

# Concept of Luminosity

(in particle colliders)

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([http://cern.ch/lhc-beam-beam/talks/Daresbury\\_luminosity.pdf](http://cern.ch/lhc-beam-beam/talks/Daresbury_luminosity.pdf))

## Why colliding beams ?

■ Two beams:  $E_1, \vec{p}_1, E_2, \vec{p}_2, m_1 = m_2 = m$

$$E_{cm} = \sqrt{(E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2}$$

■ Collider versus fixed target:

Fixed target:  $\vec{p}_2 = 0 \rightarrow E_{cm} = \sqrt{2m^2 + 2E_1 m}$

Collider:  $\vec{p}_1 = -\vec{p}_2 \rightarrow E_{cm} = E_1 + E_2$

■ LHC (pp): 14000 GeV versus  $\approx 115$  GeV

■ LEP ( $e^+e^-$ ): 210 GeV versus ?



# Collider performance issues

- Available energy
  - Number of interactions per second (useful collisions)
  - Total number of interactions
  - Secondary issues:
    - Time structure of interactions (how often and how many at the same time)
    - Space structure of interactions (size of interaction region)
    - Quality of interactions (background, dead time etc.)
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## Luminosity:

■ We want:

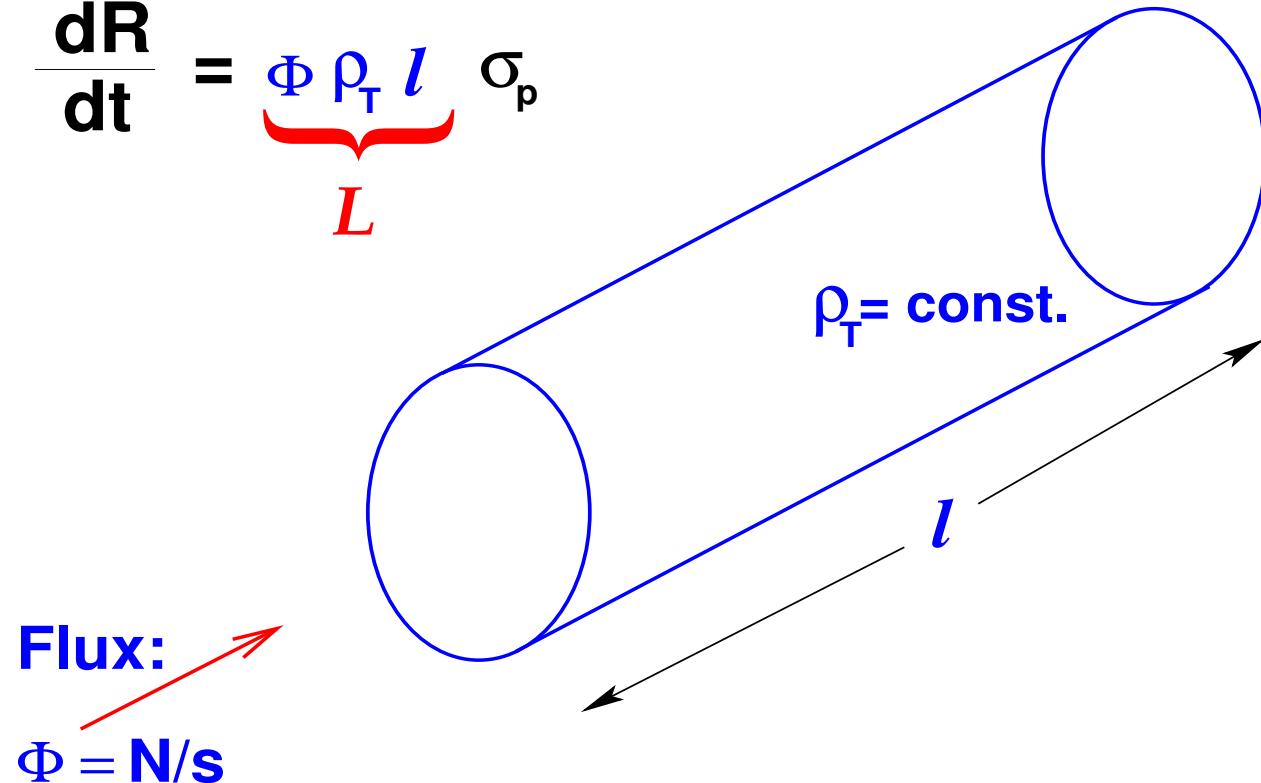
→ Proportionality factor between cross section  $\sigma_p$  and number of interactions per second  $\frac{dR}{dt}$

$$\frac{dR}{dt} = \mathcal{L} \times \sigma_p \quad (\rightarrow \text{units : cm}^{-2}\text{s}^{-1})$$

- Relativistic invariant
  - Independent of the physical reaction
  - Reliable procedures to **compute** and **measure**
-

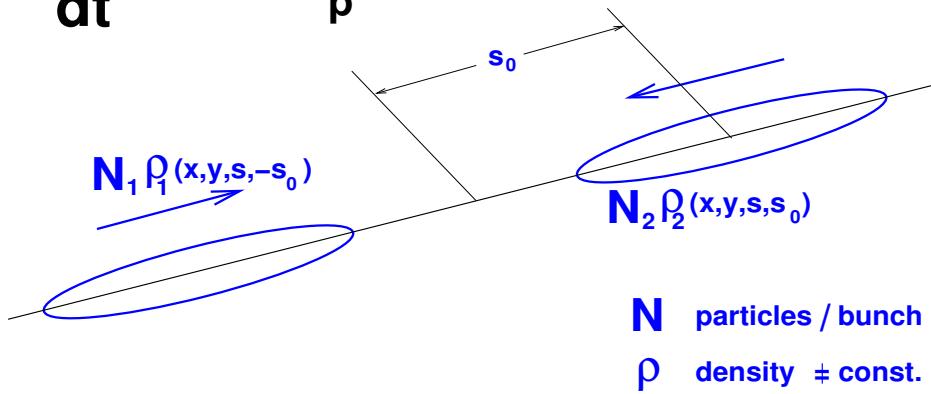
## Fixed target luminosity

$$\frac{dR}{dt} = \underbrace{\Phi \rho_T l}_{L} \sigma_p$$



# Collider luminosity (per bunch crossing)

$$\frac{dR}{dt} = L \sigma_p$$



$$\mathcal{L} \propto K N_1 N_2 \int \int \int \int_{-\infty}^{+\infty} \rho_1(x, y, s, -s_0) \rho_2(x, y, s, s_0) dx dy ds ds_0$$

$s_0$  is "time"-variable:  $s_0 = c \cdot t$

Kinematic factor:  $K = \sqrt{(\vec{v}_1 - \vec{v}_2)^2 - (\vec{v}_1 \times \vec{v}_2)^2 / c^2}$

## Collider luminosity (per beam)

■ Assume uncorrelated densities in all planes

→ factorize:  $\rho(x, y, s, s_0) = \rho_x(x) \cdot \rho_y(y) \cdot \rho_s(s \pm s_0)$

■ For head-on collisions ( $\vec{v}_1 = -\vec{v}_2$ ) we get:

$$\mathcal{L} = 2 \cdot N_1 N_2 \cdot f \cdot n_b \cdot \int \int \int \int_{-\infty}^{+\infty} dx dy ds ds_0 \\ \rho_{1x}(x) \rho_{1y}(y) \rho_{1s}(s - s_0) \cdot \rho_{2x}(x) \rho_{2y}(y) \rho_{2s}(s + s_0)$$

■ In principle: should know all distributions

→ Mostly use Gaussian  $\rho$  for analytic calculation  
(in general: it is a good approximation)

## Gaussian distribution functions

■  $\rho_{iz}(z) = \frac{1}{\sigma_z \sqrt{2\pi}} \exp\left(-\frac{z^2}{2\sigma_z^2}\right) \quad i = 1, 2, \quad z = x, y$

■  $\rho_{is}(s \pm s_0) = \frac{1}{\sigma_s \sqrt{2\pi}} \exp\left(-\frac{(s \pm s_0)^2}{2\sigma_s^2}\right)$

■ For non-Gaussian profiles not always possible to find analytic form, need a numerical integration



## Luminosity for two beams (1 and 2)

### ■ Simplest case : equal beams

→  $\sigma_{1x} = \sigma_{2x}, \quad \sigma_{1y} = \sigma_{2y}, \quad \sigma_{1s} = \sigma_{2s}$

→ but:  $\sigma_{1x} \neq \sigma_{1y}, \quad \sigma_{2x} \neq \sigma_{2y}$  is allowed

### ■ Further: no dispersion at collision point



## Integration (head-on)

for  $\sigma_1 = \sigma_2 \rightarrow \rho_1 \rho_2 = \rho^2$ :

$$\mathcal{L} = \frac{2 \cdot N_1 N_2 f n_b}{(\sqrt{2\pi})^6 \sigma_s^2 \sigma_x^2 \sigma_y^2} \int \int e^{-\frac{x^2}{\sigma_x^2}} e^{-\frac{y^2}{\sigma_y^2}} e^{-\frac{s^2}{\sigma_s^2}} e^{-\frac{s_0^2}{\sigma_s^2}} dx dy ds ds_0$$

integrating over  $s$  and  $s_0$ , using:

$$\int_{-\infty}^{+\infty} e^{-at^2} dt = \sqrt{\pi/a}$$

$$\mathcal{L} = \frac{2 \cdot N_1 N_2 f n_b}{8(\sqrt{\pi})^4 \sigma_x^2 \sigma_y^2} \int \int e^{-\frac{x^2}{\sigma_x^2}} e^{-\frac{y^2}{\sigma_y^2}} dx dy$$

finally after integration over x and y:  $\Rightarrow \mathcal{L} = \frac{N_1 N_2 f n_b}{4\pi \sigma_x \sigma_y}$

## Luminosity for two (equal) beams (1 and 2)

■ Simplest case :  $\sigma_{1x} = \sigma_{2x}, \sigma_{1y} = \sigma_{2y}, \sigma_{1s} = \sigma_{2s}$

or:  $\sigma_{1x} \neq \sigma_{2x} \neq \sigma_{1y} \neq \sigma_{2y}$ , but :  $\sigma_{1s} \approx \sigma_{2s}$

$$\Rightarrow \boxed{\mathcal{L} = \frac{N_1 N_2 f n_b}{4\pi \sigma_x \sigma_y} \left( \mathcal{L} = \frac{N_1 N_2 f n_b}{2\pi \sqrt{\sigma_{1x}^2 + \sigma_{2x}^2} \sqrt{\sigma_{1y}^2 + \sigma_{2y}^2}} \right)}$$

## Examples

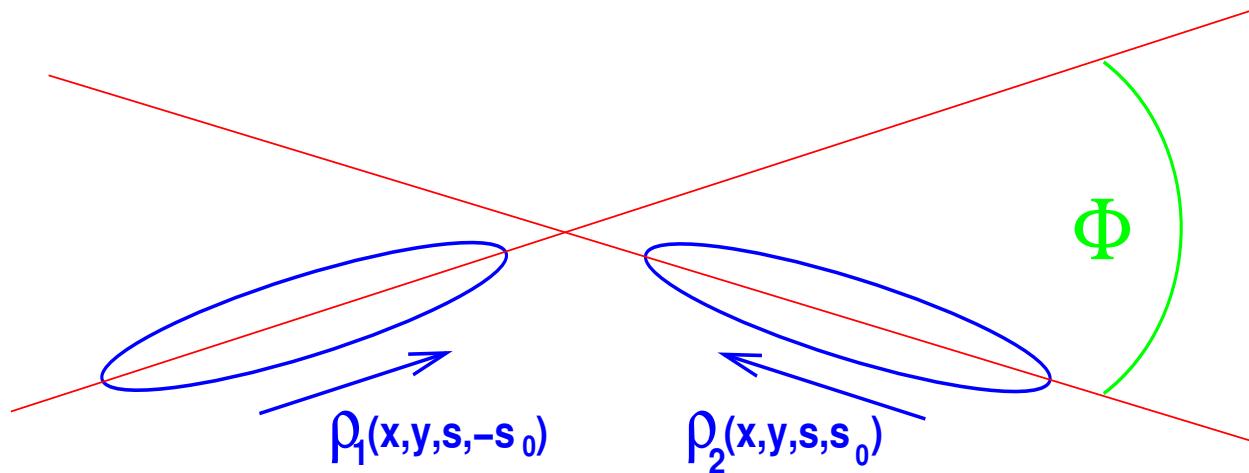
	<b>Energy (GeV)</b>	$\mathcal{L}_{max}$ $\text{cm}^{-2}\text{s}^{-1}$	<b>rate</b> $\text{s}^{-1}$	$\sigma_x/\sigma_y$ $\mu\text{m}/\mu\text{m}$	<b>Particles per bunch</b>
SPS ( $p\bar{p}$ )	<b>315x315</b>	<b><math>6 \cdot 10^{30}</math></b>	<b><math>4 \cdot 10^5</math></b>	<b>60/30</b>	$\approx 10 \cdot 10^{10}$
Tevatron ( $p\bar{p}$ )	<b>1000x1000</b>	<b><math>100 \cdot 10^{30}</math></b>	<b><math>7 \cdot 10^6</math></b>	<b>30/30</b>	$\approx 30/8 \cdot 10^{10}$
HERA ( $e^+p$ )	<b>30x920</b>	<b><math>40 \cdot 10^{30}</math></b>	<b>40</b>	<b>250/50</b>	$\approx 3/7 \cdot 10^{10}$
LHC (pp)	<b>7000x7000</b>	<b><math>10000 \cdot 10^{30}</math></b>	<b><math>10^9</math></b>	<b>17/17</b>	$\approx 11 \cdot 10^{10}$
LEP ( $e^+e^-$ )	<b>105x105</b>	<b><math>100 \cdot 10^{30}</math></b>	$\leq 1$	<b>200/2</b>	$\approx 50 \cdot 10^{10}$
PEP ( $e^+e^-$ )	<b>9x3</b>	<b><math>8000 \cdot 10^{30}</math></b>	<b>NA</b>	<b>150/5</b>	$\approx 2/6 \cdot 10^{10}$



## Complications

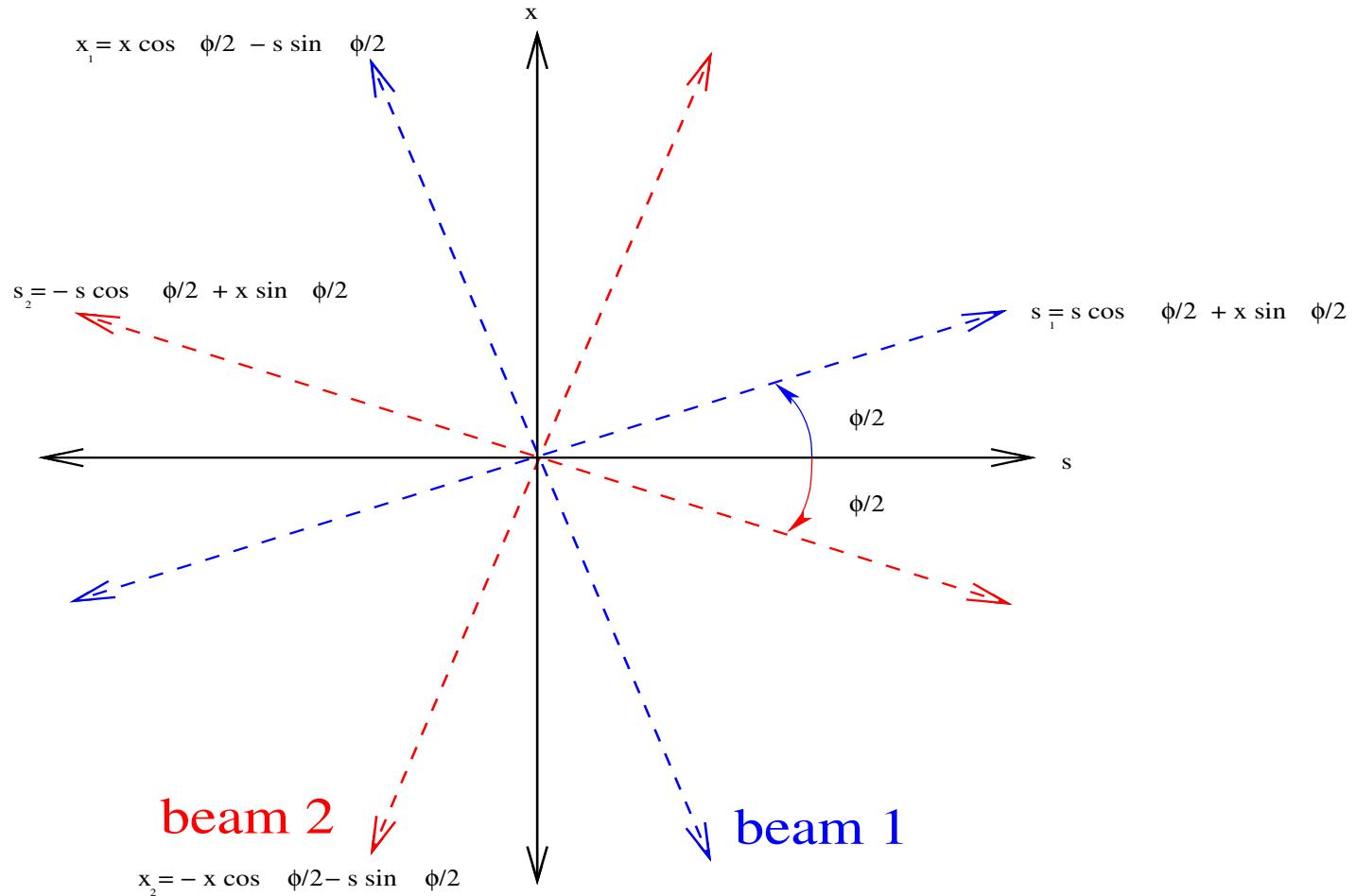
- Crossing angle
  - Hour glass effect
  - Collision offset (wanted or unwanted)
  - Non-Gaussian profiles
  - Dispersion at collision point
  - $\delta\beta^*/\delta s = \alpha^* \neq 0$
  - Strong coupling
  - etc.
-

## Collisions at crossing angle



- Needed to avoid unwanted collisions
  - For colliders with many bunches: LHC, CESR, KEKB
  - For colliders with coasting beams

# Collisions angle geometry (horizontal plane)



## Crossing angle

Assume crossing in **horizontal (x, s)- plane.**

Transform to new coordinates:

$$\begin{cases} x_1 = x \cos \frac{\phi}{2} - s \sin \frac{\phi}{2}, & s_1 = s \cos \frac{\phi}{2} + x \sin \frac{\phi}{2}, \\ x_2 = x \cos \frac{\phi}{2} + s \sin \frac{\phi}{2}, & s_2 = s \cos \frac{\phi}{2} - x \sin \frac{\phi}{2} \end{cases}$$

$$\mathcal{L} = 2 \cos^2 \frac{\phi}{2} N_1 N_2 f n_b \int \int \int \int_{-\infty}^{+\infty} dx dy ds ds_0 \\ \rho_{1x}(x_1) \rho_{1y}(y_1) \rho_{1s}(s_1 - s_0) \rho_{2x}(x_2) \rho_{2y}(y_2) \rho_{2s}(s_2 + s_0)$$



## Integration (crossing angle)

use as before:

$$\int_{-\infty}^{+\infty} e^{-at^2} dt = \sqrt{\pi/a}$$

and:

$$\int_{-\infty}^{+\infty} e^{-(at^2+bt+c)} dt = \sqrt{\pi/a} \cdot e^{\frac{b^2-ac}{a}}$$

Further:

- Since  $\sigma_x$ ,  $x$  and  $\sin(\phi/2)$  are small: drop all terms  $\sigma^k x \sin^l(\phi/2)$  or  $x^k \sin^l(\phi/2)$  for all:  $k+l \geq 4$
  - Approximate:  $\sin(\phi/2) \approx \tan(\phi/2) \approx \phi/2$
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## Crossing angle

■ Crossing Angle  $\Rightarrow$

$$\mathcal{L} = \frac{N_1 N_2 f n_b}{4\pi \sigma_x \sigma_y} \cdot S$$

■  $S$  is the reduction factor

■ For small crossing angles and  $\sigma_s \gg \sigma_{x,y}$

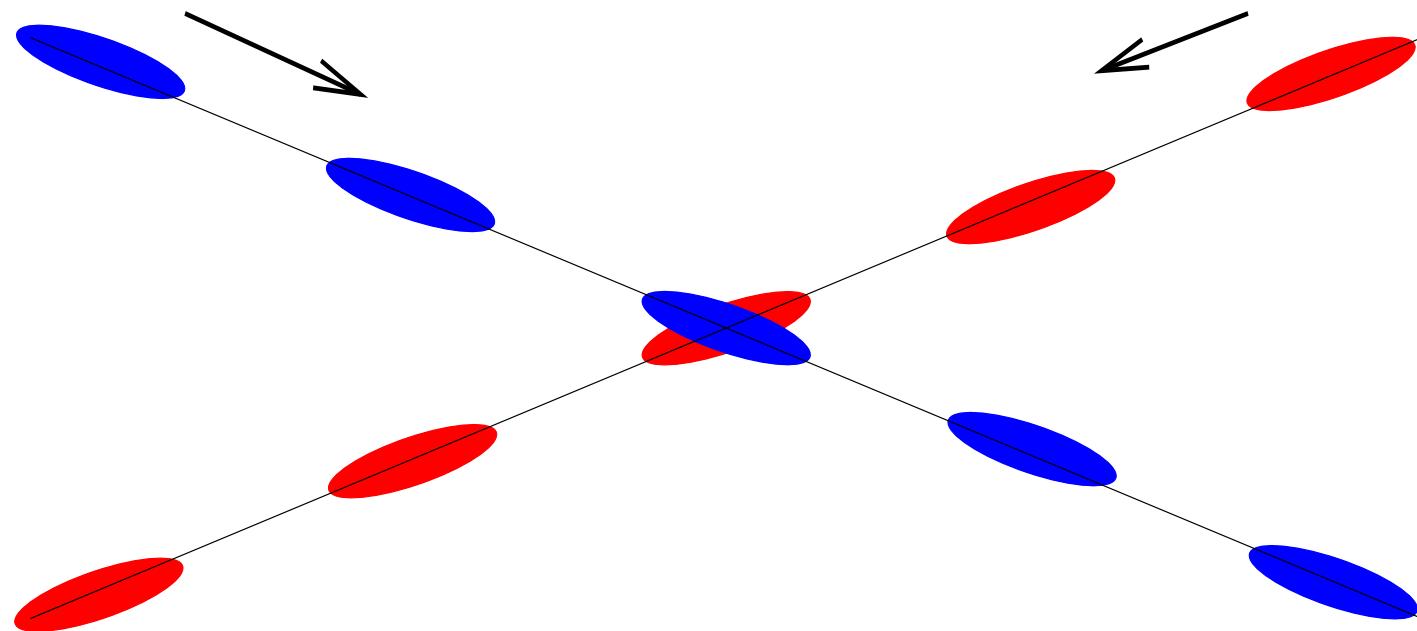
$$\Rightarrow S = \frac{1}{\sqrt{1 + \left(\frac{\sigma_s}{\sigma_x} \tan \frac{\phi}{2}\right)^2}} \approx \frac{1}{\sqrt{1 + \left(\frac{\sigma_s \phi}{\sigma_x 2}\right)^2}}$$

Example LHC:

$\Phi = 285 \text{ } \mu\text{rad}$ ,  $\sigma_s = 7.5 \text{ cm}$ ,  $S = 0.84$

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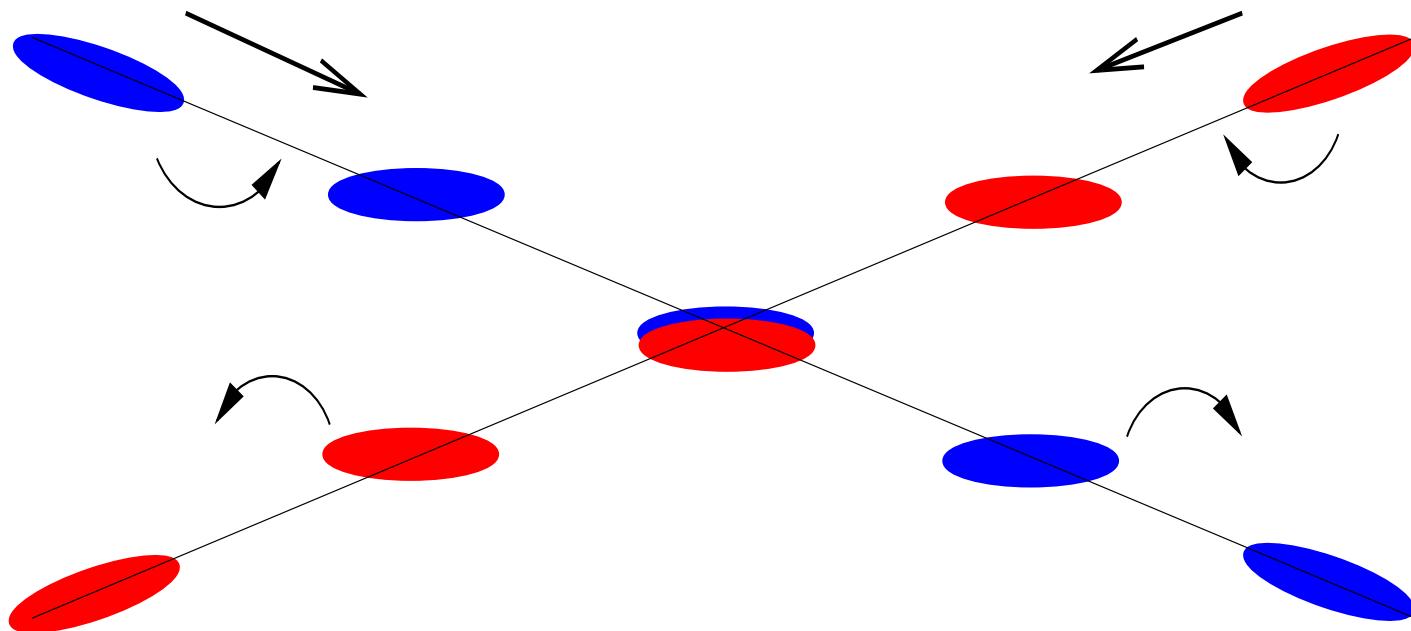
## Large crossing angle



→ Large crossing angle: large loss of luminosity



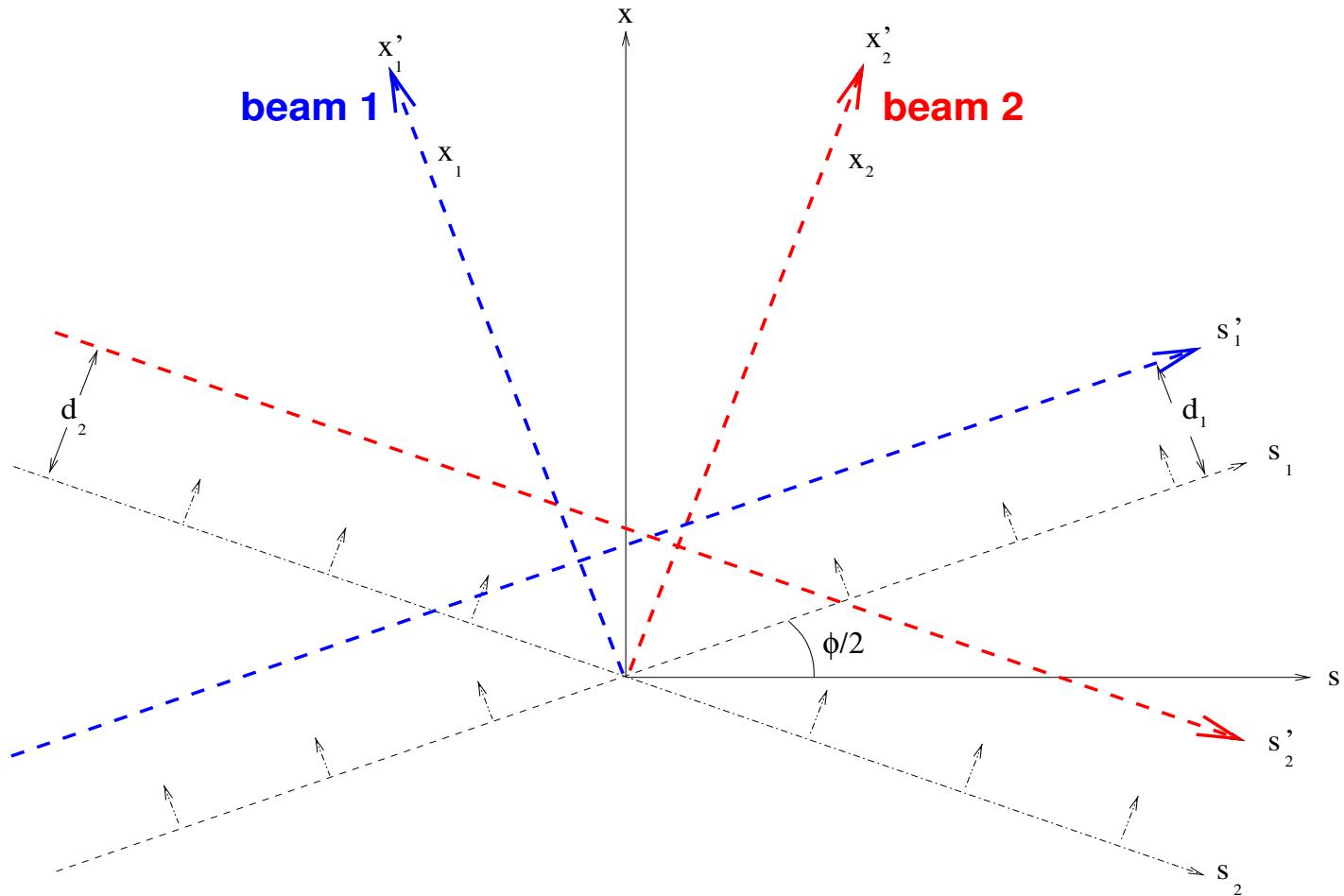
## “crab” crossing scheme



- “crab” crossing recovers geometric loss factor
- feasibility needs to be demonstrated



# Offset and crossing angle



## Offset and crossing angle

■ Transformations with offsets in crossing plane:

$$\begin{cases} x_1 = \textcolor{red}{d}_1 + x \cos \frac{\phi}{2} - s \sin \frac{\phi}{2}, & s_1 = s \cos \frac{\phi}{2} + x \sin \frac{\phi}{2}, \\ x_2 = \textcolor{red}{d}_2 + x \cos \frac{\phi}{2} + s \sin \frac{\phi}{2}, & s_2 = s \cos \frac{\phi}{2} - x \sin \frac{\phi}{2} \end{cases}$$

■ Gives after integration over  $y$  and  $s_0$ :

$$\begin{aligned} \mathcal{L} = \frac{\mathcal{L}_0}{2\pi\sigma_s\sigma_x} 2 \cos^2 \frac{\phi}{2} & \int \int e^{-\frac{x^2 \cos^2(\phi/2) + s^2 \sin^2(\phi/2)}{\sigma_x^2}} e^{-\frac{x^2 \sin^2(\phi/2) + s^2 \cos^2(\phi/2)}{\sigma_s^2}} \\ & \times e^{-\frac{d_1^2 + d_2^2 + 2(d_1 + d_2)x \cos(\phi/2) - 2(d_2 - d_1)s \sin(\phi/2)}{2\sigma_x^2}} dx ds. \end{aligned}$$

## Offset and crossing angle

After integration over x:

$$\mathcal{L} = \frac{N_1 N_2 f n_b}{8\pi^{\frac{3}{2}} \sigma_s} \cdot 2 \cos \frac{\phi}{2} \int_{-\infty}^{+\infty} W \cdot \frac{e^{-(As^2+2Bs)}}{\sigma_x \sigma_y} ds$$

with:

$$A = \frac{\sin^2 \frac{\phi}{2}}{\sigma_x^2} + \frac{\cos^2 \frac{\phi}{2}}{\sigma_s^2} \quad B = \frac{(d_2 - d_1) \sin(\phi/2)}{2\sigma_x^2}$$

and  $W = e^{-\frac{1}{4\sigma_x^2}(d_2-d_1)^2}$

⇒ Luminosity with correction factors

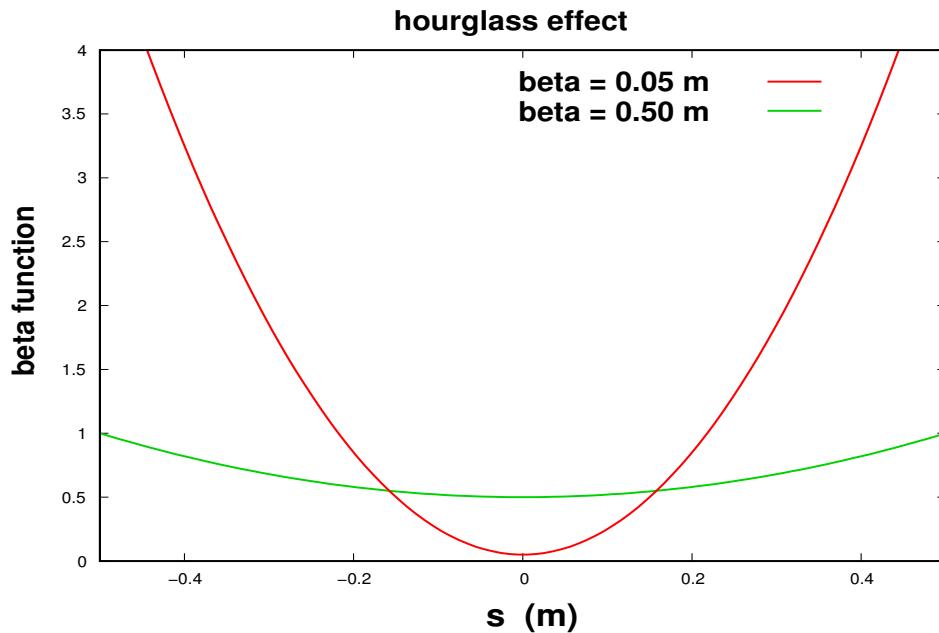
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## Luminosity with correction factors

$$\mathcal{L} = \frac{N_1 N_2 f n_b}{4\pi \sigma_x \sigma_y} \cdot W \cdot e^{\frac{B^2}{A}} \cdot S$$

- $W$ : correction for beam offset
  - $S$ : correction for crossing angle
  - $e^{\frac{B^2}{A}}$ : correction for crossing angle **and** offset
-

# Hour glass effect



- $\beta$ -functions depends on position s
- $\beta(s) \approx \beta^* \left(1 + \left(\frac{s}{\beta^*}\right)^2\right)$

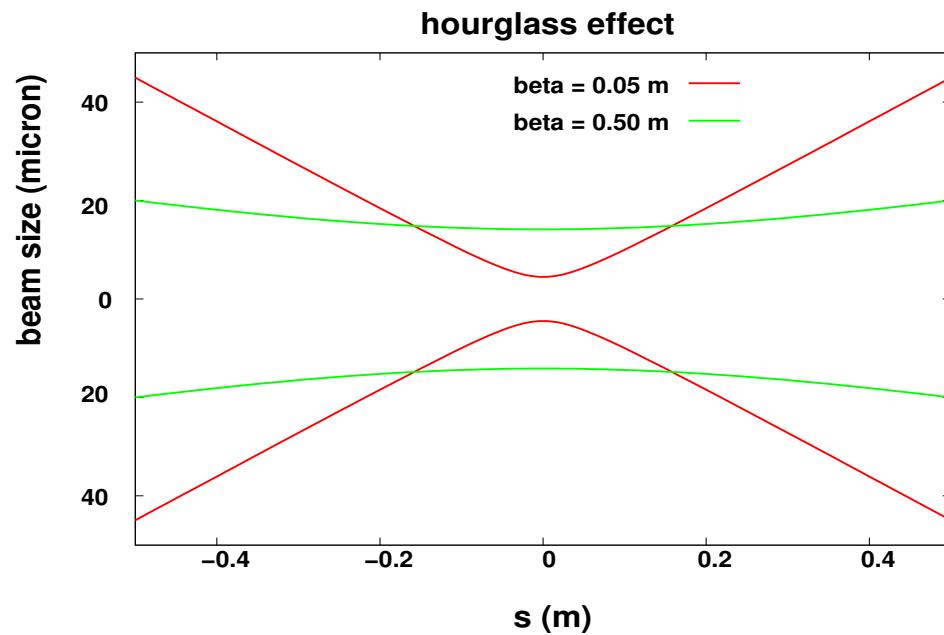
## Hour glass effect



- $\beta$ -functions depends on position s

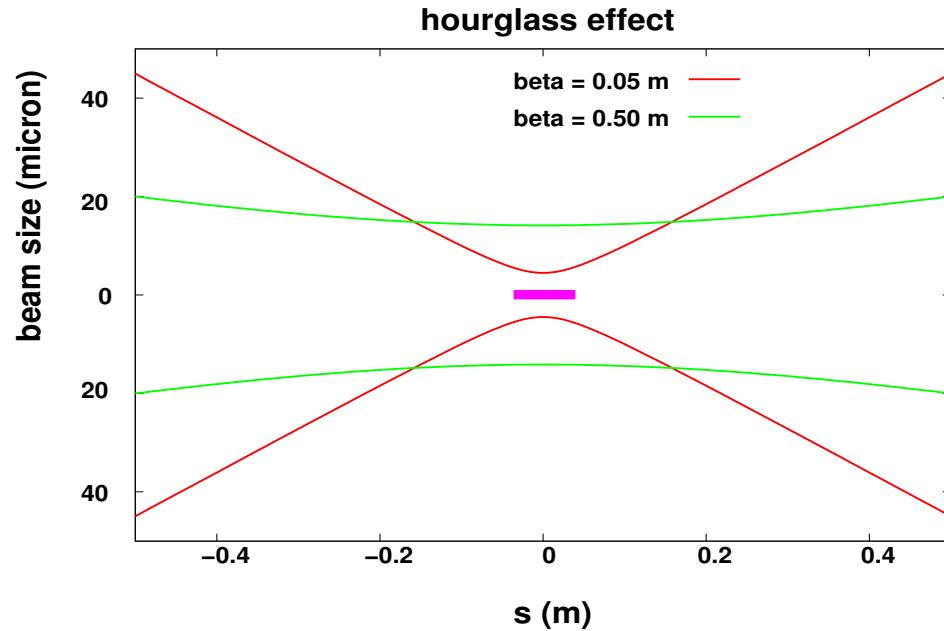


# Hour glass effect



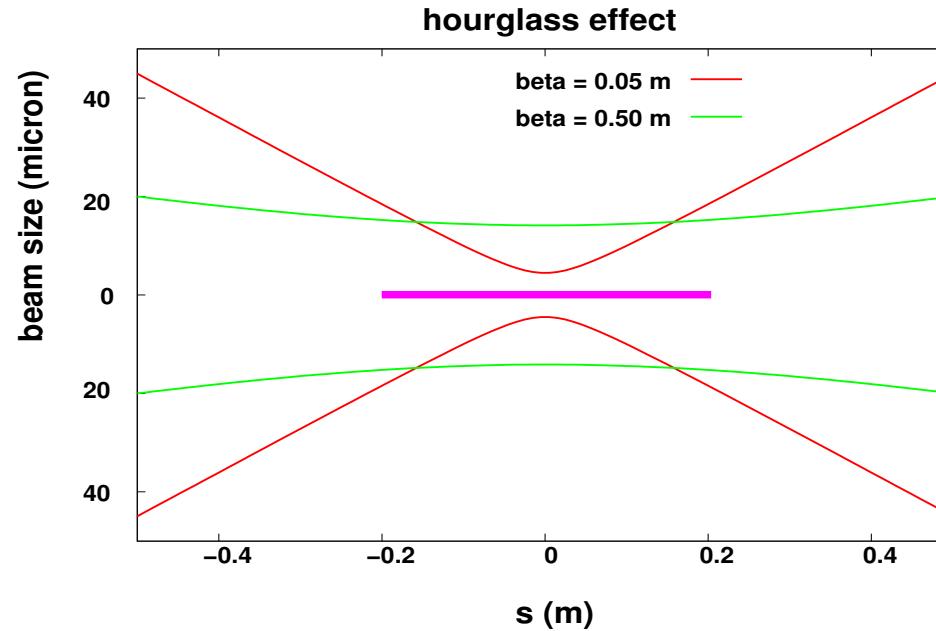
- Beam size  $\sigma \propto \sqrt{\beta^*(s)}$  depends on position s

# Hour glass effect - short bunches



- Small variation of beam size along bunch

# Hour glass effect - long bunches



- Significant effect for long bunches and small  $\beta^*$

## Hour glass effect

- $\beta$ -functions depends on position  $s$
- Usually:  $\beta(s) = \beta^*(1 + \left(\frac{s}{\beta^*}\right)^2)$ 
  - i.e.  $\sigma \implies \sigma(s) \neq \text{const.}$
  - $\sigma(s) = \sigma^* \sqrt{\left(1 + \left(\frac{s}{\beta^*}\right)^2\right)}$
- Important when  $\beta^*$  comparable to the r.m.s. bunch length  $\sigma_s$  (or smaller !)



## Hour glass effect

→ Take it easy:  $\beta_x^* = \beta_y^*$ , crossing angle, but no offset

→ Replace  $\sigma$  by  $\sigma(s)$  in standard formulae

$$\mathcal{L} = \left( \frac{N_1 N_2 f n_b}{8\pi} \right) \frac{2 \cos \frac{\phi}{2}}{\sqrt{\pi} \sigma_s} \int_{-\infty}^{+\infty} \frac{e^{-s^2 A}}{\sigma_x^* \sigma_y^* [1 + (\frac{s}{\beta^*})^2]} ds$$

$$A = \frac{\sin^2 \frac{\phi}{2}}{(\sigma_x^*)^2 [1 + (\frac{s}{\beta^*})^2]} + \frac{\cos^2 \frac{\phi}{2}}{\sigma_s^2}$$

→ Numerical Integration

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## Calculations for the LHC

- $N_1 = N_2 = 1.15 \times 10^{11}$  particles/bunch
  - $n_b = 2808$  bunches/beam
  - $f = 11.2455$  kHz,  $\phi = 285$   $\mu\text{rad}$
  - $\beta_x^* = \beta_y^* = 0.55$  m
  - $\sigma_x^* = \sigma_y^* = 16.6$   $\mu\text{m}$ ,  $\sigma_s = 7.7$  cm
-

■ Simplest case (Head on collision):

$$\mathcal{L} = 1.200 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$$

■ Effect of crossing angle:

$$\mathcal{L} = 0.973 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$$

■ Effect of crossing angle & Hourglass:

$$\mathcal{L} = 0.969 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$$



## If the beams are not Gaussian ??

Exercise:

- Assume flat distributions (normalized to 1)

$$\rho_1 = \rho_2 = \frac{1}{2a}, \quad \text{for } [-a \leq z \leq a], z = x, y$$

Calculate r.m.s. in x and y:

$$\langle (x, y)^2 \rangle = \int_{-\infty}^{+\infty} (x, y)^2 \cdot \rho(x, y) dx dy$$

and

$$\mathcal{L} = \int_{-\infty}^{+\infty} \rho_1(x, y) \rho_2(x, y) dx dy$$

- Compute:  $\mathcal{L} \cdot \sqrt{\langle x^2 \rangle \cdot \langle y^2 \rangle}$
- Repeat for various distributions and compare



## Integrated luminosity

- $\mathcal{L}_{\text{int}} = \int_0^T \mathcal{L}(t) dt$

- The figure of merit:

$$\mathcal{L}_{\text{int}} \cdot \sigma_p = \text{number of events}$$

- Experiments: continuous recording of  $\mathcal{L}$
  - For studies: assume some life time behaviour.  
E.g.  $\mathcal{L}(t) \longrightarrow \mathcal{L}_0 \exp\left(-\frac{t}{\tau}\right)$
  - Contributions to life time from: intensity decay, emittance growth etc.
-

## Integrated luminosity

- Knowledge of preparation time allows optimization of  $\mathcal{L}_{\text{int}}$



## Integrated luminosity

- Typical run times LEP:  
 $t_r \approx 8 - 10$  hours
- For LHC long preparation time  $t_p$  expected
  - Optimum combination of  $t_r$  and  $t_p$  gives maximum luminosity
  - $t_r$  is usually a "free" parameter, i.e. can be chosen



## Maximising Integrated Luminosity

- Assume exponential decay of luminosity

$$\mathcal{L}(t) = \mathcal{L}_0 \cdot e^{t/\tau}$$

- Average luminosity  $\langle \mathcal{L} \rangle$

$$\langle \mathcal{L} \rangle = \frac{\int_0^{t_r} dt \mathcal{L}(t)}{t_r + t_p} = \mathcal{L}_0 \cdot \tau \cdot \frac{1 - e^{-t_r/\tau}}{t_r + t_p}$$

- (Theoretical) maximum for:

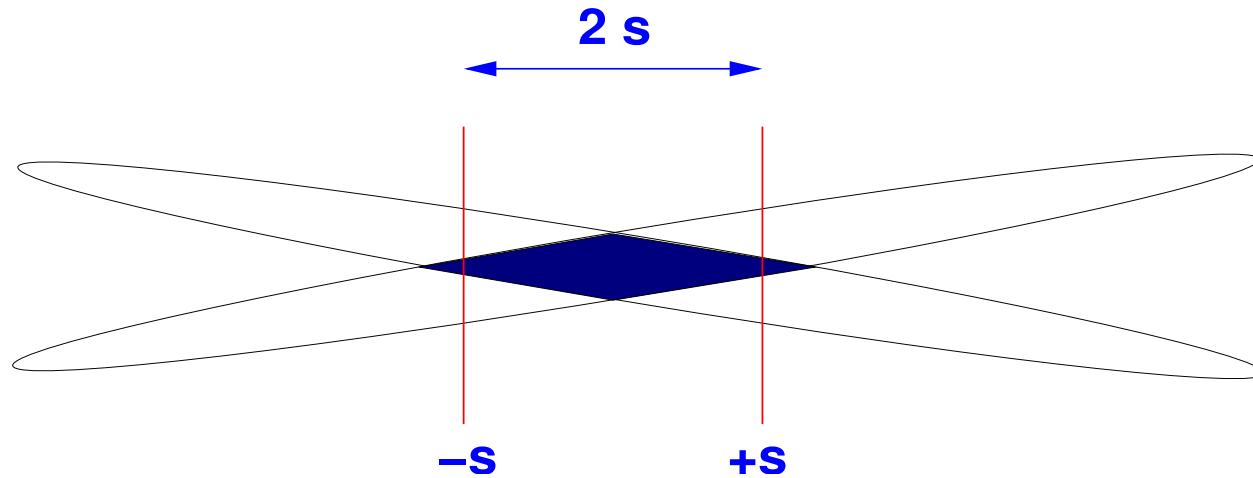
$$t_r \approx \tau \cdot \ln(1 + \sqrt{2t_p/\tau} + t_p/\tau)$$

- Example LHC:  $t_p \approx 10\text{h}$ ,  $\tau \approx 15\text{h}$ ,  $\Rightarrow t_r \approx 15\text{h}$

- Exercise: Would you improve  $\tau$  (long  $t_r$ ) or  $t_p$  ?



## Luminous region



- Density distribution of interaction vertices
- Fraction of collisions occur  $\pm s$  from the IP ?
- Important for experiments !



## Luminous region

- Depends on  $\sigma_x$ ,  $\sigma_y$ ,  $\sigma_s$  and crossing angle  $\phi$
- Integrate only along a finite longitudinal length:

$$\mathcal{L}_0 = \int_{-\infty}^{+\infty} \mathcal{L}(s') ds' \longrightarrow \mathcal{L}(\textcolor{blue}{S}) = \int_{-\textcolor{blue}{S}}^{+\textcolor{blue}{S}} \mathcal{L}(s') ds'$$

$$\mathcal{L}(\textcolor{blue}{S}) = \left( \frac{N_1 N_2 f n_b}{8\pi \sigma_x^* \sigma_y^*} \right) \frac{2 \cos \frac{\phi}{2}}{\sqrt{\pi} \sigma_s} \sqrt{\frac{\pi}{A}} \operatorname{erf} \left( \sqrt{A} \textcolor{blue}{S} \right)$$

- Real figure of merit:  $\mathcal{L}_{\text{int}}(\textcolor{blue}{S}) = \int_0^T \int_{-\textcolor{blue}{S}}^{+\textcolor{blue}{S}} \mathcal{L}(s', t) ds' dt$



## Some results for LHC

- $\sigma_s = 7.7 \text{ cm}, \beta^* = 0.55 \text{ m}, \phi = 285 \mu\text{rad}:$
  
  - 100% lumi  $\rightarrow S = \pm 12 \text{ cm}$
  
  - 90% lumi  $\rightarrow S = \pm 7 \text{ cm}$
  
  - 80% lumi  $\rightarrow S = \pm 5.5 \text{ cm}$
-

## Interactions per crossing

- Luminosity /  $fn_b \propto N_1 N_2$
  - In LHC: crossing every 25 ns
  - Per crossing approximately 20 interactions
  - May be undesirable (pile up in detector)
  - $\Rightarrow$  more bunches  $n_b$ , or smaller N ??
- Beware: maximum (peak) luminosity  
 $\mathcal{L}_{max}$  is not the whole story ... !



## Luminosity measurement

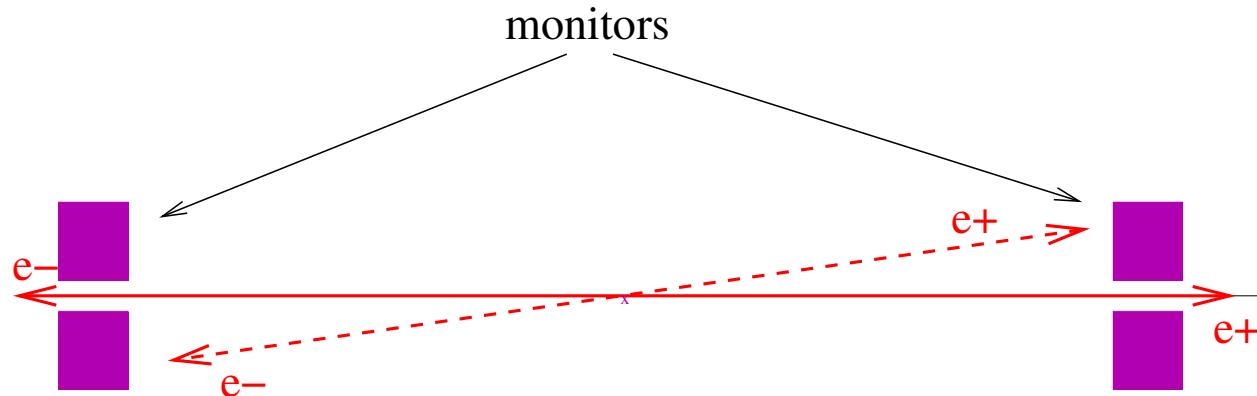
- One needs to get a signal proportional to interaction rate → Beam diagnostics
  - Large dynamic range:  
 $10^{27} \text{ cm}^{-2}\text{s}^{-1}$  to  $10^{34} \text{ cm}^{-2}\text{s}^{-1}$
  - Very fast, if possible for individual bunches
  - Used for optimization
  - For absolute luminosity need calibration
-

## Luminosity calibration

$(e^+e^-)$

- Use well known and calculable process
  - $e^+e^- \rightarrow e^+e^-$  elastic scattering (Bhabha scattering)
  - Have to go to small angles ( $\sigma_{el} \propto \Theta^{-3}$ )
  - Small rates at high energy ( $\sigma_{el} \propto \frac{1}{E^2}$ )
-

# Luminosity calibration



- Measure coincidence at small angles
  - Low counting rates, in particular for high energy !
  - Background may be problematic
-

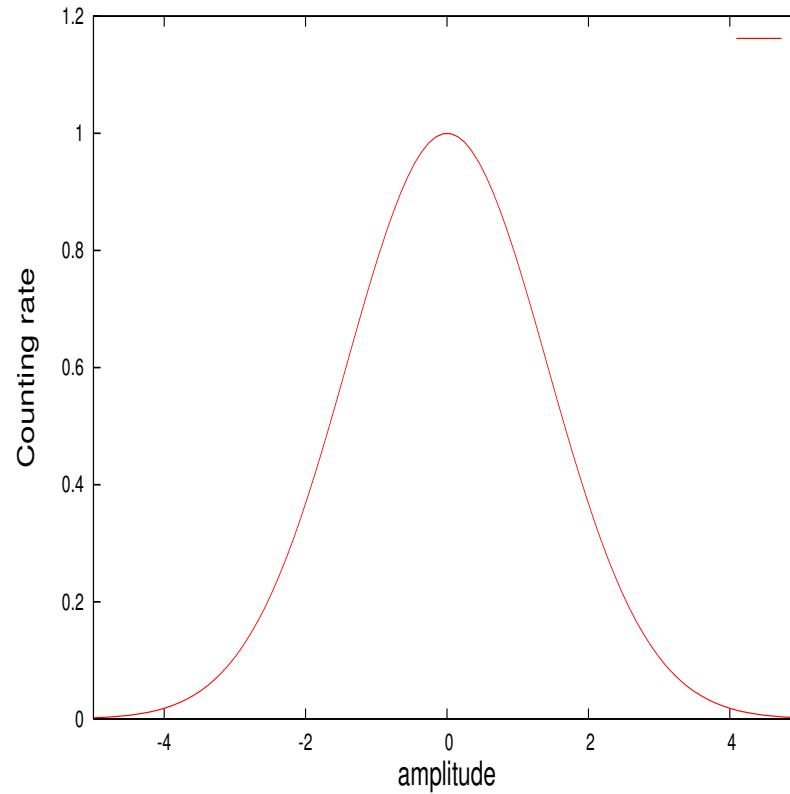
# Luminosity calibration

(hadrons, e.g.  $pp$  or  $p\bar{p}$ )

- Must measure beam current and beam sizes
  - Beam size measurement:
    - Wire scanner or synchrotron light monitors
    - Measurement with beam ... → remember luminosity with offset
    - Move the two beams against each other in transverse planes (van der Meer scan)
-

# Luminosity optimization

Van der Meer scan



Record counting rates  $R(d)$   
as function of movement  $d$

► Since  $R(d)$  is proportional to  
luminosity  $L(d)$

► Get ratio of luminosity  
 $L(d)/L(0)$

## Luminosity optimization

- From ratio of luminosity  $\mathcal{L}(d)/\mathcal{L}_0$
  - Remember:  $W = e^{-\frac{1}{4\sigma^2}(d_2-d_1)^2}$
  - Determines  $\sigma$
  - ... and centres the beams !
  - ”beam-beam deflection scans ...”  $\Rightarrow$
-

## Absolute value of $\mathcal{L}$ ( $pp$ or $p\bar{p}$ )

### ■ By total rate and optical theorem

(also: luminosity independent determination of  $\sigma_{tot}$ ):

- $\sigma_{tot} \cdot \mathcal{L} = N_{inel} + N_{el}$  (Total counting rate)

- $\lim_{t \rightarrow 0} \frac{d\sigma_{el}}{dt} = (1 + \rho^2) \frac{\sigma_{tot}^2}{16\pi} = \frac{1}{\mathcal{L}} \frac{dN_{el}}{dt}|_{t=0}$

→ 
$$\mathcal{L} = \frac{(1 + \rho^2)}{16\pi} \frac{(N_{inel} + N_{el})^2}{(dN_{el}/dt)|_{t=0}}$$

### ■ Luminosity determined from experimental rates



## Absolute value of $\mathcal{L}$ ( $pp$ or $p\bar{p}$ )

■ By Coulomb normalization:

- Coulomb amplitude exactly calculable:

$$\begin{aligned} \bullet \quad \lim_{t \rightarrow 0} \frac{d\sigma_{el}}{dt} &= \frac{1}{\mathcal{L}} \frac{dN_{el}}{dt} \Big|_{t=0} = \pi |f_C + f_N|^2 \\ &\simeq \pi \left| \frac{2\alpha_{em}}{-t} + \frac{\sigma_{tot}}{4\pi} (\rho + i) e^{b\frac{t}{2}} \right|^2 \simeq \frac{4\pi\alpha_{em}^2}{t^2} \Big|_{|t| \rightarrow 0} \end{aligned}$$

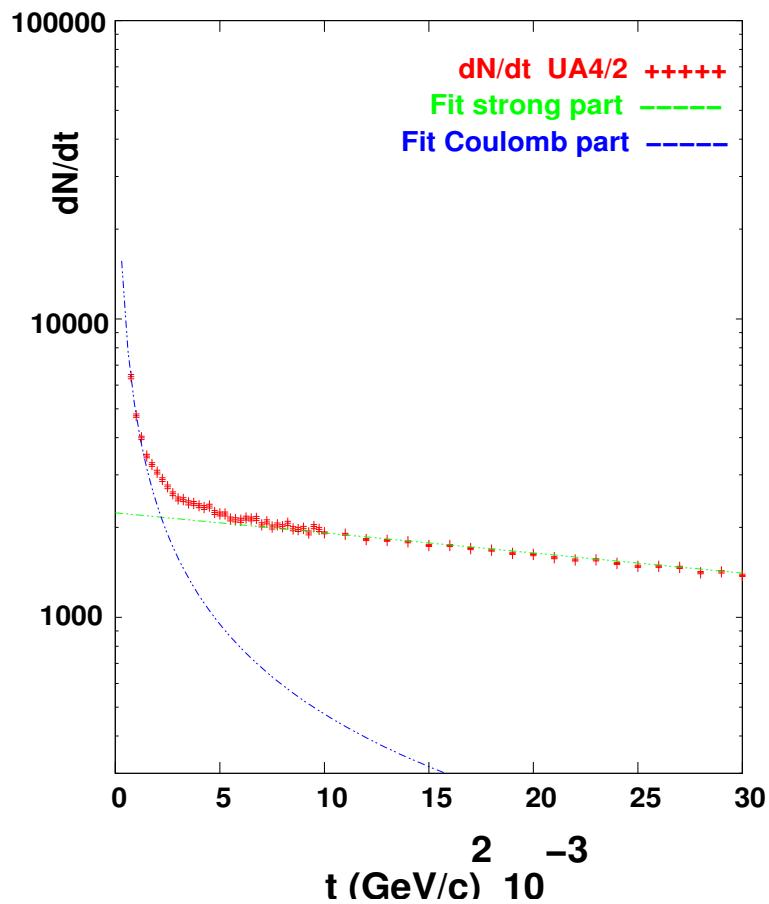
- Fit gives:  $\sigma_{tot}, \rho, b$  and  $\mathcal{L}$

■ Can be done measuring **only** elastic scattering  
(No  $N_{inel}$  needed !)

■  $\Rightarrow$  Roman pots



## Differential elastic cross section



- Measure  $dN/dt$  at small  $t$  ( $0.01 < (\text{GeV}/c)^2 < 10$ ) and extrapolate to  $t = 0.0$
- Needs special optics to allow measurement at very small  $t$
- Measure total counting rate  $N_{\text{el}} + N_{\text{inel}}$   
Needs good detector coverage
- Often use slightly modified method, precision 1 – 2 %



## Not treated :

- Coasting beams (e.g. ISR)
  - Asymmetric colliders (e.g. PEP, HERA)
  - Linear colliders (SLC, TESLA etc.)
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# How to cook high Luminosity ?

- Get high intensity
- Get small beam sizes (small  $\epsilon$  and  $\beta^*$ )
- Get many bunches
- Get small crossing angle (if any)
- Get exact head-on collisions
- Get short bunches

