



# Accelerators for Newcomers

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D. Brandt, CERN

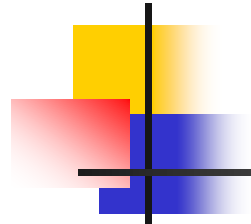


# Why this Introduction?

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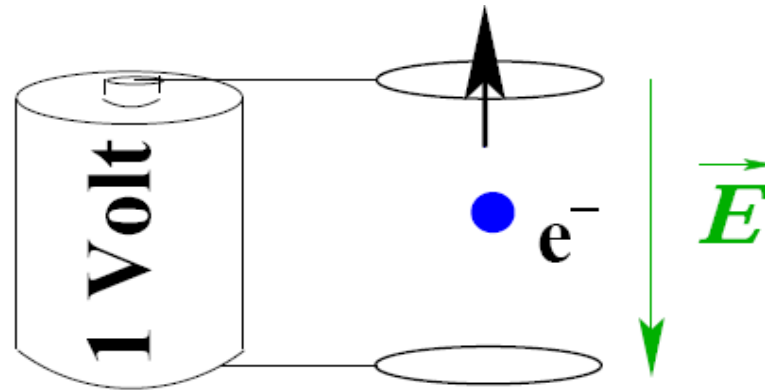
- During this school, you will learn about **beam dynamics** in a rigorous way...
- but some of you are completely new to the field of accelerator physics.
- It seemed therefore justified to start with the introduction of a few very **basic concepts**, which will be used throughout the course.

This is a completely **intuitive approach** (no mathematics) aimed at highlighting the **physical concepts**, without any attempt to achieve any scientific derivation.



## Some generalities ...

# Units: the electronvolt (eV)



The **electronvolt (eV)** is the energy gained by an electron travelling, in vacuum, between two points with a voltage difference of 1 Volt.

$$1 \text{ eV} = 1.602 \cdot 10^{-19} \text{ Joule}$$

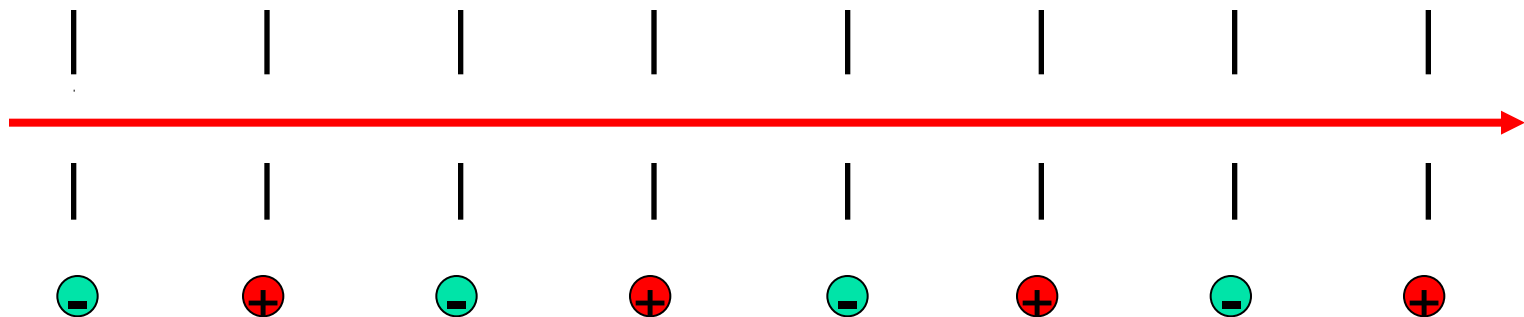
We also frequently use the electronvolt to express masses from  $E=mc^2$ :

$$1 \text{ eV}/c^2 = 1.783 \cdot 10^{-36} \text{ kg}$$

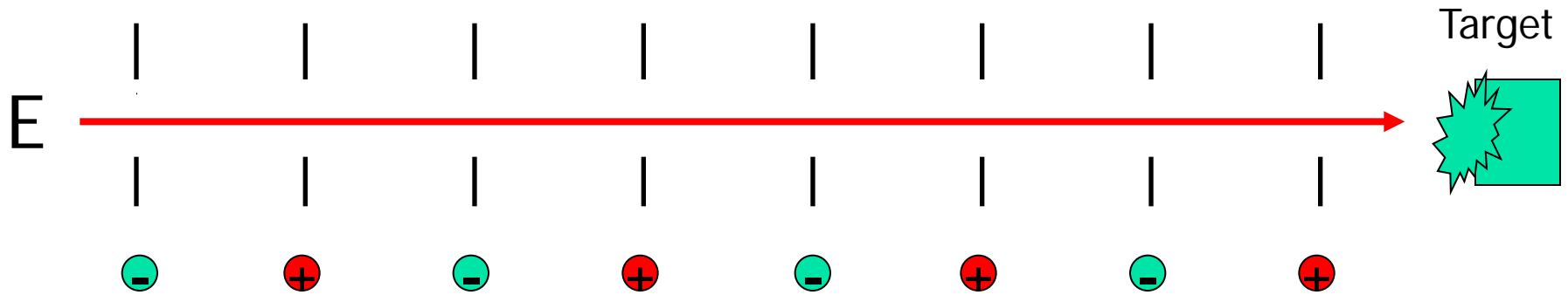
# What is a Particle Accelerator?

➤ a machine to accelerate some particles ! **How is it done ?**

➤ Many different possibilities, but rather easy from the general principle:



# Ideal linear machines (linacs)



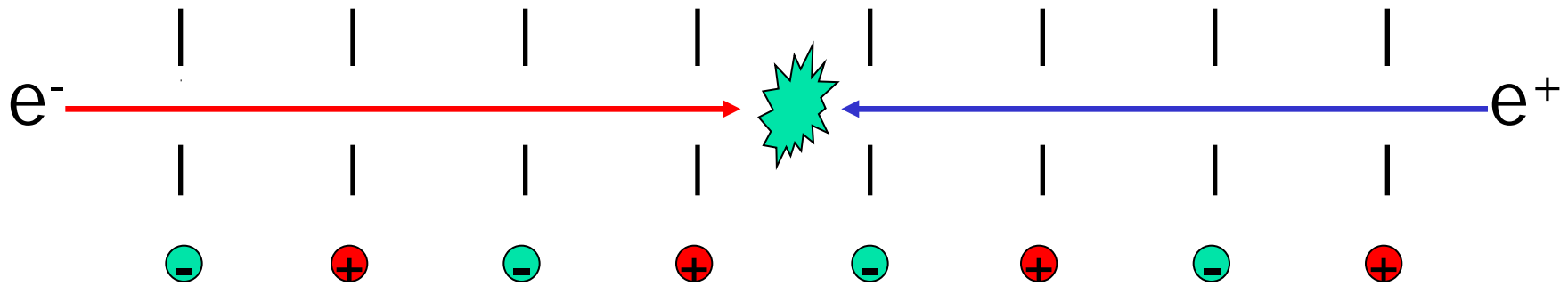
$$\text{Available Energy : } E_{\text{c.m.}} = m \cdot (2+2\gamma)^{1/2} = (2m \cdot (m+E))^{1/2}$$

with  $\gamma = E/E_0$

Advantages: Single pass  
High intensity

Drawbacks: Single pass  
Available Energy

# Improved solution for $E_{c.m.}$



Available Energy :  $E_{c.m.} = 2m\gamma = 2E$

with  $\gamma = E/E_0$

Advantages: High intensity

Drawbacks: Single pass  
Space required

# Watch out !

The difference between fixed target and colliding mode deserves to be considered in some detail:

Fixed target mode:

$$E_{\text{c.m.}} \propto (2mE)^{1/2}$$



Colliding mode:

$$E_{\text{c.m.}} \propto 2E$$

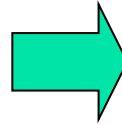
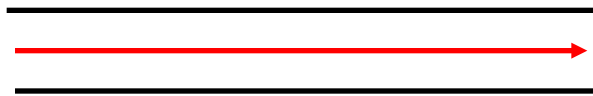
What would be the required beam energy to achieve  $E_{\text{c.m.}} = 14 \text{ TeV}$  in fixed target mode ?



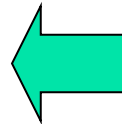
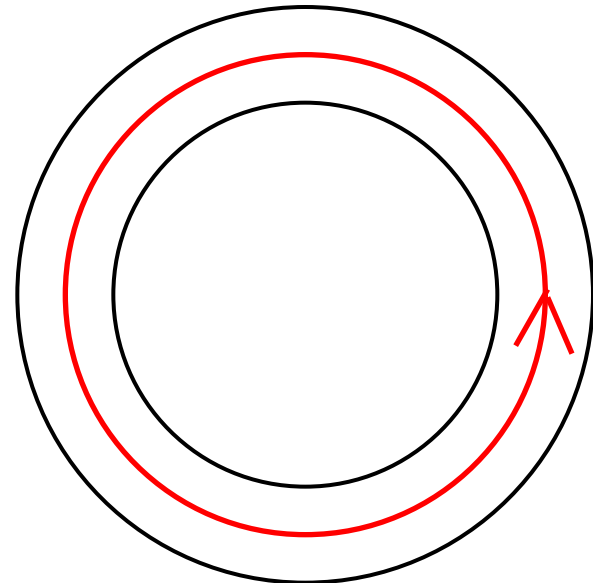
# Keep particles: circular machines

Basic idea is to keep the particles in the machine for many turns.

Move from the linear design



To a circular one:



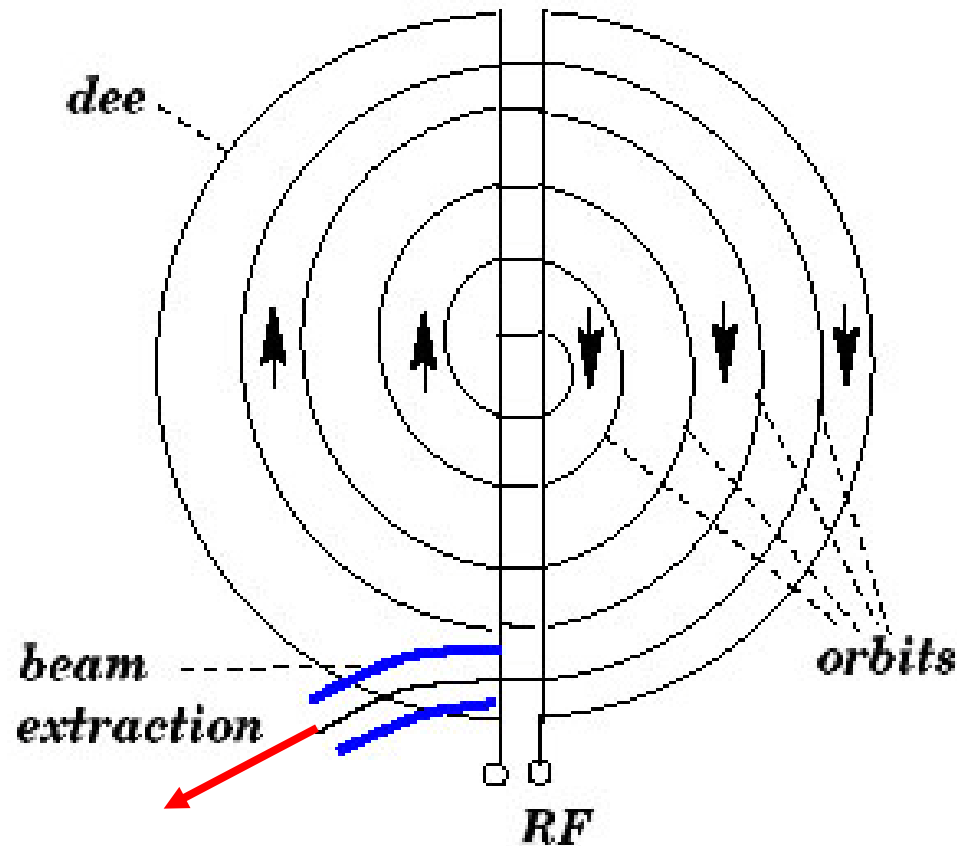
➤ Need Bending

➤ Need **Dipoles!**

# Circular machines 1 ( $E_{c.m.} \sim (mE)^{1/2}$ )

fixed target:

cyclotron



huge dipole, compact design,  $B = \text{constant}$ , low energy, single pass.

# Circular machines 2 ( $E_{\text{c.m.}} \sim (mE)^{1/2}$ )

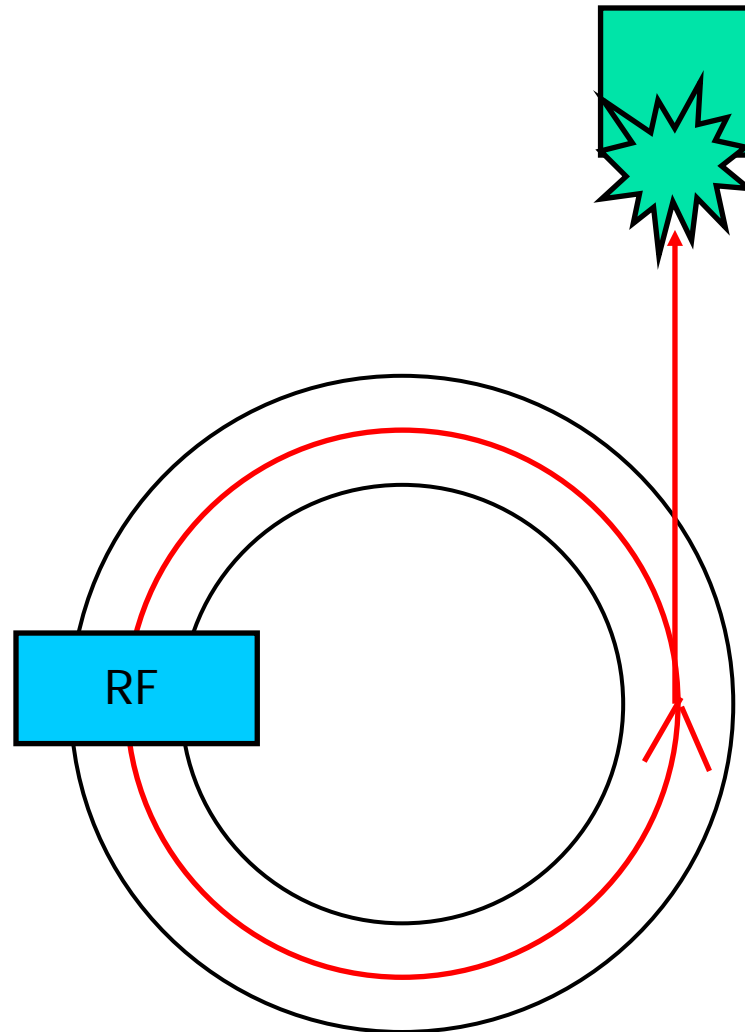
fixed target:

synchrotron

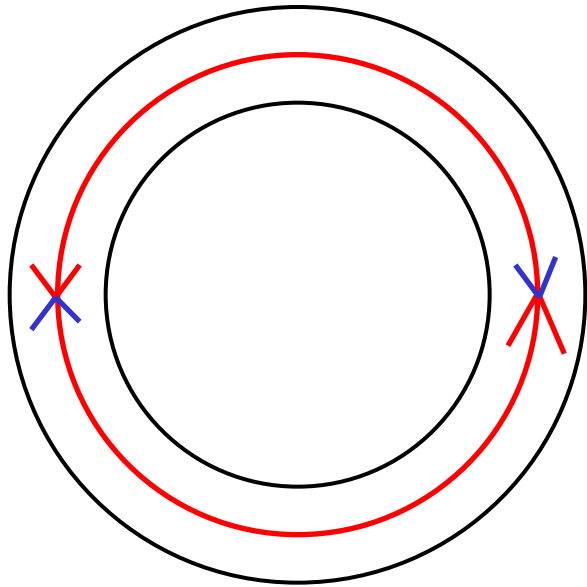
varying B

small magnets,

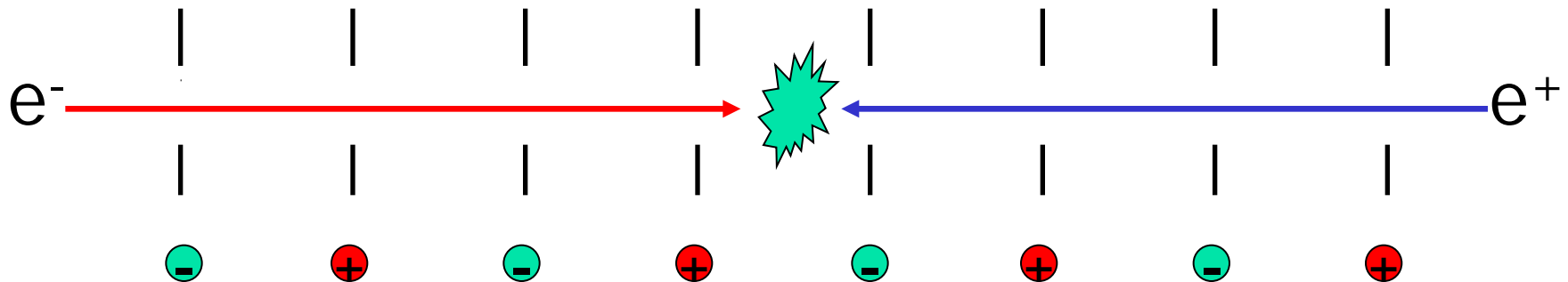
high energy



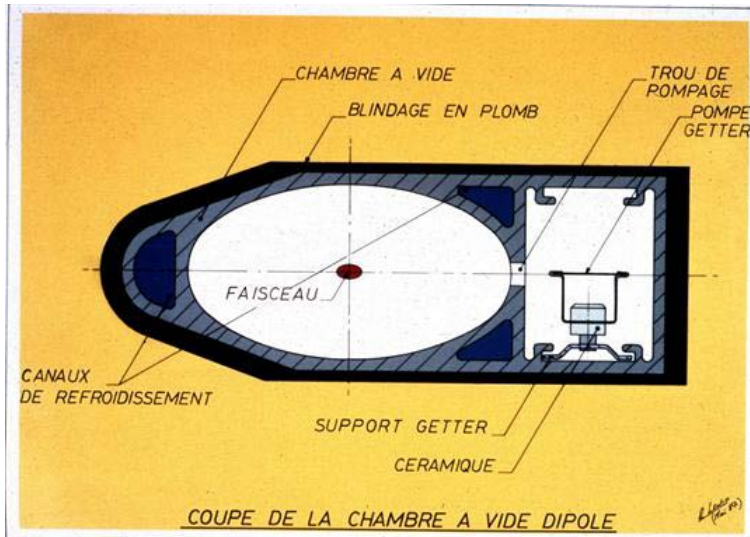
# Colliders ( $E_{c.m.}=2E$ )



Colliders with the same type of particles (e.g. p-p) require two separate chambers. The beam are brought into a common chamber around the interaction regions



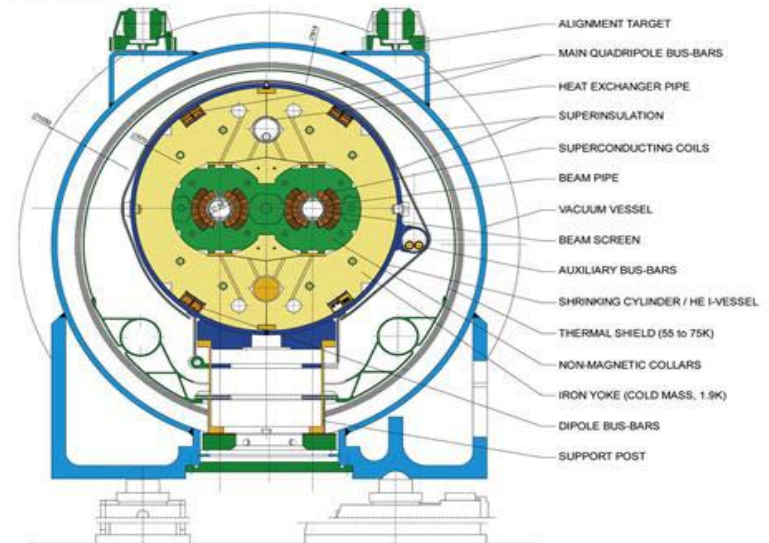
# Colliders ( $e^+ - e^-$ ) et ( $p - p$ )



LEP

LHC

LHC DIPOLE : STANDARD CROSS-SECTION





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## Transverse Dynamics

$$F = e (E + v \times B)$$




# Beam Dynamics (1)

In order to describe the motion of the particles, each particle is characterised by:

- Its azimuthal position along the machine:  $s$
- Its momentum:  $p$  (or Energy  $E$ )
- Its horizontal position:  $x$
- Its horizontal slope:  $x'$
- Its vertical position:  $y$
- Its vertical slope:  $y'$

i.e. a sixth dimensional vector

$(s, p, x, x', y, y')$



# Beam Dynamics (2)

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- In an accelerator designed to operate at the energy  $E_{nom}$ , all particles having  $(s, E_{nom}, 0, 0, 0, 0)$  will happily fly through the center of the vacuum chamber without any problem. These are “ideal particles”.
- The difficulties start when:
  - one introduces **dipole magnets**
  - the energy  $E \neq E_{nom}$  or  $(p - p_{nom}/p_{nom}) = \Delta p/p_{nom} \neq 0$
  - either of  $x, x', y, y' \neq 0$





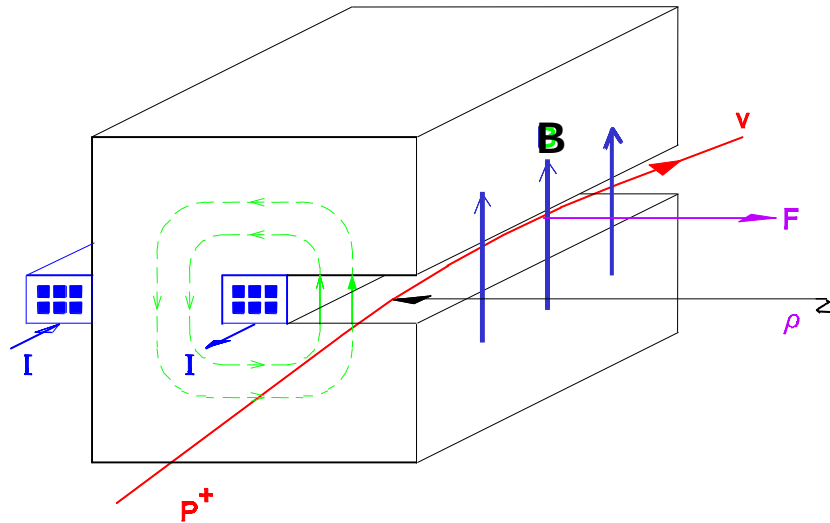
# Basic problem:

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With more than  $10^{10}$  particles per bunch, most of them will **not** be **ideal particles**, i.e. they are going to be lost !

Purpose of this lecture: how can we keep the particles in the machine ?

# Circular machines: Dipoles



$$p = m_0 \cdot c \cdot (\beta\gamma)$$



Magnetic rigidity:  
 $B\rho = mv/e = p/e$

Classical mechanics:

Equilibrium between two forces

Lorentz force

Centrifugal force

$$F = e \cdot (\underline{v} \times \underline{B})$$

$$F = mv^2/\rho$$

$$evB = mv^2/\rho$$

Relation also holds for relativistic case provided the classical momentum  $mv$  is replaced by the relativistic momentum  $p$



# Why fundamental ?

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Constraints:

**E** and  $\rho$  given  $\Rightarrow$  Magnets defined (**B**)

Constraints:

**E** and **B** given  $\Rightarrow$  Size of the machine ( $\rho$ )

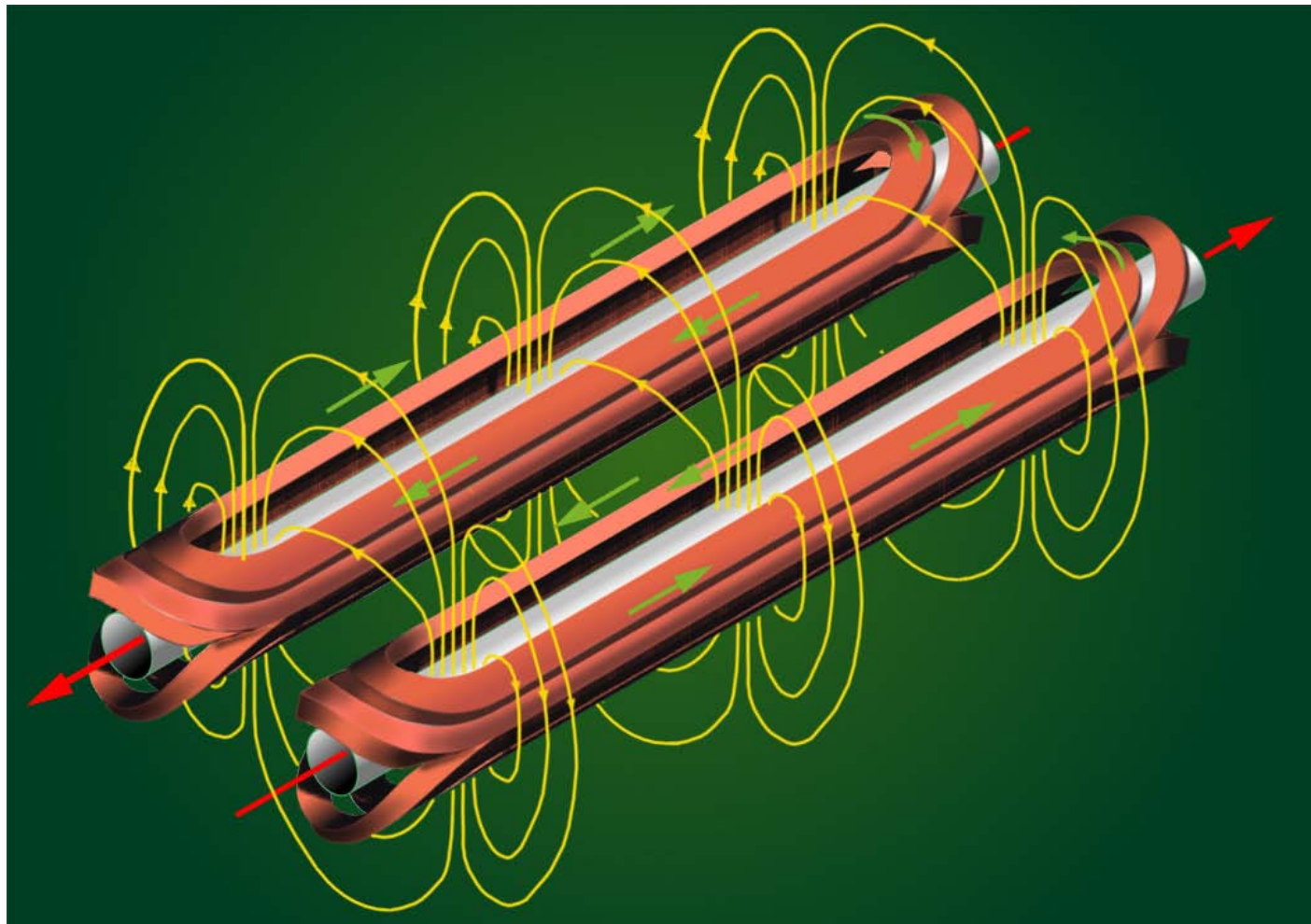
Constraints:

**B** and  $\rho$  given  $\Rightarrow$  Energy defined (**E**)

# Dipoles (1):



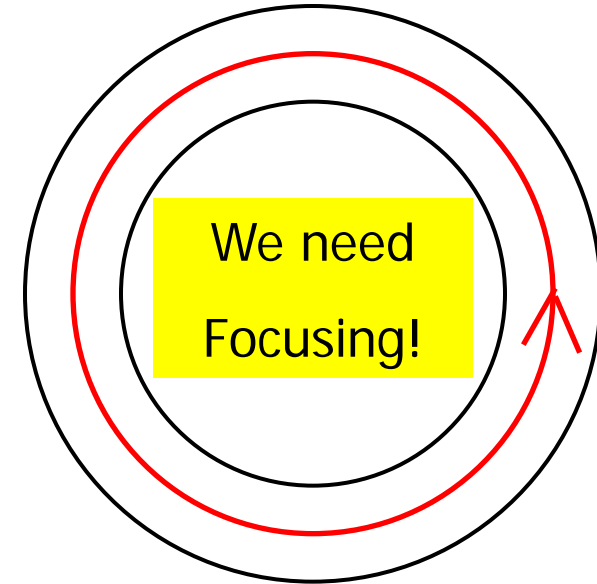
# Dipoles (2):



# Ideal circular machine:

- Neglecting radiation losses in the dipoles
- Neglecting gravitation

ideal particle would happily circulate on axis in the machine for ever!



**Unfortunately: real life is different!**

Gravitation:  $\Delta y = 20$  mm in 64 msec!

Alignment of the machine

Limited physical aperture

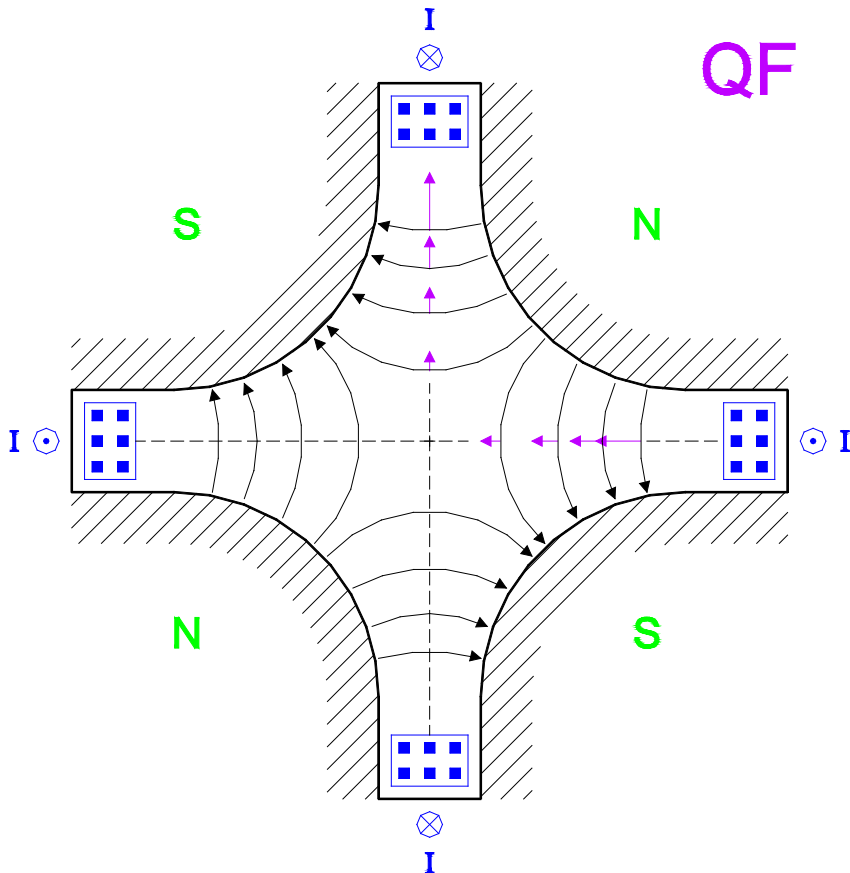
Ground motion

Field imperfections

Energy error of particles and/or  $(x, x')_{inj} \neq (x, x')_{nominal}$

Error in magnet strength (power supplies and calibration)

# Focusing with quadrupoles



$$F_x = -g.x$$

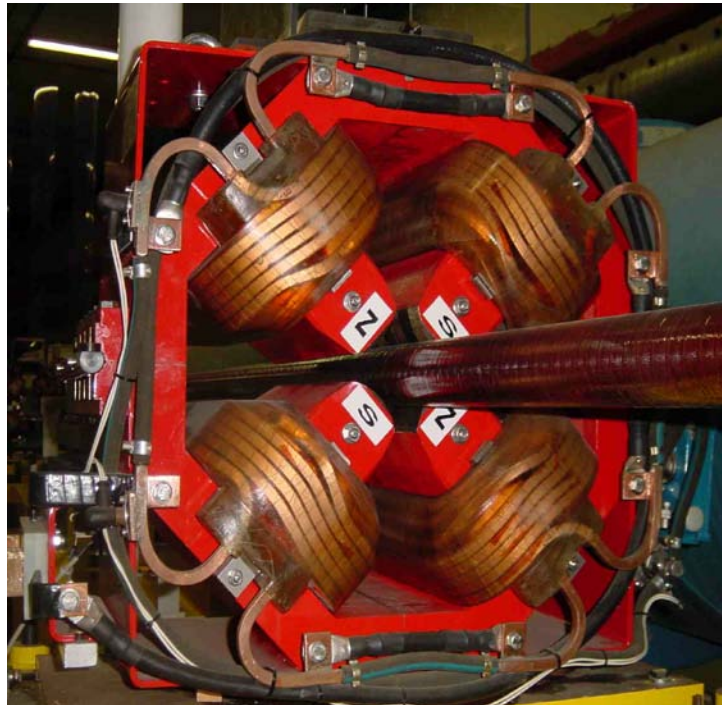
$$F_y = g.y$$

Force increases **linearly** with displacement.

Unfortunately, effect is **opposite** in the two planes (H and V).

Remember: **this** quadrupole is **focusing** in the **horizontal** plane but **defocusing** in the **vertical** plane!

# Quadrupoles:







# Focusing properties ...

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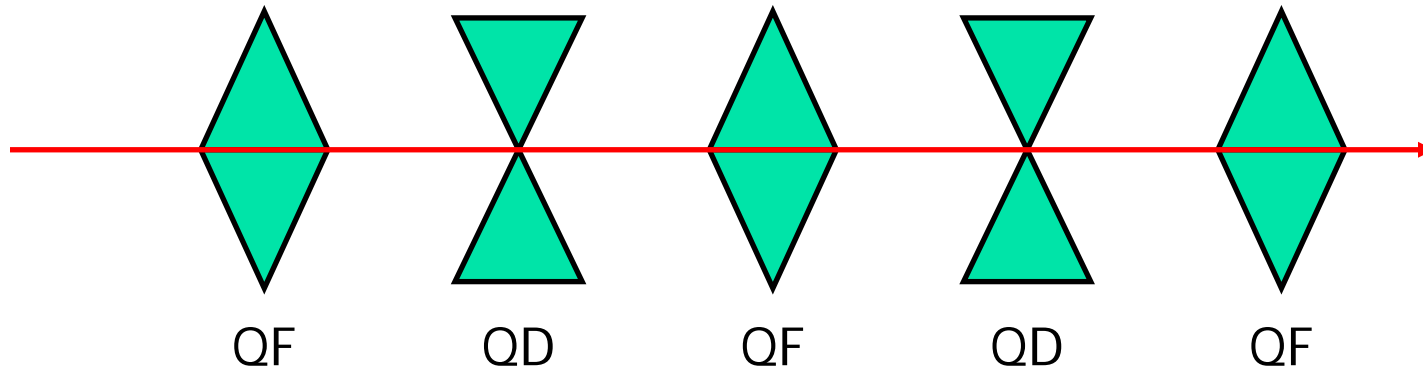
A quadrupole provides the required effect in one plane...

but the opposite effect in the other plane!

Is it really interesting ?

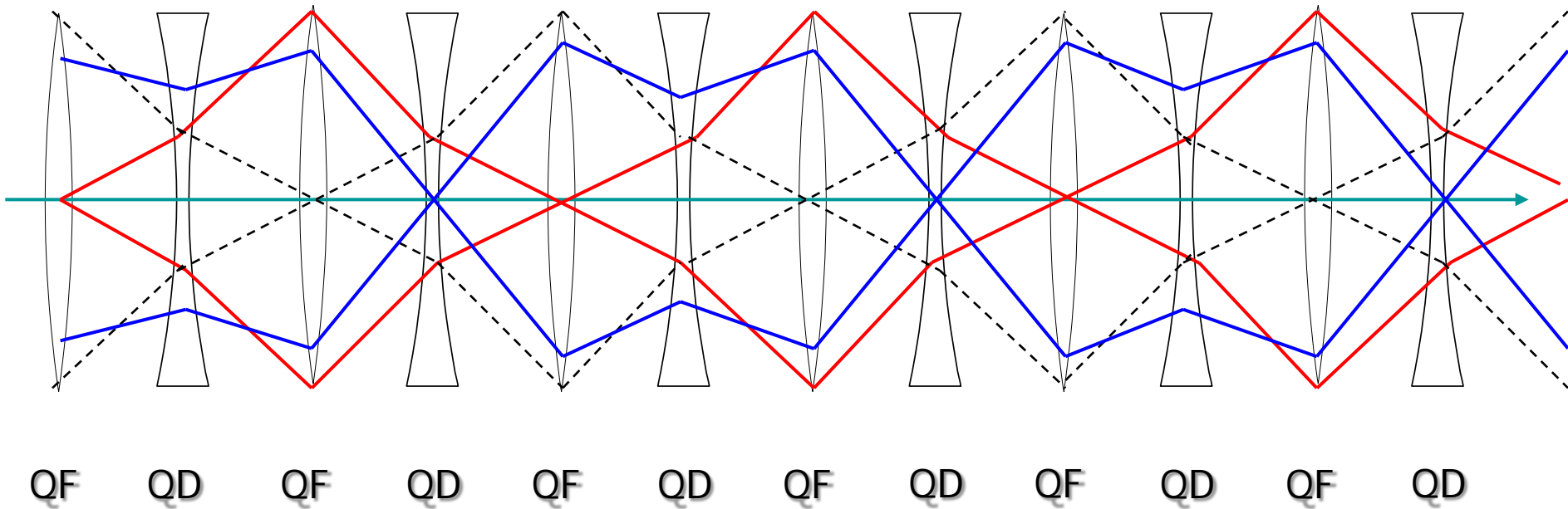
# Alternating gradient focusing

Basic new idea:  
Alternate QF and QD



valid for one plane only (H or V) !

# Alternating gradient focusing





# Alternating gradient focusing:

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Particles for which  $x, x', y, y' \neq 0$  thus oscillate around the ideal particle ...

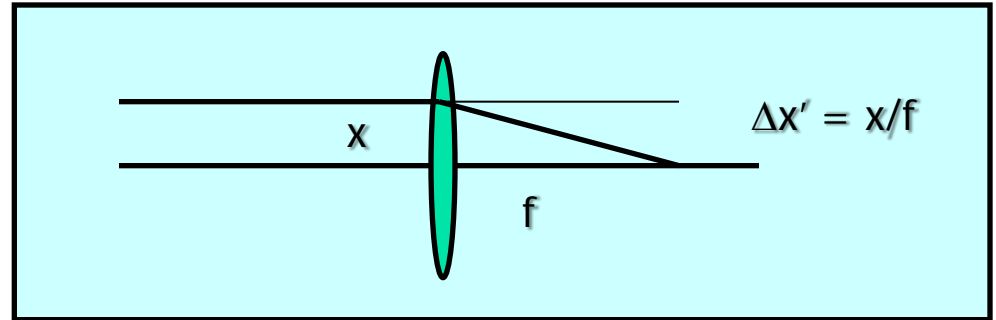
but the trajectories remain inside the vacuum chamber !

# Thin lens analogy of AG focusing

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{\text{out}} = \begin{pmatrix} 1 & 0 \\ -1/f & 1 \end{pmatrix} \begin{pmatrix} x \\ x' \end{pmatrix}_{\text{in}}$$

$$X_{\text{out}} = X_{\text{in}} + 0 \cdot X'_{\text{in}}$$

$$X'_{\text{out}} = (-1/f) \cdot X_{\text{in}} + X'_{\text{in}}$$

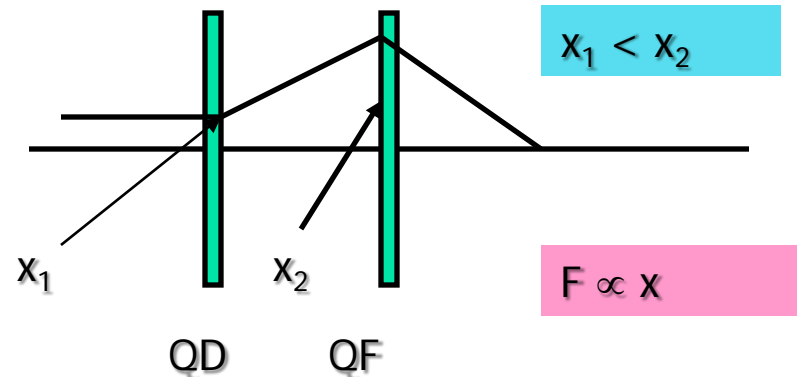


$$\text{Drift} = \begin{pmatrix} 1 & L \\ 0 & 1 \end{pmatrix}$$

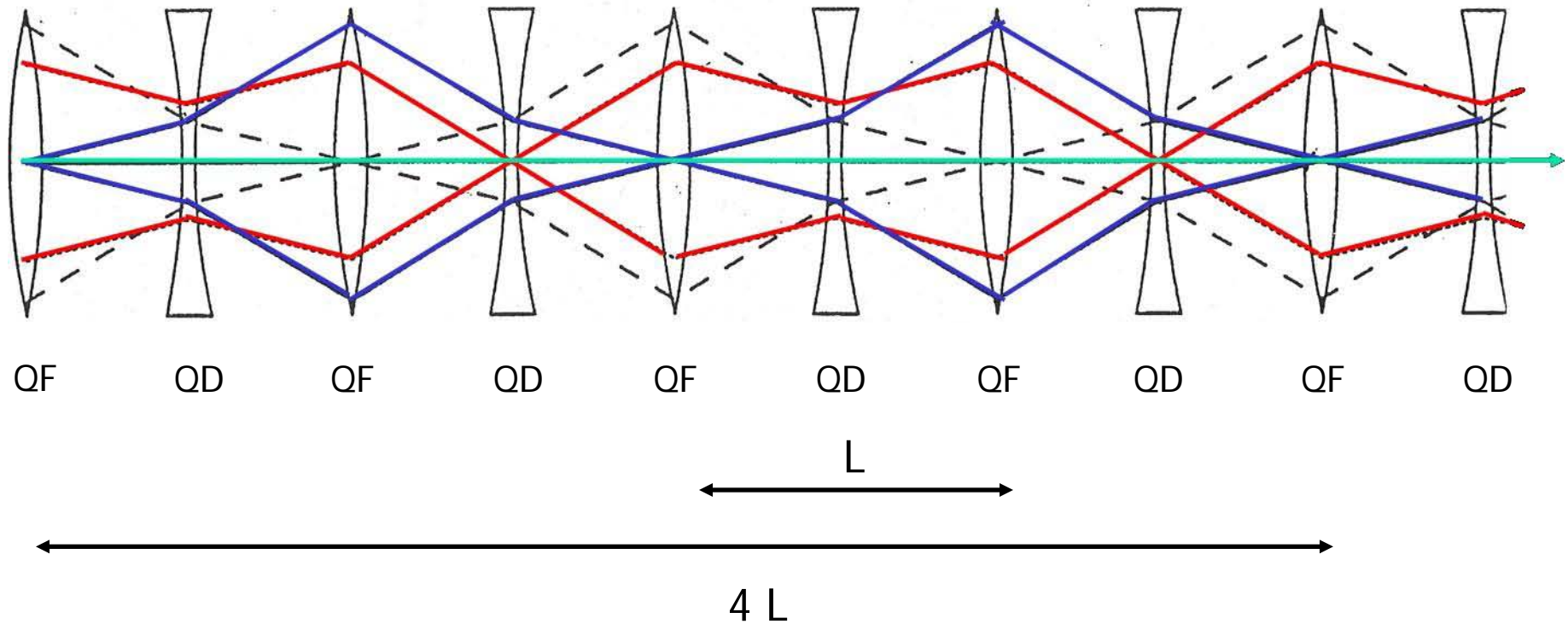
$$\text{QF-Drift-QD} = \begin{pmatrix} 1-L/f & L \\ -L/f^2 & 1+L/f \end{pmatrix}$$

Initial:  $x = x_0$  and  $L < f$   
 $x' = 0$

More intuitively:



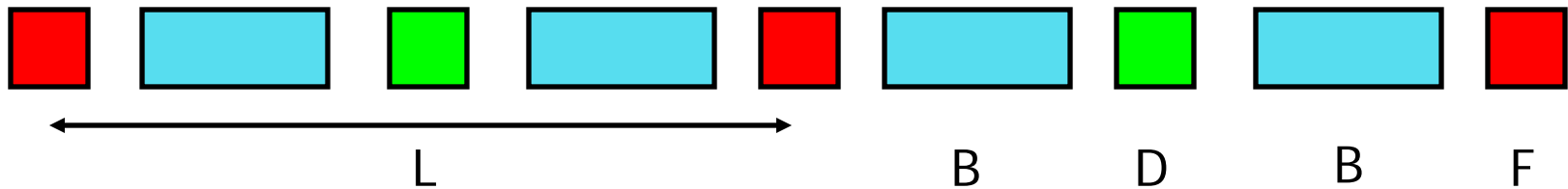
# The concept of the « FODO cell »



One complete oscillation in 4 cells  $\Rightarrow 90^\circ / \text{cell} \Rightarrow \mu = 90^\circ$

# Circular machines (no errors!)

The accelerator is composed of a **periodic** repetition of **cells**:



➤ The phase advance per cell  $\mu$  can be modified, in each plane, by varying the strength of the quadrupoles.

➤ The ideal particle will follow a **particular** trajectory, which **closes on itself** after one revolution: **the closed orbit**.

➤ The real particles will perform oscillations **around the closed orbit**.

➤ The number of **oscillations for a complete revolution** is called the **Tune Q** of the machine ( $Q_x$  and  $Q_y$ ).

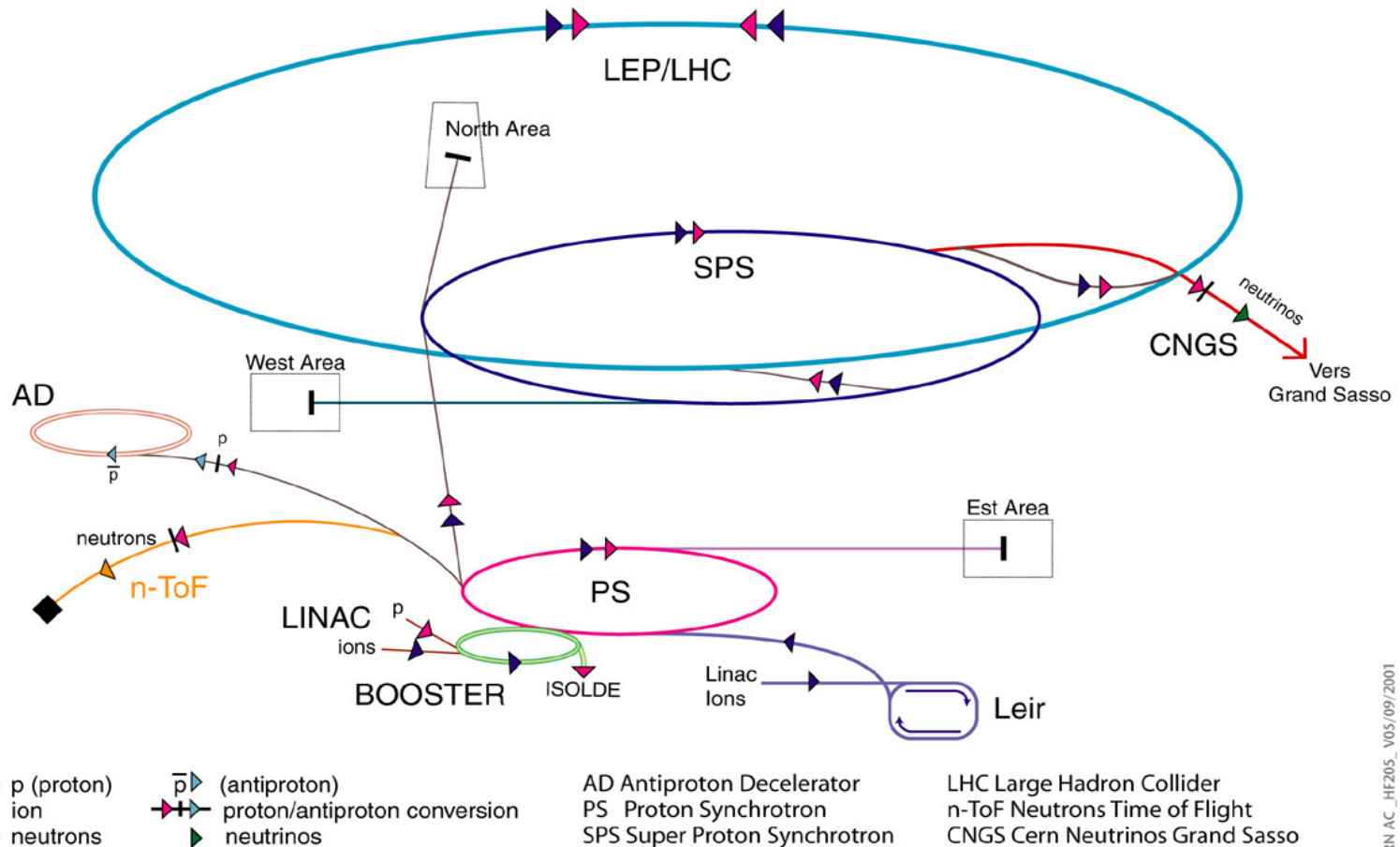
# Regular periodic lattice: The Arc





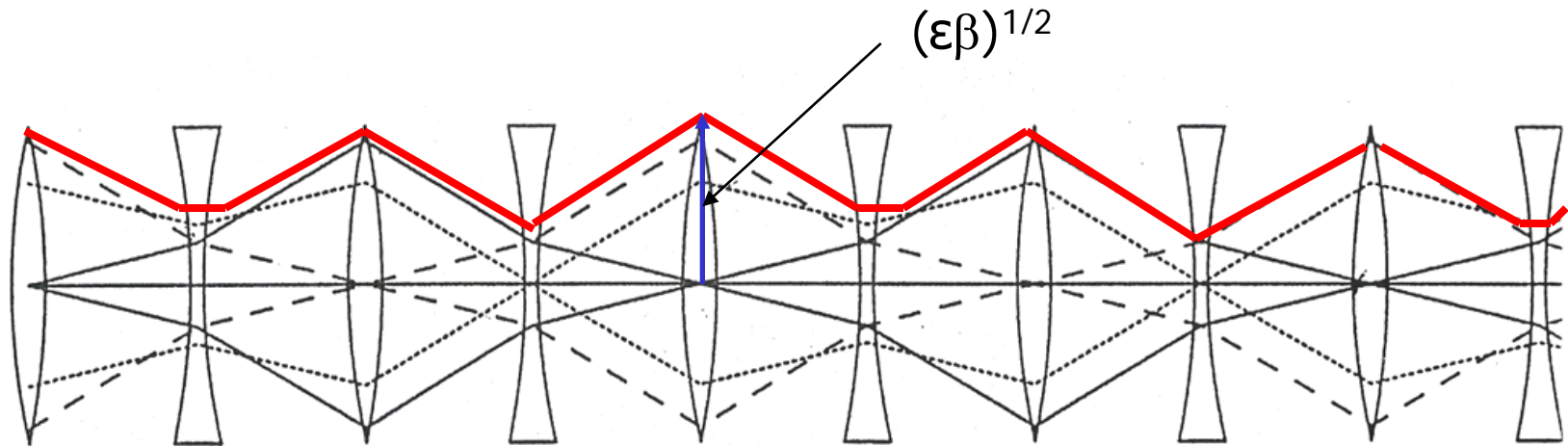
# Synchrotrons ...

## Accelerator chain of CERN (operating or approved projects)



CERN AC\_HF205\_V05/09/2001

# The beta function $\beta(s)$



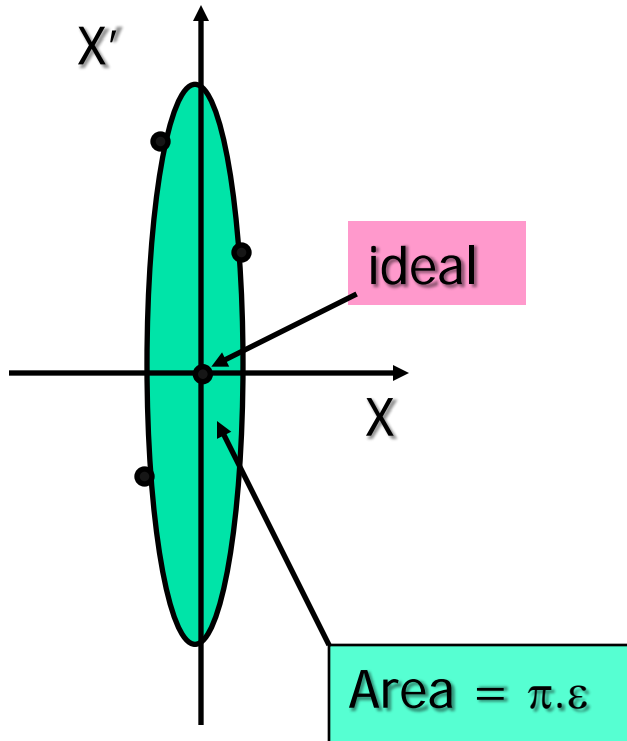
The  $\beta$ -function is the **envelope** around all the trajectories of the particles circulating in the machine.

The  $\beta$ -function has a **minimum at the QD** and a **maximum at the QF**, ensuring the net focusing effect of the lattice.

It is a **periodic function** (repetition of cells). The oscillations of the particles are called **betatron motion** or **betatron oscillations**.

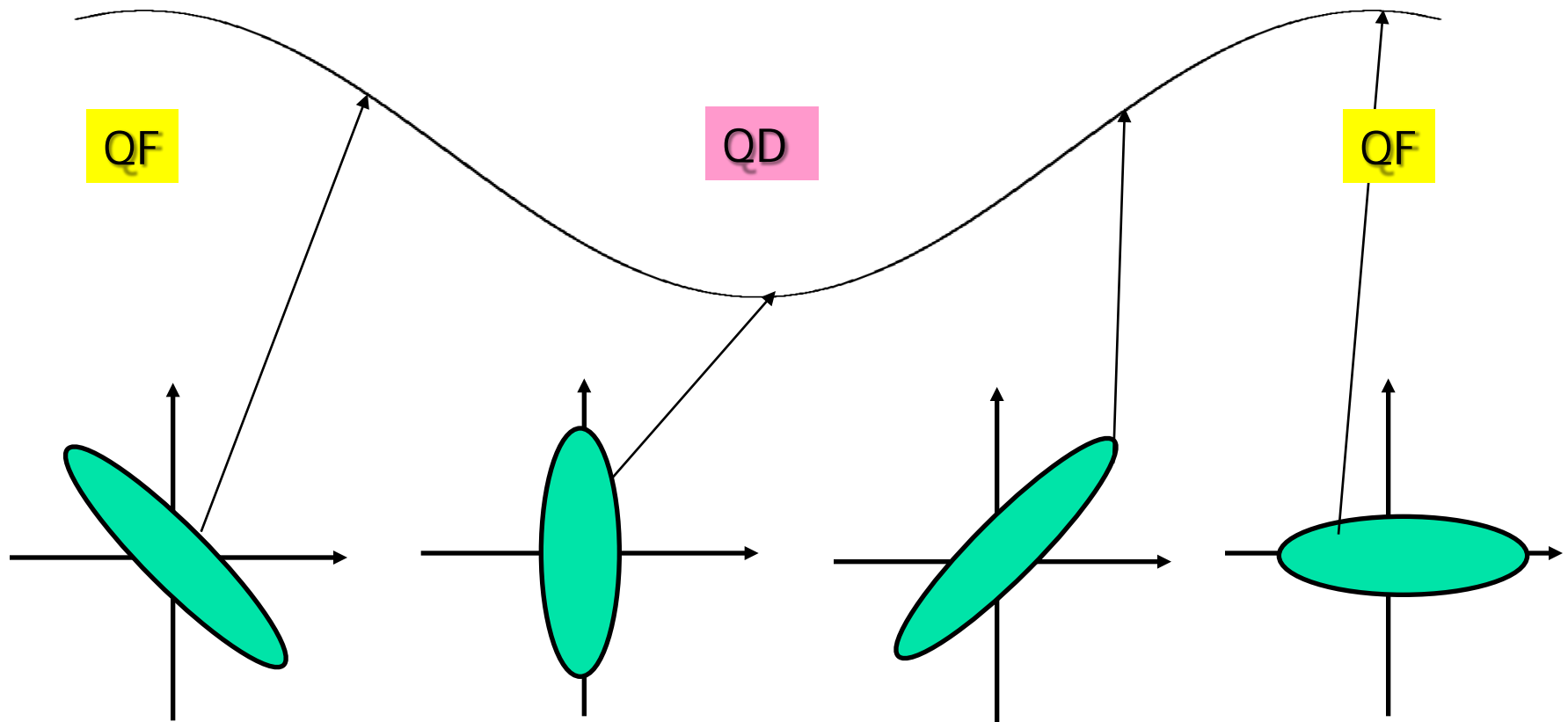
# Phase space

- Select a particle in the beam being at 1 sigma (68%) of the distribution and plot its **position vs. its phase** ( $x$  vs.  $x'$ ) at some location in the machine for many turns.



- $\epsilon$  Is the emittance of the beam [mm mrad]
- $\epsilon$  describes the quality of the beam
- Measure of how much particle depart from ideal trajectory.
- $\beta$  is a property of the machine (quadrupoles).

# Emittance conservation



The shape of the ellipse varies along the machine, but its area (**the emittance  $\epsilon$** ) remains constant at **a given energy**.

# Why introducing these functions?

The  $\beta$  function and the emittance are fundamental parameters, because they are directly related to the beam size (**measurable quantity** !):

Beam size [m]

$$\sigma_{x,y}(s) = (\varepsilon \cdot \beta_{x,y}(s))^{1/2}$$

$$\sigma(\text{IP}) = 17 \mu\text{m}$$

at 7 TeV ( $\beta=0.55$  m)

The emittance  $\varepsilon$  characterises the quality of the injected beam (kind of measure how the particles depart from ideal ones). It is an **invariant** at a given energy.

$\varepsilon$  = beam property

$\beta$  = machine property (quads)



# Recapitulation 1

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- The **fraction** of the oscillation performed in a periodic cell is called the **phase advance  $\mu$  per cell** (x or y).
- The total number of oscillations over **one full turn of the machine** is called the **betatron tune  $Q$**  (x or y).
- The **envelope** of the betatron oscillations is characterised by the **beta function  $\beta(s)$** . This is a **property of the quadrupole settings**.
- The quality of the (injected) beam is characterised by the **emittance  $\epsilon$** . This is a **property of the beam** and is **invariant** around the machine for a given energy.
- The r.m.s. beam size (measurable quantity) is  **$\sigma = (\beta \cdot \epsilon)^{1/2}$** .



# Off momentum particles:

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- These are “non-ideal” particles, in the sense that they do not have the right energy, i.e. all particles with  $\Delta p/p \neq 0$

What happens to these particles when traversing the magnets ?

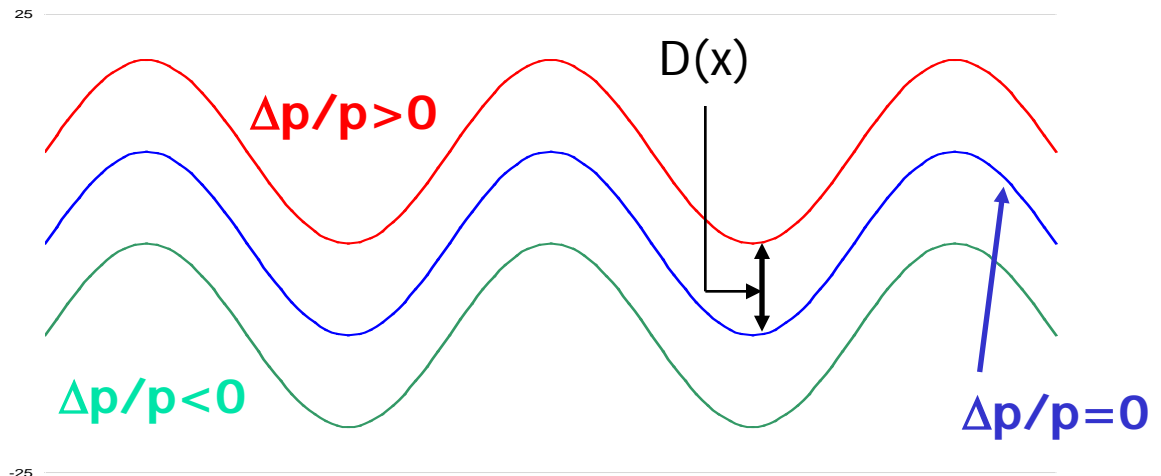
# Off momentum particles ( $\Delta p/p \neq 0$ )

## Effect from Dipoles

- If  $\Delta p/p > 0$ , particles are **less** bent in the dipoles → should spiral out !
- If  $\Delta p/p < 0$ , particles are **more** bent in the dipoles → should spiral in !

**No!**

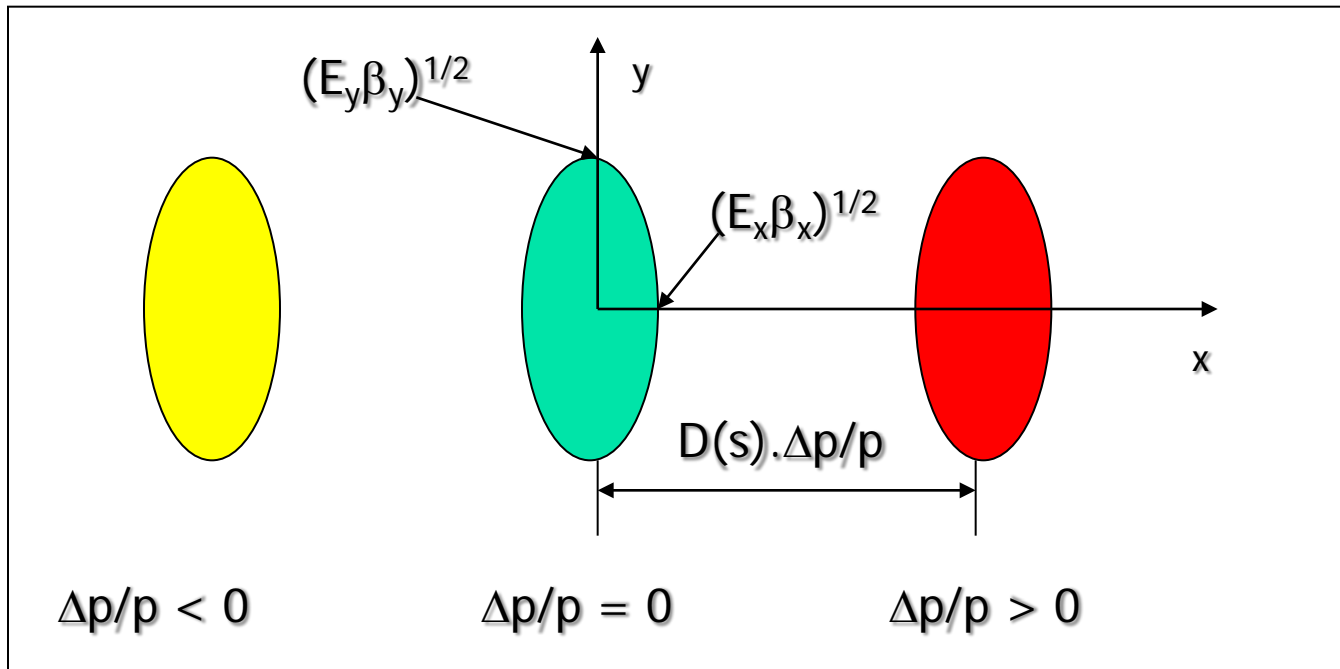
There is an equilibrium with the restoring force of the quadrupoles





# Dispersion

In general:



Only extreme values of  $\Delta p/p$  are shown.

The vacuum chamber must accommodate the full width.

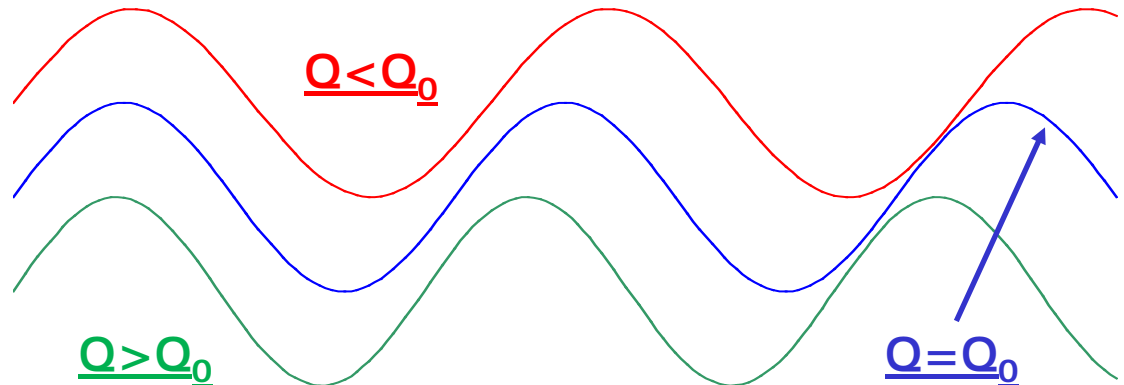
VH:  $A_y(s) = (E_y \beta_y(s))^{1/2}$     and    HW:  $A_x(s) = (E_x \beta_x(s))^{1/2} + D(s) \cdot \Delta p/p$

# Off momentum particles ( $\Delta p/p \neq 0$ )

## Effect from Quadrupoles

- If  $\Delta p/p > 0$ , particles are **less** focused in the quadrupoles → **lower Q !**
- If  $\Delta p/p < 0$ , particles are **more** focused in the quadrupoles → **higher Q !**

Particles with different momenta would have a different **betatron tune**  $Q=f(\Delta p/p)$ !





# The chromaticity $Q'$

Particles with different momenta ( $\Delta p/p$ ) would thus have different tunes  $Q$ .  
So what ?

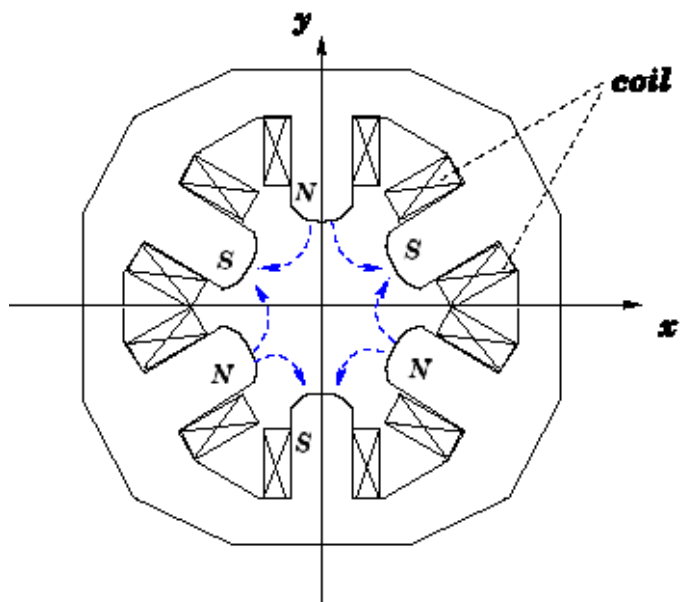
unfortunately

- The tune dependence on momentum is of **fundamental** importance for the **stability** of the machine. It is described by the **chromaticity** of the machine  $Q'$ :

$$Q' = \Delta Q / (\Delta p/p)$$

The chromaticity has to be carefully **controlled and corrected** for stability reasons.

# The sextupoles (SF and SD)

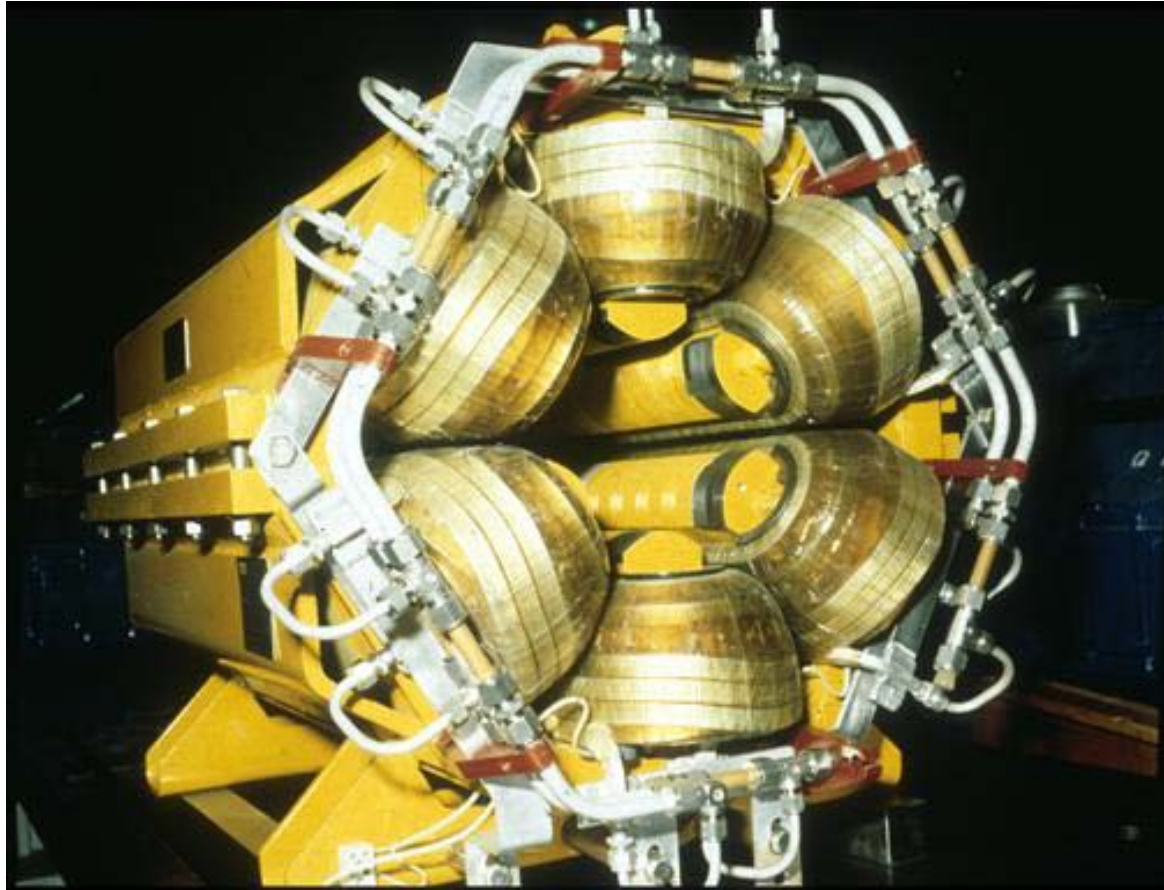


$$\triangleright \Delta X' \propto X^2$$

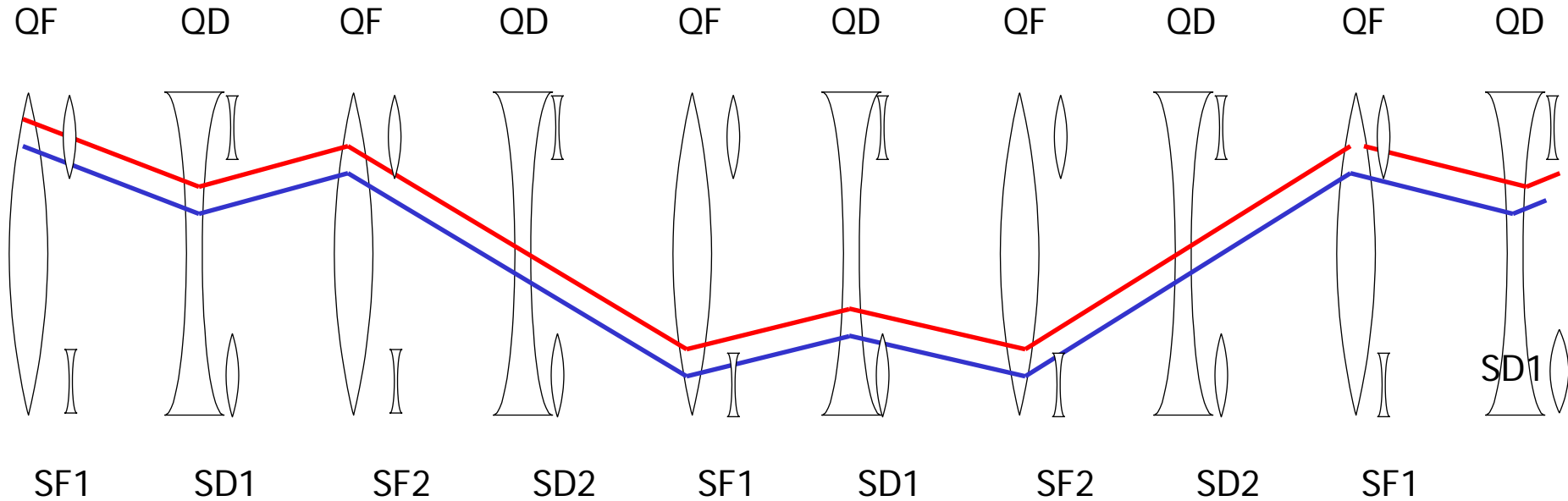
- A SF sextupole basically « adds » focusing for the particles with  $\Delta p/p > 0$ , and « reduces » it for  $\Delta p/p < 0$ .
- The chromaticity is corrected by adding a sextupole after each quadrupole of the FODO lattice.

# Sextupoles:

SPS



# Effect of sextupoles





# Recapitulation 2

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- For off momentum particles ( $\Delta p/p \neq 0$ ), the magnets induce other important effects, namely:
  - The dispersion (dipoles)
  - The chromaticity (quadrupoles)



# Longitudinal plane

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➤ So far, we considered only the motion in the transverse planes from an intuitive point of view. The corresponding rigorous treatment will be given in the lectures on "**Transverse Beam Dynamics**".

➤ The lectures on "**Longitudinal Beam Dynamics**" will explain the details of the corresponding longitudinal motion as well as the RF acceleration of the particles.





# The course:

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Beam Dynamics is certainly a “core” topic of accelerator physics, but the objective of this school is to give you a broader introduction covering:

- Relativity and E.M. Theory                      History, physics and applications
- Particle sources                                      Injection, Extraction
- Transfer Lines                                        Magnets
- Beam Diagnostics                                  Radio Protection
- Linear Imp. and Resonances                    Vacuum
- Synchrotron Radiation, Electron Dynamics, SLS, FELs
- Multi particle Effects                              Power Converters