

# Nonlinearities in light sources

**or:** *How to correct the chromaticity without destroying the dynamic aperture in a high brightness light source lattice*

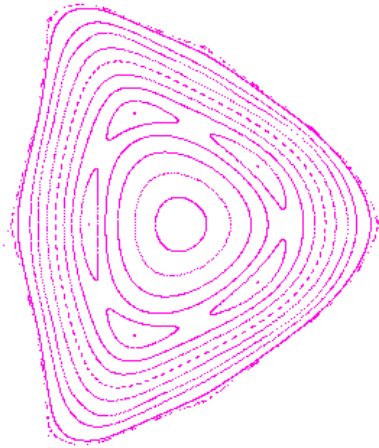
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## Contents:

- The source of chromaticity
- Chromaticity correction with sextupoles
- Breakdown of dynamic acceptance: investigation of sextupole effects
- First order optimization of the sextupole pattern
- Particular light source problem
- Second order optimization: octupole terms

## Nonlinearities in Light Sources



by design:

**Sextupoles**

for chromaticity correction

**The RF bucket**

→ **how to deal with?**



parasitic:

multipolar errors:

→ ring magnets

→ insertion devices

→ **define tolerances**

Nonlinearities → **CHAOS** → separatrix for bounded motion (= dynamic acceptance)

### Light source problems:

- Transverse acceptance **bad** [i.e. separatrix < beam pipe] → **no injection**
- Lattice energy acceptance (= energy dependant transverse acceptance) **bad** [i.e. < RF energy acceptance] → **low beam lifetime**

# Methods for nonlinear optimization

## analytical

1<sup>st</sup> & 2<sup>nd</sup> order perturbation theory  
→ *maps, resonance drive terms, tune shifts with amplitude ...*

- ⊕ quick calculation
- ⊕ interactive optimization  
(semi-analytic minimization)
- ⊖ no prediction of performance  
(perturbation!)  
→ “*the art of weighting...*”

⇒ **Design**

## numerical

Particle tracking  
→ *Poincaré plots, dynamic aperture scans, particle spectra...*

- ⊖ slow calculation
- ⊖ difficult to use in minimizer  
(fractal parameter space!)
- ⊕ valid prediction of performance  
→ complete & correct model

⇒ **Proof**

⇔  
**Iterate**  
⇔

# Chromaticity

## Quadrupole:

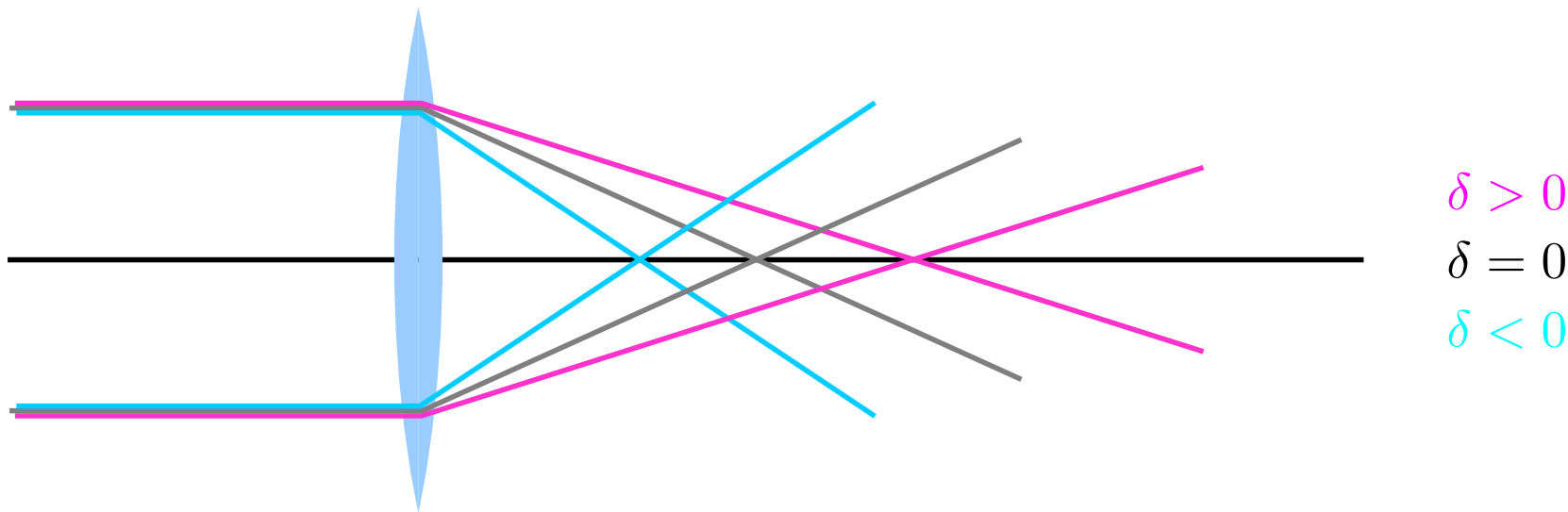
Length  $L$

Strength  $b_2 = \frac{1}{(B\rho)} \frac{dB_y}{dx}$

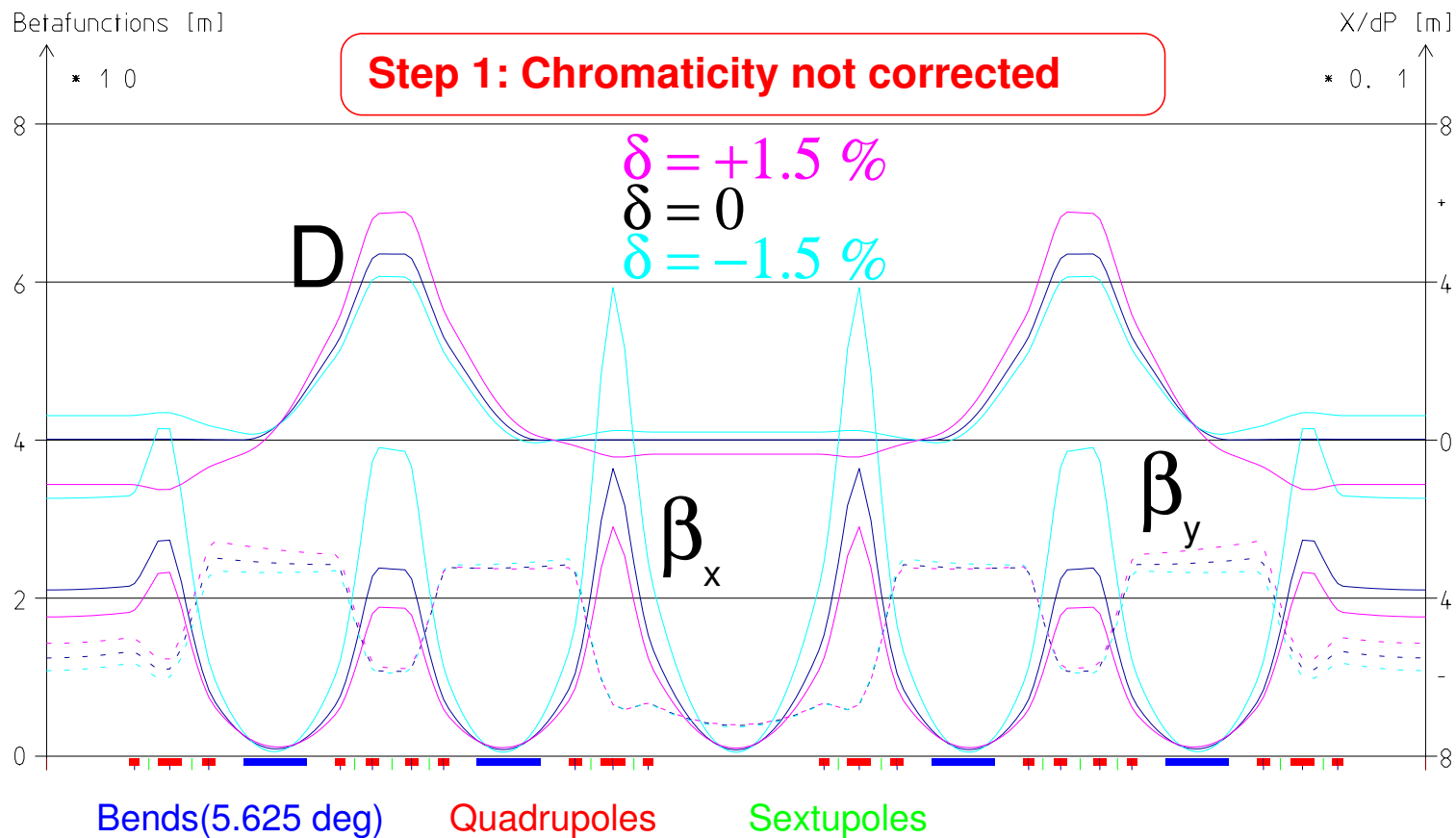
$$(B\rho) := \frac{p}{e} = 3.3356 \text{ Tm} \cdot E[\text{GeV}]$$

Kicks on particle:  $\Delta x' = -b_2 L x$     $\Delta y' = b_2 L y$    ( $b_2 > 0 \rightarrow$  horiz.foc.)

Chromatic aberration:  $b_2(\delta) = \frac{b_2}{(1+\delta)} \approx b_2 (1 - \delta)$     $\delta := \frac{\Delta p}{p}$



# Chromaticity



16*L = 845.99 m (periodic)	$\delta = 0$	Qx = 35.4503	Qz = 11.3976	Ex = 8.17E-0009 @ 6 GeV
dp/p = -1.50 %		Cx = -96.1574	Cz = -30.4910	Al = 0.00028
dp/p = 1.50 %		Qx = 37.2591	Qz = 11.8640	$\delta = -1.5\%$
		Qx = 34.1891	Qz = 10.9465	$\delta = +1.5\%$

Test lattice: ESRF with original optics for dispersion free straights

# Chromaticity

## Chromaticity

(Gradient error  $\Delta b_2 ds$ )  $\times$  (one turn matrix  $\mathcal{M}$ ) = (new one turn matrix  $\tilde{\mathcal{M}}$ )<sup>a</sup>

Gradient error due to chromatic aberration:  $\pm \Delta b_2 = \mp b_2 \delta$  (hor./vert.)

$$\begin{pmatrix} 1 & 0 \\ \pm b_2 \delta ds & 1 \end{pmatrix} \times \begin{pmatrix} \cos 2\pi Q & \beta \sin 2\pi Q \\ -\frac{\sin 2\pi Q}{\beta} & \cos 2\pi Q \end{pmatrix} = \begin{pmatrix} \cos 2\pi \tilde{Q} & \beta \sin 2\pi \tilde{Q} \\ -\frac{\sin 2\pi \tilde{Q}}{\beta} & \cos 2\pi \tilde{Q} \end{pmatrix}$$

$$\frac{1}{2} \text{Tr}(\tilde{\mathcal{M}}) = \cos 2\pi \tilde{Q} = \cos 2\pi(Q + \Delta Q) = \cos 2\pi Q \pm \frac{1}{2} b_2 \delta \beta \sin 2\pi Q ds$$

$$\Delta Q \ll 1 \longrightarrow \Delta Q = \mp \frac{1}{4\pi} b_2 \delta \beta ds \quad \xi = \frac{\Delta Q}{\delta} = \mp \frac{1}{4\pi} \oint_C b_2(s) \beta(s) ds$$

**Light source:** Low emittance:  $\epsilon \downarrow \rightarrow \int b_2 \uparrow \rightarrow \xi_x \uparrow$  ( $\xi_x \approx -50 \dots -100$ )

$\Rightarrow$  **Head tail instability**<sup>b</sup>

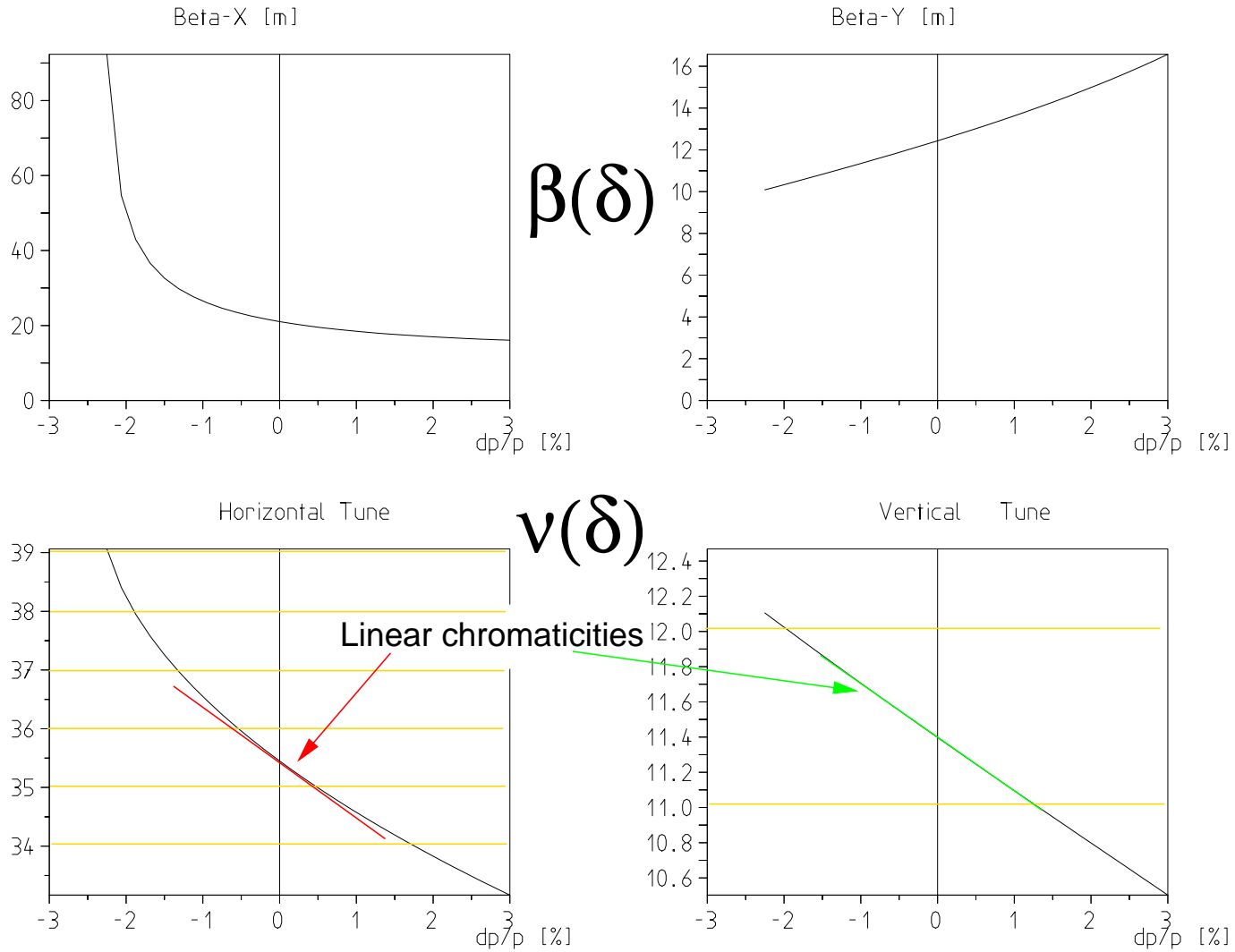
$\Rightarrow$  **Low energy acceptance**

<sup>a</sup>E. D. Courant and H. S. Snyder, The alternate gradient synchrotron, Ann. Phys. 3

<sup>b</sup>A.W.Chao, Coherent instabilities of a relativistic bunched beam, AIP Conf. Proc. 105 (1982) 353

# Chromaticity

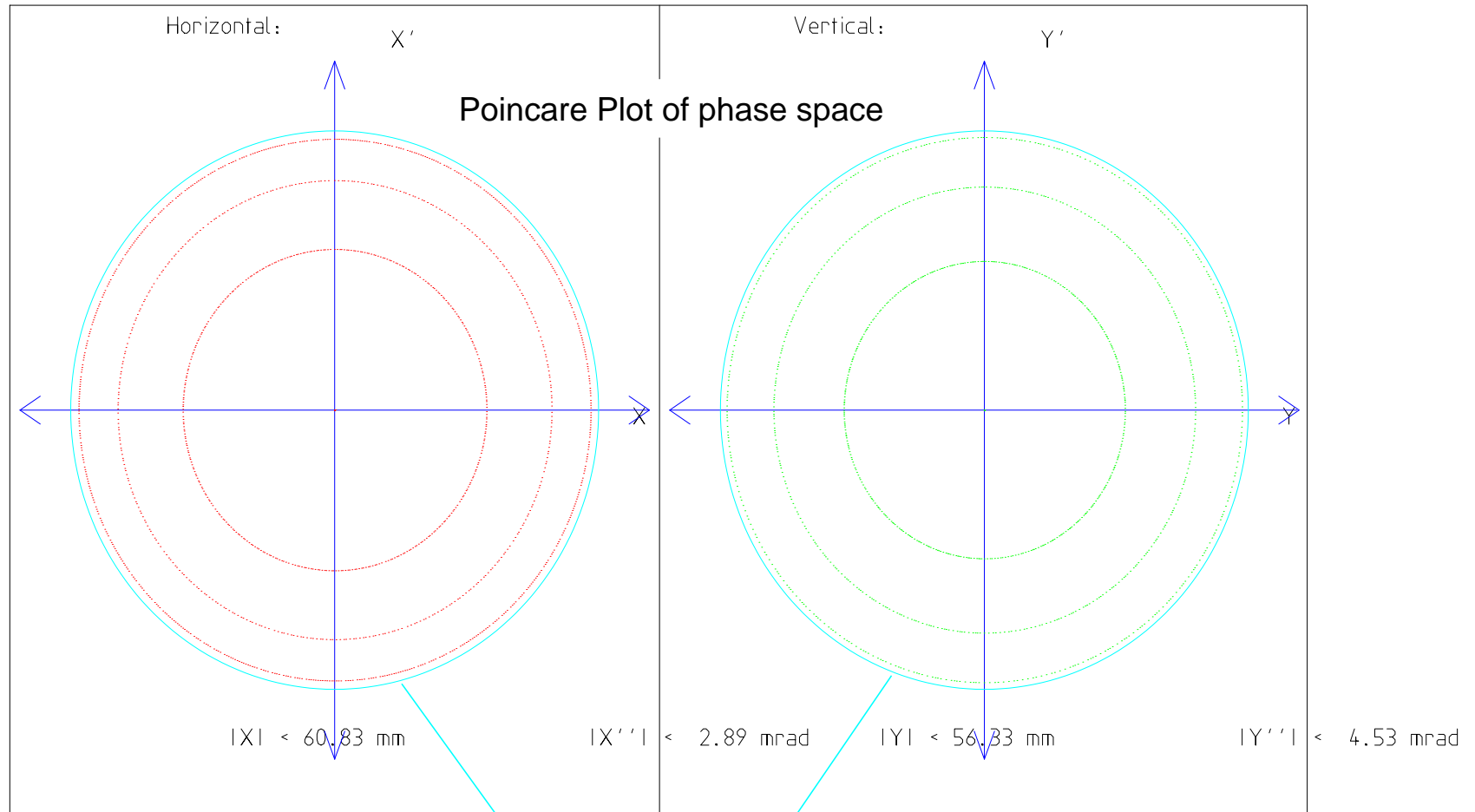
## Step 1: Chromaticity not corrected



Nonlinearities in Light Sources

# Chromaticity

## Step 1: No chromaticity correction: Linear betatron motion



Limitation from 160 mm x 160 mm vacuum chamber

Nonlinearities in Light Sources



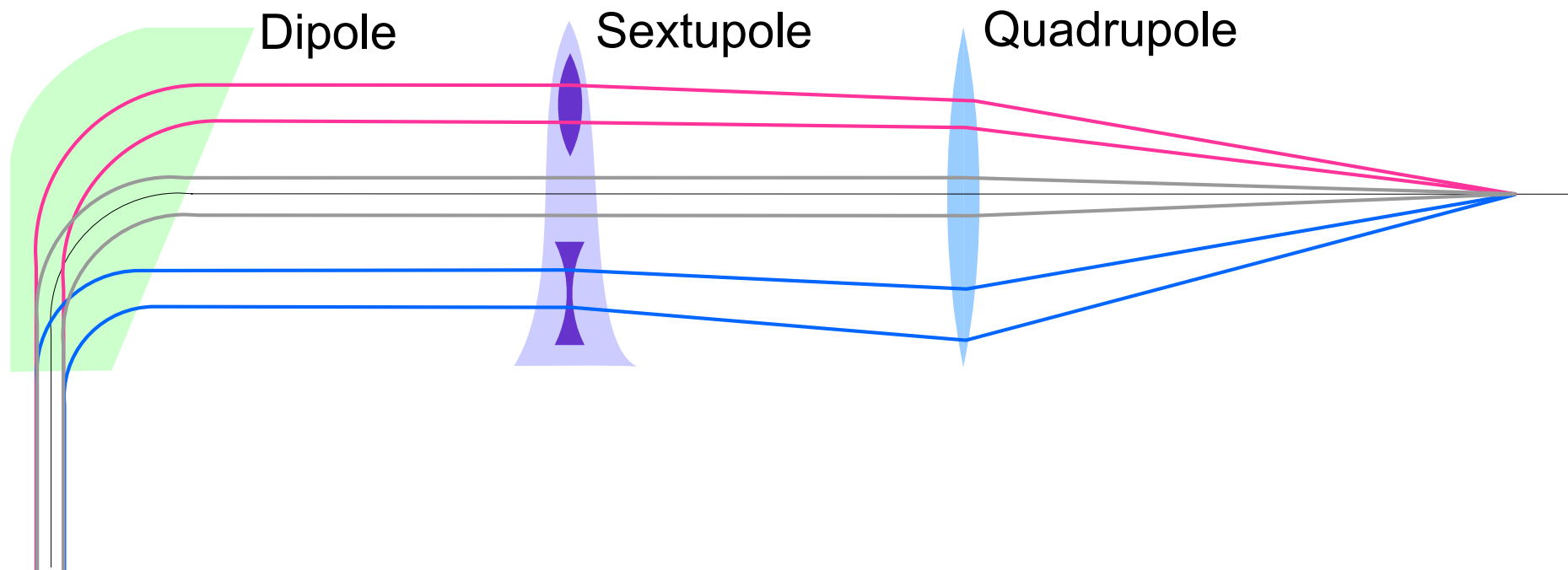
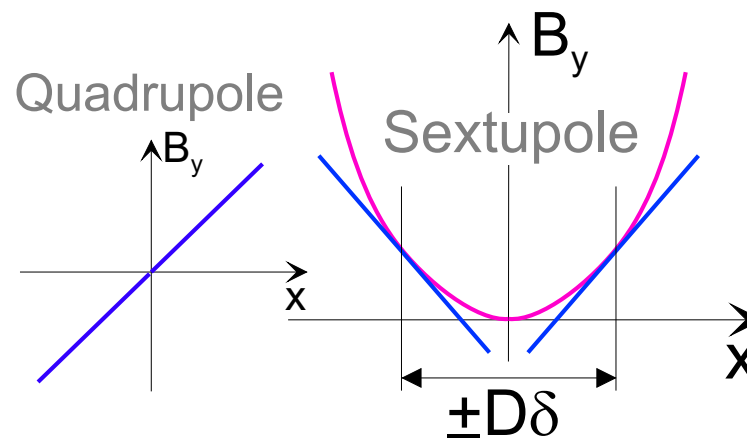
# Chromaticity correction

## Chromaticity correction

Sextupole:  $B_y(x) = \frac{1}{2} B'' x^2$

local gradient:  $B'_y(x) = B'' x$

“Order” by momentum:  $x(\delta) = D\delta$



Nonlinearities in Light Sources

## Chromaticity correction

**Quadrupole:**  $b_2 = \frac{1}{(B\rho)} \frac{dB_y}{dx}$

**Sextupole:**  $b_3 = \frac{1}{2} \frac{1}{(B\rho)} \frac{d^2 B_y}{dx^2}$

$$\Delta x' = -b_2 L x$$

$$\Delta x' = -b_3 L (x^2 - y^2)$$

$$\Delta y' = b_2 L y$$

$$\Delta y' = 2b_3 L x y$$

Chromatic aberrations:  $b_n(\delta) = b_n / (1 + \delta) \approx b_n (1 - \delta)$

Sextupoles in dispersive regions:  $x \rightarrow D\delta + x \quad y \rightarrow y$

Kicks on a particle (keep up to second order in products of  $x, y, \delta$ ):

Quadrupole:  $\Delta x' = -b_2 L x + [b_2 L] \delta x \quad \Delta y' = +b_2 L y - [b_2 L] \delta y$

Sextupole:  $\Delta x' = -[2b_3 L D] \delta x - b_3 L (x^2 - y^2) - b_3 L D^2 \delta^2$

$$\Delta y' = +[2b_3 L D] \delta y + 2b_3 L x y$$

$\Rightarrow$  ☺ Chromaticity correction for  $(2b_3 L D \stackrel{!}{=} b_2 L)$ :

$\Rightarrow$  ∞ nonlinear kicks...

## Chromaticity correction

$$\begin{aligned}\xi_{x/y} &= \pm \frac{1}{4\pi} \oint_C [2b_3(s)D(s) - b_2(s)] \beta_{x/y}(s) ds \\ &= \frac{1}{4\pi} \left( \pm \sum_{\text{sext}} 2(b_3L)_n \beta_{(x/y)n} D_n \mp \sum_{\text{quad}} (b_2L)_n \beta_{(x/y)n} \right) \stackrel{!}{=} 0 \quad [+ \Delta\xi_{x/y}]\end{aligned}$$

Linear system: **2 families** of sextupoles **SF, SD**

$$\begin{aligned}& \frac{1}{2\pi} \begin{pmatrix} + \sum_{n \in \text{SF}} \beta_{xn} D_n & + \sum_{n \in \text{SD}} \beta_{xn} D_n \\ - \sum_{n \in \text{SF}} \beta_{yn} D_n & - \sum_{n \in \text{SD}} \beta_{yn} D_n \end{pmatrix}_{2 \times 2} \times \begin{pmatrix} (b_3L)_{\text{SF}} \\ (b_3L)_{\text{SD}} \end{pmatrix}_{1 \times 2} \\ &= \frac{1}{4\pi} \begin{pmatrix} + \sum_{\text{Quad}} (b_2L) \beta_{xn} \\ - \sum_{\text{Quad}} (b_2L) \beta_{yn} \end{pmatrix}_{1 \times 2} \begin{bmatrix} + \Delta\xi_x \\ \Delta\xi_y \end{bmatrix} \Rightarrow (b_3L)_{\text{SF}}, (b_3L)_{\text{SD}}\end{aligned}$$

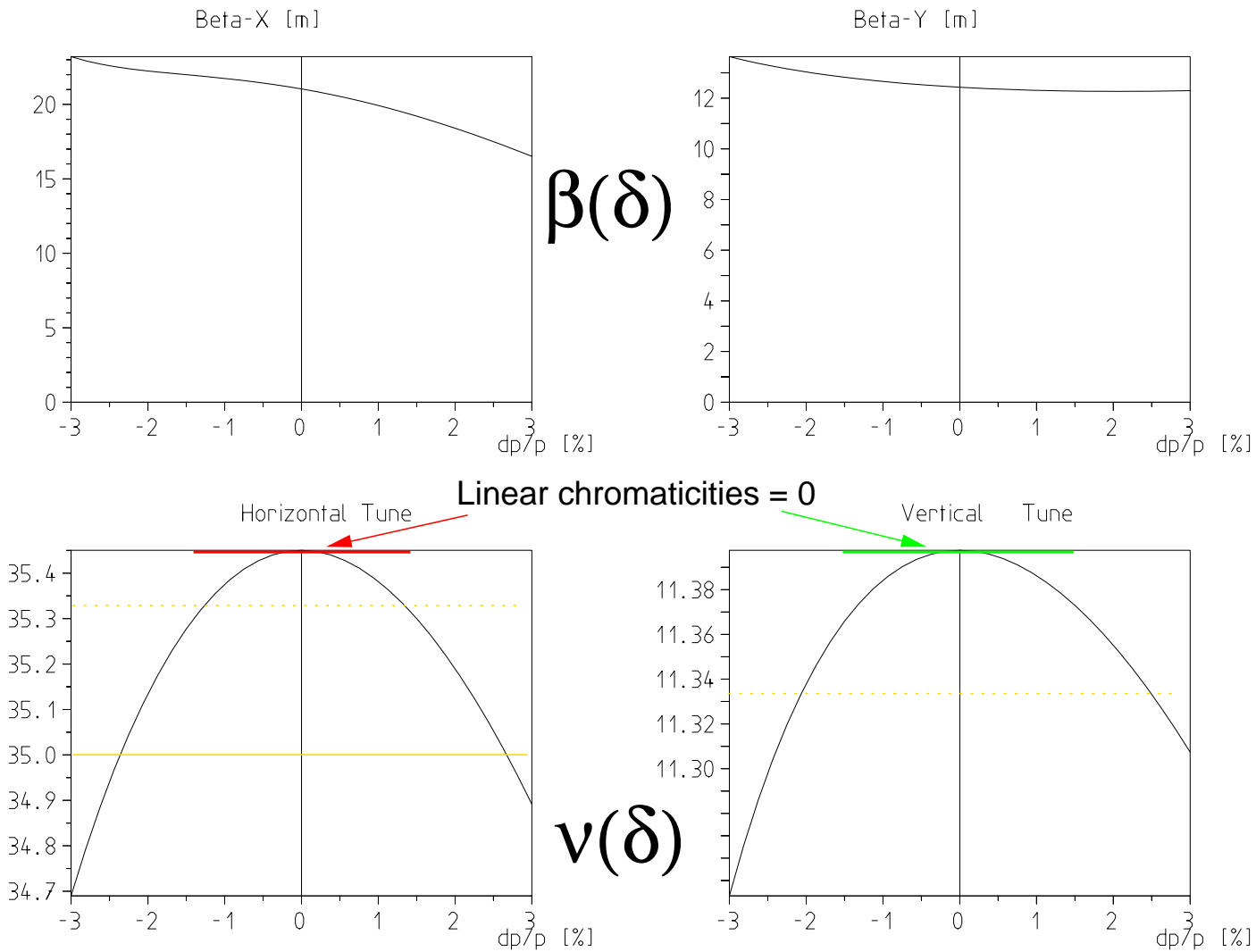
**Decoupling:**

→ SF locations:  $D \uparrow, \beta_x \uparrow, \beta_y \downarrow$

→ SD locations:  $D \uparrow, \beta_x \downarrow, \beta_y \uparrow$

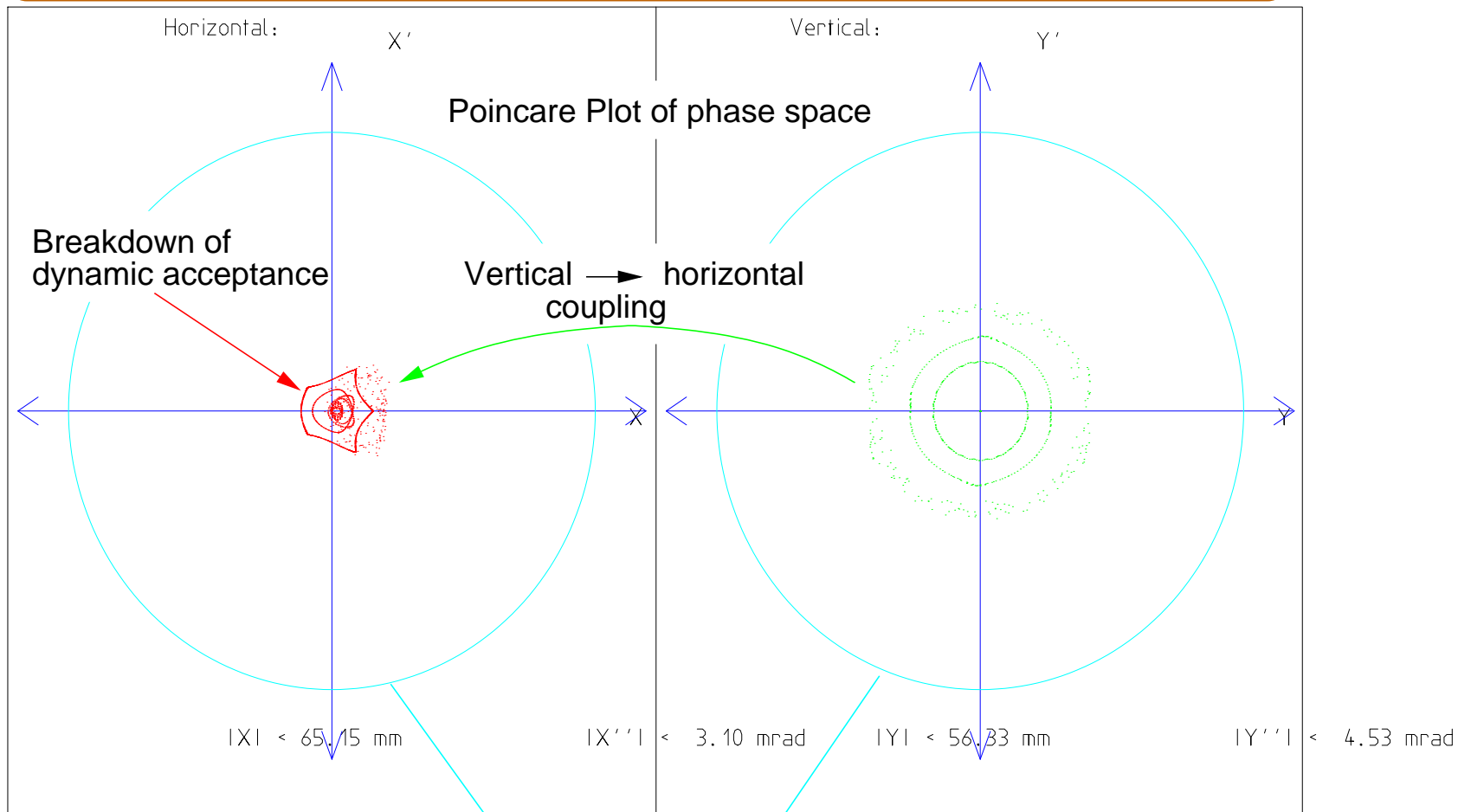
# Chromaticity correction

## Step 2: Chromaticity correction with 2 sextupole families



# Chromaticity correction

## Step 2: Chromaticity correction with 2 sextupole families



Limitation from 160 mm x 160 mm vacuum chamber

## The Hamiltonian

$$H(s) = \underbrace{\frac{p_x^2 + p_y^2}{2(1 + \delta)}}_{\text{kinetic}} - \underbrace{b_1 x \delta}_{\text{dispersive}} + \underbrace{\frac{b_1^2}{2} x^2}_{\text{focussing}} + \underbrace{\frac{b_2}{2} (x^2 - y^2)}_{H_2(s)} + \underbrace{\frac{b_3}{3} (x^3 - 3xy^2)}_{H_3(s)} + \dots$$

Hamiltonian equations of motion (local):

$$x'(s) = \frac{\partial H}{\partial p_x} = \frac{p_x}{1 + \delta} \quad p'_x(s) = -\frac{\partial H}{\partial x} = b_1 \delta - (b_1^2 + b_2)x + b_2 y - b_3(x^2 - y^2)$$

$$\rightarrow \text{e.g. kick from quad. (length } L\text{): } \Delta x' = \int_Q x'' ds = \int_Q \frac{p'_x}{1 + \delta} ds = -\frac{b_2}{1 + \delta} L x$$

**Goal:** Global Quad. & Sextupole contribution:  $\int_{\text{cell}} [H_2(s) + H_3(s)] ds$

→ independant of  $\delta$ : chromaticity corrected

→ linear and uncoupled ( $\sim x^2, y^2$ ): cancellation of nonlinear kicks

## First order sextupole optimization

Insert betatron oscillations (global, linear, flat lattice):

$$x(s) = \sqrt{2J_x \beta_x(s)} \cos \phi(s) + D(s) \delta \quad y(s) = \sqrt{2J_y \beta_y(s)} \cos \phi(s)$$

$$\Rightarrow \int_{\text{cell}} [H_2(s) + H_3(s)] ds = \sum h_{jklmp} \text{ with}$$

$$h_{jklmp} \propto \sum_n^{N_{\text{sext}}} (b_3 L)_n \beta_{xn}^{\frac{j+k}{2}} \beta_{yn}^{\frac{l+m}{2}} D_n^p e^{i\{(j-k)\phi_{xn} + (l-m)\phi_{yn}\}} - \left[ \sum_n^{N_{\text{quad}}} (b_2 L)_n \beta_{xn}^{\frac{j+k}{2}} \beta_{yn}^{\frac{l+m}{2}} e^{i\{(j-k)\phi_{xn} + (l-m)\phi_{yn}\}} \right]_{p \neq 0}$$

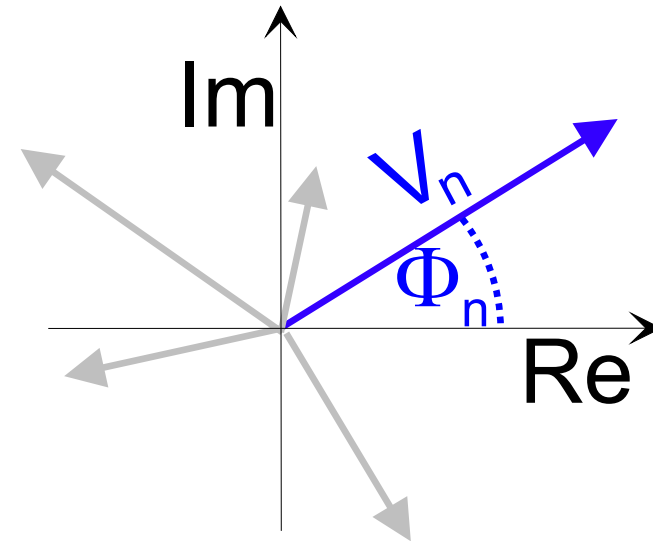
$$h = \sum_n^{N_{\text{sext}}} V_n e^{i\Phi_n} [ + \dots \text{quads for } p \neq 0 \dots ]$$

Sextupole<sub>n</sub> ↔ complex vector:

Length  $V_n = V_n(b_3, L, \beta_x, \beta_y, D)$

Angle  $\Phi_n = \Phi_n(\phi_x + \phi_y)$

- $\Phi_n = 0 \quad \forall n \rightarrow$  tune shifts
- $\Phi_n \neq 0 \rightarrow$  resonances



## First order sextupole [+quadrupole] Hamiltonian

- 2 phase independent terms  $\rightarrow$  chromaticities:

$$h_{11001} = +J_x \delta \left[ \sum_n^{N_{sext}} (2b_3 L)_n \beta_{xn} D_n - \sum_n^{N_{quad}} (b_2 L)_n \beta_{xn} \right] \rightarrow \xi_x$$

$$h_{00111} = -J_y \delta \left[ \sum_n^{N_{sext}} (2b_3 L)_n \beta_{yn} D_n - \sum_n^{N_{quad}} (b_2 L)_n \beta_{yn} \right] \rightarrow \xi_y$$

- 7 phase dependant terms  $\rightarrow$  resonances:  $h^N := h$  for  $N$  cells,  $N \rightarrow \infty \implies$

$$|h_{jklmp}^\infty| = \frac{|h_{jklmp}|}{2 \sin \pi [a_x Q_x^{\text{cell}} + a_y Q_y^{\text{cell}}]}$$

$$a_x = (j - k) \quad a_y = (l - m)$$

$$h_{21000} = h_{12000}^* \rightarrow Q_x$$

$$h_{30000} = h_{03000}^* \rightarrow 3Q_x$$

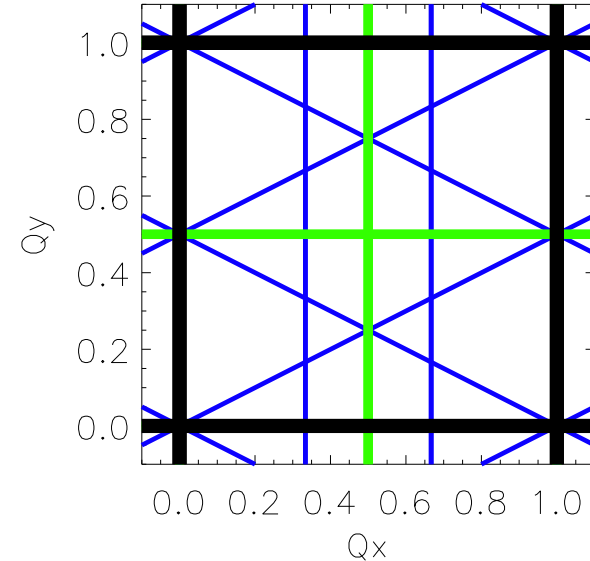
$$h_{10110} = h_{01110}^* \rightarrow Q_x$$

$$h_{10200} = h_{01020}^* \rightarrow Q_x + 2Q_y$$

$$h_{10020} = h_{01200}^* \rightarrow Q_x - 2Q_y$$

$$h_{20001} = h_{02001}^* \rightarrow 2Q_x$$

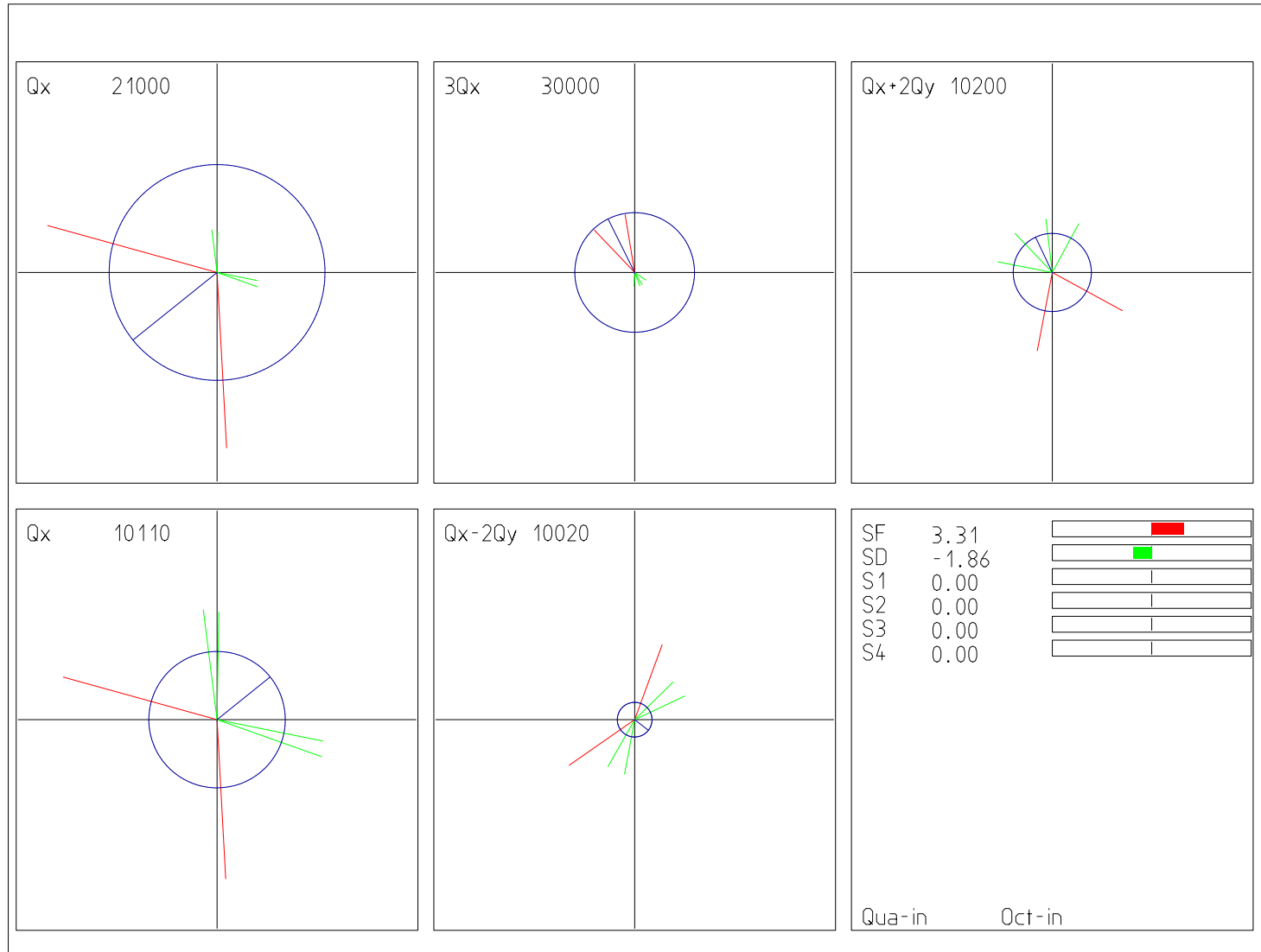
$$h_{00201} = h_{00021}^* \rightarrow 2Q_y$$





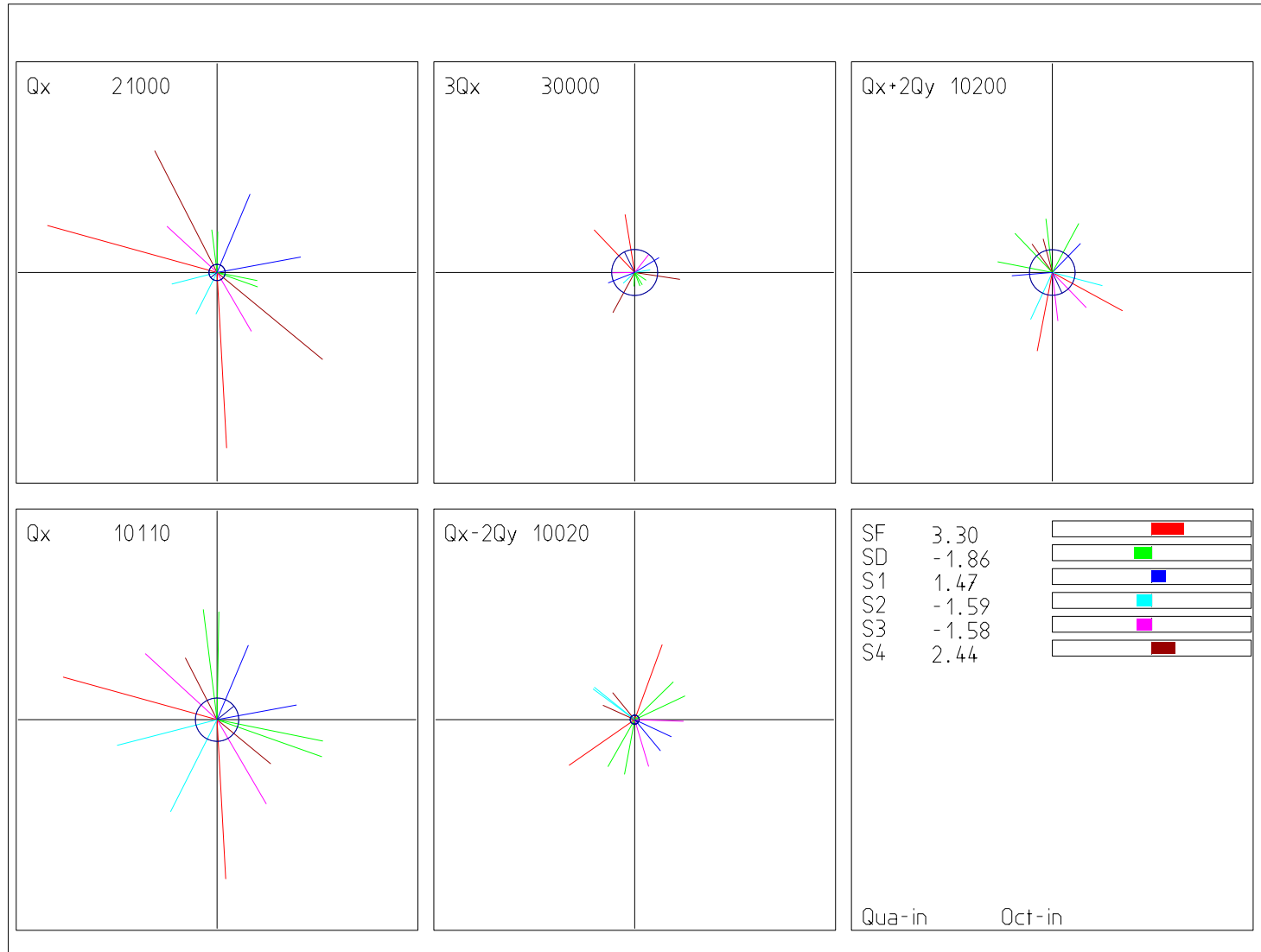
# First order sextupole optimization

## 2 Sextupole families (ESRF standard cell)



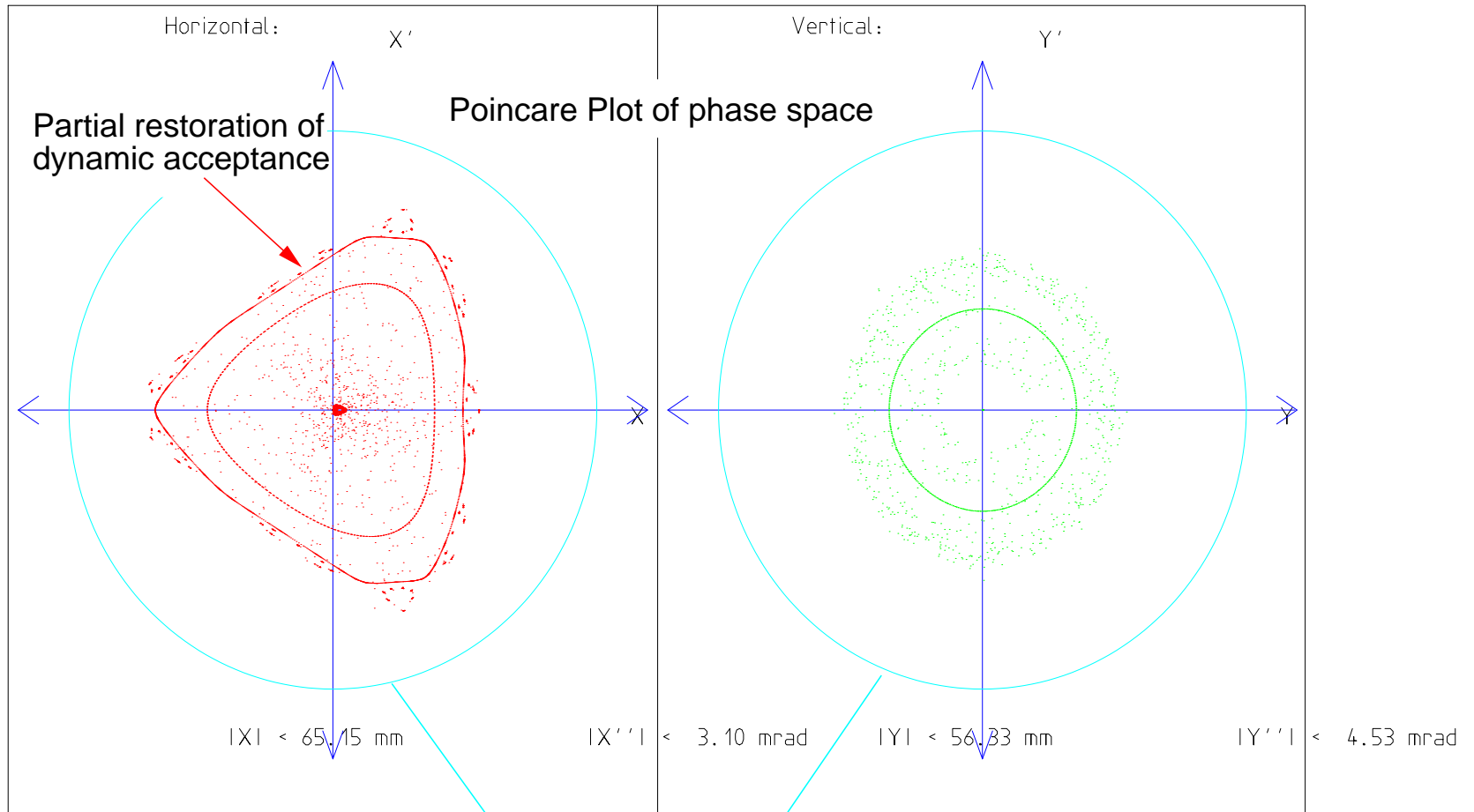
# First order sextupole optimization

## 6 Sextupole families (4 *harmonic* families in straight sections)



# First order sextupole optimization

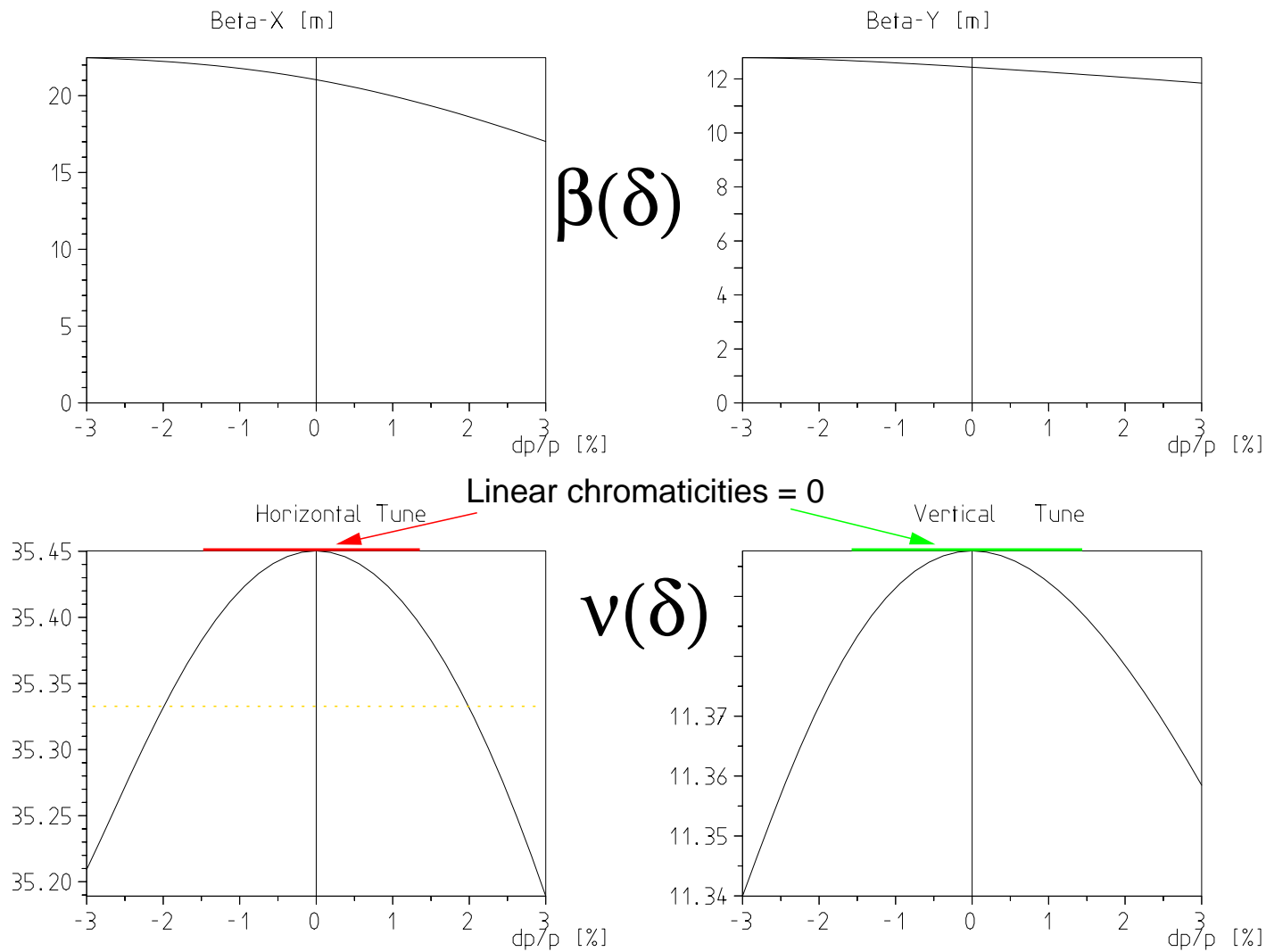
## Step 3: First order sextupole optimization with 6 families



Limitation from 160 mm x 160 mm vacuum chamber

# First order sextupole optimization

## Step 3: First order optimization with 6 sextupole families



# First order sextupole optimization

## Fourier spectrum

### Suppression of sextupolar resonances

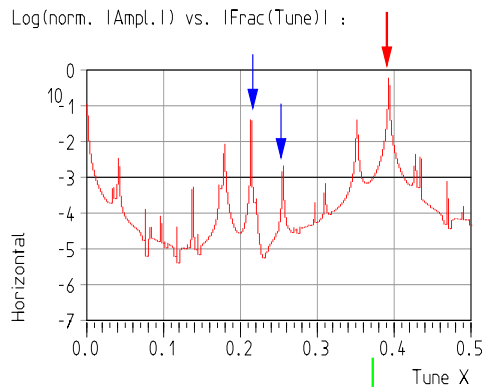
2 sextupole families

6 sextupole families

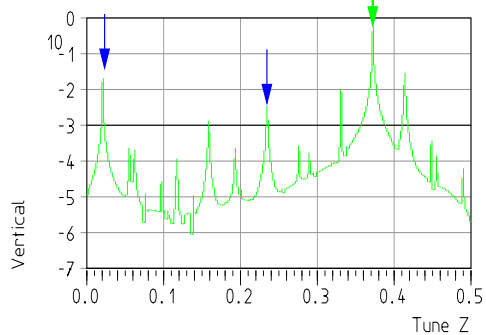
Tracking of test particle ( $x_o = y_o = 4$  mm,  $x'_o = y'_o = 0$ ,  $\delta = 0$ ) and FFT:

Fourier spectrum of Test-particle

Log(norm. |Ampl.|) vs. |Frac(Tune)| :



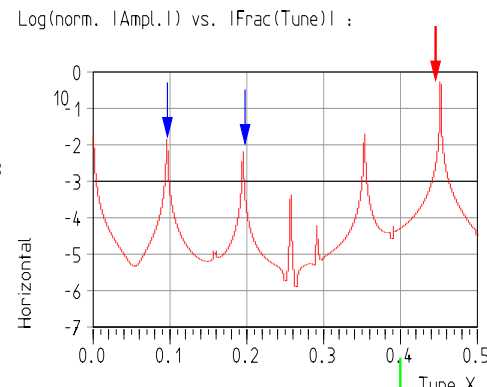
Tune	Amp. [mm]	Guess
35.4503:		
0.3931: → Qx	3.5591	-
-0.0003: -	0.6327	-
0.2138: 2Qx	0.2809	→ 3Qx=112
0.3517: Qx-2Qz	0.2256	2Qx-2Qz=48
0.1793: 3Qx	0.0491	4Qx=144
0.4275: ?	0.0251	-
0.4345: ?	0.0248	-
0.0415: ?	0.0201	-
0.2552: 2Qz	0.0128	→ Qx-2Qz=16
0.4687: ?	0.0043	-



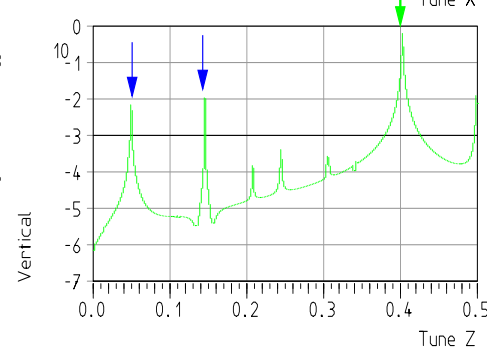
11.3976:		
0.3724: → Qz	3.9946	-
0.4138: 2Qx-Qz	0.1829	2Qx-2Qz=48
0.0207: Qx-Qz	0.1454	→ Qx-2Qz=16
0.3309: ?	0.0719	-
0.2345: Qx+Qz	0.0217	→ Qx+2Qz=64
0.1586: 2Qx+Qz	0.0082	2Qx+2Qz=96

Fourier spectrum of Test-particle

Log(norm. |Ampl.|) vs. |Frac(Tune)| :



Tune	Amp. [mm]	Guess
35.4503:		
0.4520: → Qx	4.1865	-
0.3531: Qx-2Qz	0.1378	2Qx-2Qz=48
-0.0003: -	0.1166	-
0.0960: 2Qx	0.0955	→ 3Qx=112
0.1949: 2Qz	0.0445	→ Qx-2Qz=16



11.3976:		
0.4025: → Qz	3.9480	-
0.4986: 2Qx-Qz	0.0818	2Qx-2Qz=48
0.1455: Qx+Qz	0.0830	→ Qx+2Qz=64
0.0495: Qx-Qz	0.0472	→ Qx-2Qz=16

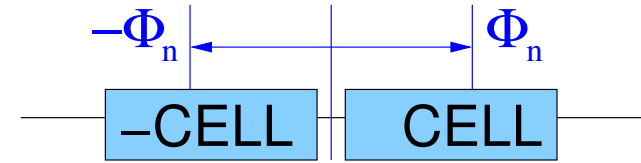
## First order sextupole optimization

Systematic first order optimization:<sup>†</sup>

9 terms  $h_{jklmp}$  (7 complex, 2 real)

→ **16** sextupole families

⇒ Symmetry:



$Im(h_{jklmp}) = 0 \rightarrow$  **9** sextupole families.

Linear system for  $M$  families of sextupoles:

$$\left\{ \sum_{n \in \{Sm\}} \beta_n^{(\dots)} D_n^{(\dots)} e^{i\{(\dots)\phi_n\}} \dots \right\}_{9 \times M} \times \left\{ (b_3 L)_m \right\}_{M \times 1} = \left\{ \sum_{\text{Quad}} (b_2 L) \dots \right\}_{1 \times 9}$$

**Light source problem:**  $\epsilon \downarrow \Rightarrow \Delta\phi_x^{\text{cell}} \rightarrow 180^\circ \Rightarrow e^{i2\phi_x} \approx 1$

$2Q_x$  resonance driving term  $h_{20001}$  proportional to chromaticity  $\xi_x \propto h_{11001}$

$$2Q_x \rightarrow \frac{\partial \beta_x}{\partial \delta} \rightarrow \xi_x^{(2)} = \frac{\partial^2 Q_x}{\partial \delta^2} \rightarrow \text{energy acceptance } \downarrow$$

No solution for  $\{(b_3 L)_m\} \Rightarrow$  suppression by  $\Delta\phi_x^{\text{straight}}$

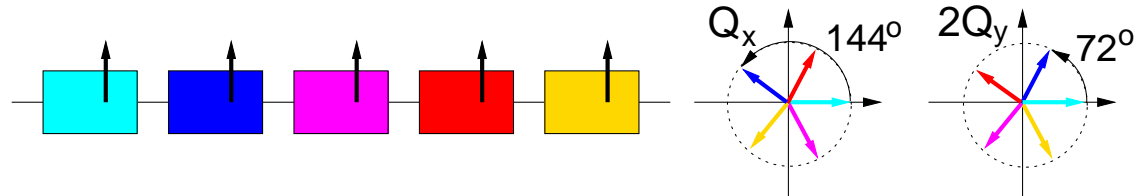
<sup>†</sup> J.Bengtsson et al., Increasing the energy acceptance of synchrotron light storage rings, NIM A 404 (1998) 237

## Phase cancellation schemes

**Periodicity:**  $N$  cells

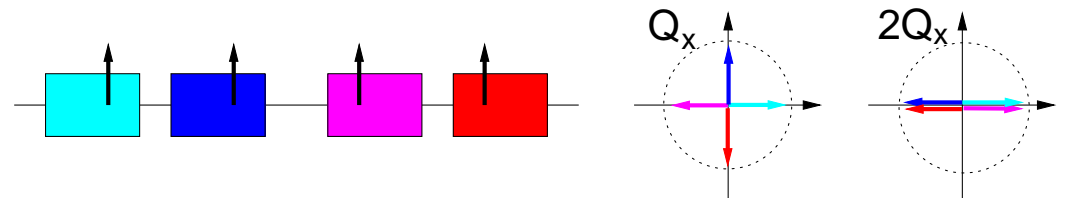
$$N\Delta Q_x^{\text{cell}}, [3N\Delta Q_x^{\text{cell}}, 2N\Delta Q_x^{\text{cell}}], 2N\Delta Q_y^{\text{cell}} \longrightarrow \text{integer!}$$

e.g.  $N = 5$ ,  $\Delta Q_x^{\text{cell}} = 0.4 (= 144^\circ)$ ,  $\Delta Q_y^{\text{cell}} = 0.1 (= 36^\circ)$



**Symmetry:** Lattice section vs. mirror image

$$\Delta Q_x^{\text{cell}} = \frac{2n_x+1}{4}, \Delta Q_y^{\text{cell}} = \frac{2n_y+1}{4} \quad (n_x, n_y \text{ integers}):$$



→ phase advances over straight sections →  $\beta_x, \beta_y$  in straights.

⇒ Iterate: **linear**  $\iff$  **nonlinear** lattice design

### Second order sextupole [+first order octupole] Hamiltonian

$$\sum_n \sum_m (b_3 L)_n (b_3 L)_m \times (\beta_n, \phi_n \beta_m, \phi_m \dots) + \left[ \sum_q (b_4 L)_q \times (\beta_q, \phi_q \dots) \right]$$

- **3 phase independent terms** → **amplitude dependant tune shifts:**

$$\frac{\partial Q_x}{\partial J_x} \quad \frac{\partial Q_x}{\partial J_y} = \frac{\partial Q_y}{\partial J_x} \quad \frac{\partial Q_y}{\partial J_y}$$

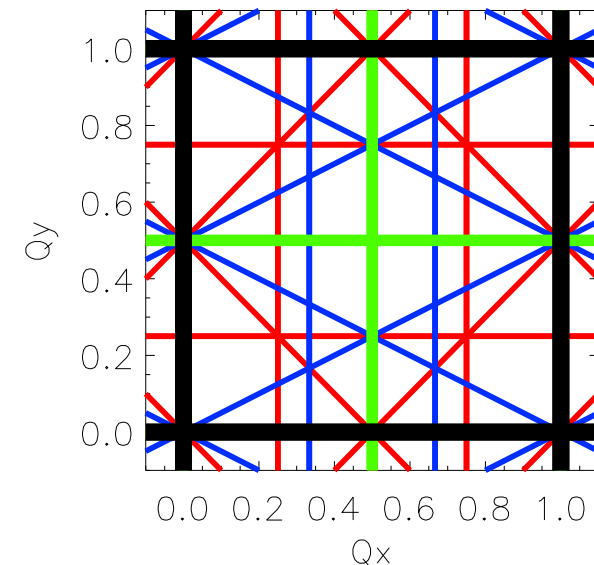
- **2 phase independent off-momentum terms** → **second order chromaticities:**

$$\xi_{x/y}^{(2)} = \frac{\partial^2 Q_{x/y}}{\partial \delta^2}$$

- **8 phase dependant terms**

→ **octupolar resonances:**

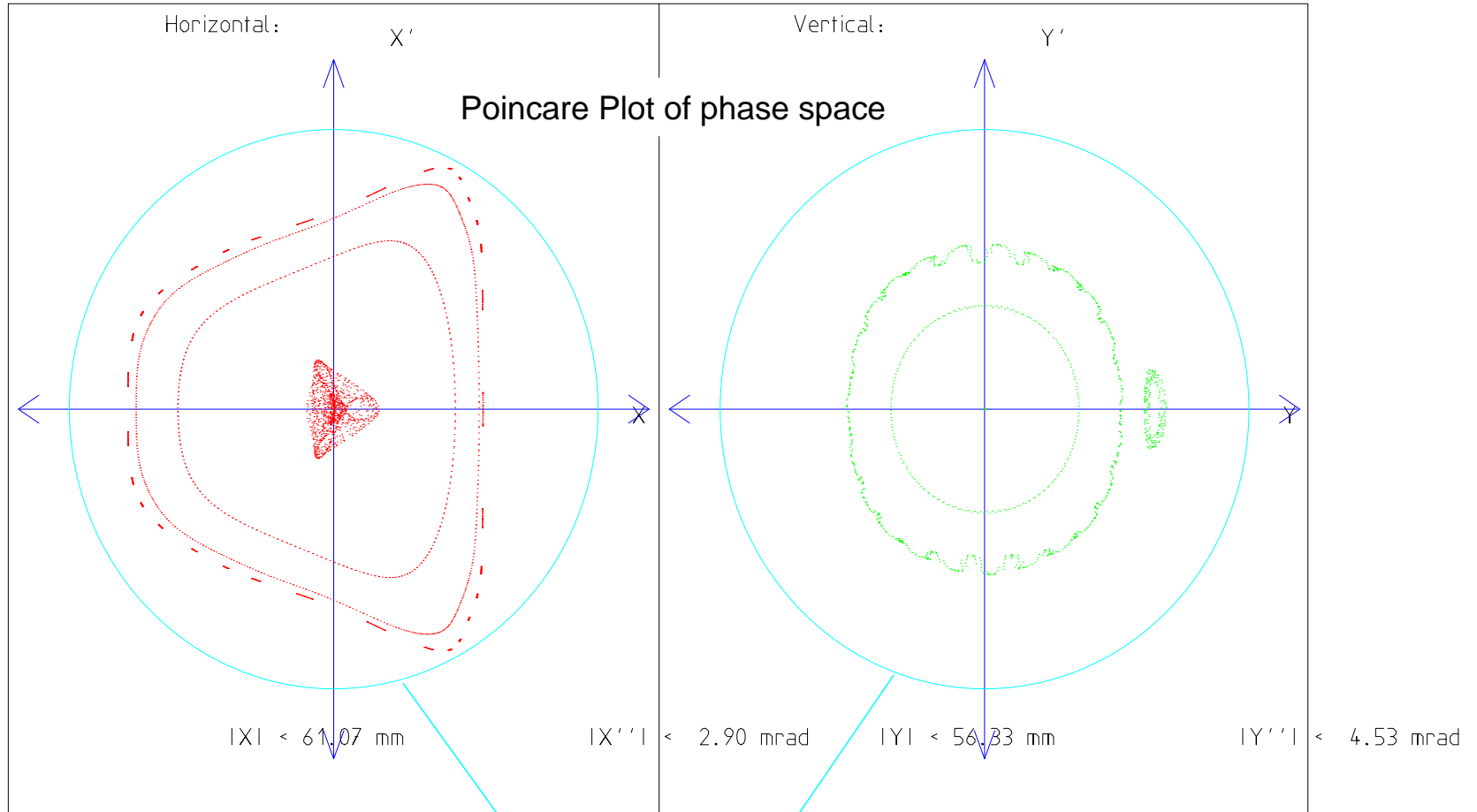
$h_{40000} \rightarrow 4Q_x$	$h_{31000} \rightarrow 2Q_x$
$h_{00400} \rightarrow 4Q_y$	$h_{20110} \rightarrow 2Q_x$
$h_{20200} \rightarrow 2Q_x + 2Q_y$	$h_{00310} \rightarrow 2Q_y$
$h_{20020} \rightarrow 2Q_x - 2Q_y$	$h_{01110} \rightarrow 2Q_y$





# Second order sextupole optimization

## Step 4: Second order sextupole optimization with 6 families

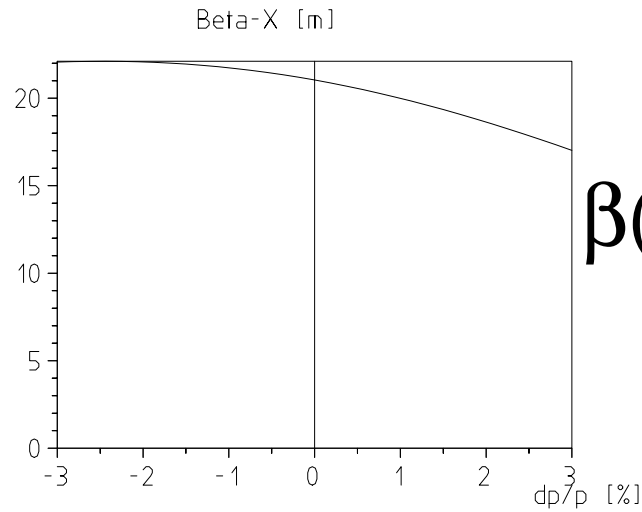


Limitation from 160 mm x 160 mm vacuum chamber

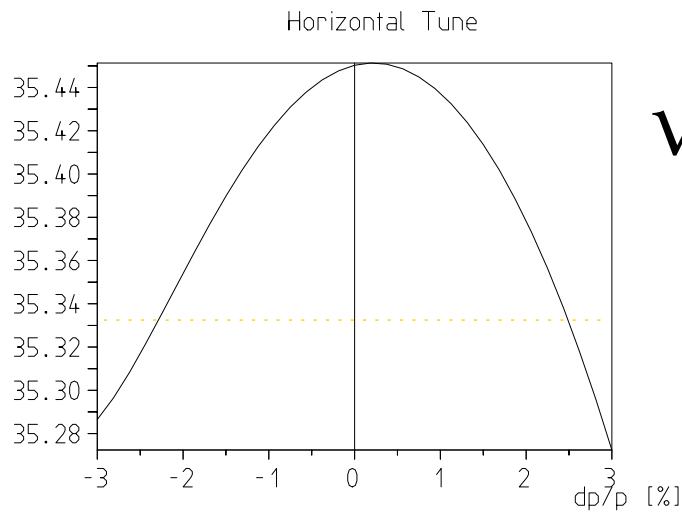
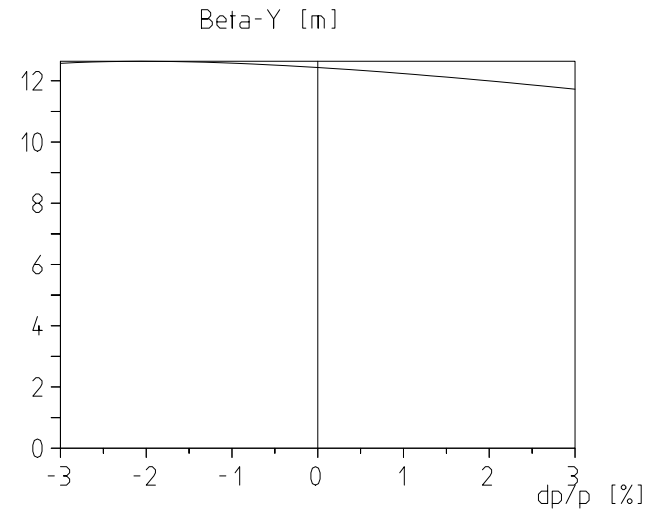
Nonlinearities in Light Sources

# Second order sextupole optimization

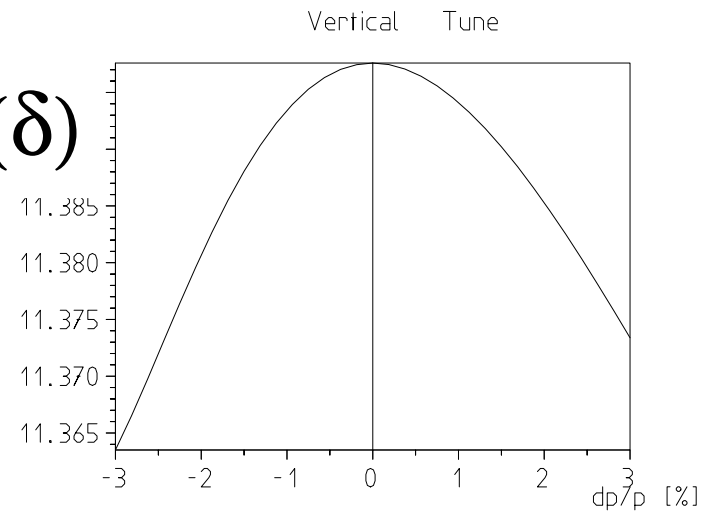
## Step 4: Second order optimization with 6 sextupole families



$\beta(\delta)$

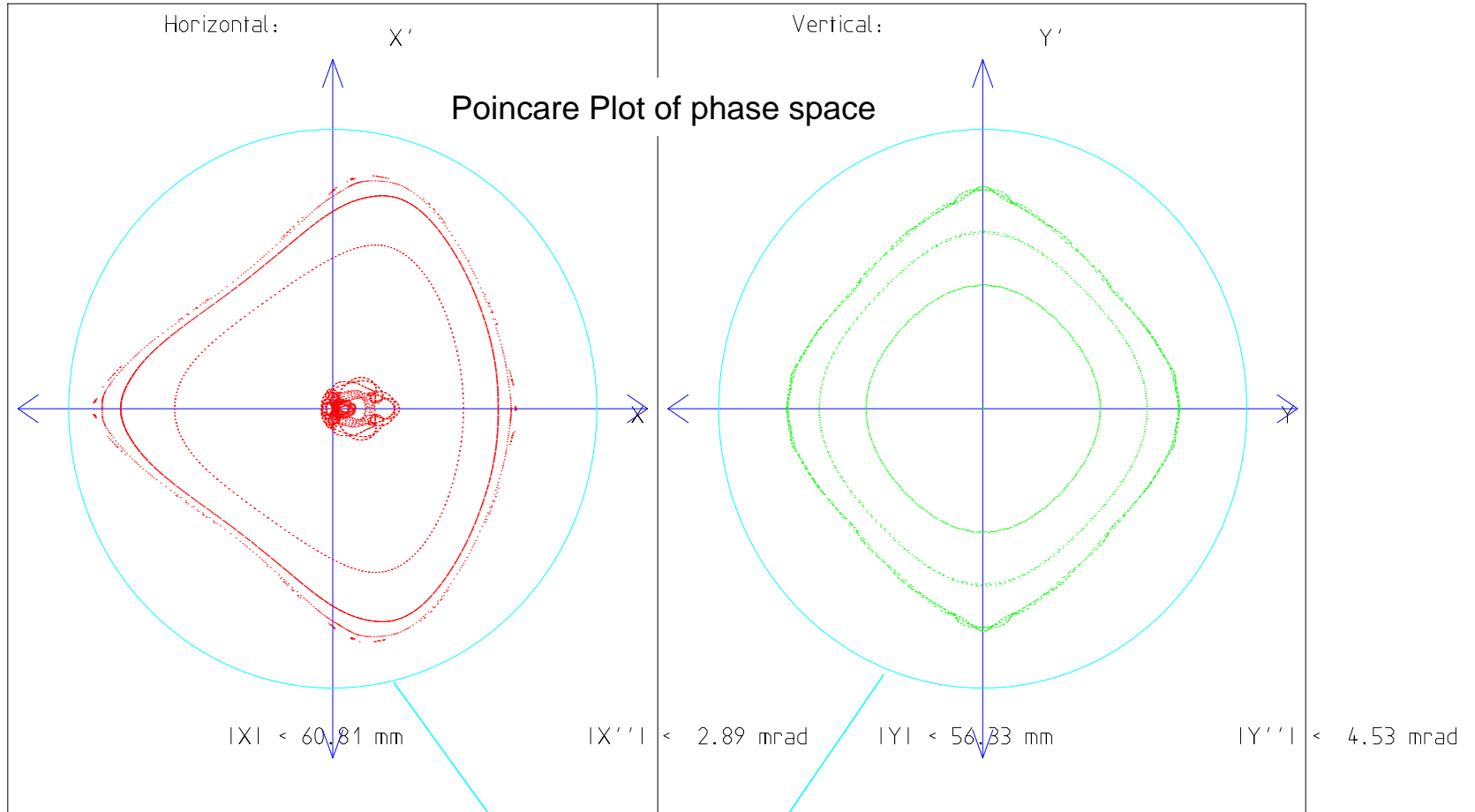


$\nu(\delta)$



# Second order sextupole optimization

[ Step 5: Further optimization by adding 8 octupole families ]

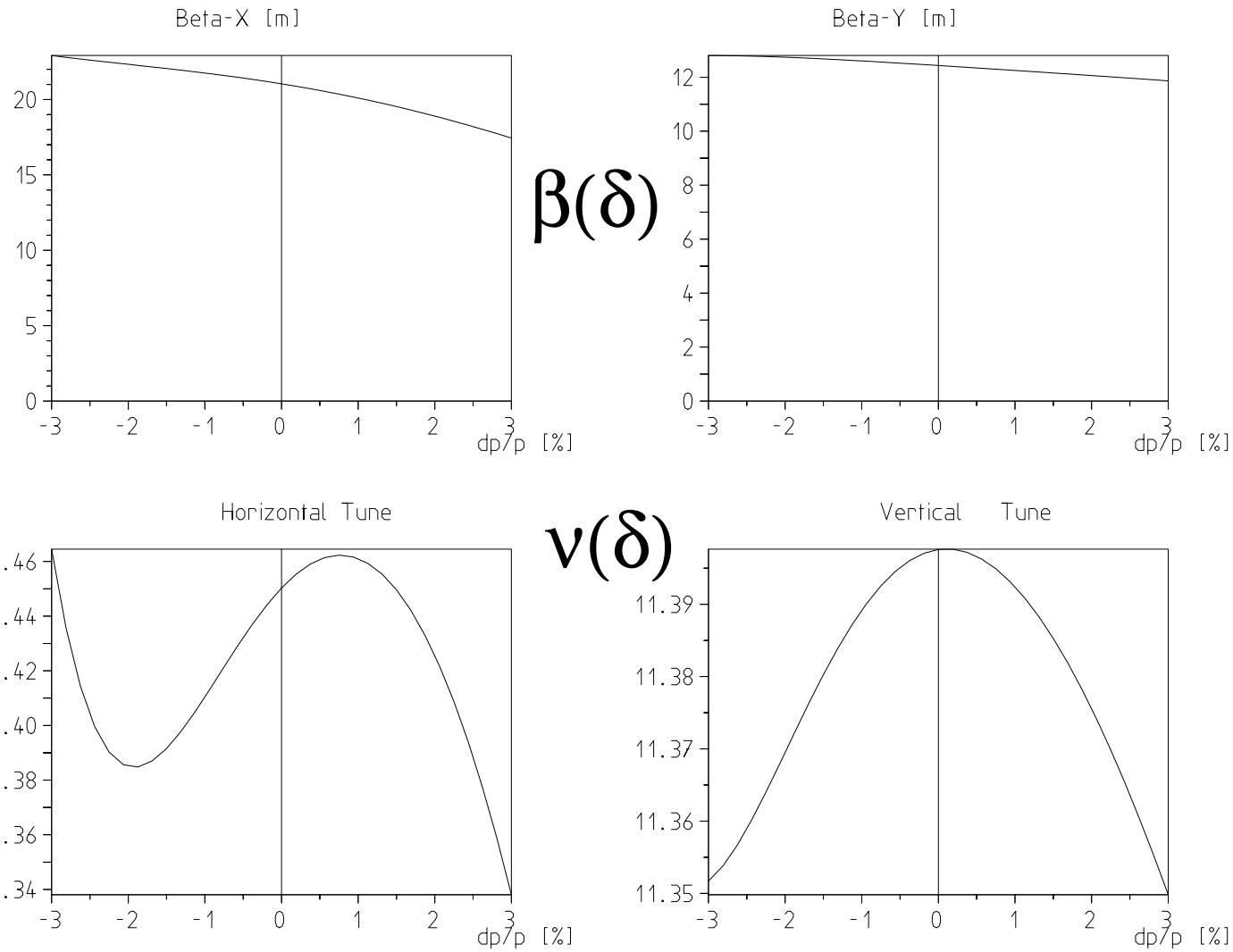


Limitation from 160 mm x 160 mm vacuum chamber

Nonlinearities in Light Sources

# Second order sextupole optimization

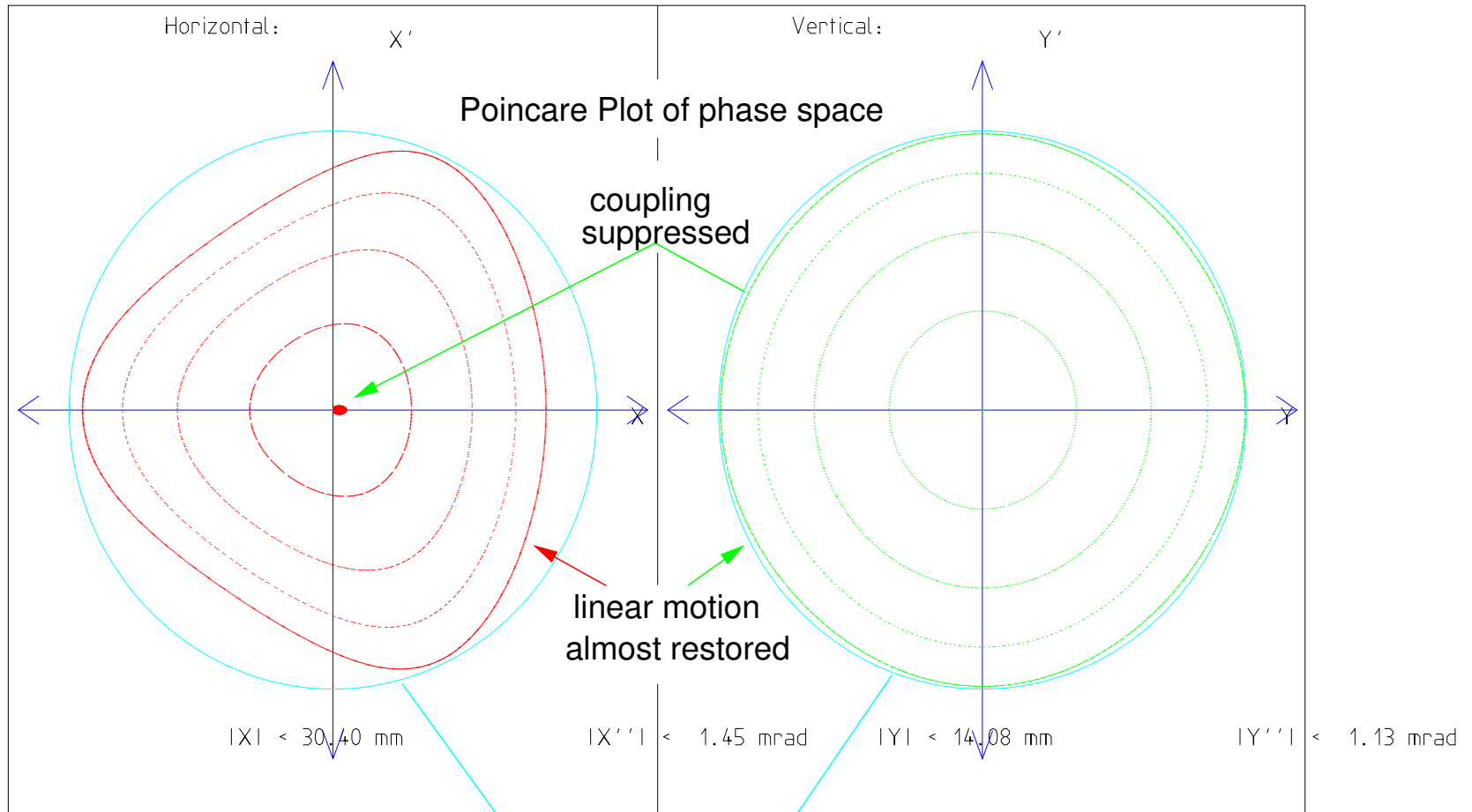
[ Step 5: Octupoles for suppression of 2nd order chromaticities ]



Nonlinearities in Light Sources

# Second order sextupole optimization

Including real beam pipe limitations ( $\pm 40$  mm wide,  $\pm 20$  mm high)



Limitation from 80 mm x 40 mm vacuum chamber (linear)

### Sextupole pattern optimization strategy

- **decouple chromatic sextupoles**  
→ careful placement of SD, SF families
- **exploit symmetry and periodicity**  
→ back to *linear* lattice design and machine layout
- **install “harmonic” sextupoles**  
→  $\approx 4 \dots 8$  additional families [in straight sections]
- **minimize all 1<sup>st</sup> order and some 2<sup>nd</sup> order terms**  
→ sextupolar resonances, tune shifts with amplitude and 2<sup>nd</sup> order chromaticities
- **check by tracking**  
→ dynamic aperture scans – also include misalignments and multipolar errors
- **iterate...**