

## Magnets

Gijs de Rijk CERN

CAS - Basics of Accelerator Physics and Technology

Chavannes-de-Bogis, 7th February 2017



### Contents

- Introduction
  - magnetic field and magnet principles for resistive and superconducting magnets
  - Field description and magnet types
- Resistive magnets
- Superconducting magnets
- Literature on Magnets



## The 2 Maxwell equations relevant for magnets



With:

$$\vec{B} = \mu \vec{H} = \mu_0 \mu_{\rm r} \vec{H}$$



## Magnetic field quality: multipole description

$$B_{y}(z) + iB_{x}(z) = 10^{-4}B_{1}\sum_{n=1}^{\infty} (b_{n} + ia_{n}) \left(\frac{x + iy}{R_{ref}}\right)^{n-1}$$

with:

z = x + iy,

 $B_x$  and  $B_y$  the flux density components in the x and y direction,  $R_{ref}$  the radius of the reference circle,

 $B_1$  the dipole field component at the reference circle,

 $b_n$  the normal nth multipole component,

 $a_n$  the skew nth multipole component.

In a ring shaped accelerator, where the beam does multiple passes, one typically demands :

 $a_n, b_n \leq 1 \text{ unit } 10^{-4}$ 



## **Magnetic fields**

From Ampere's law with no time dependencies, We can derive the law of Biot and Savart





If you wanted to make a B = 1.5 T magnet with just two infinitely thin wires placed at 100 mm distance in air one needs : I = 187500 A

- To get reasonable fields (B > 1 T) one needs large currents
- Moreover, the field homogeneity will be poor

And If you wanted to make a B = 8 T one needs :  $I = 5 \cdot 10^5$  A (LHC dipole coil 80 turns of 11850 A at 8.3 T = 9.48 \cdot 10^5 A )

For LHC dipole@8.3 T ~1 MA in 3300 mm<sup>2</sup> : ~300 A/mm<sup>2</sup> (overall current density in the coil area)







## Iron dominated magnets (aka: Classical, Resistive)

With the help of an iron yoke we can get fields with less current

Example: C shaped dipole for accelerators



 $[\tilde{\mathbf{0}}]_{\mathcal{C}}\vec{H} \times d\vec{l} = N \times I$  $N \times I = H_{iron} \times l_{iron} + H_{airgap} \times l_{airgap} \triangleright$  $N \times I = \frac{B}{M_0 M_r} \times l_{iron} + \frac{B}{M_0} \times l_{airgap} \bowtie$  $N \times I = \frac{l_{airgap} \times B}{M_0}$  This is valid as  $\mu_r \gg \mu_0$  in the iron : limited to B < 2 T coil B = 1.5 T**M**, as function of B for low carbon magnet steel (Magnetil BC) Gap = 50 mm9000 8000 *N* . *I* = 59683 A 7000 6000 2 x 30 turn coil m<sub>r 5000</sub> 4000 I = 994 A3000 2000 1000 @5 A/mm<sup>2</sup>, 200 mm<sup>2</sup> 0 2 0.5 1.5 0 1 14 x 14 mm Cu

B (T)



## Comparison : iron magnet and air coil

Imagine a magnet with a 50 mm vertical gap (horizontal width ~100 mm) Iron magnet wrt to an air coil:

- Up to 1.5 T we get ~6 times the field
- Between 1.5 T and 2 T the gain flattens of : the iron saturates
- Above 2 T the slope is like for an air-coil: currents become too large to use resistive coils







## Types of magnet fields for accelerators



Courtesy D. Tommasini, CERN SKEW : horizontal field on mid-plane



## **Basic resistive magnet types**

Magnet		Pole shape	<b>Transfer function</b>
NI/	w : pole width g : vertical gap	parallel	$B=\mu_0 NI/g$
w NI/2	w : pole width g : vertical gap	parallel	$B=\mu_0 NI/g$
d E Ni Ni total ampertures = 2N	w : pole width g : pole gap t : coil width	parallel	B=µ0NI/g
"	w : pole width g : pole gap t : coil width	parallel	$B=\mu_0 NI/g$
	R : aperture radius d : coil distance t : coil width	2xy=R <sup>2</sup>	$\begin{array}{c} B(r)=G \cdot r\\ G=2\mu_0 NI/R^2 \end{array}$
	R : aperture radius d : coil distance t : coil width	$3x^2y-y^3=R^3$	$\begin{array}{c} B(r) = S \cdot r^2 = \frac{1}{2}B'' \cdot r^2 \\ S = 3\mu_0 NI/R^3 \end{array}$



## "Free" Coils for generating the Perfect Dipole Field

- Used for magnets where the field is fully determined by the coil only (not by the poles), mostly the case for superconducting magnets
  Conductors 2 solid Intercepting ellipses (or circles)
  - A uniform, opposite polarity, current density in the area of two intersecting ellipses produces a pure dipolar field, but:
    - The aperture is not circular
    - Not easy to simulate with a flat cable

Thick conductor shell with a  $cos\theta$  current distribution  $J = J_0 \cos\Theta$ 

- Pure dipolar field
- Easier to reproduce with a flat rectangular cable









## Magnet types and higher orders

a "pure" multipolar field can be generated by a specific coil geometry





## What is specific about sc accelerator magnets ?

- Cylindrical volume with perpendicular field
- Dipoles, quadrupoles, etc



Field quality:







Artist view of a dipole, from M. N. Wilson « Superconducting Magnets »

 $\cos\Theta$  coil :  $J = J_0 \cos\Theta$ 



Field quality formulated and measured in a multipole expansion,

$$B_{y} + iB_{x} = 10^{-4} B_{1} \overset{\stackrel{\times}{a}}{\underset{n=1}{\overset{}}} (b_{n} + ia_{n}) \overset{\stackrel{\circ}{b}}{\underset{e}{\overset{}}} \frac{x + iy}{R_{ref}} \overset{\stackrel{\circ}{o}^{n-1}}{\underset{o}{\overset{}}}$$

- $b_n, a_n \in few \times units$
- Long magnets: dipoles from 6 m (Tevatron) to 15 m (LHC)
- Often magnets are bend (9.14 mm sagitta for the LHC dipoles)



## **Resistive magnets**

Also called: Classical magnets, Iron dominated magnets



## Practical resistive magnet design & manufacturing

Steps in the process:

- 1. Specification
- 2. Conceptual design
- 3. Raw materials choice
- 4. Detailed design
  - 1. Coil cross-section geometry: cooling
  - 2. Yoke shape, pole shape: FE model optimization
  - 3. Yoke ends, coil ends design
- 5. Yoke manufacturing, tolerances, alignment, structure
- 6. Coil manufacturing, insulation, impregnation type
- 7. Magnetic field measurements



## **Specification**

Before you start designing you need to get from the accelerator designers:

- B(T) or G (T/m) (higher orders:  $G_3(T/m^2)$ , etc)
- Magnet type: C-type, H-type, DC (slow ramp) or AC (fast ramp)
- Aperture:
  - Dipole : "good field region"  $\rightarrow$  airgap height and width
  - quads and higher order: "good field region"  $\rightarrow$  aperture inscribed circle
- Magnetic length and estimated real length
- Current range of the power convertor (and the voltage range: watch out for the cables )
- Field quality:

dipole:  $\frac{\Delta B}{B}$  (ref volume), quadrupole:  $\frac{\Delta G}{G}$  (reference circle)

or  $b_n, a_n$  for n = 1, 2, 3, 4, 5, ...

- Cooling type: air, water ( $\mathsf{P}_{max}$  ,  $\Delta p_{max}$  and  $\mathsf{Q}_{max}$  (I/min)
- Jacks and Alignment features
- Vacuum chamber to be used  $\rightarrow$  fixations, bake-out specifics

These need careful negotiation and often iteration after conceptual (and detailed) design



## **Conceptual design**

• From *B* and *I* you get *NI* (A)

$$NI = \frac{l_{airgap}B}{\mu_0}$$

- From NI (A) and the power convertor I<sub>max</sub> you get N
- Then you decide on a coil X-section using:

 $j_{coil} = 5 \ ^{A}/_{mm^{2}}$  for water cooled or  $j_{coil} = 1 \ ^{A}/_{mm^{2}}$  for air cooled

- This defines the coil cavity in the yoke (you add 0.5 mm insulation around each conductor and 1 mm ground insulation around the coil) and select the best fitting rectangular
- You can the draw the draft X-section using:

$$W_{yoke} = W_{pole} \frac{B}{B_{sat}}$$
 with  $1.5 T < B_{sat} < 2 T$ 

- Decide on the coil ends: racetrack, bedstead
- You now have the rough magnet cross section and envelope









## **Power and cooling**

Power generated by coil

• DC: from the length of the conductor  $N \cdot L_{turn}$ , the cross section  $\sigma$  and the specific resistivity  $\rho$  of the material one gets the spent Power in the coil

$$P/l[W/m] = \frac{\rho}{S}I^2 \quad \text{with:} \qquad \qquad \rho_{Cu} = 1.72(1+0.0039(T-20))10^{-8}\Omega m$$
$$\rho_{Al} = 2.65(1+0.0039(T-20))10^{-8}\Omega m$$

For pulsed magnets: take the average I<sup>2</sup> for the duty cycle

### Coil Cooling

- □  $\Delta T$  (entrance-exit of the coil) is usually chosen as 20°C or 30°C depending on the  $T_{cooling water}$ )
- The cooling water needs to be in moderately turbulent regime (with laminar flow the flow speed is zero on the wall !) (*Reynolds* > 2000)
  - In general a few bars (eg. 4 bar) pressure drop is chosen



## **Theoretical pole shapes**

The ideal poles for dipole, quadrupole, sextupole, etc. are lines of constant scalar potential

Dipole  $y = \pm h/2$  straight line

quadrupole  $2xy = \pm r^2$  hyperbola

sextupole  $3x^2y - y^3 = \pm r^3$ 



## Practical pole shapes: shims and alignment features

 Dipole example: below a lamination of the LEP main bending magnets, with the pole shims well visible



- Quadrupoles: at the edge of the pole one can put a combination a shim and alignment feature (examples: LHC-MQW, SESAME quads, etc)
- This then also allows to measure the pole distances : special instrumentation can be made for this







## Finite Element electromagnetic models

- Aim of the electromagnetic FE models:
  - The exact shape of the yoke needs to be designed
    - Optimize field quality: adjust pole shape, minimize high saturation zones
    - Minimize the total steel amount (magnet weight, raw material cost)
  - Calculate the field: needed for the optics and dynamic aperture modelling
    - transfer function  $B_{xsection}(I)$ ,  $\int Bdl$ , magnetic length
    - multipoles (in the centre of the magnet and integrated)  $b_n$  and  $a_n$
  - Some Electromagnetic FE software packages that are often used:
    - Opera from Cobham: 2D and 3D commercial software see: <u>http://operafea.com/</u>
    - "Good old" Poisson, 2D: now distributed by LANL-LAACG see: http://laacg.lanl.gov/laacg/services/download\_sf.phtml
    - ROXIE (CERN) 2D and 3D, specialized for accelerator magnets; single fee license for labs & universities see: <u>ttps://espace.cern.ch/roxie/default.aspx</u>
    - ANSYS Maxwell: 2D and 3D commercial software see: <u>http://www.ansys.com/Products/Electronics/ANSYS-Maxwell</u>



## Yoke shape, pole shape: FE model optimization

Use symmetry and the thus appropriate boundary conditions to model only  $\frac{1}{4}$ <sup>th</sup> (dipoles, quadrupoles) or even 1/6<sup>th</sup> sextupoles.

Meshing needs attention in the detailed areas like

poles, slits, etc

Magnet type	MQWA	MQWB
Magnetic length	3	3.1 m
Beam separation	22	24 mm
Aperture diameter	4	6 mm
Operating temperature	< 65° C	
Nominal gradient	35 T/m	30 T/m
Nominal current	710 A	600 A
Inductance	2	8 mH
Resistance	3	$7 \text{ m}\Omega$
Conductor V section	20.5 x 18.0	mm <sup>2</sup> inner poles
Conductor X-section	17.0 x 17.0	mm <sup>2</sup> outer poles
Cooling hole diameter	7 mm i	nner poles,
Cooling note diameter	8 mm outer poles	
Number of turns per magnet	8	x 11
Minimum water flow	28	3 1/min
Dissipated power at Inom	19 kW	14 kW
Mass	11	700 kg







# 7-Febr-2017, Chavannee-de-Bogis, Basics, CAS

GdR

Magnets,

Btot (T)

2.069 1.960 1.852 1.743 1.634 1.525

1.416 1.307 1.198 1.089 0.980 0.871 0.762 0.653 0.544 0.436 0.327 0.218 0.109



## Yoke manufacturing

- Yokes are nearly always laminated to reduce eddy currents during ramping
- Laminations can be coated with an inorganic (oxidation, phosphating, Carlite) or organic (epoxy) layer to increase the resistance
- Magnetic properties: depend on chemical composition + temperature and mechanical history
- Important parameters: coercive field  $H_c$  and the saturation induction.
  - $H_c$  has an impact on the remnant field at low current
    - $H_c < 80$  A/m typical or  $H_c < 20$  A/m for magnets at low field (B < 0.05 T)
- Iow carbon steel (C content < 0.006%) is best for higher fields B > 1 T
- Yokes are either :
  - Glued, using epoxy coated laminations
  - Welded, full length plates are welded on the outside
  - Compressed by tie rods in holes
  - or a combination of all these
  - To be able to keep the yoke (or yoke stack) stable you probably need end plates (can range from  $\pm$  1 cm to 5 cm depending on the size)
  - The end plates have pole chamfers and often carry end shims



### Yoke manufacturing

Stacking an MBW dipole yoke stack



Double aperture LHC quad. MQW Stacking on a precision table





#### Welding the structural plates



#### Half magnet handling



#### Glued yoke (MCIA LHC TL)





## Coil manufacturing, insulation, impregnation type

- Winding Cu conductors is an well established technique
- When the Cu conductor is thick it is best to use "dead soft" Cu (T treatment)
- Insulation of the coil
  - Glass fibre epoxy impregnated
    - Individual conductor 0.5 mm glass fibre, 0.25 mm tape wound half lapped
    - Impregnated with radiation resistant epoxy, total glass volume ratio >50%
  - For thin conductors: Cu emanel coated, possibly epoxy impregnated afterwards

## For dipoles some main types are racetrack or bedstead







#### Quadrupoles





## **Coil manufacturing**

MQW Glass fibre tape wrapping.







## Winding the hollow Cu conductor



#### coil electrical test (under water !)



Finished MBXW coil





## **Magnetic field measurements**

Several Magnetic Measurements techniques can be applied, e.g.:

- Rotating coils: multipoles and integrated field or gradient in all magnets
- Stretched wire: magnetic centre and integrated gradient for n > 1 magnets
- Hall probes: field map
- Pickup coils: field on a current ramp
- Example below : MQW : double aperture quadrupole for the LHC.



Mole Assembly Top View



### **Permanent magnets**

#### Linac4 @ CERN permanent magnets , quadrupoles



Pictured : Cell-Coupled Drift Tube Linac module.

- Permanent magnet because of space between DTL tanks
- Sm<sub>2</sub>Co<sub>17</sub> permanent magnets
- Integrated gradient of 1.3 to 1.6 Tesla
- 15 magnets
- Magnet length 0.100 m
- Field quality/amplitude tuning blocks









## Examples of resistive magnets;

## Some history, some modern regular magnets and some special cases



GdR

## Some early magnets (early 1950-ies)

Bevatron (Berkeley) 1954, 6.2 GeV











## **PS combined function dipole**

Magnetic field:			Equation of hyperbolic part: $(243.00+r) z = 12150 mm^2$
	at injection	147 G	
	for 24.3 GeV	1.2 T	
<u>s</u> dR	maximum	1.4 T	
ن م	Weight of one magnet unit	38 t	
let	Gradient @1.2 T : 5 T/m		
lagi			192.0
Σ	Equipped with pole-face		57*28'
017	windings for higher order	Water cooled AI race-	Toterance within the shaded area 10,02 mm
2(	CONTECTIONS Connection of the PFW main windings	track coils	n <sub>nem</sub> = 268,4 R = 70079mm
ebr	for R type magnet Defocusing upstream Ring center Ring center Focusing downstream		FINAL POLE PROFILE.
7-F(			Fig. 9: Final pole profile.
S, I			
ogi		100 to 100	Refinaçe cerars Dove-teller state
4 B			
-de			Coll window
Jee	Connection of the PFW main windings		
anı	for magnet type S Focusing upstream ring center pring center Defocusing downstream		
Jav		A A REAL RANK	
Ù			
ics			a) "OPEN" BLOCK, b) "CLOSED PLOCK, b) "CLOSED PLOCK, c) "CLOSED PLOCK, c) b) "C) "CLOSED PLOCK, c) b) "C) "C) "CLOSED PLOCK, c) "C) "C) "C) "C) "C) "C) "C)
3as		K K C	Fig. 12: Final form of the magnet blocks.
S			
Q			Courtesy D. Tommasini, CERN 30



## dipole magnet : SPS dipole

H magnet type MBB B = 2.05 T Coil : 16 turns  $I_{max} = 4900$  A Aperture = 52 × 92 mm<sup>2</sup> L = 6.26 m Weight = 17 t







## **Quadrupole magnet : SPS quadrupole**







type MQ G = 20.7 T/m Coil : 16 turns  $I_{max} = 1938$  A Aperture inscribed radius = 44 mm  $L_{coil} = 3.2$  m Weight = 8.4 t



## **MBW LHC** warm separation dipole (1)





## MQW: LHC warm double aperture quadrupole







## Superconducting magnets



## Superconducting accelerators magnets; the state of the art

- Maximum attainable field slowly approaches 16 T
  - 20% margin needed (80% on the load line):
    - for a 16 T nominal field we need to design for 20 T




# **Forces and Stored Energy**

Scaling of **forces** on coil quadrant vs. Field, for recent production and R&D dipoles

Scaling of the **stored energy** per unit length of magnet vs. Field, for recent production and R&D dipoles



The electromagnetic loads in a 20 T dipole would be a factor 5 to 8 larger than in the LHC dipoles



# Existing Superconducting Accelerator dipole magnets (1)





Courtesy P. Ferracin, CERN



# Existing Superconducting Accelerator dipole magnets (2)





GdR

7-Febr-2017, Magnets,

e-Bogis,

Chavannee-d

Basic

AS

# **Type II Superconductors**

Below a the critical surface the material is "superconducting". Above the surface it is "normal conducting"

- $\Theta_c$  Critical Temperature (at zero field and current density)
- *B<sub>c2</sub>* Critical Field (at zero temperature and current density)
- $J_c$  Critical Current Density ( at zero temperature and field)

The Critical surface depends on the material type Nb-Ti, Nb<sub>3</sub>Sn, etc) and the processing

Superconducting means: R = 0

J: few x 10<sup>3</sup> A/mm<sup>2</sup> inside the superconductor



Courtesy L. Bottura

Quantized fluxoids

in a superconductor



40



#### Superconductivity

Typical operational conditions (0.85 mm diameter strand)





## **Available Superconductors**



CAS



# Superconducting strands: Nb-Ti

• Nb-Ti is the workhorse for present accelerators, medical magnets, cyclotrons, etc

#### **Strands and Cables for LHC Dipole Magnets**



Performance specification		
STRAND	Type 01	Type 02
Diameter (mm)	1.065	0.825
Cu/NbTi ratio	$1.6-1.7 \pm 0.03$	$1.9-2.0 \pm 0.03$
Filament diameter (µm)	7	6
Number of filaments	8800	6425
$Jc (A/mm^2) @1.9 K$	1530 @ 10 T	2100 @ 7 T
$\mu_0 M(mT) @1.9 K, 0.5 T$	30 ±4.5	23 ±4.5
CABLE	Type 01	Type 02
Number of strands	28	36
Width (mm)	15.1	15.1
Mid-thickness (mm)	1.900 ±0.006	1.480 ±0.006
Keystone angle (degrees)	$1.25 \pm 0.05$	$0.90 \pm 0.05$
Cable Ic (A) @ 1.9 K	13750 @ 10T	12960 @ 7T
Interstrand resistance $(\mu \Omega)$	10-50	20-80





#### Cable compaction ~ 91 %





## Multifilament wires Fabrication of Nb<sub>3</sub>Sn multifilament wires

- Since Nb<sub>3</sub>Sn is brittle, it cannot be extruded and drawn like Nb-Ti.
  - The process requires several steps:
    - Assembly multifilament billets from  $Nb_3Sn$  precursor: Nb and Sn
    - Fabrication of the wire through extrusion-drawing
    - Fabrication of the cable
    - Fabrication of the coil
    - "reaction": the Cu, Sn and Nb are heated to 600-700 C and the Sn diffuses in Nb and reacts to form Nb<sub>3</sub>Sn











# Superconducting strands and tapes: BSCO



- Available in strands (OST)
- Can reach 400 A/mm<sup>2</sup> (overall)
- Is fragile under stress and strain
- Powder in a silver tube
- Has to be reacted at 850° C with a temperature oprecision of 1° C in an oxygen atmosphere
  - Can be cabled in high current Rutherford cables Nexa





OST wire 0.8 mm using Nexans precursor





Difficult technology but could be promising for high field magnets in >20 T region  $_{45}$ 



# Superconducting tapes: YBCO

YBCO: Yttrium barium copper oxide

- Available in tapes : YBCO deposited on a substrate to impose the texture (1-2 μm)
- Can reach > 600 A/mm<sup>2</sup> (overall)
- Is strong under axial stress and strain
- Limited cabling possibilities:







 Difficult technology but could be promising for high field magnets in >20 T region.





YBCO SUPERPOWER Record field (25 T), adding 3 T NHMFL - Florida

Potted racetrack coils



# Superconducting cables for magnets

#### We need multi-strand cables

- Superconducting accelerators are ramped up in time spans 100 s to 1000 s
- Coils are designed for voltages to ground of around 1000 V
- With the number of turns and the current the inductance is to be limited to keep the voltage below 1000 V
- Dipoles and Current:
  - Tevatron B = 4.4 T ; I ~ 4000 A
  - Hera B = 5 T ;  $I \sim 6000 A$
  - LHC B = 8.3 T ; I ~ 12000 A
- For magnets 10 T < B < 15 T the current has to be 10kA < I < 15 kA
- For stability reasons strands are
  0.6 mm < strand diameter < 1 mm</li>
- With a Cu-nonCu ratio (stability) around 1 and a Jc ~ 1000 A/mm<sup>2</sup>
  - ➔ a 1 mm diameter strand can carry ~400 A
  - ➔ so we need a 30 strand cable to get up to 12 kA



## **Rutherford cables**

- Compact cables giving high overall current density
- Easy rectangular geometry for convenient winding









Basics, Chavannee-de-Bogis, 7-Febr-2017, Magnets, GdR

CAS

# How to get high fields in accelerator dipole and quadrupole magnets ?

From Ampere's law one can derive the field resulting from the current in a line conductor and integrate this over the surface of a coil

- Dipole 60° sector coil [see ref 10, 14]
  - The field is proportional to the current density j
  - The field is proportional to coil width
  - The field is independent of aperture

$$B_{1} = -4\frac{jm_{0}}{2\rho} \overset{\rho/3}{\overset{\rho}{}_{0}} \overset{r}{\overset{\rho}{}_{0}} \frac{\cos q}{r} r dr dq = -\frac{\sqrt{3}m_{0}}{\rho} jw$$

- with: r: inner radius coil
  - w : coil width
  - $\boldsymbol{\rho}$  : radial coordinate
  - J : current density
- Quadrupole 30<sup>o</sup> sector coil [see ref 11, 14]
  - The gradient is proportional to the current density j
  - The gradient depends on w/r

#### → by having very high current density close to the beam pipe

See: E. Todesco et al. ref[10] and indirectly : N. Wilson ref[1], K-H Mess et al. ref[2] For a in depth study of magnetic field calculations: S. Russenschuck ref[4]



Cross-section of a dipole based on 60° sector coils



Cross-section of a quadrupole based on 30° sector coils



## **Electro-magnetic forces**

The e.m. forces in a dipole magnet tend to push the coil

- Towards the mid plane in the vertical-azimuthal direction
  - $(F_{y}, F_{\theta} < 0)$
- Outwards in the radial-horizontal direction ( $F_x$ ,  $F_r > 0$ )



The e.m. forces in a quadrupole magnet tend to push the coil

- Towards the mid plane in the vertical-azimuthal direction  $(F_{y}, F_{\theta} < 0)$
- Outwards in the radial-horizontal direction ( $F_x$ ,  $F_r > 0$ )





# **Conductor stability and AC behaviour**

- Pure massive superconductor is not stable as they (Nb-Ti, Nb<sub>3</sub>Sn) are poor normal conductors
  - To 'cryogenically stabilize' the conductor one surrounds it in Cu:
    - good electrical conductivity
    - good heat transfer to the He
- During current ramping the filaments will magnetize
  - → make them thinner
- Filaments will have magnetic coupling twist the strand





Courtesy M. Wilson

Chavannee-de-Bogis, 7-Febr-2017, Magnets, GdR Basics, CAS



# Quench: a thermal runaway effect

Due to perturbations locally the conductor can get  $T > T_c (J_l, B_l)$ 

A thermal runaway can then occur, called a **Quench** 

With stored energies > MJ the coils can overheat if nothing is done (T = 3000K is possible !) What to do ?

- Detect the quench : SC: R=0 → V=0, quench V>0 (typically 100mV threshold)
- Switch power convertor off
- Heat up the whole coil with quench heaters
- Dump energy of the circuit into a dump resistor





# Practical accelerator magnet design: Dipoles

Two types of coils are in use for high field magnets:

Cos(@) coil and Block coil

- Cos( $\Theta$ ) coil (the traditional solution)
  - Allows a very good field quality ( $b_n < 1.10^{-4}$ ) in thin coils
    - all (but one) existing accelerators use this type of coil
  - Is very efficient wrt the quantity of superconductor used
  - The EM forces cause a stress buildup at the midplane where also high fields are located
  - Wedges are needed in the straight part ('Keystoned' cable)
  - The ends are short, special geometry for which there is a large experience but not it is easy







saddle' coils make

better field

shapes

Courtesy M. Wilson

simplest winding

uses racetrack

coils

53



# Practical accelerator magnet design: Dipoles

- Block coil (used on development magnets)
  - Used with thick coils the field quality is good
    - Not yet used in accelerators
  - Is less efficient (~10%) wrt to  $cos(\Theta)$  for quantity of superconductor used
  - The EM forces cause a stress buildup at the outside edge of the coil where the fields are lower
  - The straight part is very easy : rectangular cable and wedges (field quality)
  - 'flared ends' look easy but there is little experience exists to make them





#### **Pre-stress**

- Why pre-stress ?
  - Field quality is determined by the cable positioning (be precise to ~0.02 mm)
  - Under the MN forces the coils will move
    - $\rightarrow$  Apply pre-stress to fix the positioning
  - Very small amounts of heat can quench the coil: limit the movement (avoid stick-slip effects on ~10  $\mu$ m movements)
    - →Apply pre-stress to fix the positioning
- How to put pre-stress ?

Three methods:

- 1. Compress at room temperature: collar system
- 2. Use room temperature pre-stress plus differential shrinkage at cool-down: Al or stainless steel shrinking cylinder and/or a (shrinking) key
- 3. Compress a bit at room temperature and use differential shrinkage at cooldown: Al shrinking cylinder + bladder and key system
- Order of magnitudes: LHC @ 8.34 T: 70 MPa warm, 30 MPa cold

Fresca2 @ 13 T: 60 MPa warm, 130 MPa cold



#### **Prestress: collars**

"The classical solution"

- Thin collars put around the coil
- The coil is well contained in a fixed cavity
- Pressed together and locked with pins or keys
- At 300K apply a prestress 2-3 times of what is needed as part of the stress is lost during cooldown: for very high field tends to be too high (LHC:70 MPa at 300 K and 40 MPa at cold)
- Field quality is in good part determined by collar shape
- If the coils size is not so well controlled, the stress can be too high or too low
  - Nb<sub>3</sub>Sn is stress sensitive and this could be a problem







LHC dipole CERN



#### TQC quadrupole LARP-FNAL



# Prestress: Al shrinking cylinder + bladder and keys

Developed at LBNL, example: TQS a LARP model quadrupole

300 K: Bladders pressurized with water (<600 bar) , then insert keys  $\rightarrow$  load between 10 MPa and 80 MPa

Cooldown: differential shrinkage between AL shell and Fe yoke load another ~100 MPa

Needs careful mechanical FE modeling before and strain measurements during bladder operations and cooldown







# Looking in the kitchen of future SC magnet development

#### What is happening after the 8T magnets for LHC? At CERN

- 1) Upgrade the LHC luminosity: HL-LHC (HILUMI)
  - use large aperture  $Nb_3Sn$  triplet quadrupoles (12T class)
  - improve collimation: use a few 11T dipoles to make space
- 2. Go to higher energies
  - 16 T Nb<sub>3</sub>Sn dipoles in the LHC ring for  $E_{com}$ =26 TeV : HE-LHC
  - 16 T Nb<sub>3</sub>Sn dipoles in a 100 km new ring for  $E_{com}$ =100 TeV : FCC (Future Circular Collider)

But even !

- 20 T HTS hybrid dipoles in the LHC ring: for  $E_{com}$ =33 TeV : HE-LHC
- 20 T HTS hybrid dipoles in a 80 km new ring for  $E_{com}$ =100 TeV : FCC

#### In China

A similar completely new project is being studied in China: SPPC (C=54 km, B=20 T,  $E_{com}$ =71 TeV)

For these, basic High Field Magnet development programs are since many years running in the US and Europe and recently in China



### **CERN-European development evolution**





# Basic magnet technology development for HILUMI and beyond (2004-2013) ; US development evolution





### **Basic HFM development :** Some achievements at LBNL (1995-2004)

Since 20 years LBNL is running a high field dipole development program Some achievements:

- D20, 50 mm aperture, cosQ 4 layer dipole, reached 13.5 T@1.9K
- HD1, flat block coil, 8 mm aperture, reached 16 T
- HD2, flared end block coil, 36 mm aperture, reached 13.8 T

These pose a clear breakthrough above 10 T with a new coil layout (block coil) and a mechanical structure aimed (shell-bladder and keys) at high fields



Fig. 2. HD2 cross-section.





# Basic HFM development : EuCARD high field dipole (Fresca2):

- Fresca2 : CERN, CEA construction phase
- First tests 2014





- 156 turns per pole
- Iron post
- $B_{center} = 13.0 T$
- I<sub>13T</sub> = 10.7 kA
- B<sub>peak</sub> = 13.2 T
- $E_{mag} = 3.6 \text{ MJ/m}$
- L = 47mH/m





- Diameter Aperture = 100 mm
- L coils = 1.5 m
- L straight section = 700 mm
- L yoke = 1.6 m
- Diameter magnet = 1.03 m



Courtesy Attilio Milanese, Pierre Manil



# **Fabrication of Fresca2 coils**

Straightforward technology to wind block coils with flared ends:

This is a lesson for FCC magnets !











#### HILUMI IT magnet zoo







# HL-LHC: MQXF low beta Nb<sub>3</sub>Sn quadrupole

Spring 2016 the first model achieved the nominal and ultimate field at FNAL !



By courtesy of G. Ambrosio (FNAL), P. Ferracin (CERN et al)

A CERN LARP collaboration. Nominal Gradient 132.6 T/m Aperture diameter 150 mm Peak Field 12.1 T Current 17.5 A Loadline Margin 20% @ 1.9 K Stored Energy 1.32 MJ/m











# **HL-LHC: 11 T Dispersion suppressor magnet**









- First Nb3Sn magnet to go into an accelerator (2019) !
- Present model program (CERN and FNAL)
  - demonstrated the required performance (11.25 T at 11850 A) and Achieved accelerator field quality

Nominal Field 11 T Aperture diameter 60 mm Peak Field 11.35 T Current 11.85 kA Loadline Margin 19.7% @ 1.9 K Stored Energy 0.96 MJ/m







# **FCC development**





#### FCC: 16T dipole options





S. Farinon, P. Fabbricatore (INFN)



CAS Basics, Chavannee-de-Bogis, 7-Febr-2017, Magnets, GdR

# HTS: First attempt towards 20 T



6 T HTS (YBCO) insert for test in FReSCa2 To get to 19 T But without bore

CEA + CRNS Grenoble





# **EuCARD2 5T accelerator quality ReBCO magnet**

5 Tesla stand alone, (18 T in 13 T background), @ 4.5K, 40 mm aperture, 10 kA class cable, Accelerator Field quality





# EuCARD2 5T : Feather0 - Feather-M2.0

- Feather0: First coil in the test station
- Feather2: winding of first coil with dummy cable in progress











### Literature on warm accelerator magnets

- Books
  - G.E.Fisher, "Iron Dominated Magnets" AIP Conf. Proc., 1987 -- Volume 153, pp. 1120-1227
  - J. Tanabe, "Iron Dominated Electromagnets", World Scientific, ISBN 978-981-256-381-1, May 2005
  - P. Campbell, Permanent Magnet Materials and their Application, ISBN-13: 978-0521566889
  - S. Russenschuck, Field computation for accelerator magnets : analytical and numerical methods for electromagnetic design and optimization / Weinheim : Wiley, 2010. - 757 p.
  - Schools
    - CAS Bruges, 2009, specialized course on magnets, 2009, CERN-2010-004
    - CAS Frascati 2008, Magnets (Warm) by D. Einfeld
    - CAS Varna 2010, Magnets (Warm) by D. Tommasini
  - Papers and reports
    - D. Tommasini, "Practical definitions and formulae for magnets," CERN, Tech. Rep. EDMS 1162401, 2011


## Literature on Superconducting Magnets

- Books
- 1) M. Wilson, Superconducting magnets / Oxford : Clarendon Press, 1983 (Repr. 2002). 335 p
- 2) K-H. Mess, P. Schmüser, S. Wolff, Superconducting Accelerator Magnets, Singapore, World Scientific, 1996. 218 p.
- Y. Iwasa, Case studies in superconducting magnets : design and operational issues . 2nd ed. Berlin : Springer, 2009. -682 p.
- 4) S. Russenschuck, Field computation for accelerator magnets : analytical and numerical methods for electromagnetic design and optimization / Weinheim : Wiley, 2010. 757 p.
- 5) CERN Accelerator school, Magnets, Bruges, Belgium 16 25 June 2009, Editor: D. Brandt, CERN–2010–004

- Conference proceedings and reports
- 21st International Conference on Magnet Technology, Hefei, China, 18 23 Oct 2009, IEEE Trans. Appl. Supercond. 20 (2010)
- 7) The 2010 Applied Superconductivity Conference, Washington DC, US, 1-6 Aug 2010, , IEEE Trans. Appl. Supercond. 21 (2011)



## Literature on Superconducting Magnets (2)

- Papers and reports
- 8) S. Caspi, P. Ferracin, "Limits of Nb3Sn accelerator magnets", *Particle Accelerator Conference* (2005) 107-11.
- 9) S. Caspi, P. Ferracin, S. Gourlay, "Graded high field Nb3Sn dipole magnets", *19th Magnet Technology Conference, IEEE Trans. Appl. Supercond.*, (2006) in press.
- 10) E. Todesco, L. Rossi, "Electromagnetic Design of Superconducting Dipoles Based on Sector Coils", Phys. Rev. Spec. Top. Accel. Beams 10 (2007) 112401
- 11) E. Todesco, L. Rossi, AN ESTIMATE OF THE MAXIMUM GRADIENTS IN SUPERCONDUCTING QUADRUPOLES, CERN/AT 2007-11(MCS),
- 12) P. Fessia, et al., Parametric analysis of forces and stresses in superconducting dipoles, IEEE, trans. Appl, Supercond. Vol 19, no3, June 2009.
- 13) P. Fessia, et al., Parametric analysis of forces and stresses in superconducting quadrupole sector windings, sLHC Project Report 0003

## Websites

15) http://www.magnet.fsu.edu/magnettechnology/research/asc/plots.html

٠



## Acknowledgements

For this lecture I used material from lectures, seminars, reports, etc. from the many colleagues. Special thanks goes to:

Davide Tommasini, Attilio Milanese, Antoine Dael, Stephan Russenschuck, Thomas Zickler, Giorgio Ambrosio (FNAL), Luca Bottura (CERN), Shlomo Caspi (LBNL), Arnaud Devred (ITER), Paolo Ferracin (LBNL), Attilio Milanese (CERN), Jeroen van Nugteren, Juan-Carlos Perez (CERN), Lucio Rossi (CERN), Ezio Todesco (CERN), Martin Wilson

And to the people who taught me, years ago, all the fine details about magnets !



www.cern.ch