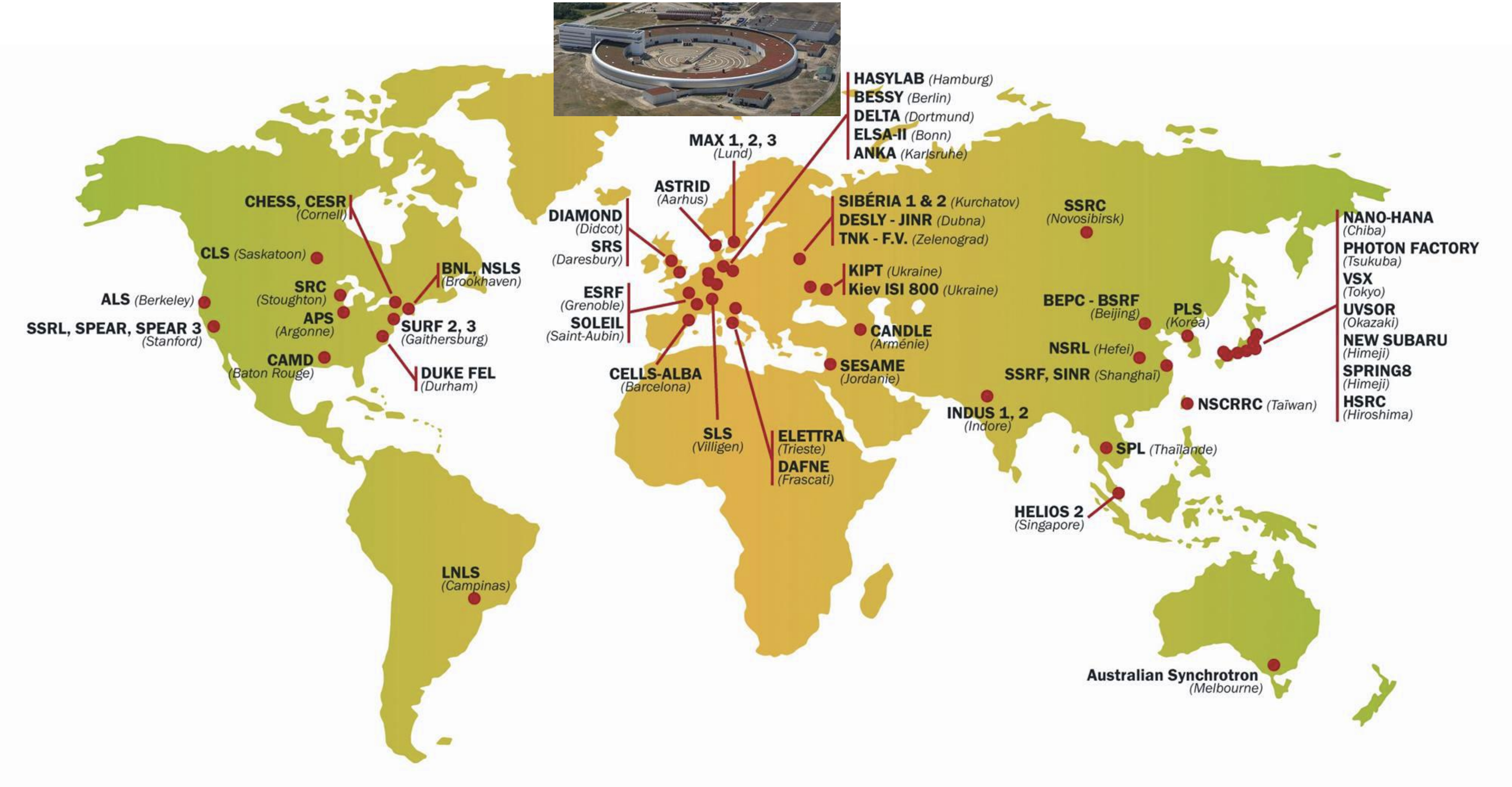
Light Sources
based on Storage Rings

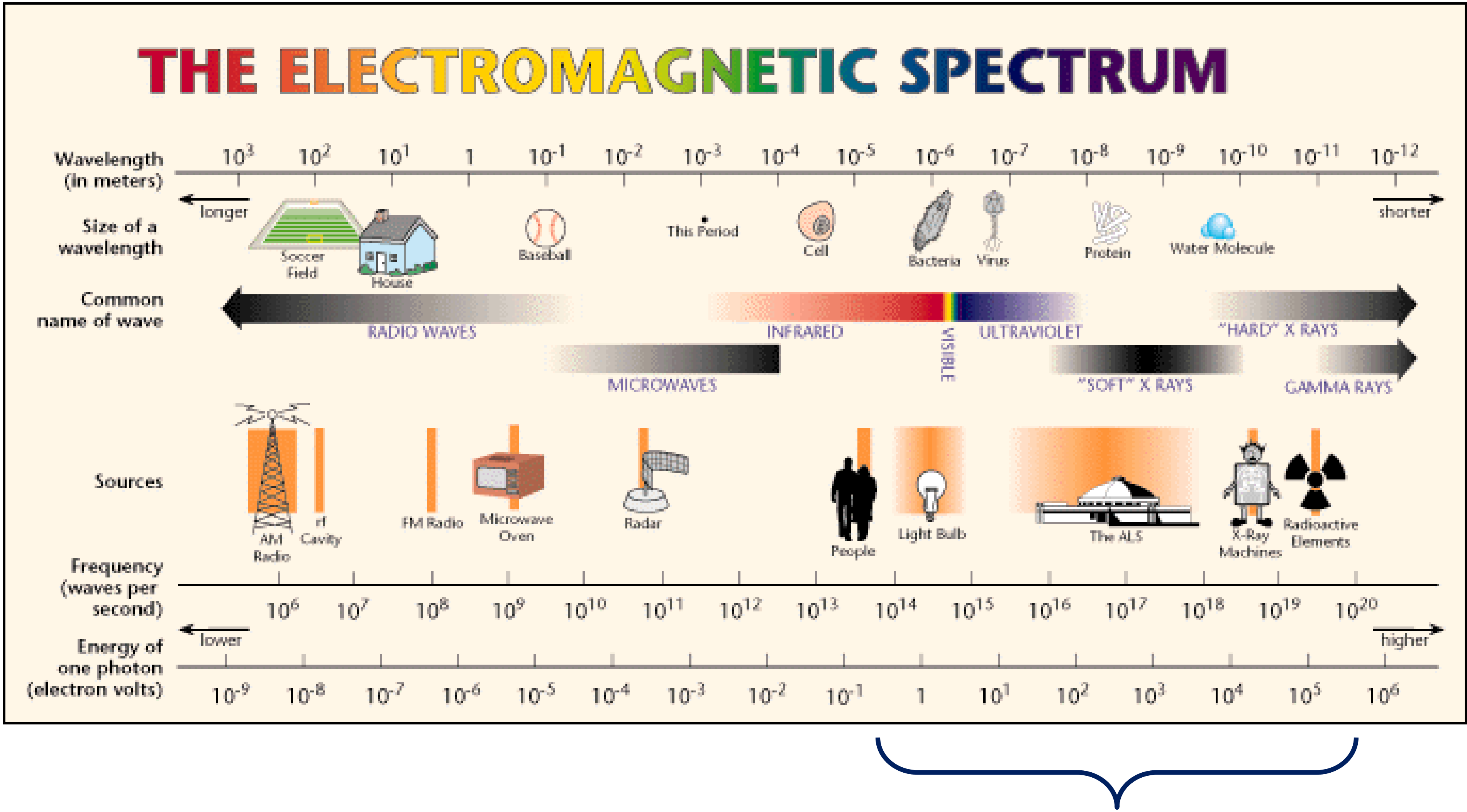
Lenny Rivkin

Paul Scherrer Institute (PSI)
and
Swiss Federal Institute of Technology Lausanne (EPFL)

Light sources: > 50 producing synchrotron light

60‘000 users world–wide

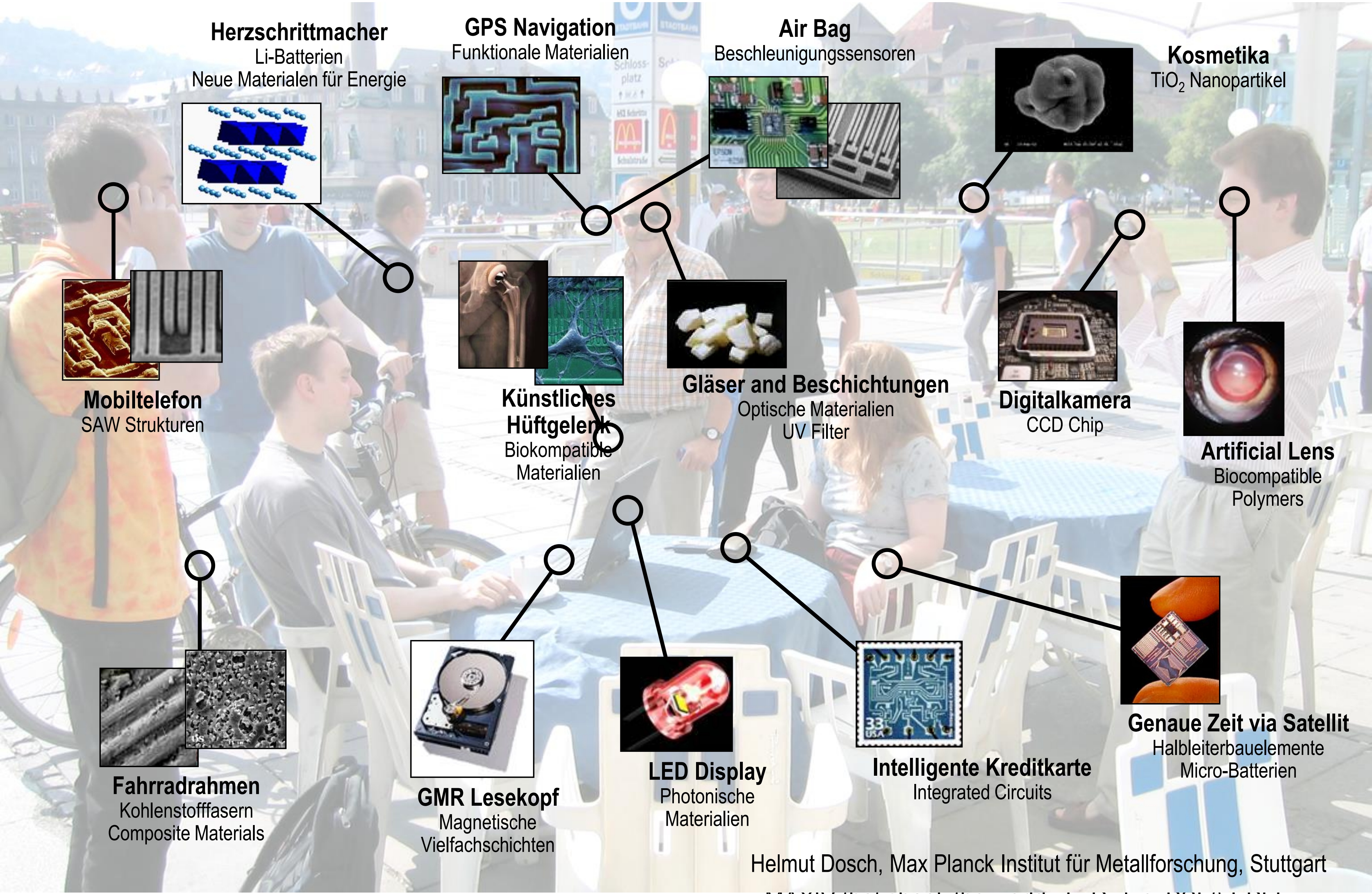




Wavelength continuously tunable !

Materials – key to our technologies

Materials – key to our technologies



Herzschrittmacher
Li-Batterien
Neue Materialen für Energie

GPS Navigation
Funktionale Materialien

Air Bag
Beschleunigungssensoren

Kosmetika
TiO_2 Nanopartikel

Mobiltelefon
SAW Strukturen

Künstliches Hüftgelenk
Biokompatible Materialien

Gläser and Beschichtungen
Optische Materialien
UV Filter

Digitalkamera
CCD Chip

Artificial Lens
Biocompatible Polymers

Fahrradrahmen
Kohlenstofffasern
Composite Materials

GMR Lesekopf
Magnetische Vielfachschichten

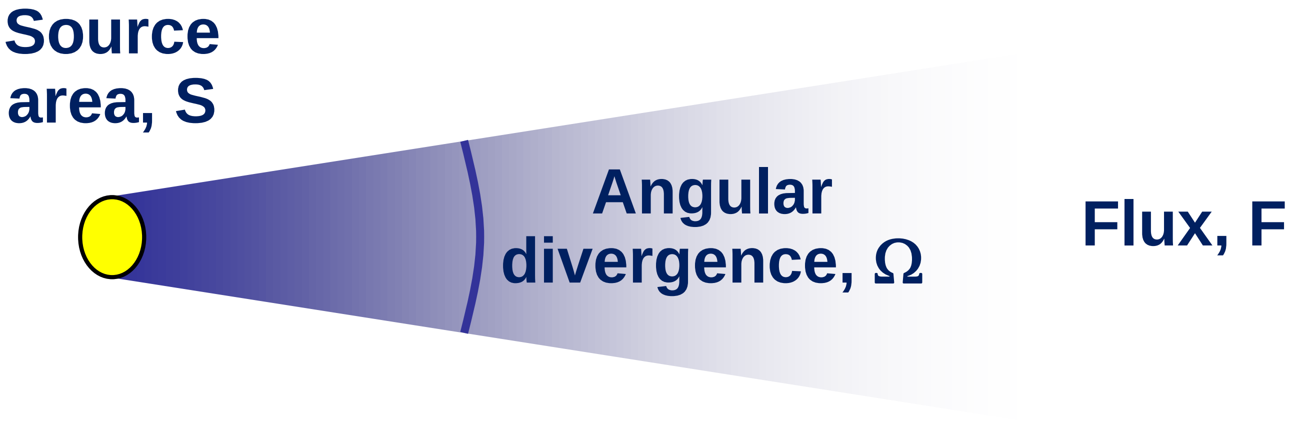
LED Display
Photonische Materialien

Intelligente Kreditkarte
Integrated Circuits

Genaue Zeit via Satellit
Halbleiterbauelemente
Micro-Batterien

Helmut Dosch, Max Planck Institut für Metallforschung, Stuttgart

The “brightness” of a light source:



Source area, S

Angular divergence, Ω

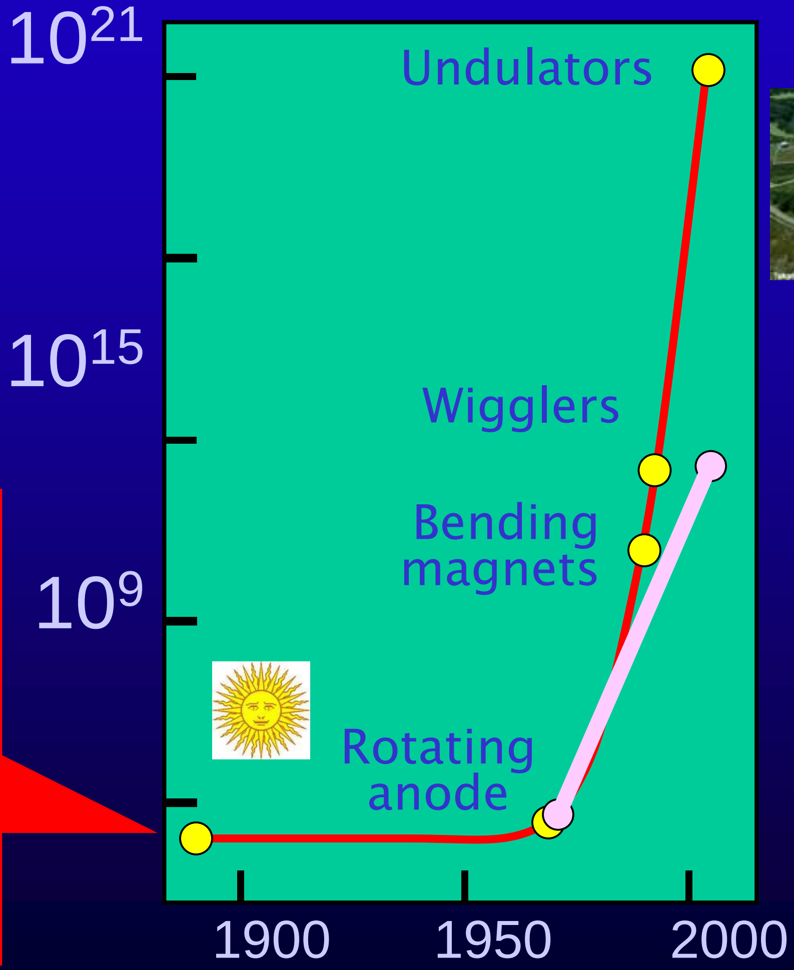
Flux, F

$$\text{Brightness} = \text{constant} \times \frac{F}{S \times \Omega}$$

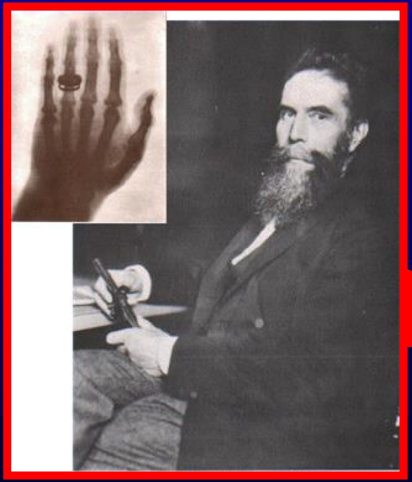
Steep rise in brightness

the second wave

XFEL



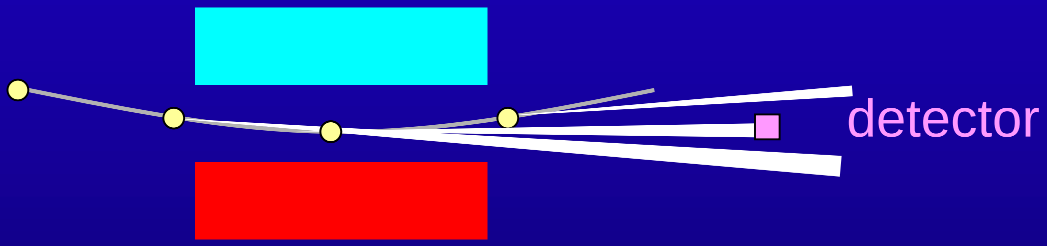
Moore’s Law for semiconductors

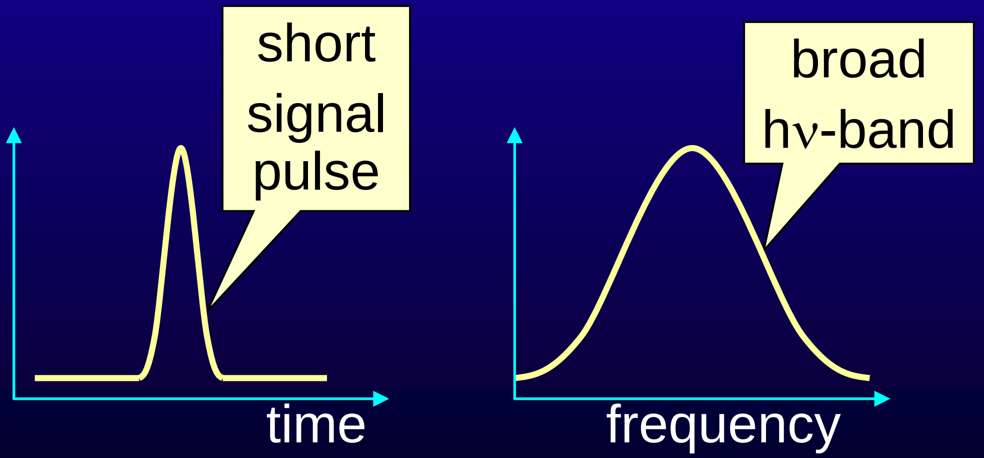


Bertha Roentgen’s hand
(exposure: 20 min)

3 types of storage ring sources:

1. Bending magnets: $\mathbf{B \sim N_e}$

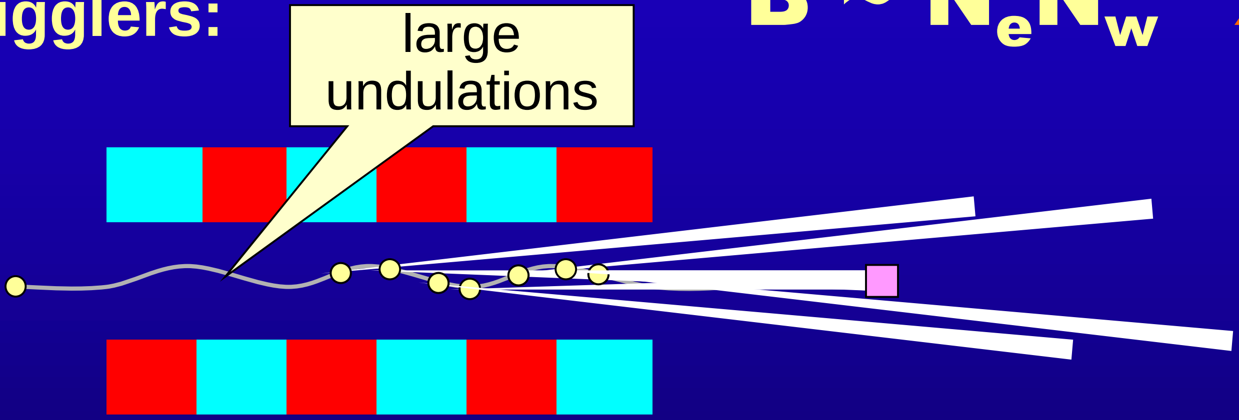




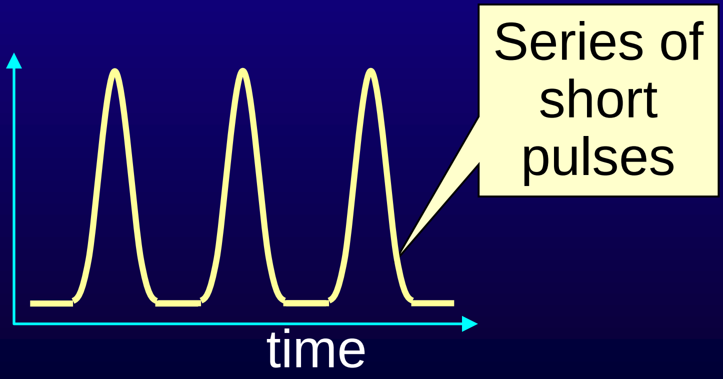
3 types of storage ring sources:

2. Wigglers:

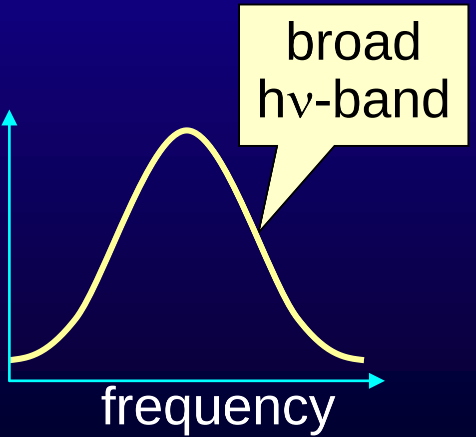
$B \sim N_e N_w$ x10



large undulations



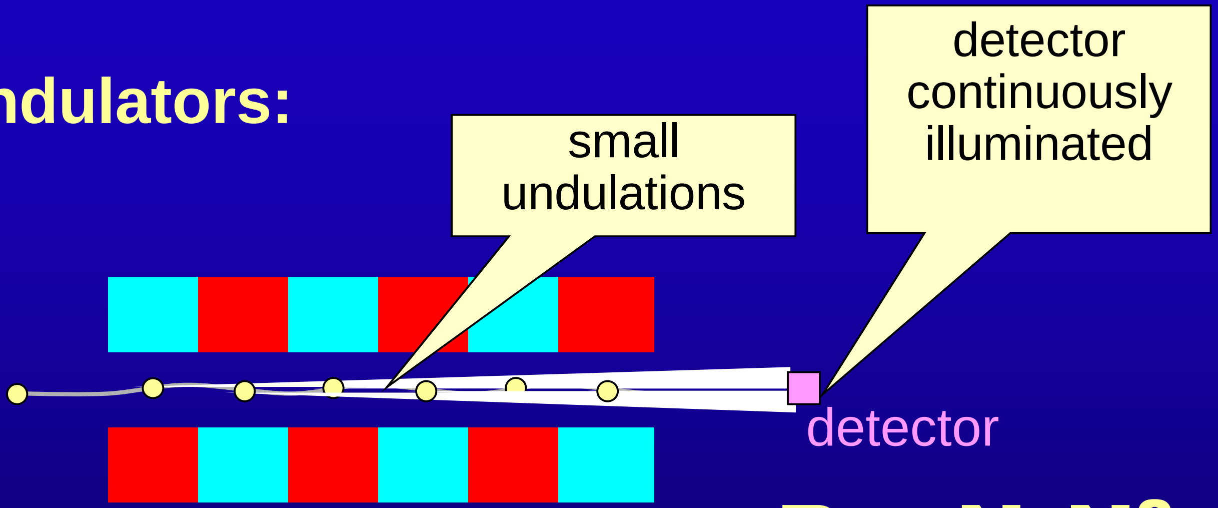
Series of short pulses



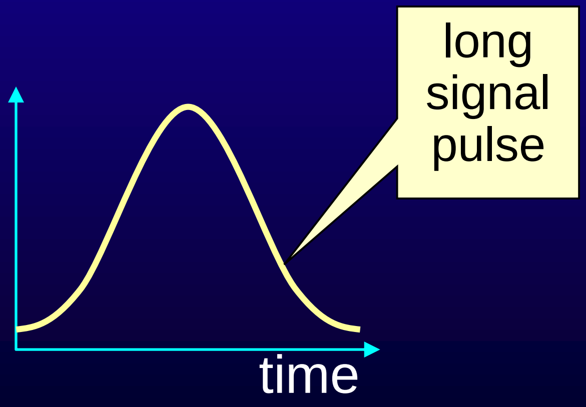
broad hν-band

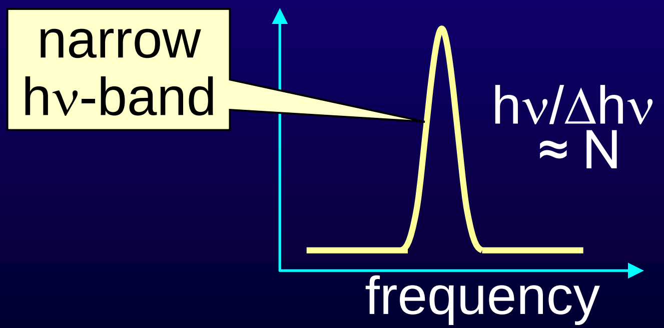
3 types of storage ring sources:

3. Undulators:

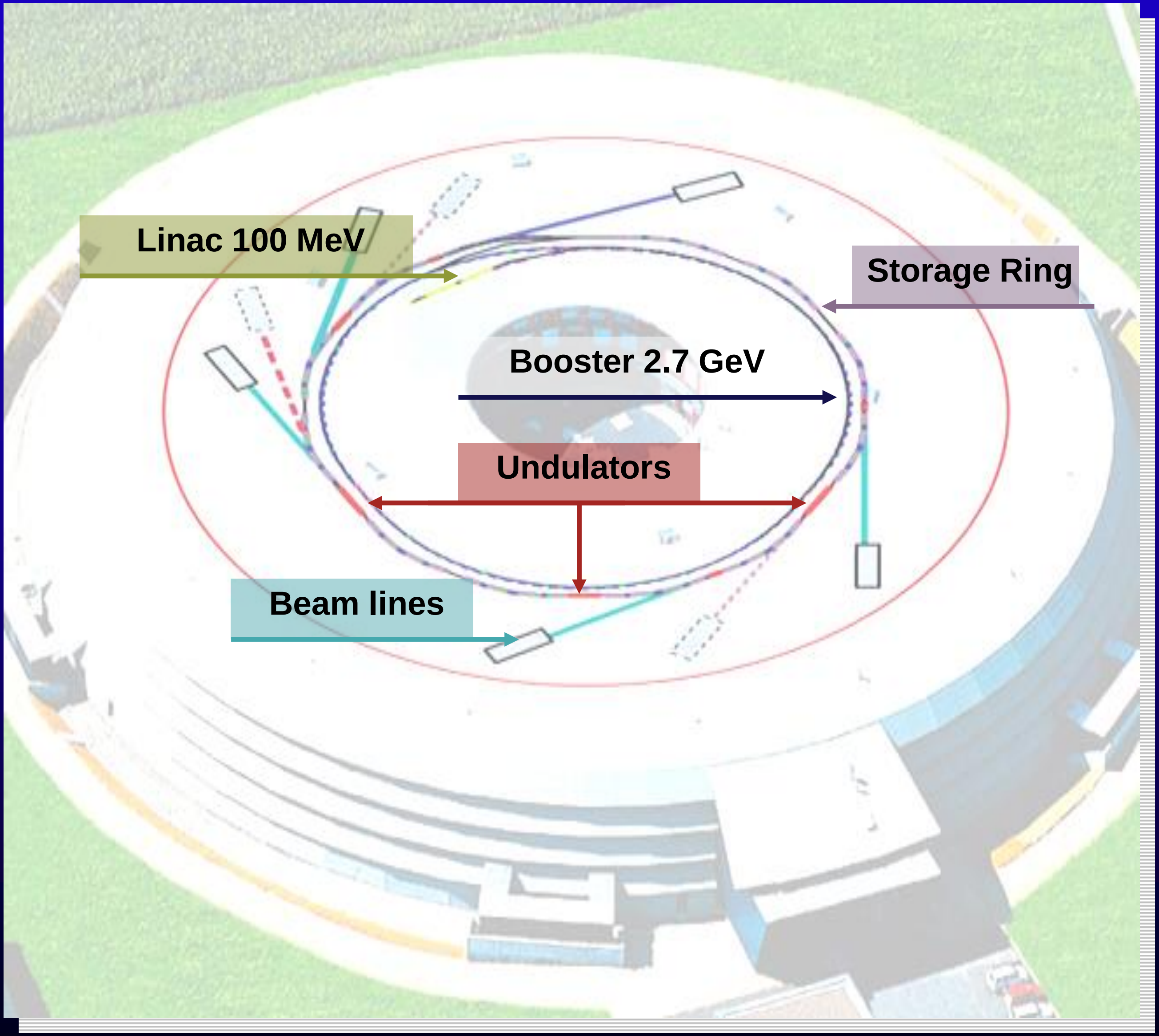


$B \sim N_e N^2_u \quad x10^3$

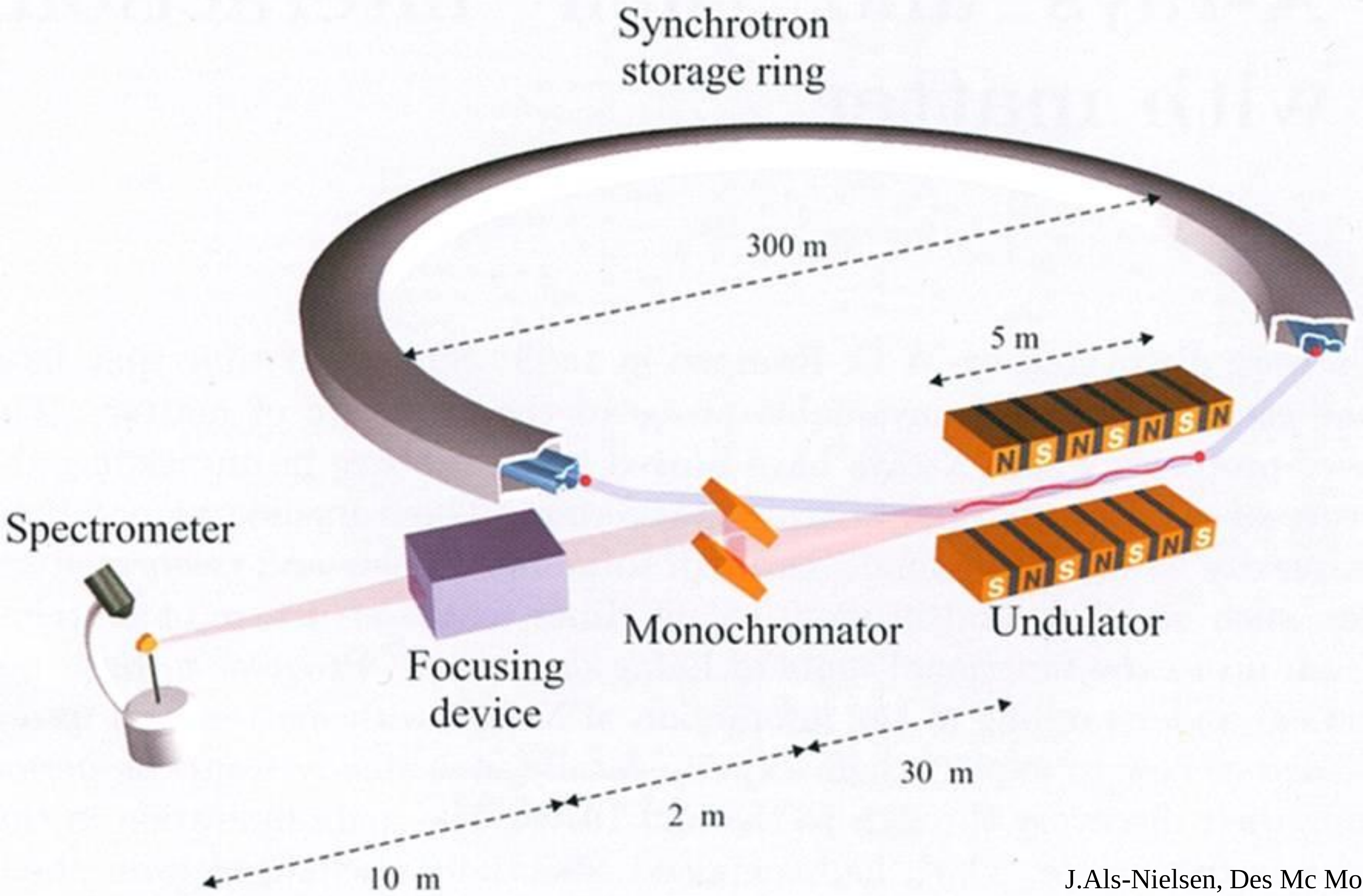




Anatomy of a light source



Undulator based beamline



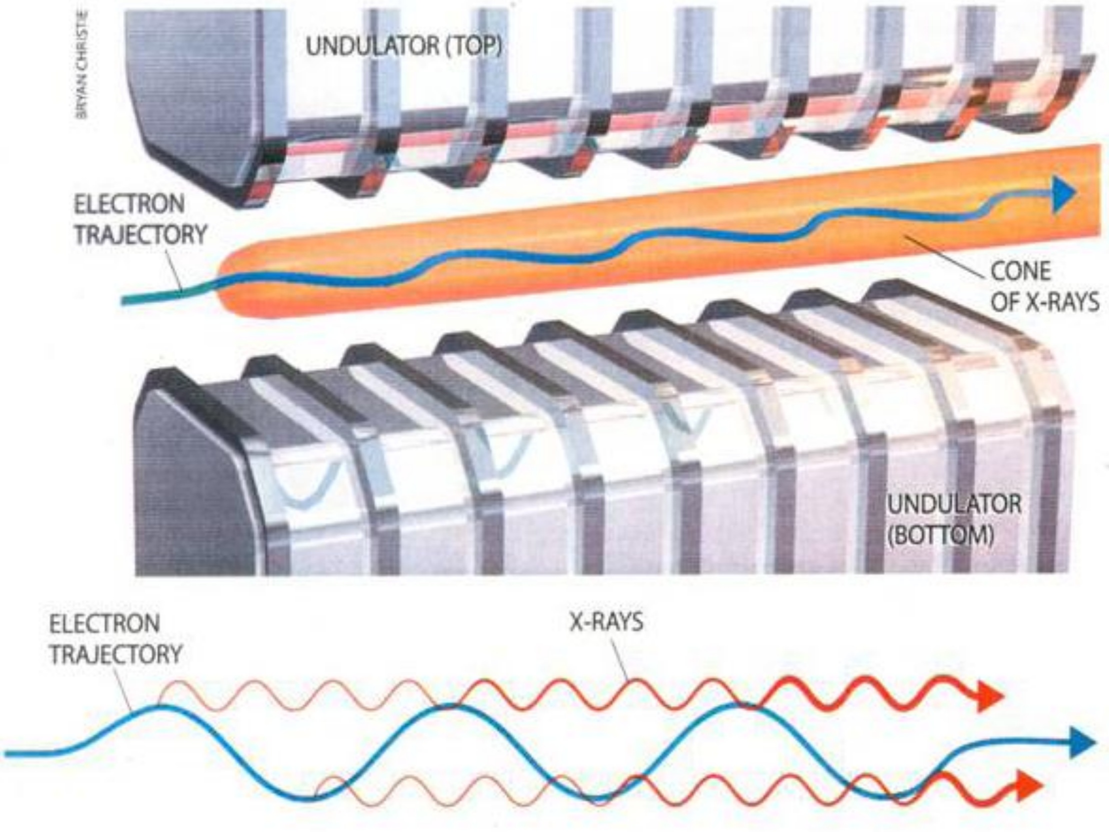
J.Als-Nielsen, Des Mc Morrow

Bright beams of particles: phase space density

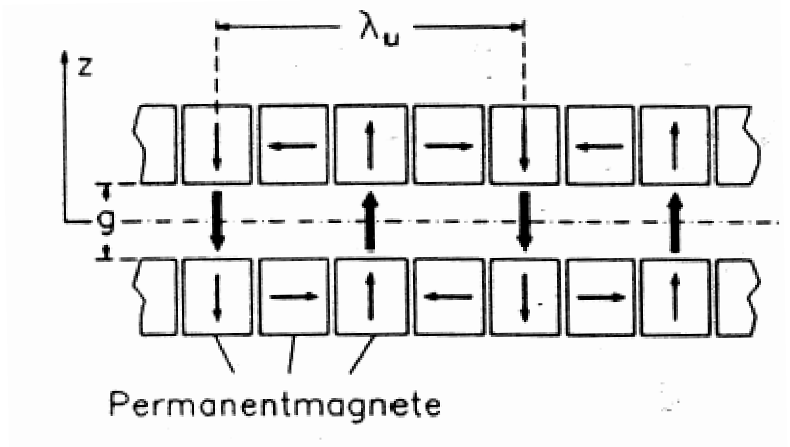
Incoherent, spontaneous emission of light:

Large phase space

Coherent, stimulated emission of light



Permanent magnet undulators

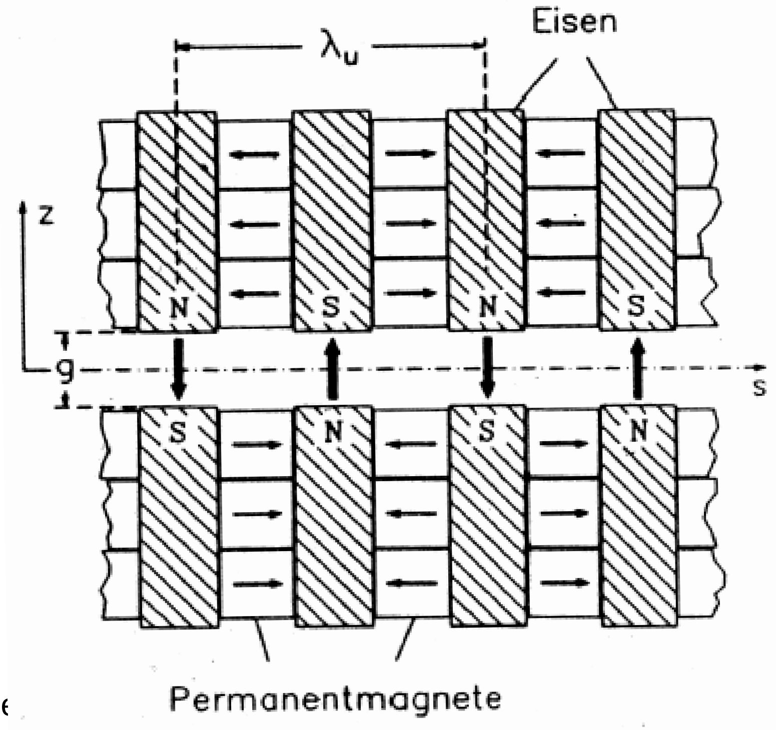


Permanent magnet materials: $SmCo_5$, NdFeB

e.g. a pencil made of such material corresponds to 15‘000 A-turns!

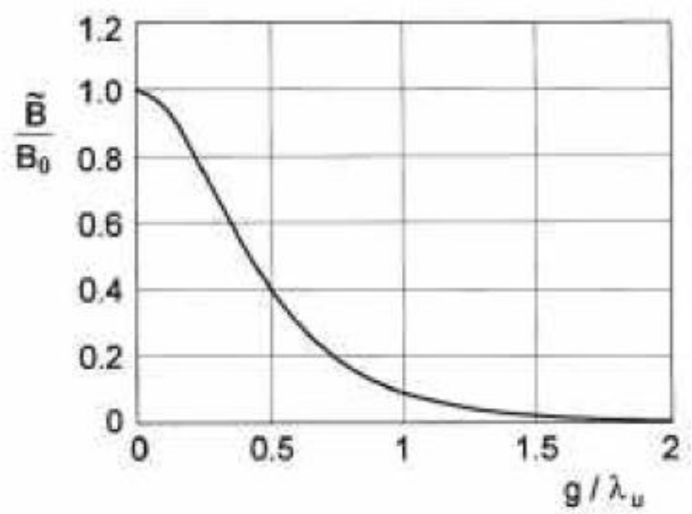
Hybrid undulator:

permanent magnets and iron



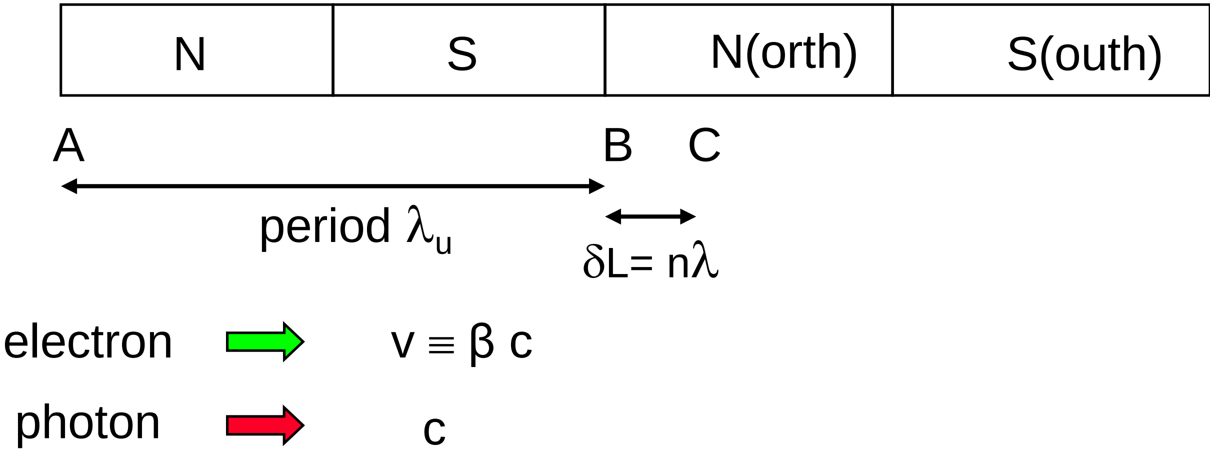
Field tuning with gap

$$B \approx 1.8 \cdot B_r \cdot e^{-\pi \cdot \frac{gap}{\lambda_u}}$$



Permanent magnet material	Remanent field [T]
$SmCo_5$	0.9 – 1.0
Sm_2Co_{17}	1.0 – 1.1
NdFeB	1.0 – 1.4

Selection of wavelength in an undulator II



The path difference

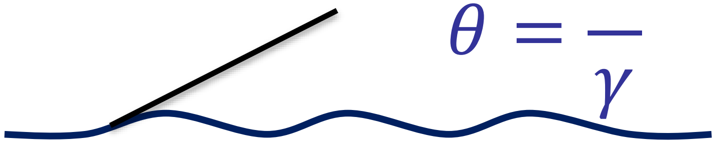
$$\boxed{\delta L \equiv n\lambda \approx (1-\beta)\,\lambda_u\,, \quad 1-\beta \approx \frac{1}{2\gamma^2}}$$

$$\boxed{\lambda = \frac{\lambda_u}{2n\gamma^2}\left(1+\frac{K^2}{2}\right)}$$

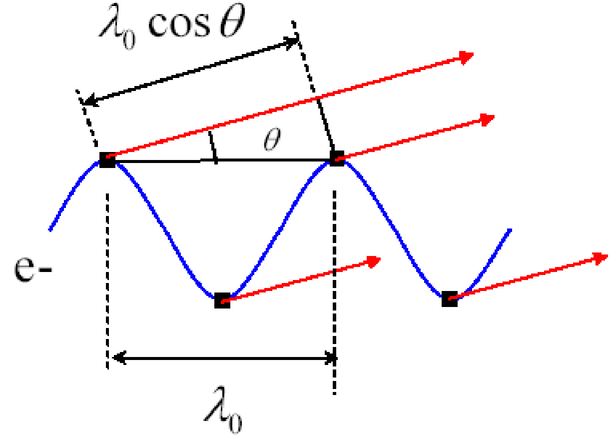
detour through slalom

$$\boxed{K = 0.0934\cdot\lambda_u[mm]\cdot B[T]}$$

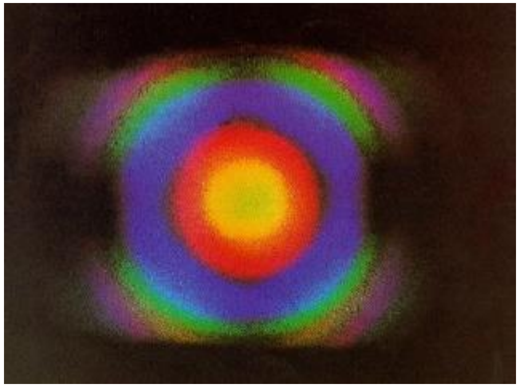
$$\theta = \frac{K}{\gamma}$$

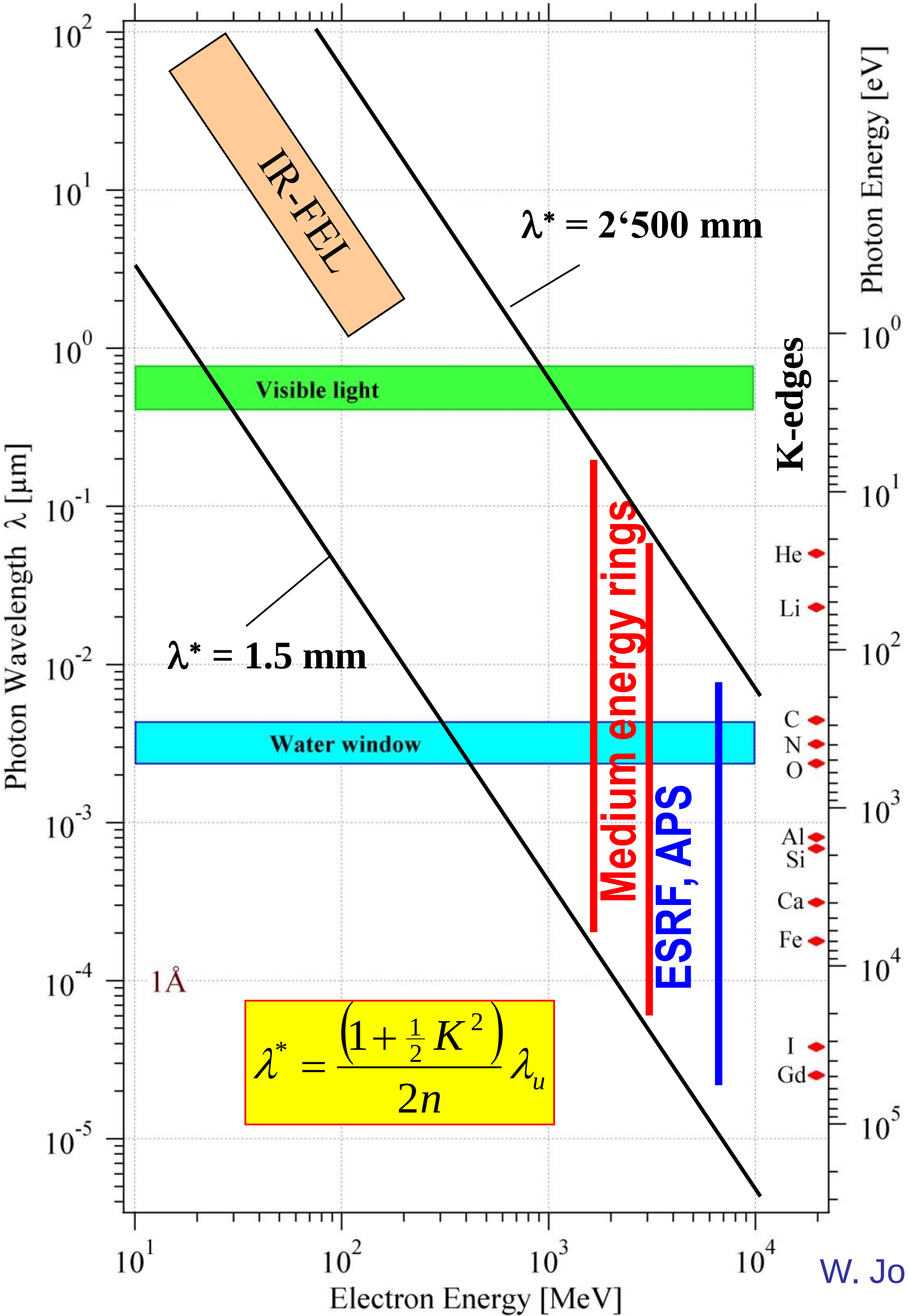


Undulator radiation



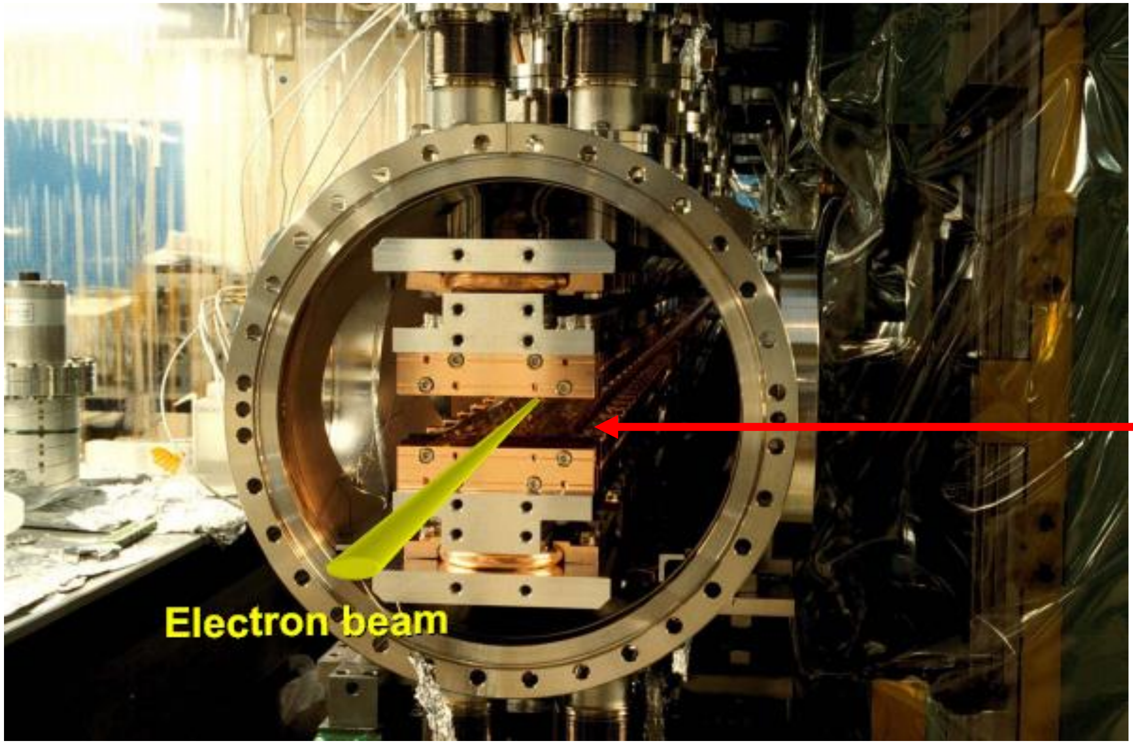
$$\boxed{\lambda = \frac{\lambda_u}{2n\gamma^2}\left(1+\frac{K^2}{2}+\gamma^2\theta^2\right)}$$





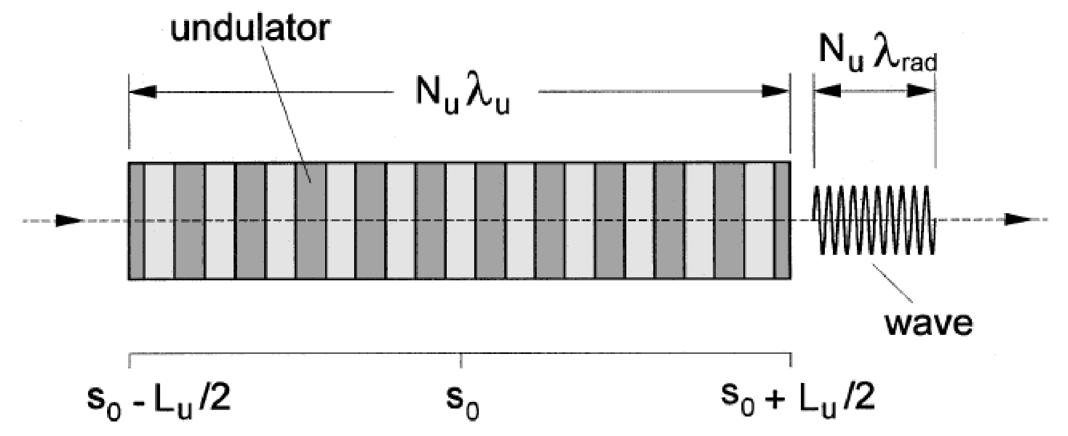
W. Joho

In-vacuum undulators / s.c. undulators



Gaps
down
to
3 mm

Undulator line width



Undulator of infinite length

$$N_u = \infty \quad \Rightarrow \quad \frac{\Delta\lambda}{\lambda} = 0$$

Finite length undulator

- radiation pulse has as many periods as the undulator
- the line width is

$$\frac{\Delta\lambda}{\lambda} \sim \frac{1}{N_u}$$

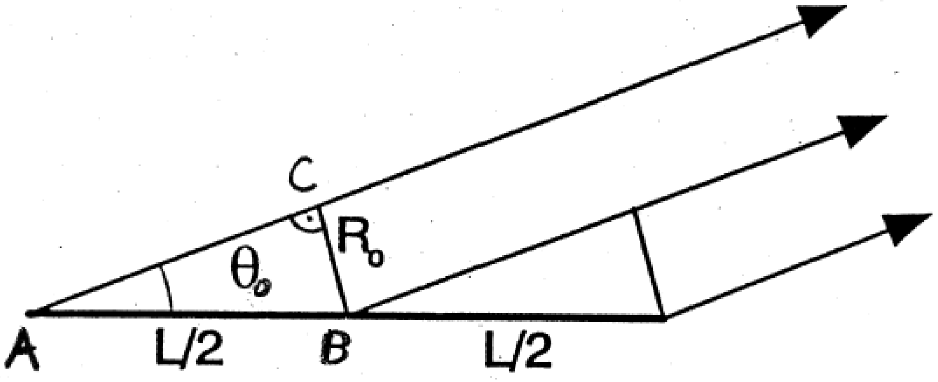
Due to the electron energy spread

$$\frac{\Delta\lambda}{\lambda} = 2\frac{\sigma_E}{E}$$

Radiation cone of an undulator

Undulator radiates from the whole length L into a narrow cone.

Propagation of the wave front BC is suppressed under an angle θ_0,



if the path length AC is just shorter by a half wavelength compared to AB (negative interference). This defines the central cone.

$$\Delta L = AB - AC = \frac{1}{2}L(1-\cos\theta_0) \approx \frac{1}{4}L\theta_0^2$$

Negative interference for $\Delta L = \dfrac{\lambda}{2}$

$$\boxed{\theta_0 = \sqrt{\frac{2\lambda}{L}}} \quad \boxed{R_0 = \sqrt{\frac{\lambda \cdot L}{2}}} \quad \boxed{\varepsilon_0 = \theta_0 R_0 = \lambda}$$

W. Joho

WHAT DO USERS EXPECT FROM A HIGH PERFORMANCE LIGHT SOURCE ?

- **PROPER PHOTON ENERGY FOR THEIR EXPERIMENTS**
- **BRILLIANCE** →
- **STABILITY**

$$B = \frac{\Phi}{(2\pi)^2 \Sigma_x \Sigma_{x'} \Sigma_y \Sigma_{y'}}$$

FIGURE OF MERIT

$$\Sigma^2 = \sigma_e^2 + \sigma_\gamma^2$$

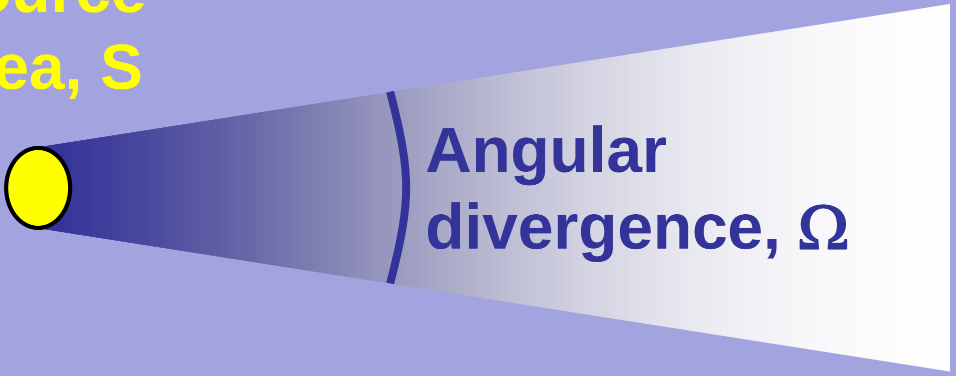
$$\Sigma_x \Sigma_{x'} \approx \sigma_x \sigma_x^{'} \sim \varepsilon_x$$

Photon beam size (U):

$$\sigma_\gamma^{'} = \sqrt{\frac{\lambda}{L}}$$

$$\sigma_\gamma = \frac{\sqrt{\lambda L}}{4\pi}$$

The electron beam “emittance”:



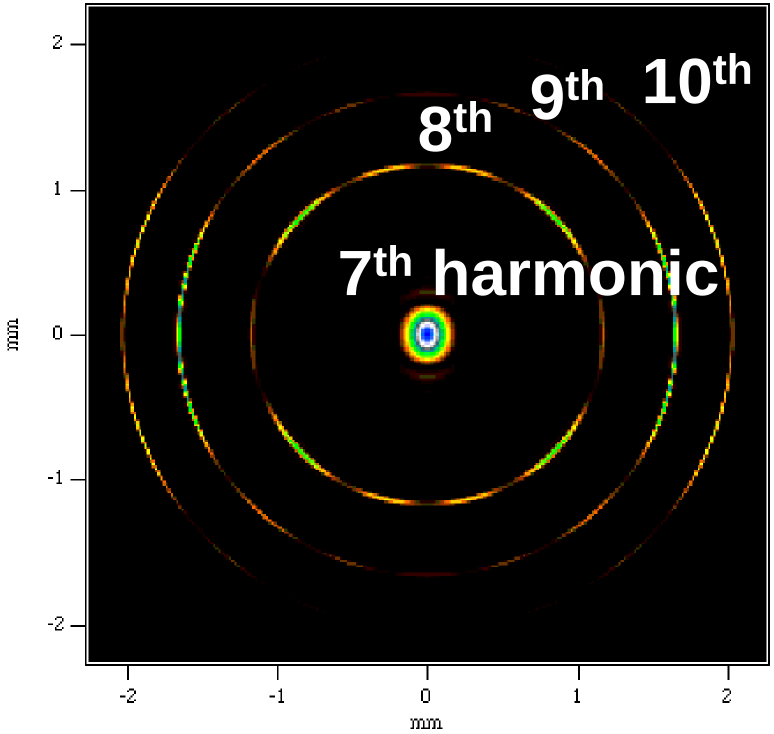
Source area, S

Angular divergence, Ω

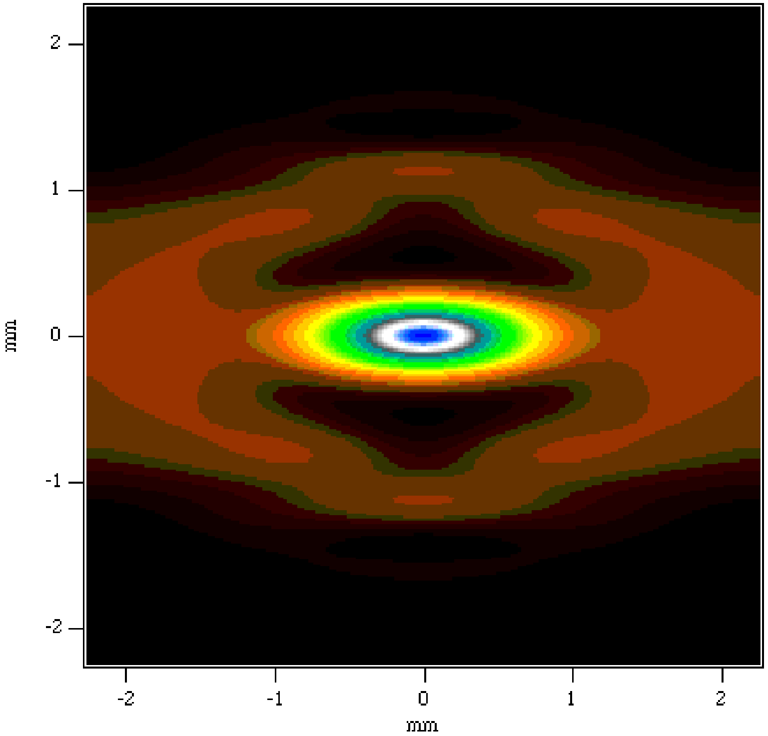
The brightness depends on the geometry of the source, i.e., on the electron beam emittance

Emittance = S x Ω

Undulator radiation from 6 GeV beam with zero emittance, energy spread (example ESRF)



Emittance 4 nm·rad, 1% coupling, finite energy spread



Electron beam phase space

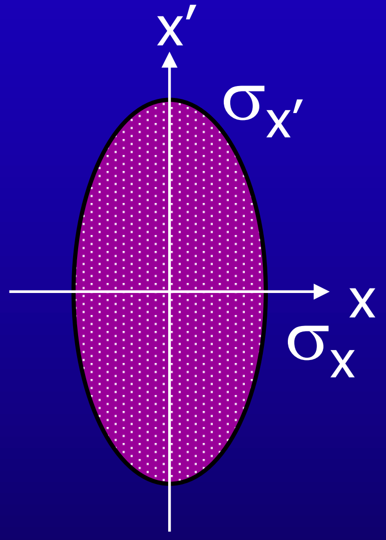
$\text{Area} = \pi \cdot \varepsilon$

Emittance

$\text{Units of } \varepsilon \quad [m \cdot rad]$

Transverse electron beam distribution

- Gaussian
- "Typical" particle: 1 - σ ellipse
  (in a place where $\alpha = \beta' = 0$)



$$\sigma_x = \sqrt{\varepsilon\,\beta}$$
$$\sigma_{x'} = \sqrt{\varepsilon / \beta}$$

$$\varepsilon = \sigma_x \cdot \sigma_{x'}$$

$$\beta = \frac{\sigma_x}{\sigma_{x'}}$$

Radiation effects in electron storage rings

Average radiated power restored by RF

- Electron loses energy each turn to synchrotron radiation
- RF cavities accelerate electrons back to the nominal energy

Radiation damping

- Average rate of energy loss produces **DAMPING** of electron oscillations in all three degrees of freedom (if properly arranged!)

Quantum fluctuations

- Statistical fluctuations in energy loss (from quantized emission of radiation) produce **RANDOM EXCITATION** of these oscillations

Equilibrium distributions

- The balance between the damping and the excitation of the electron oscillations determines the equilibrium distribution of particles in the beam

Small emittance lattices

Equilibrium horizontal emittance

$$\boxed{\varepsilon_{x0} \equiv \frac{\sigma_{x\beta}^2}{\beta} = \frac{C_q E^2}{J_x} \cdot \frac{\langle \mathcal{H} \rangle_{mag}}{\rho}}$$

- one tries to optimize the $\mathcal{H}$ function in **bending magnets**

$$\boxed{\mathcal{H} = \gamma D^2 + 2\alpha D D' + \beta D'^2}$$

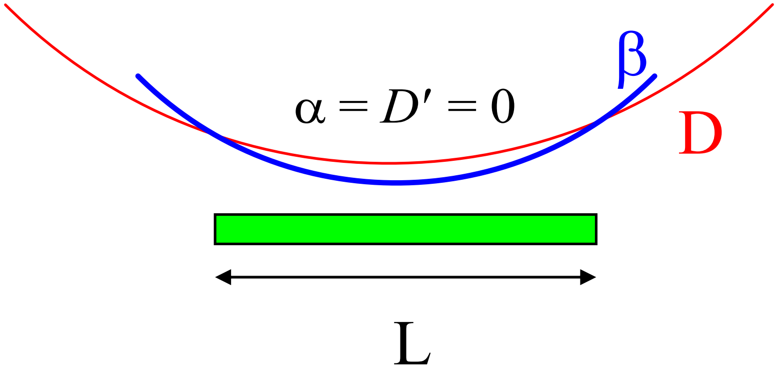
- the equilibrium emittance can be written as:

$$\boxed{\varepsilon_{x0} = \frac{C_q E^2}{J_x} \cdot \theta^3 \cdot \mathbf{F_{latt}}}$$

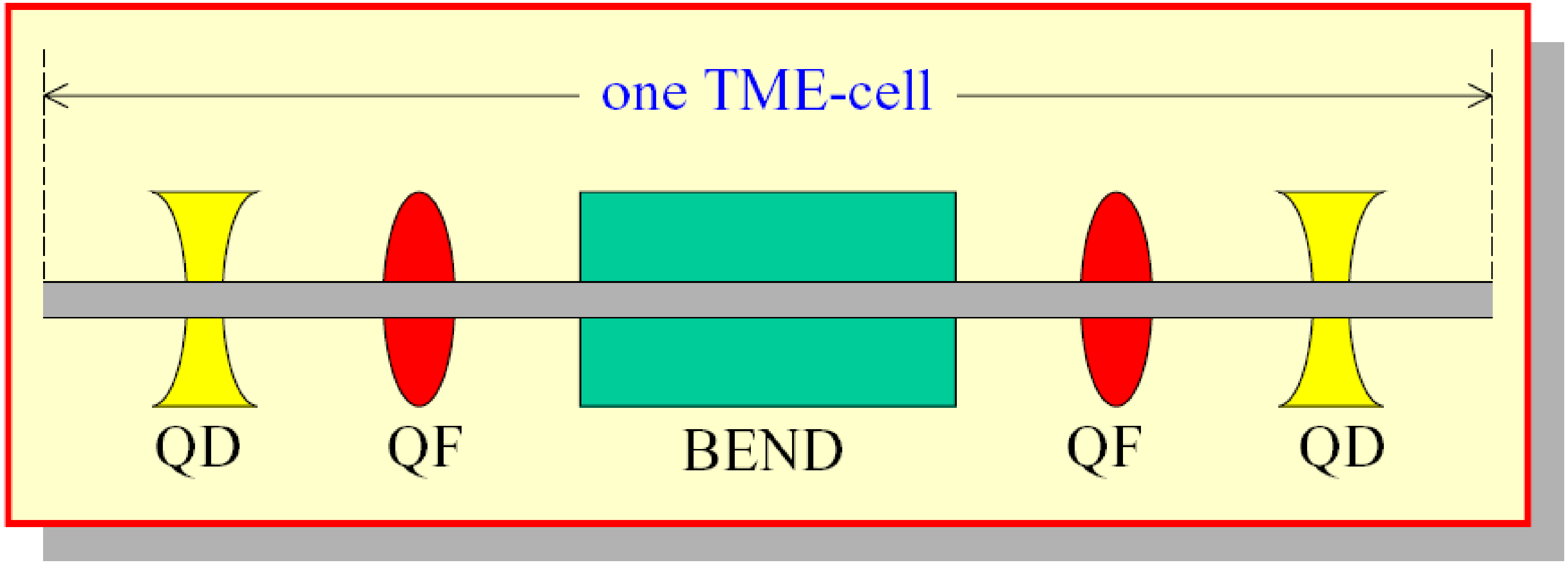
$$\boxed{\mathrm{F}_{\min} = \frac{1}{12\sqrt{15}}}$$

there exists a minimum

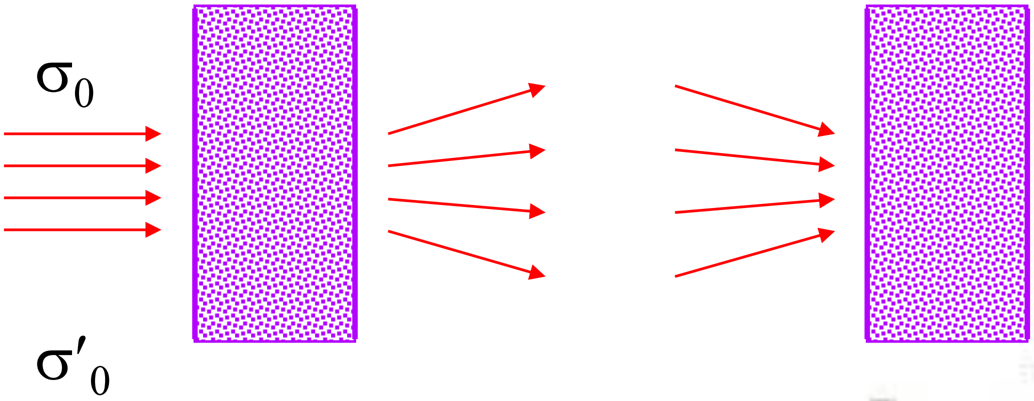
$$\boxed{\beta^* = \frac{L}{2\sqrt{15}}, \quad D^* = \frac{L\theta}{24}}$$



Theoretical minimum emittance



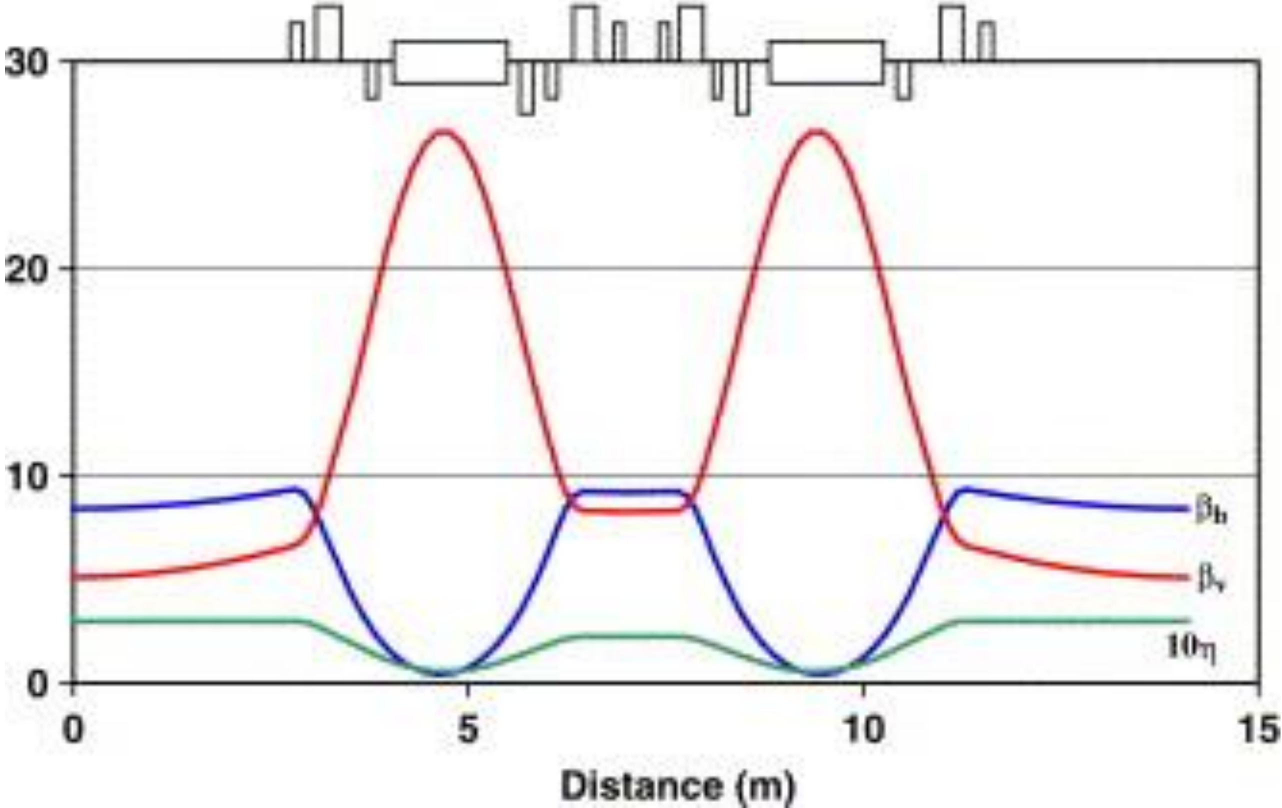
Minimum emittance lattices



$$\varepsilon_{x0} = \frac{C_q E^2}{J_x} \cdot \theta^3 \cdot \mathrm{F}_{\mathrm{latt}}$$

$$\mathrm{F}_{\mathrm{min}} = \frac{1}{12\sqrt{15}}$$

Double Bend TME

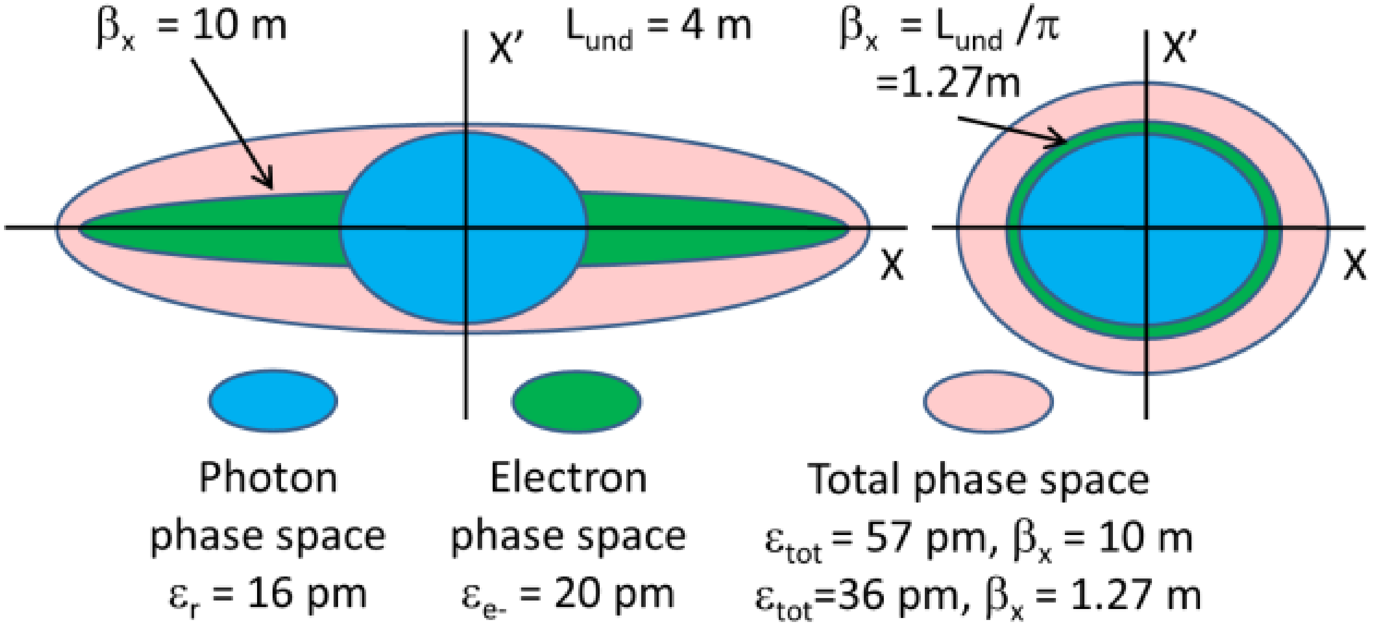


Tight focus in the middle of the bending magnets – need space!

Many bending magnets – need space!

X-ray emittance from electron source: a convolution of electron and photon phase space

$$\text{Brightness} = \frac{\Phi}{(2\pi)^2 \Sigma_x \Sigma_{x'} \Sigma_y \Sigma_{y'}} \qquad \Sigma^2 = \sigma_e^2 + \sigma_\gamma^2$$



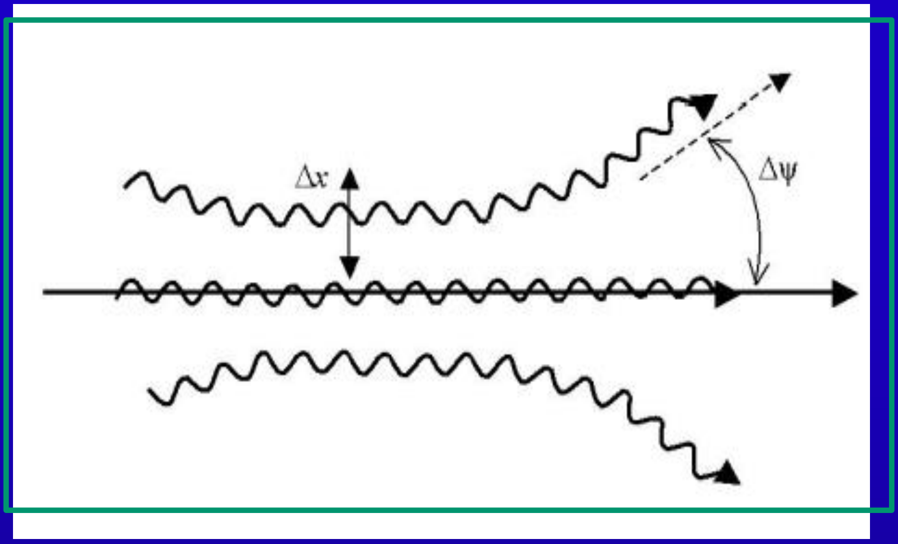
Photon phase space ε_r = 16 pm

Electron phase space ε_{e-} = 20 pm

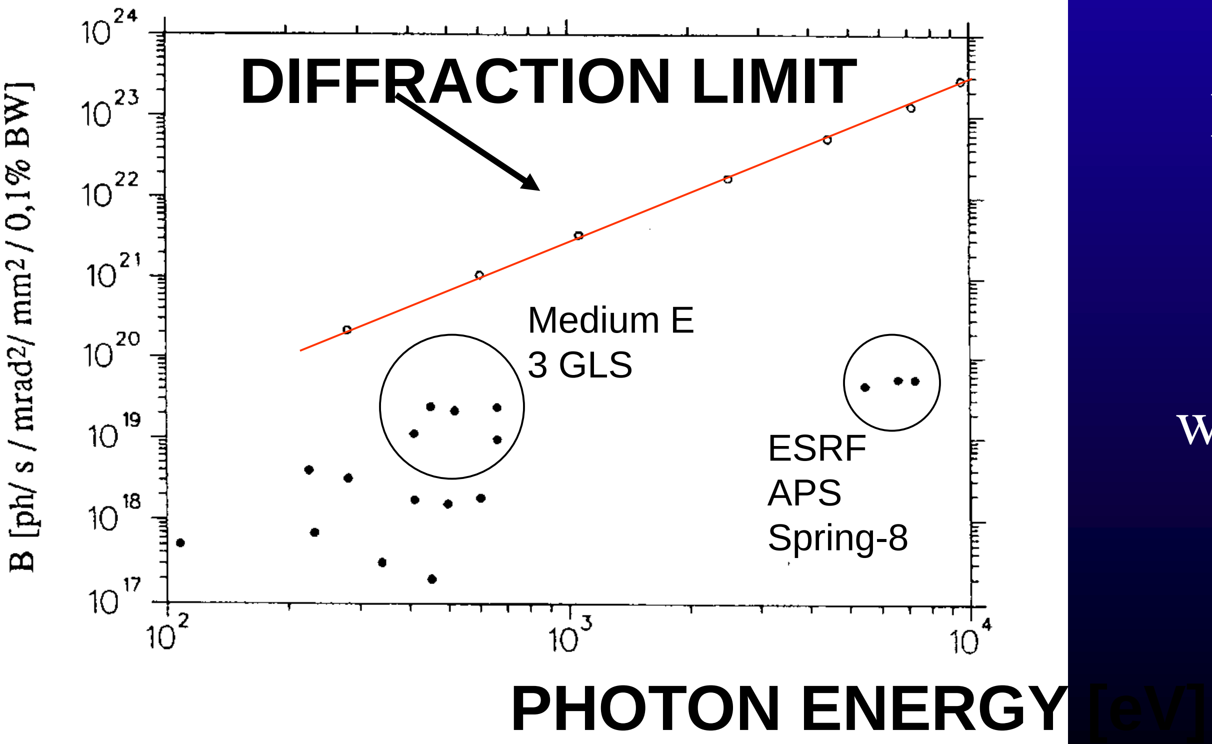
Total phase space
ε_{tot} = 57 pm, β_x = 10 m
ε_{tot}=36 pm, β_x = 1.27 m

R. Hettel

HITTING THE DIFFRACTION LIMIT



BRIGHTNESS:



Light of wavelength

λ

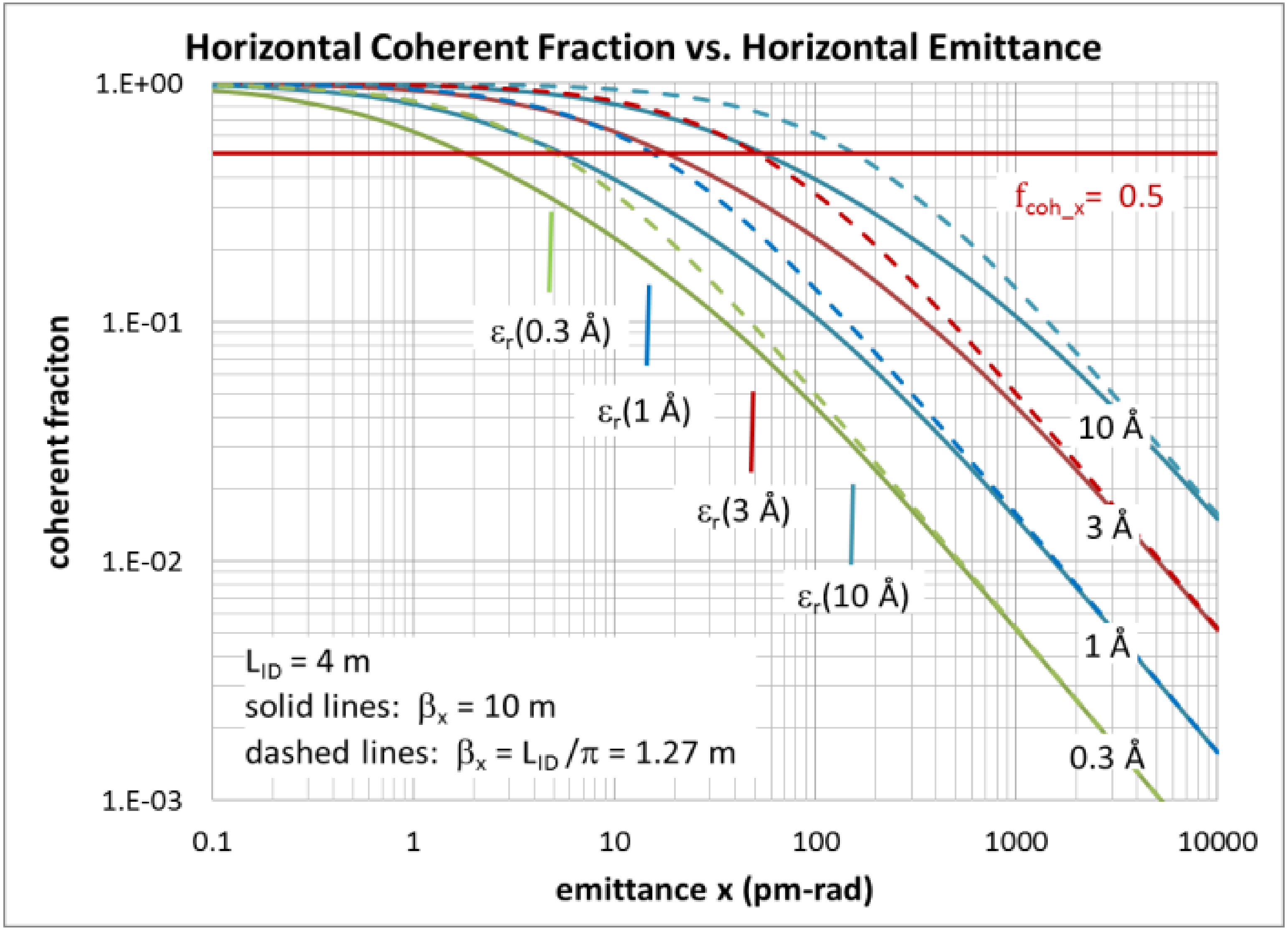
focused to spot size

Δx

will diffract with angle

$\Delta\psi = \sim \lambda/\Delta x$

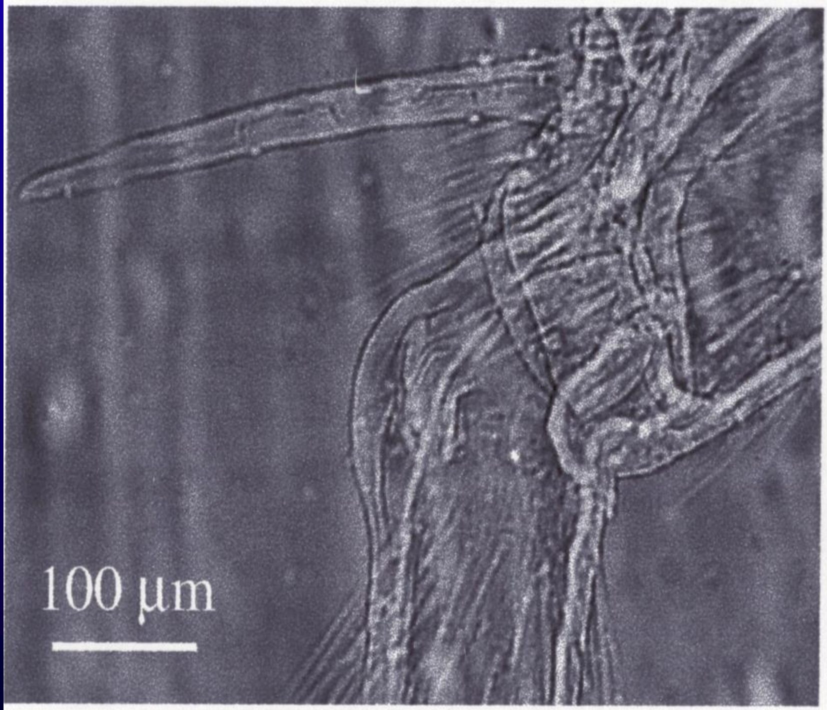
Coherence fraction



Transverse coherence

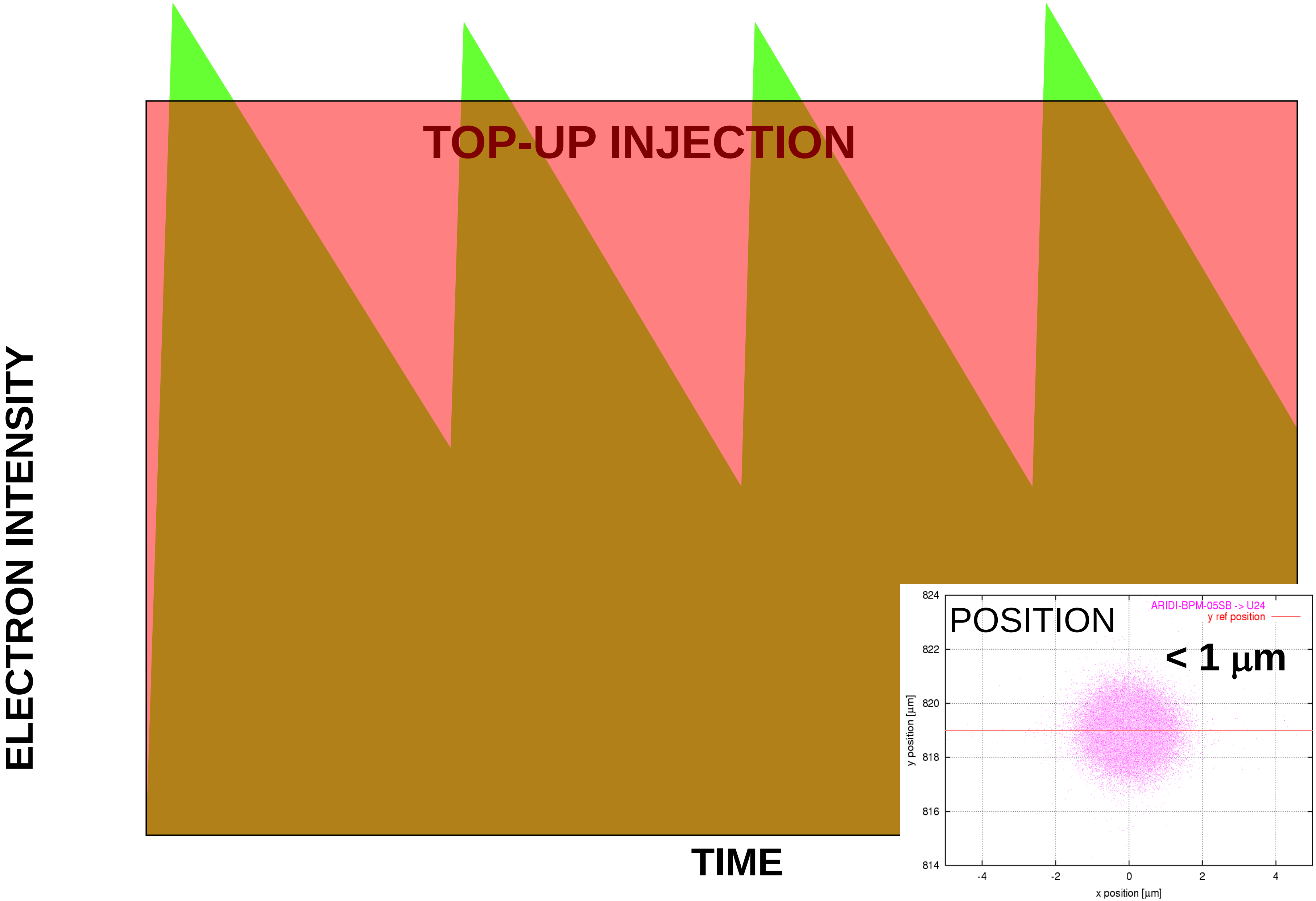
- High brightness gives coherence
- Wave optics methods for X-rays (all chapters in Born & Wolf)
- Holography

The knee of a spider



phase contrast imaging

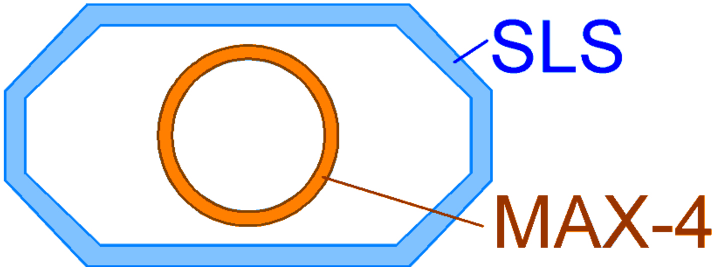
Top-up injection: key to stability

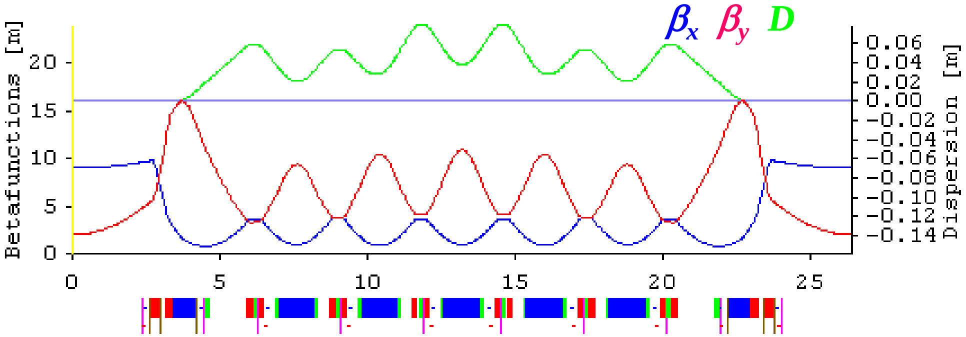


A revolution in storage ring technology

Pioneer work: MAX IV (Lund, Sweden)

Aperture reduction ↔ **Multi–Bend Achromat (MBA)**





Technological achievement:
NEG coating of small vacuum chambers

- ⇨ short lattice cells
- ⇨ many lattice cells
- ⇨ low angle per bend

- ⇨ Small magnet bore
- ⇨ High magnet gradient

emittance $\varepsilon \propto (\text{energy})^2 \times (\text{bend angle})^3$

⇨ Emittance reduction from nm to 10...100 pm range

The MAX IV Laboratory in Lund, Sweden

The world is moving to ever brighter ring sources

2-bend achromat

BNL**: NSLS-II** (2014): 3 GeV,
<1000pm x 8 pm, 500 mA
(**New**)

7- bend achromat

Sweden: **MAX-4** (2016): 3 GeV,
230 pm x 8 pm, 500 mA (**New**)

5- bend achromat

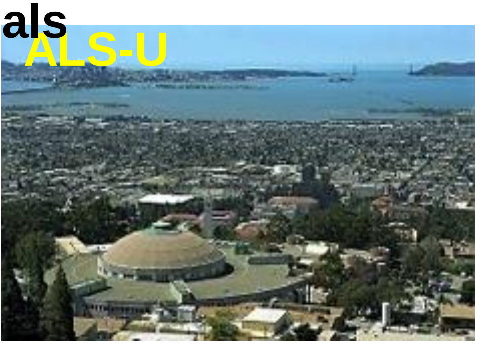
Brazil**: SIRIUS** (2016/17): 3 GeV,
280 pm x 8 pm, 500 mA (**New**)

1st multi-bend achromat ring upgrade

France: **ESRF-II** (2020): 6 GeV,
160 pm x 3 pm, 200 mA (**New**)

U.S. Proposals

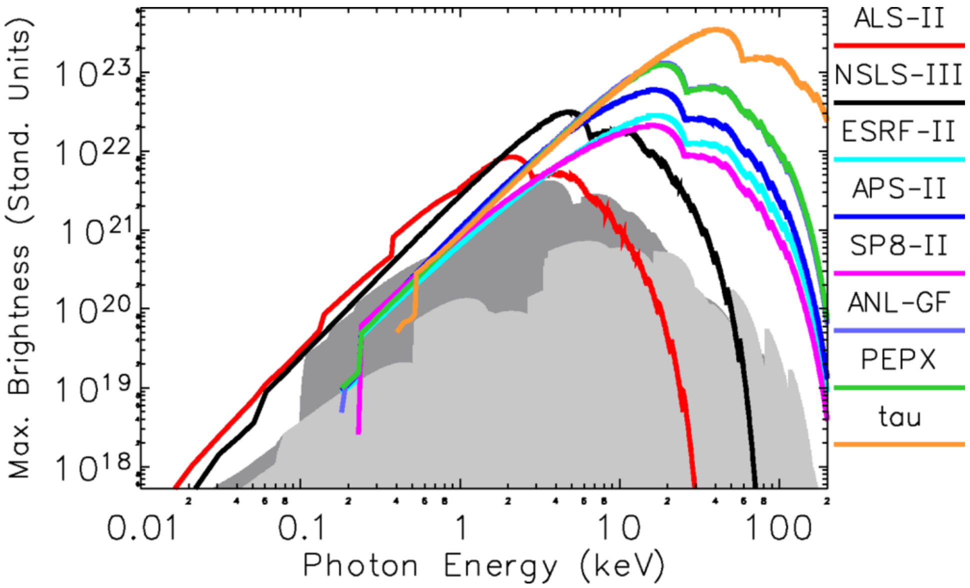
APS-U: 6 GeV, 60 pm x 8 pm,
200 mA (**Upgrade Proposal**)

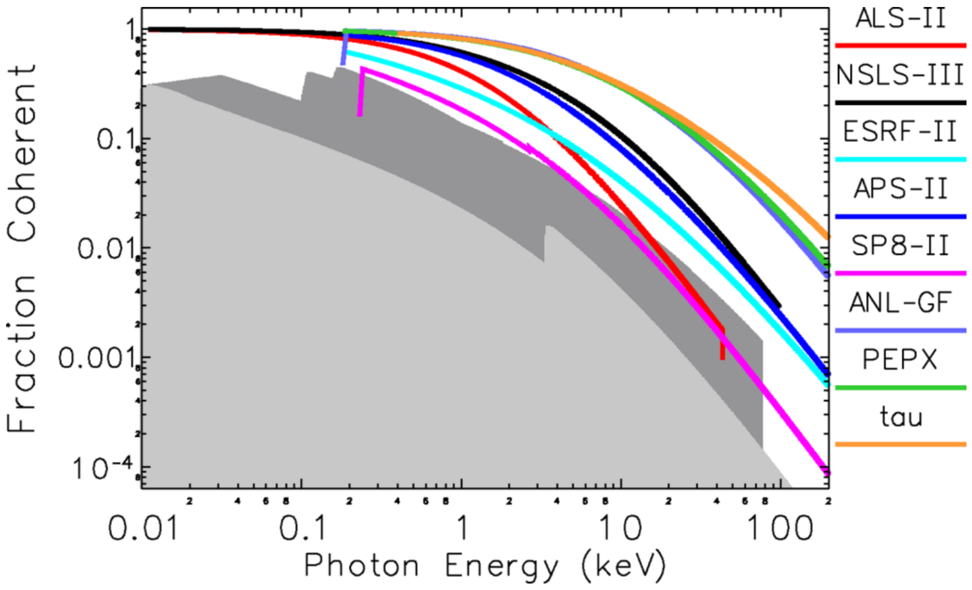


ALS-U: 2 GeV, 50 pm x 50 pm,
500 mA (**Upgrade proposal**)

Other international upgrades: ***Japan (Spring 8, 6 GeV), China (BAPS, 5 GeV), Germany (PETRA-IV, France (SOLEIL), Switzerland (SLS, 2.4 GeV), Italy (ELETTRA) and others are developing plans***

Brightness and coherence of 3rd and 4th gen rings





Legend:

🟥 0.2km/2GeV: ALS-II, 52 pm
⬛ 0.8km/3GeV: NSLS-III, 30 pm
🟦 1.1km/6GeV: APS-II, 80 pm
🟩 2.2km/6GeV: PEP-X, 5 pm
🟦 6.2km/9GeV: tauUSR, 3 pm

For 6 GeV rings, a common set of IDs, including advanced SCUs, was assumed

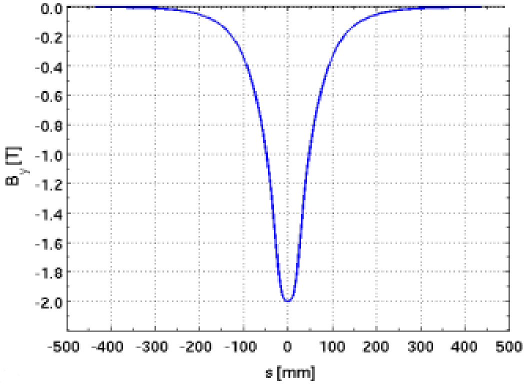
M. Borland

Sirius: 5 Bend Achromat

$\varepsilon_{x/y}$ = 190-270/3 pm.rad @ 3 GeV, 350 mA, C = 518 m

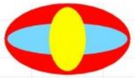
- New optics with lower β at the 6 m straight sections to match the electron beam phase-space to the photon beam phase-space from IDs.

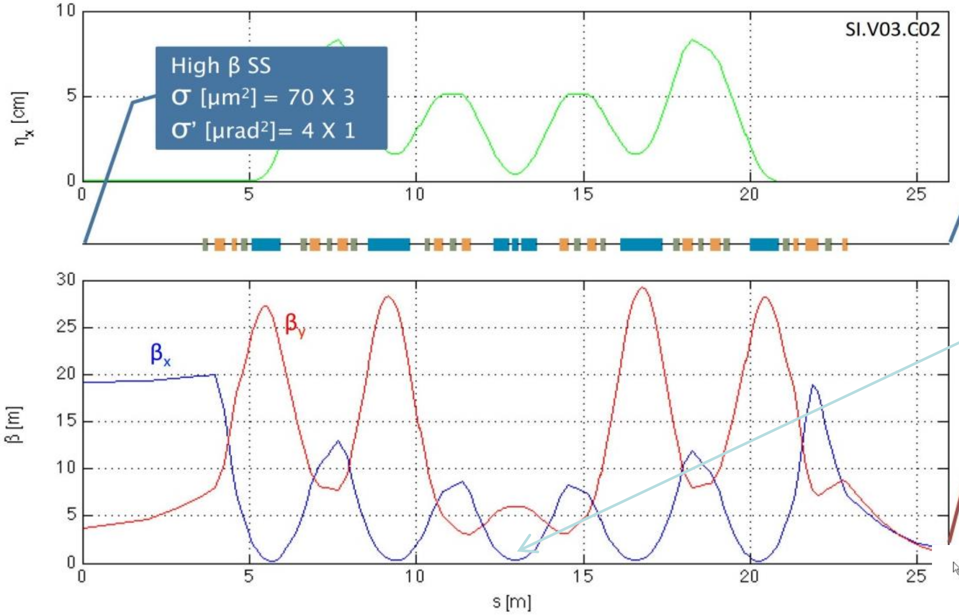
- PM Superbend
- B = 2 T
- Gap = 28 mm
- Field flexibility = 6% (using lateral control gap)
- 20 units total



κ = 1%

ε_y = 2.7 pm.rad



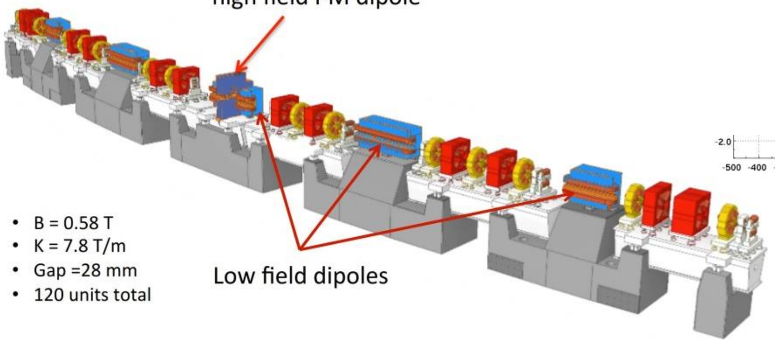


High β SS
σ [μm²] = 70 X 3
σ' [μrad²]= 4 X 1

Low β SS
σ [μm²] = 20 X 1.8
σ' [μrad²]= 13 X 1.5

superbend
σ [μm²] = 10 X 4
σ' [μrad²]= 29 X 1
ε_c [keV] = 12

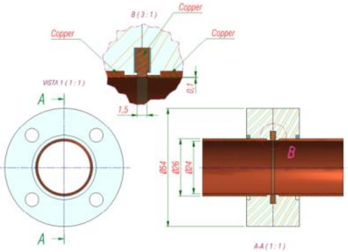
$\beta_x \approx \beta_y \approx 1\ \mathrm{m} \approx L/\pi$



high field PM dipole

- B = 0.58 T
- K = 7.8 T/m
- Gap =28 mm
- 120 units total

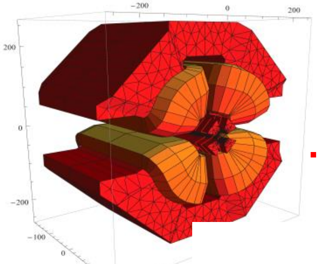
Low field dipoles



- gapless flanges
- NEG coating capability

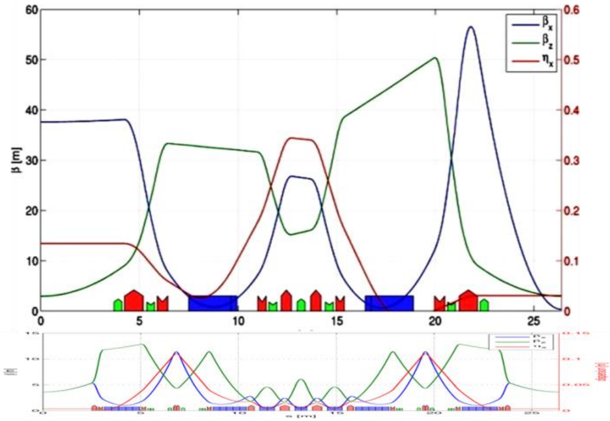
ESRF-II – hybrid 7BA

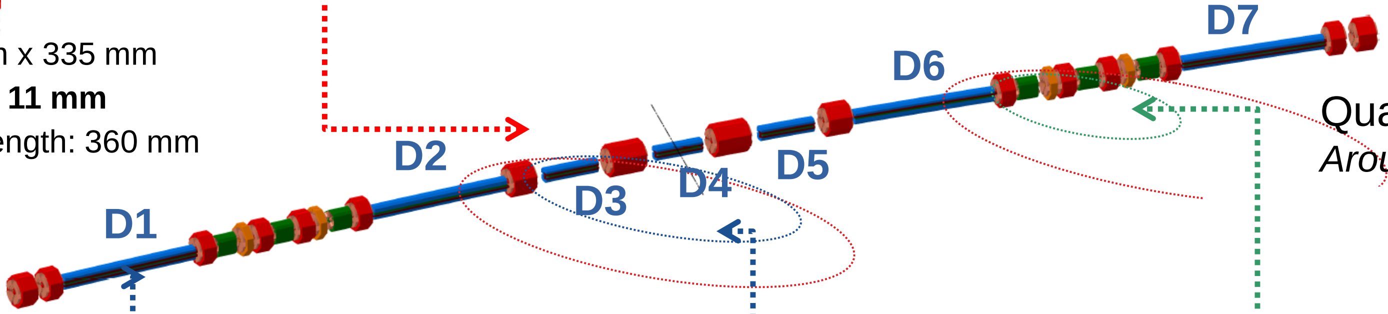
6 GeV, 844 m, 4 nm → 0.15 nm



High gradient quadrupoles
85 T/m

- Spec:100 T/m x 335 mm
- **Bore radius: 11 mm**
- Mechanical length: 360 mm
- 1 kW

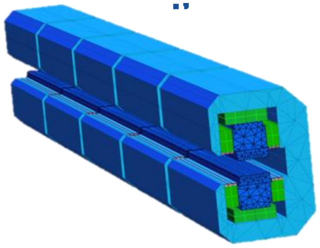


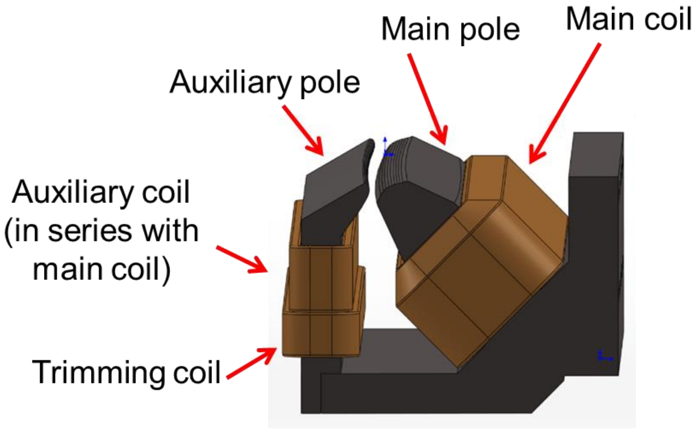


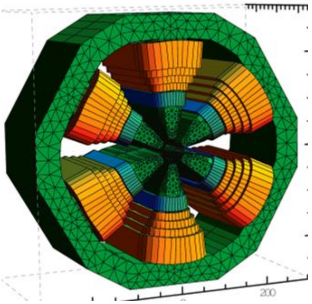
Quadrupole
Around 50 T/m

Combined dipole quadrupoles
0.85 T / 45 T/m & 0.34 T / 50 T/m

Sextupoles
1700 T/m^{-2}



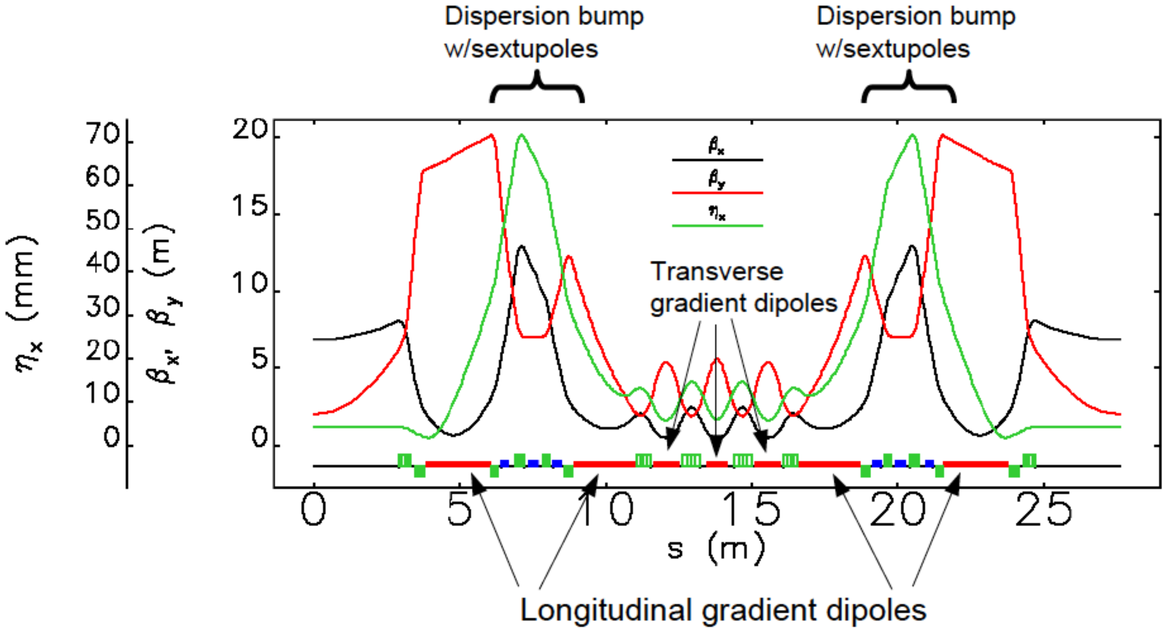




Permanent magnet dipoles
longitudinal gradient 0.16 – 0.6 T,
magnetic gap 22 mm
2 metre long, 5 modules
With a small tuning coil 1%

APS-U – hybrid 7BA

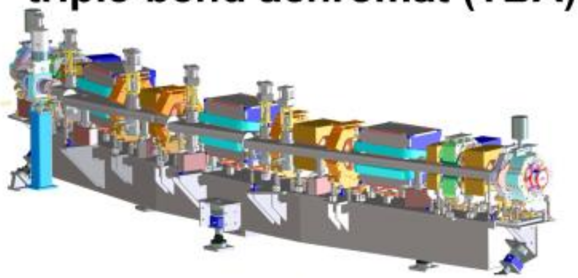
- $\varepsilon_{x/y}$ = 67/8 pm.rad @ 6 GeV, 200 mA
- C = 1.1 km
- 41-pm.rad option using reverse bend lattice

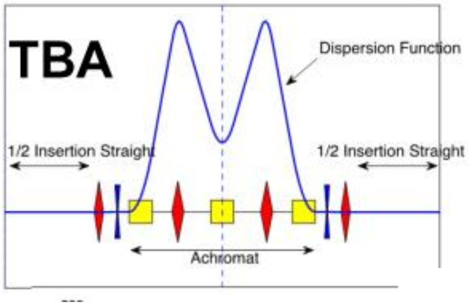


ALS-U – hybrid 9 Bend Achromat

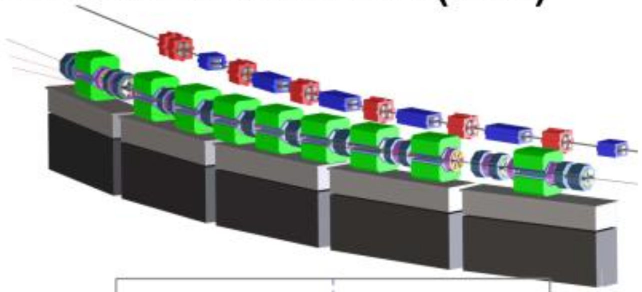
$\varepsilon_{x/y}$ = 50/50 pm.rad @ 2 GeV, 500 mA

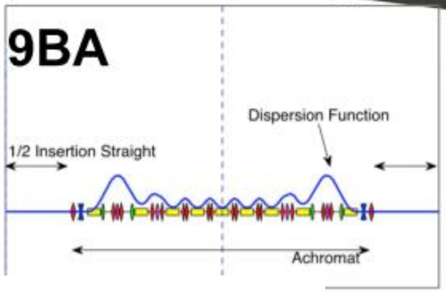
ALS today
triple-bend achromat (TBA)

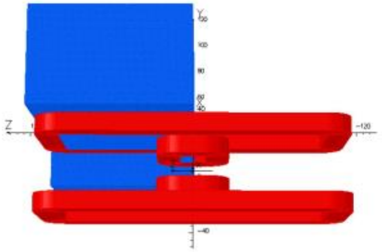


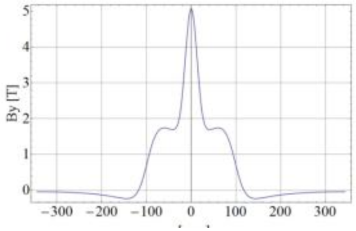


ALS-U
multi-bend achromat (9BA)

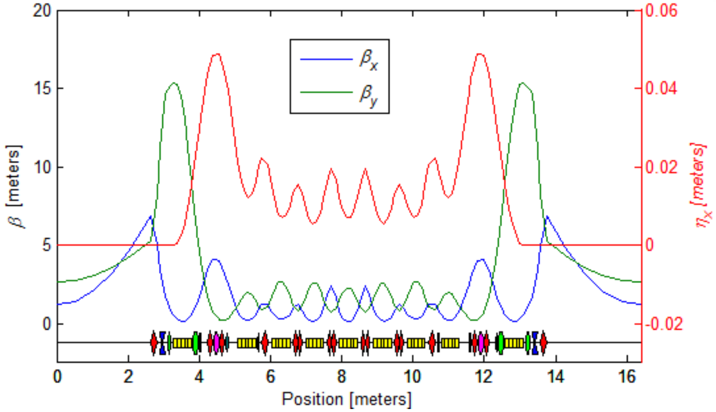








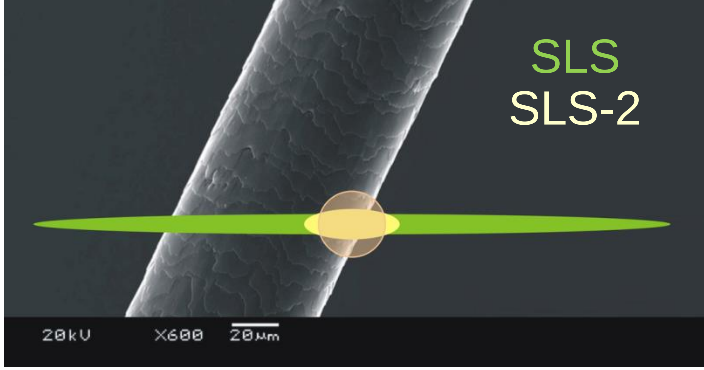
superbend option

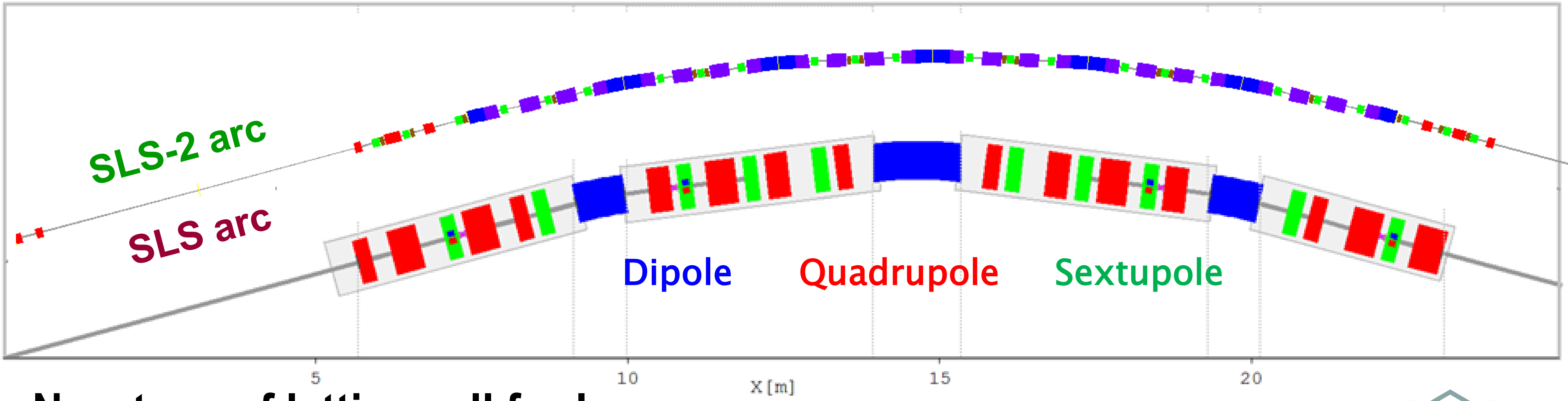


includes octupoles

SLS 2.0 upgrade plans

ε_x = **140 pm.rad @ 2.4 GeV**

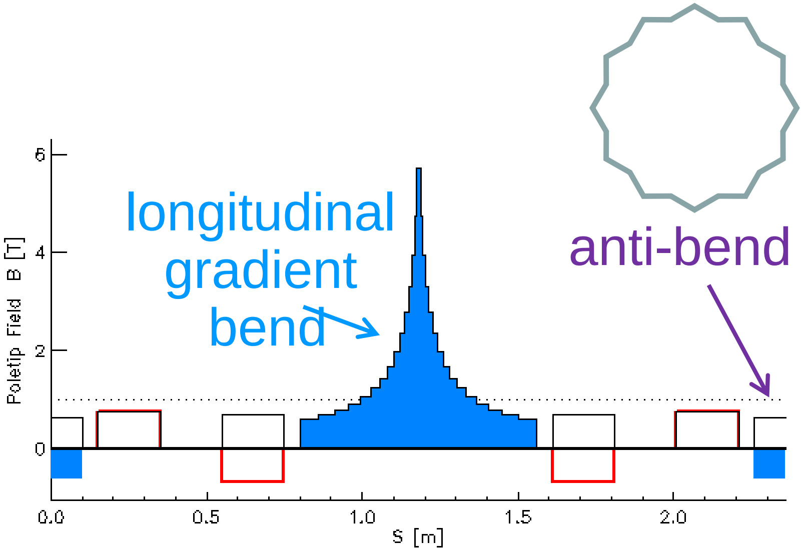




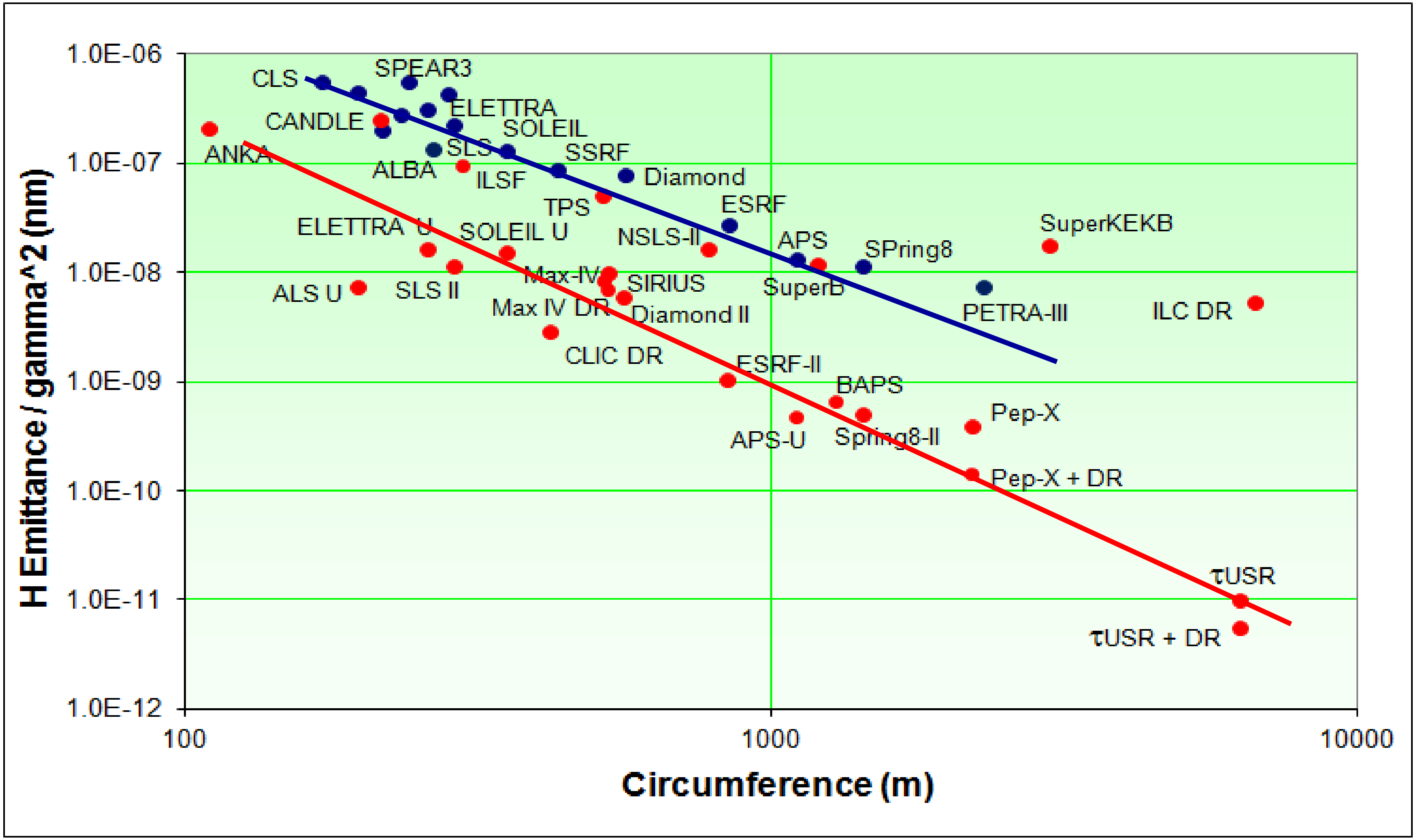
New type of lattice cell for low emittance

- bending magnets with longitudinal field variation (2 T peak)
- options for 5-6 T peak field superbends
- anti-bends for beam dynamics

⇒ star-shaped lattice



Diffraction limited rings: the new wave

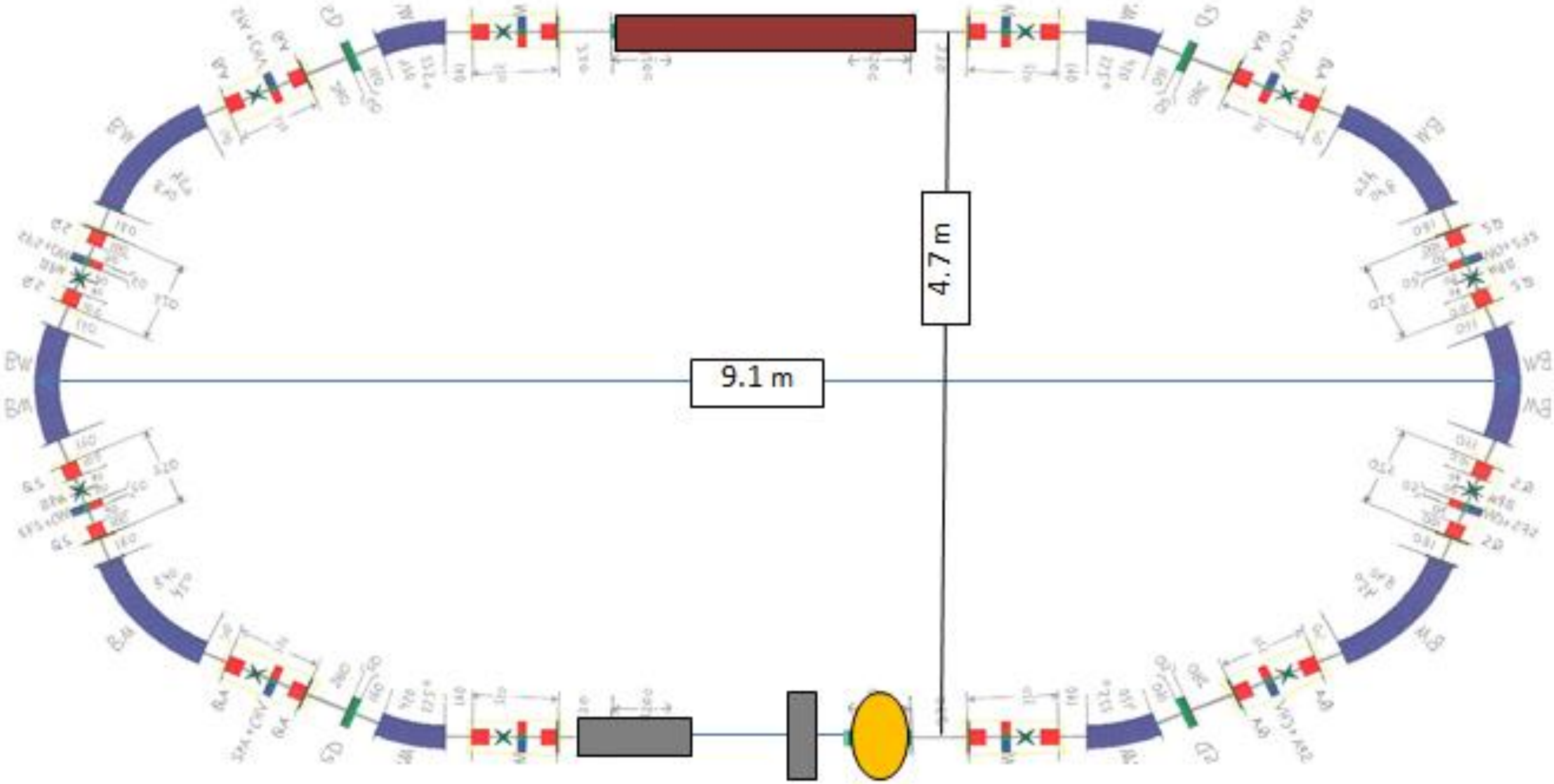


Storage rings in operation (●) and planned (●).
The old (—) and the new (—) generation.

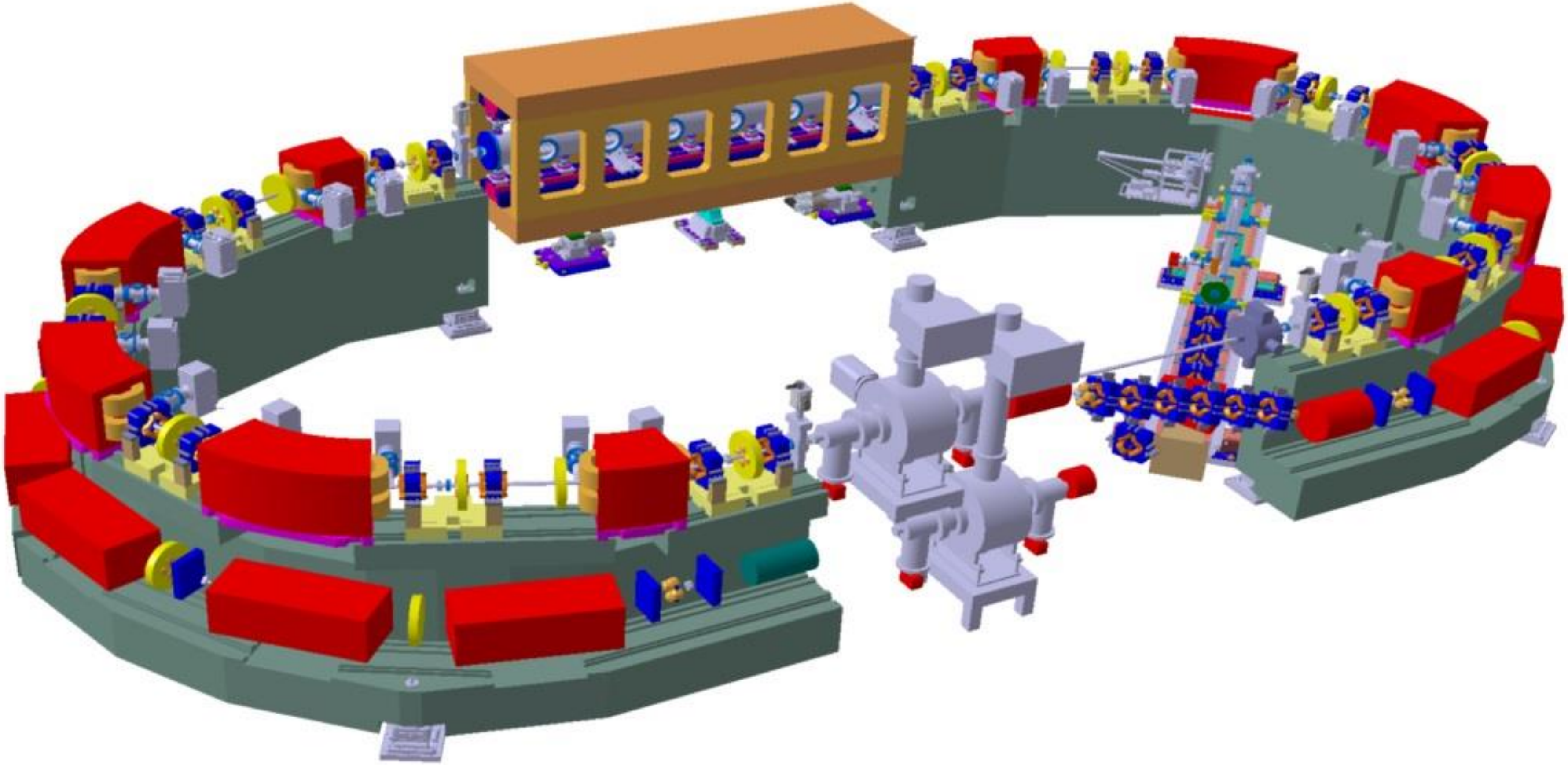
R. Bartolini

Compact light source COSAMI

Conventional, normal conducting magnetic structure

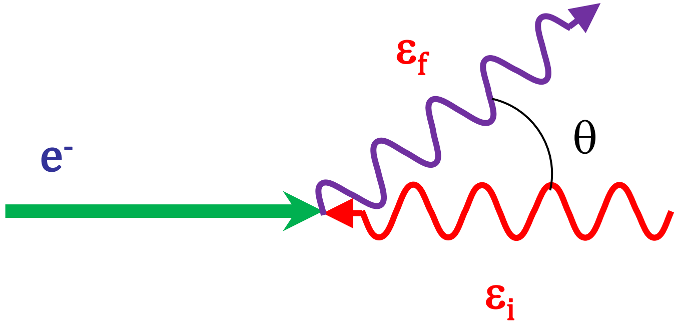


Injector and storage ring integration



When an electron collides with a photon…

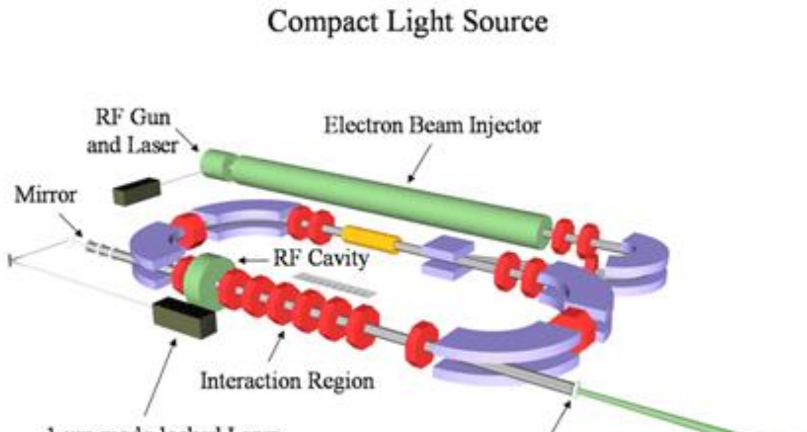
Also known as **Compton** or Thomson scattering

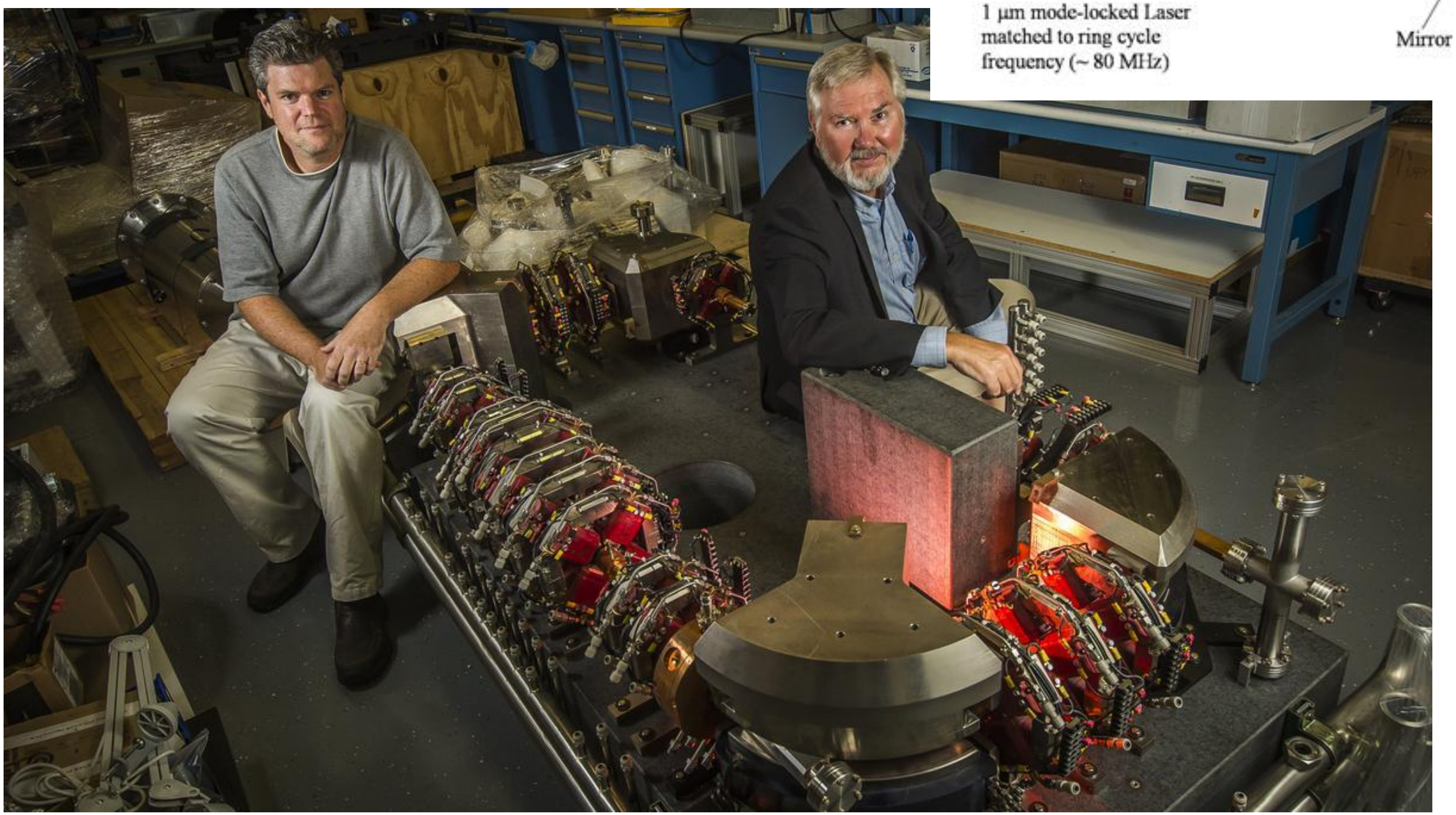


$$\varepsilon_f = \frac{4\gamma^2 \varepsilon_i}{1+\gamma^2\theta^2}$$

- backscattered photon has the maximum energy
- at an angle of $1/\gamma$ the energy drops by a factor of 2
- undulator’s periodic magnetic field could be viewed as a «photon», with useful parallels between the two cases

Compact light source based on Compton scattering





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