<sup>11</sup>p/b ( $\epsilon_l$ =0.35eVs,  $\tau$ =3.8ns) with low vertical chromaticity  $N_{\rm th} \propto \frac{\varepsilon_l}{\beta_{\rm re}} \eta$

 $\square$  Threshold higher than  $1.2 \times 10^{11} \text{p/b}$ 

$$N_{
m th} \propto Q_s \propto \sqrt{\eta}$$

#### Longitudinal instabilities

- Single bunch and coupled bunch
  - Threshold at  $2x10^{10}$ p/b for single harmonic RF (800 MHz cavity use is mandatory)

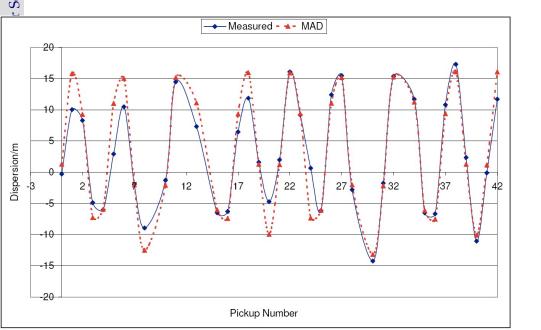
$$N_{th} \propto \epsilon_l^{5/2} \eta$$



#### Resonant tune



- By setting the SPS integer tune to a multiple of 6, large dispersion wave can be introduced (dispersion becomes even negative) by overall reducing transition energy
- Successfully establishing cycle in the SPS and measuring dispersion very close to the one of MAD
- 3-fold increase of the slippage factor can be achieved (model)
- "Difficult" beam conditions (especially for injection)
- Need optics were dispersion is suppressed in straight section



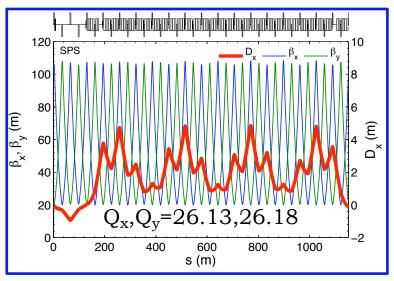
G. Arduini et al., CERN/SL-Note 98-001 (MD), 1998

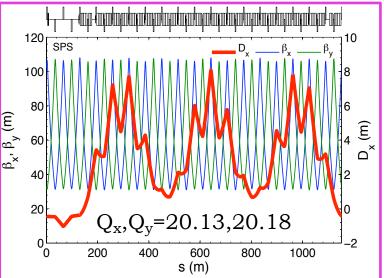
$Q_h$	$Q_v$	$\gamma_{tr}$	$\eta \ (10^{-3})$
24.18	24.22	18.54	1.61
24.29	24.32	19.59	1.30
26.62	26.58	23.23	0.551

### Q20 optics

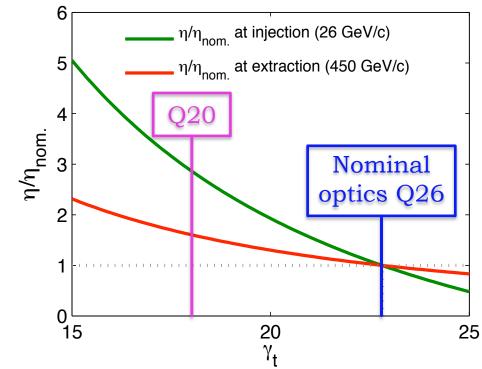


Moving FODO phase advance from  $4/16*2\pi$  ( $\pi/2$ ) to  $3/16*2\pi$  ( $3\pi/8$ )

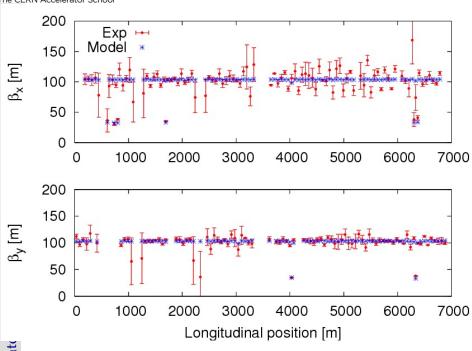


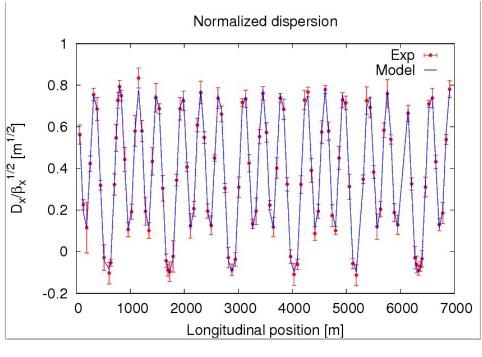


Slippage factor increased by a factor of **2.8** at **injection** and **1.6** at **flat top** 



Slip factor relative to nominal SPS optics 3/





- Measurement of the optics functions of the new lattice
  - **Beta beating** around 20% in horizontal and 10% in vertical plane
  - Normalized **dispersion** in striking agreement with the model

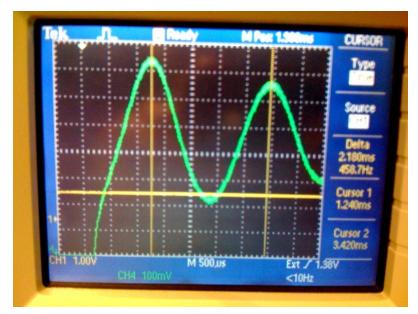
Designing a Synchrotron, CERN Accelerat



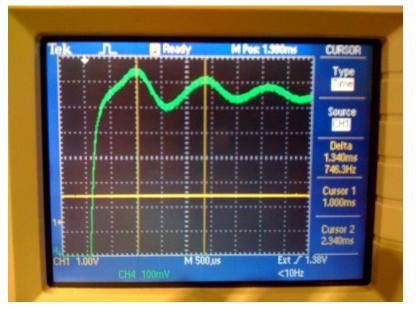
### Synchrotron frequency



- ☐ Measured synchrotron frequency from "quadrupole" oscillations at injection
  - Same RF-voltage for both optics
- □ Ratio of Synchrotron frequencies ~ **1.63** corresponds to an **increase** in slippage factor  $\eta$  by **factor 2.65** (MADX prediction: 2.86)



Q26: Fs=458/2=229Hz, Qs=0.0106/2=0.0053



Q20: Fs=746/2=373Hz, Qs=0.0172/2=0.0086

### TMCI threshold



0.00 10.00 I growth rate (1/turns)

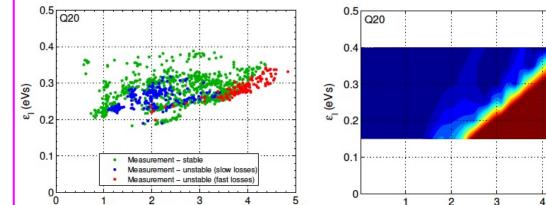
N (p/b)

- ☐ In **nominal optics**, measured/simulated threshold at 1.6x10<sup>11</sup>p/b for low chromaticity
  - High-chromaticity helps increasing threshold, but also losses along the cycle become excessive
- $\square$  Measured/simulated threshold in  $Q_{20} > 4x10^{11}p/b!!!$

N(p/b)

$$N_{
m th} \propto rac{arepsilon_l}{eta_y} \eta$$

Q26 0.0 10.00 I growth rate (1/turns) x 10<sup>11</sup> N(p/b)x 10<sup>11</sup> N(p/b)



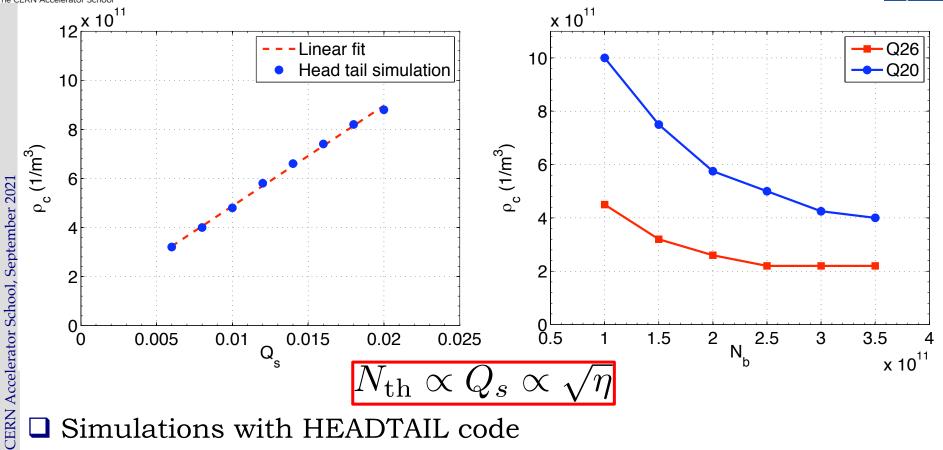
x 10<sup>11</sup>

H. Bartosik et al, **IPAC 2014** 



### E-cloud instability





- Simulations with HEADTAIL code
  - Injection energy, uniform cloud distribution, located in dipole regions
- Linear scaling with Synchrotron tune demonstrated
  - Clearly higher thresholds predicted for **Q20**

More margin with Q20 if e-cloud becomes issue for high intensity

H. Bartosik et al, IPAC2011

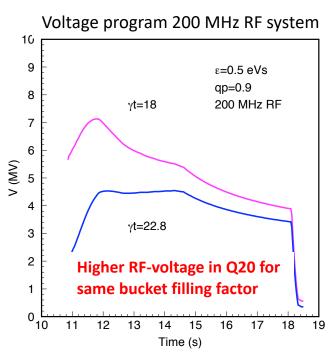


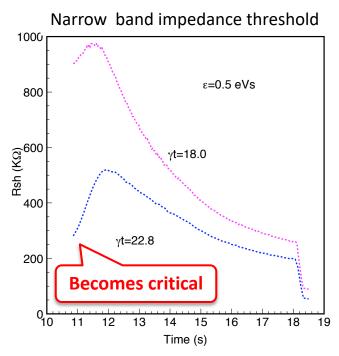
#### Longitudinal impedance threshold



$$N_{th} \propto \epsilon_l^{5/2} \eta$$

#### E. Shaposhnikova





- Impedance threshold has minimum at flat top
  - ☐ Controlled longitudinal emittance blow-up during ramp for **Q26**
  - ☐ Less (or no) longitudinal emittance blow-up needed in Q20
- ☐ Instability limit at flat bottom
  - Critical with Q26 when pushing intensity
  - Big margin with Q20 (factor of 3)

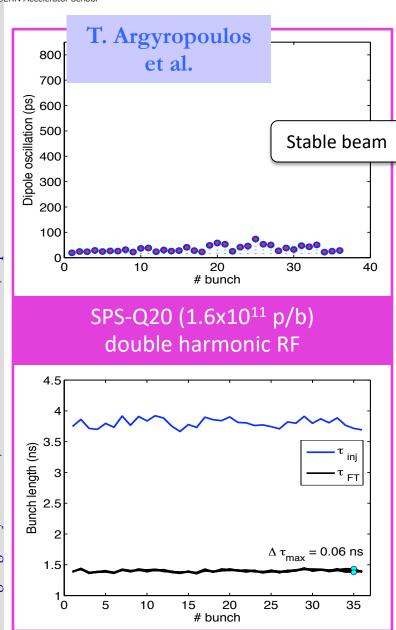


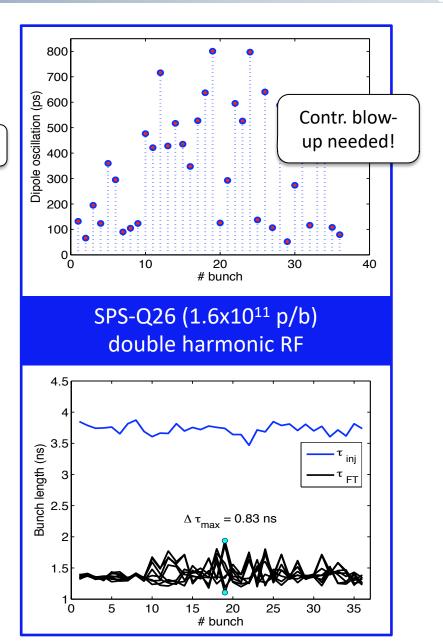
a Synchrotron, CERN Accelerator School, September 2021

Designing

### Congitudinal beam stability







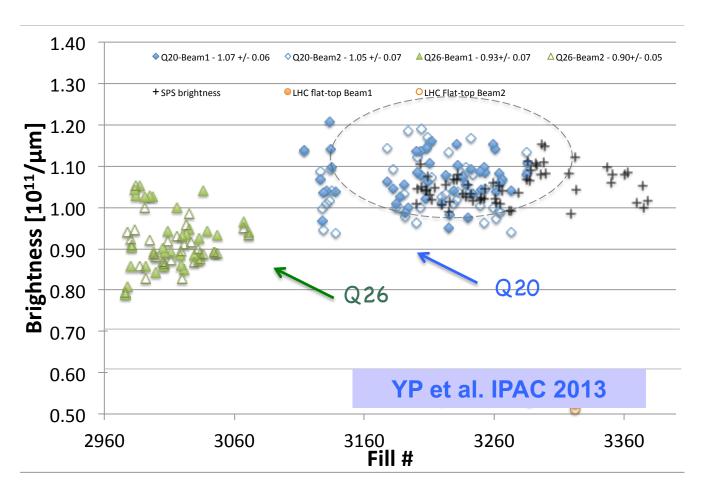




#### LHC brightness with SPS Q20



- Operational deployment of Q20 optics for LHC beams since 2012 allowing around 20% brighter beams on LHC flat bottom
- Opened way for ultra-high brightness beams of HL-LHC era





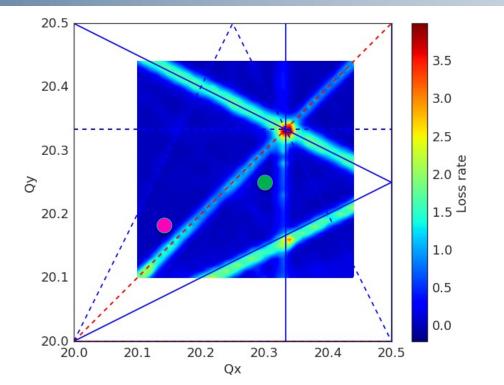


# Non-linear dynamics



### Loss map for low brightness beam





- proton working point
- ion working point

#### resonances:

red: systematic

blue: non-systematic

- upright
- - skew

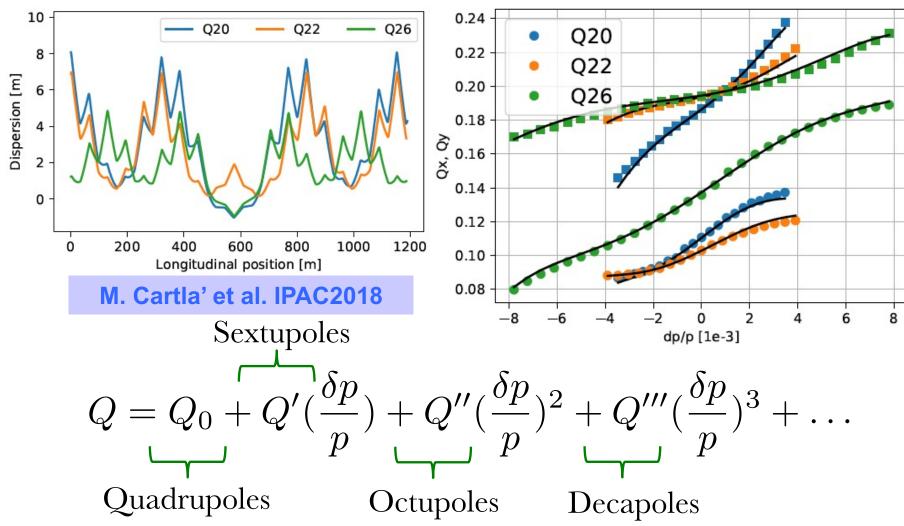
H. Bartosik et al. HB2018

- Dynamic tune scan for identification of resonances
  - □ Losses around 3<sup>rd</sup> order (normal) resonances and the diagonal clearly observed
  - Faint traces of 4<sup>th</sup> order resonances
  - Operational working point for protons 20.13/20.18 (moved up for high brightness beams)

### COP

#### Non-linear model through chromaticity





Estimate "effective" magnet multi-poles that reproduce non-linear chromaticity measurement for three different optics





# Space-charge

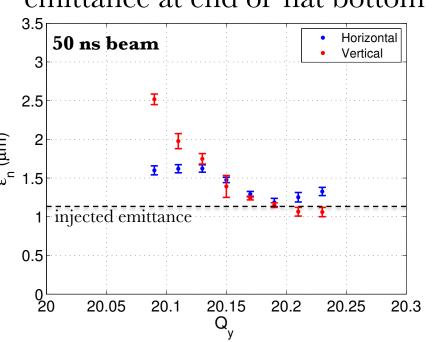
# The CERN Accelerator School

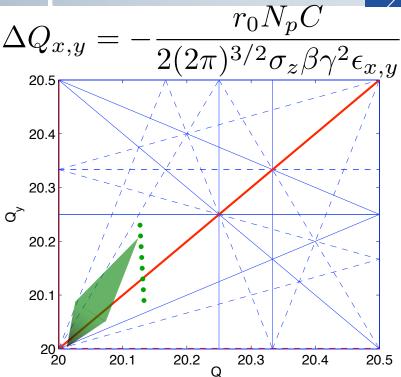
Designing a Synchrotron, CERN Accelerator School, September 2027

#### Space-charge tune spread



emittance at end of flat bottom





- **Vertical tune scan** with high brightness beam for 10 s storage time
  - $\square$  N = 1.95x10<sup>11</sup> p/b (at injection)
  - $\supseteq \varepsilon \sim 1.1 \ \mu m \ (at injection)$

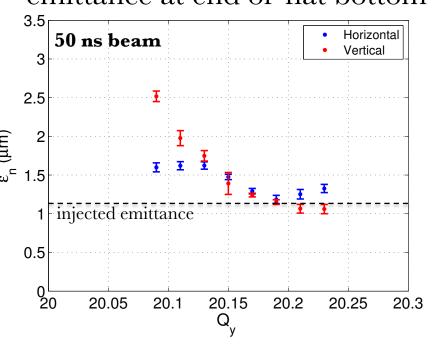
- $\Delta Q_x/\Delta Q_y \sim 0.10/0.20$
- □ Transmission to flat top around 94% (very small losses on flat bottom)
- Budget of 10% losses and 10% blow-up allows for tune spread of  $\Delta Qy$ =0.21

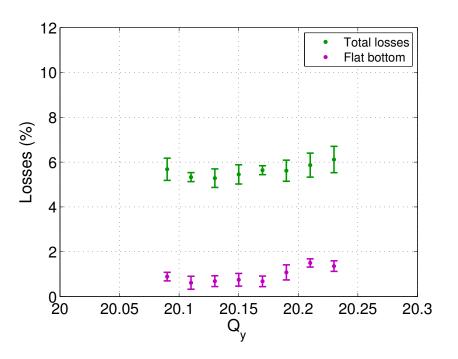
Designing a Synchrotron, CERN Accelerator School, September 2027

### Space-charge tune spread



emittance at end of flat bottom





- Vertical tune scan with high brightness beam for 10 s storage time
  - $\square$  N = 1.95x10<sup>11</sup> p/b (at injection)
  - $\epsilon \sim 1.1 \, \mu \text{m} \text{ (at injection)}$

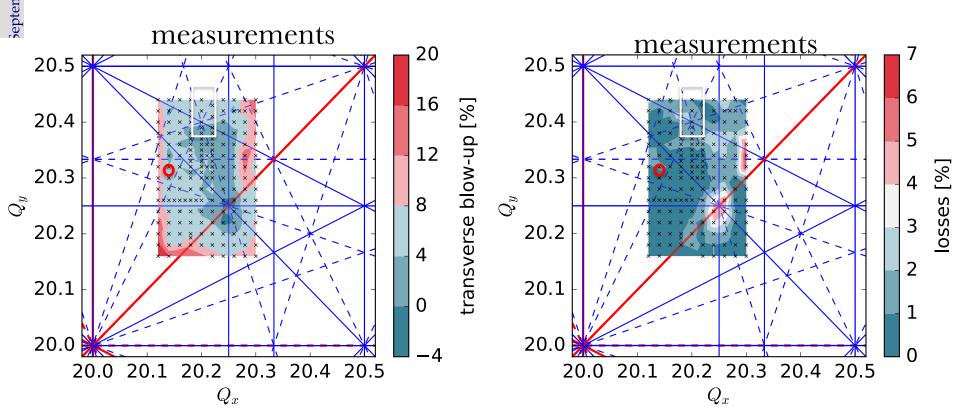
- $\Delta Q_x/\Delta Q_y \sim 0.10/0.20$
- □ Transmission to flat top around 94% (very small losses on flat bottom)
- $\blacksquare$  Budget of 10% losses and 10% blow-up allows for tune spread of  $\Delta Qy{=}0.21$



#### Exploration of tune diagram with SC



- Tune scan with high brightness single bunch beam for 3 s storage time
  - □ Blow-up at integer resonances as expected (tune spread  $\Delta Qx$ ,  $\Delta Qy \sim 0.10,0.19$ )
  - Margin for higher brightness for working points in white box (enhanced losses only close to Qx + 2Qy = 61 normal 3<sup>rd</sup> order resonance and around 4Qx = 81 normal 4<sup>th</sup> order resonance)







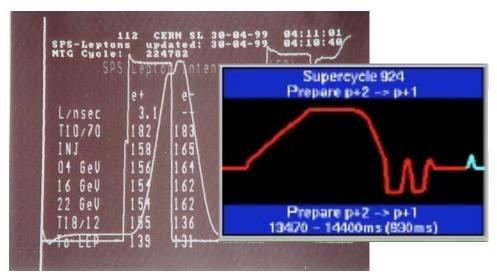
# Electron-positron dynamics



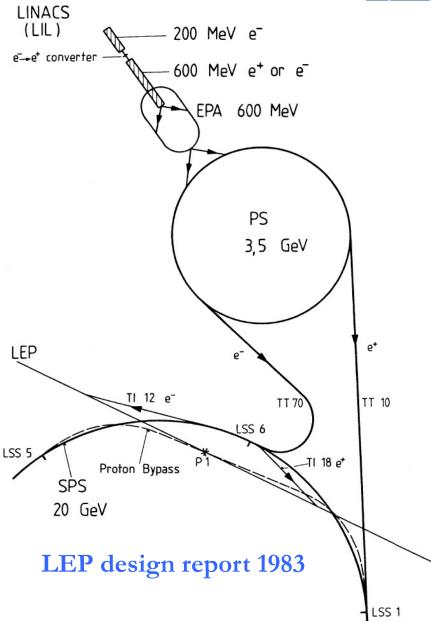
#### SPS as LEP Injector



#### P. Collier – Academic Training 2005

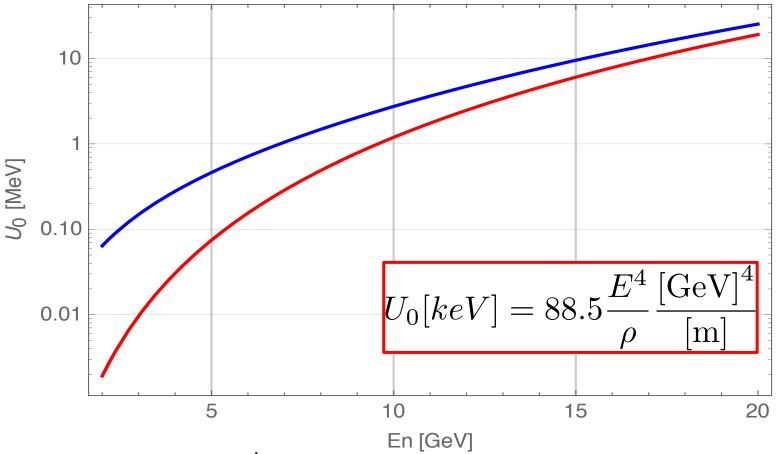


- LEP filling interleaved with proton operation
- 4 cycles with 4 bunches (2e<sup>+</sup>, 2e<sup>-</sup>) evolved to 2 cycles with 8 bunches (~2.5x10<sup>10</sup> p/b)
- Energy to LEP:  $18 \rightarrow 20 \rightarrow 22$  GeV
- Lots of RF for leptons (200MHz SWC, 100MHz SWC, 352MHz SC),
- 2 Extractions in Point 6 towards LEP



## Energy loss/turn

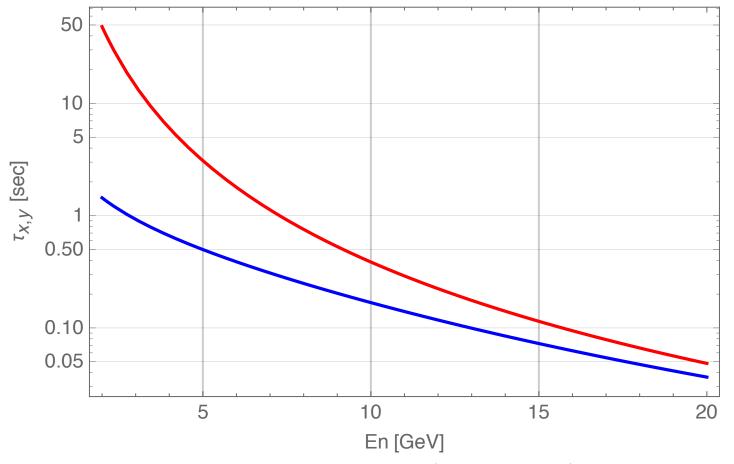




- Energy loss/turn necessitate large RF voltage (30 MV) at high energy
- Impact of a 2-m **3.5T** damping wiggler is mild at high energies

### Con Damping time





- Damping time at injection (3.5 GeV) very large (9 s)
- A 2-m **3.5T damping wiggler** could enhance damping for low energies to below 1 s (good for instabilities)

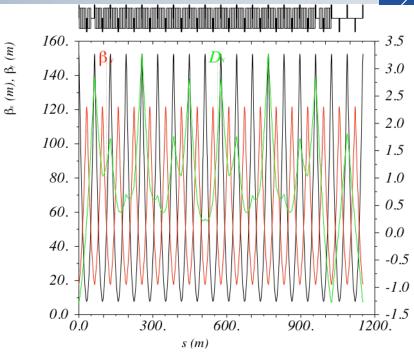
### The CERN Accelerator School

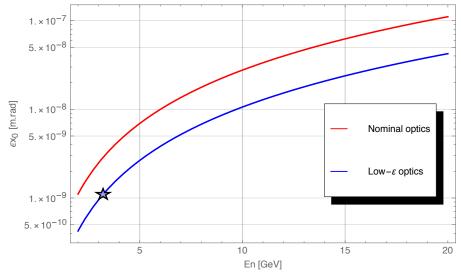
### SPS low emittance optics



Move horizontal phase advance to 135 deg. i.e.  $3\pi/4$  (Q40 optics) which is optimal for low emittance in a FODO cell

- Emittance with nominal optics @ 3.5 GeV of 3.4 nm drops to 1.3nm
- Further reduction can be achieved with damping wiggler





# Summary The CERN Accelerator School



- Using the 40+ years experience since the design and operation of the Super Proton Synchrotron (SPS), reviewed several beam dynamics concepts
  - Choice of basic parameters
    - ■Energy, bending field and circumference
  - Optics design
    - ■Cell optics, insertions, transition energy
  - Collective effects
    - ■Instabilities, Space-charge
  - □ Electron/Positron beam dynamics
    - Equilibrium beam properties, energy loss/turn, damping time