

Concept of Luminosity

(in particle colliders)

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(http://cern.ch/lhc-beam-beam/talks/Zeuthen_luminosity.pdf)

Why colliding beams ?

- Two beams:

$$E_1, \vec{p}_1, E_2, \vec{p}_2, m_1 = m_2 = m$$

- $E_{cm} = \sqrt{(E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2}$

- Collider versus fixed target:

- Fixed target: $\vec{p}_2 = 0$

- $\implies E_{cm} = \sqrt{2m^2 + 2E_1m}$

- Collider: $\vec{p}_1 = -\vec{p}_2$ (usually)

- $\implies E_{cm} = E_1 + E_2$

- LHC (pp):

14000 GeV versus 118 GeV

- LEP (e^+e^-):

210 GeV versus ?

Colliding beams

Collider performance qualified as:

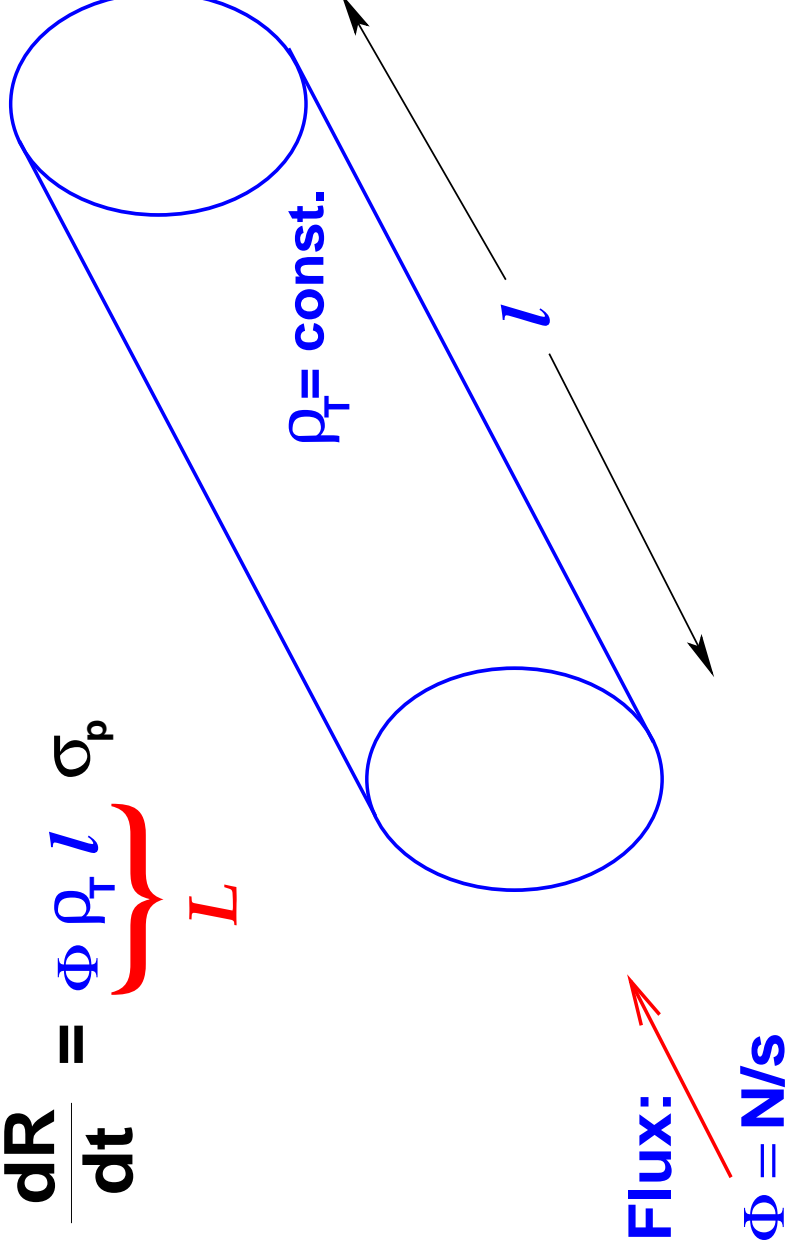
- Available energy
- Number of interactions per second (useful collisions)
- Total number of interactions
- Secondary issues:
 - Time structure of interactions (how often and how many at the same time)
 - Space structure of interactions (size of interaction region)
 - Quality of interactions (background, dead time etc.)

Examples

	Energy (GeV)	\mathcal{L} $\text{cm}^{-2}\text{s}^{-1}$	rate s^{-1}	σ_x/σ_y $\mu\text{m}/\mu\text{m}$	Particles per bunch
SPS ($p\bar{p}$)	315x315	6 10^{30}	4 10^5	60/30	$\approx 10 \ 10^{10}$
Tevatron ($p\bar{p}$)	1000x1000	50 10^{30}	4 10^6	30/30	$\approx 30/8 \ 10^{10}$
HERA (e^+p)	30x920	40 10^{30}	40	250/50	$\approx 3/7 \ 10^{10}$
LHC (pp)	7000x7000	10000 10^{30}	10^9	17/17	11 10^{10}
LEP (e^+e^-)	105x105	100 10^{30}	≤ 1	200/2	$\approx 5 \ 10^{11}$
PEP (e^+e^-)	9x3	3000 10^{30}	NA	150/5	$\approx 2/6 \ 10^{10}$
KEKB (e^+e^-)	8x3.5	10000 10^{30}	NA	77/2	$\approx 1.3/1.6 \ 10^{10}$

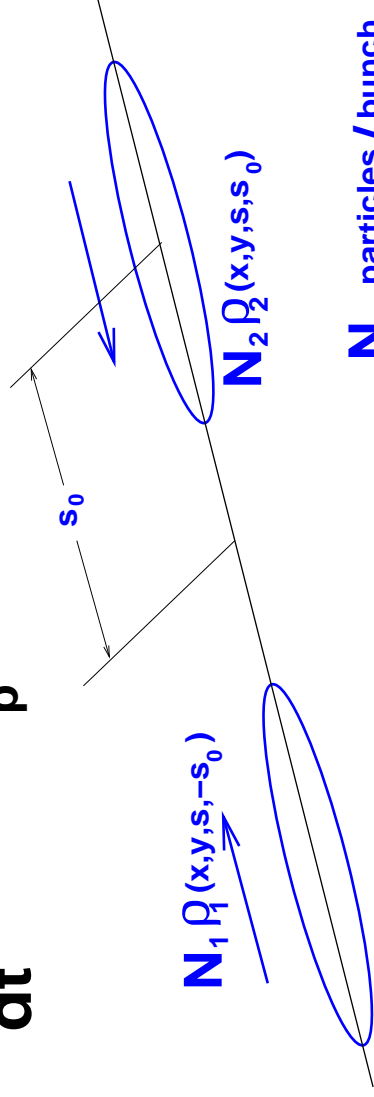
Fixed target luminosity

$$\frac{dR}{dt} = \underbrace{\Phi \rho_T l}_{L} \sigma_p$$



Collider luminosity

$$\frac{dR}{dt} = L \sigma_p$$



N particles / bunch
ρ density ≠ const.

- $\mathcal{L} \propto K \cdot \int \int \int \int_{-\infty}^{+\infty} \rho_1(x, y, s, -s_0) \rho_2(x, y, s, s_0) dx dy ds ds_0$
- s_0 is "time"-variable: $s_0 = c \cdot t$
- Kinematic factor : $K = \sqrt{(\vec{v}_1 - \vec{v}_2)^2 - (v_1 \times v_2)^2} / c^2$

Luminosity

- Assume densities are uncorrelated in all planes
- For head-on collisions ($\vec{v}_1 = -\vec{v}_2$) we get:

$$\mathcal{L} = 2N_1N_2fB \int \int \int_{-\infty}^{+\infty} dx dy ds s_0 \rho_{1x}(x) \rho_{1y}(y) \rho_{1s}(s - s_0) \rho_{2x}(x) \rho_{2y}(y) \rho_{2s}(s + s_0)$$

- In principle: should know all distributions
- Gaussian distributions for analytic calculation
- In general: they are a good approximation for most realistic distributions

Gaussian distribution functions

- $\rho_{iz}(z) = \frac{1}{\sigma_z \sqrt{2\pi}} \exp\left(-\frac{z^2}{2\sigma_z^2}\right) \quad i = 1, 2, \quad z = x, y$

- $\rho_{is}(s \pm s_0) = \frac{1}{\sigma_s \sqrt{2\pi}} \exp\left(-\frac{(s \pm s_0)^2}{2\sigma_s^2}\right)$

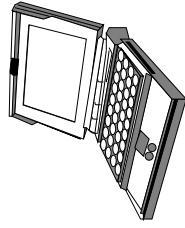
- **For non-Gaussian profiles not always possible to find analytic form, need a numerical integration**

Luminosity for two (equal) beams

- Simplest case : $\sigma_{1x} = \sigma_{2x}, \sigma_{1y} = \sigma_{2y}, \sigma_{1s} = \sigma_{2s}$
or: $\sigma_{1x} \neq \sigma_{2x}, \sigma_{1y} \neq \sigma_{2y}, \sigma_{1s} \approx \sigma_{2s}$

- Further: no horizontal or vertical dispersion at collision point

$$\Rightarrow \mathcal{L} = \frac{N_1 N_2 f B}{4\pi \sigma_x \sigma_y} \left(\mathcal{L} = \frac{N_1 N_2 f B}{2\pi \sqrt{\sigma_{1x}^2 + \sigma_{2x}^2} \sqrt{\sigma_{2y}^2 + \sigma_{2y}^2}} \right)$$



Hints for the integration in the handout

Integration (head-on)

for $\sigma_1 = \sigma_2 \rightarrow \rho_1 \rho_2 = \rho^2$:

$$\mathcal{L} = \frac{2 \cdot N_1 N_2 f N_b}{(\sqrt{2\pi})^6 \sigma_s^2 \sigma_x^2 \sigma_y^2} \int \int e^{-\frac{x^2}{\sigma_x^2} - \frac{y^2}{\sigma_y^2} - \frac{s^2}{\sigma_s^2} - \frac{s_0^2}{\sigma_s^2}} dx dy ds ds_0$$

integrating over s and s_0 , using:

$$\int_{-\infty}^{+\infty} e^{-at^2} dt = \sqrt{\pi/a}$$

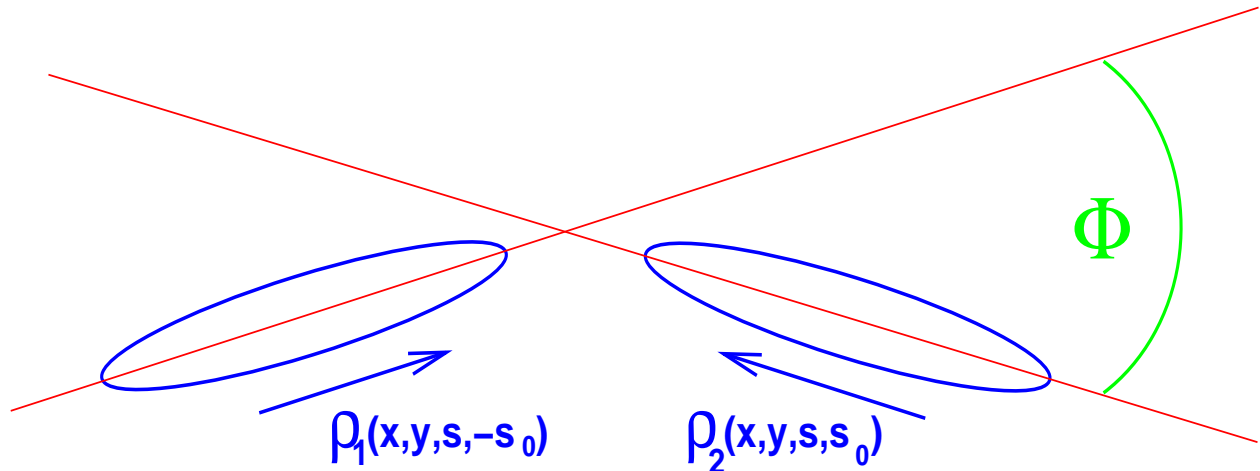
$$\mathcal{L} = \frac{2 \cdot N_1 N_2 f N_b}{8(\sqrt{\pi})^4 \sigma_x^2 \sigma_y^2} \int \int e^{-\frac{x^2}{\sigma_x^2} - \frac{y^2}{\sigma_y^2}} dx dy$$

finally after integration over x and y : $\Rightarrow \mathcal{L} = \frac{N_1 N_2 f B}{4\pi \sigma_x \sigma_y}$

Complications

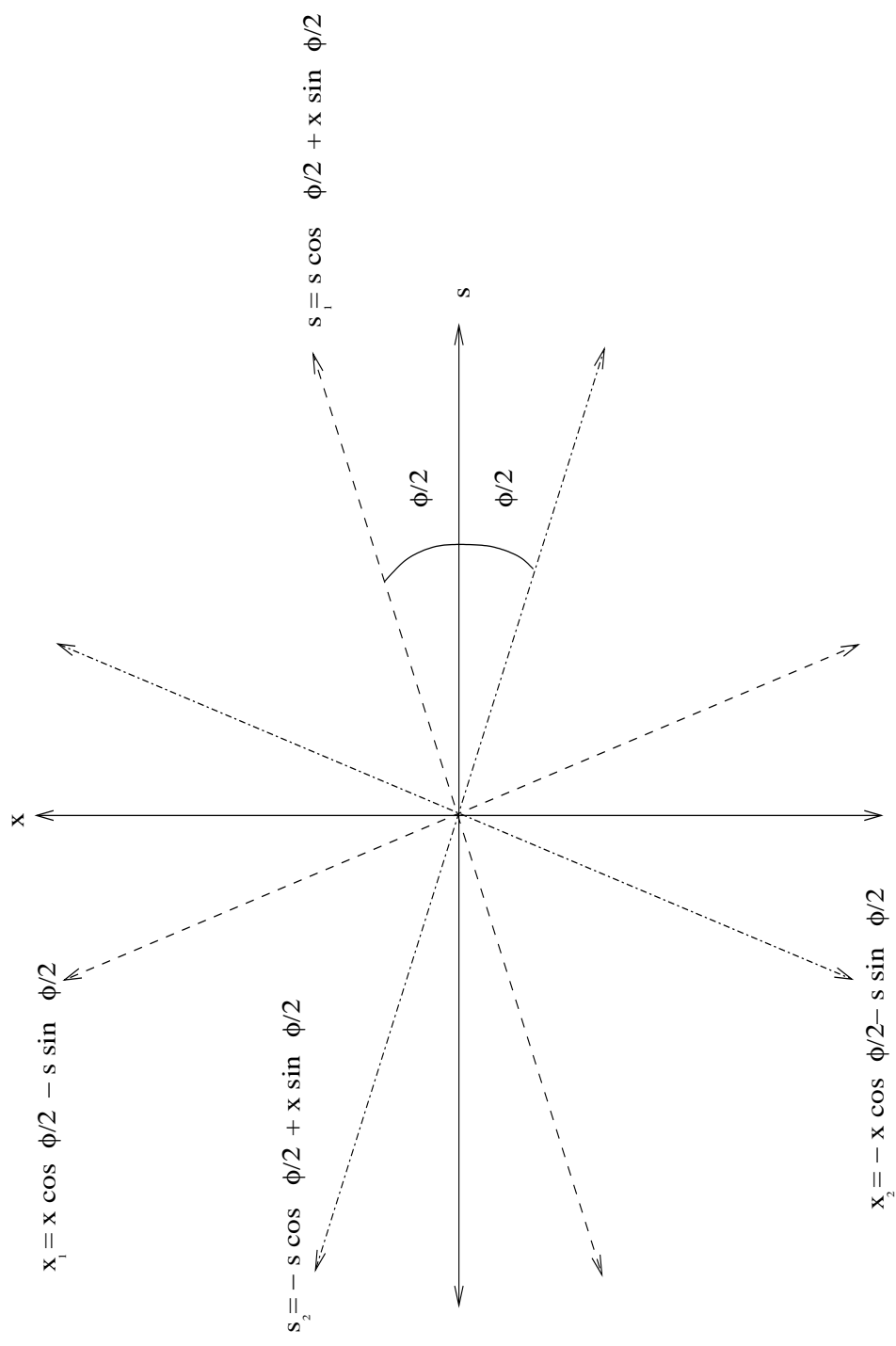
- Crossing angle
- Hour glass effect
- Collision offset (wanted or unwanted)
- Non-Gaussian profiles
- Dispersion at collision point
- $\delta\beta^*/\delta s = \alpha^* \neq 0$
- Strong coupling
- etc.

Collisions at crossing angle



- Needed to avoid unwanted collisions
- For colliders with many bunches: LHC, CESR, KEKB
- For colliders with coasting beams

Collisions angle geometry



Crossing angle

Assume crossing in horizontal plane. Transform to new coordinates:

$$\begin{cases} x_1 = x \cos \frac{\phi}{2} - s \sin \frac{\phi}{2}, & s_1 = s \cos \frac{\phi}{2} + x \sin \frac{\phi}{2}, \\ x_2 = x \cos \frac{\phi}{2} + s \sin \frac{\phi}{2}, & s_2 = s \cos \frac{\phi}{2} - x \sin \frac{\phi}{2} \end{cases}$$

$$\mathcal{L} = 2 \cos^2 \frac{\phi}{2} N_1 N_2 f B \int \int \int_{-\infty}^{+\infty} dx dy ds s_0$$

$$\rho_{1x}(x_1) \rho_{1y}(y_1) \rho_{1s}(s_1 - s_0) \rho_{2x}(x_2) \rho_{2y}(y_2) \rho_{2s}(s_2 + s_0)$$

Integration (crossing angle)

use as before:

$$\int_{-\infty}^{+\infty} e^{-at^2} dt = \sqrt{\pi/a}$$

and:

$$\int_{-\infty}^{+\infty} e^{-(at^2+bt+c)} dt = \sqrt{\pi/a} \cdot e^{\frac{b^2-ac}{a}}$$

Further:

- Since both x and $\sin(\phi/2)$ are small: drop all terms $\sigma_x^k \sin^l(\phi/2)$ or $x^k \sin^l(\phi/2)$ for all $k+l \geq 4$
- Approximate: $\sin(\phi/2) \approx \tan(\phi/2) \approx \phi/2$

Crossing angle

- Crossing Angle $\Rightarrow \mathcal{L} = \frac{N_1 N_2 f B}{4\pi \sigma_x \sigma_y} S$
- S is the reduction factor
- $S = \frac{1}{\sqrt{1 + \left(\frac{\sigma_x}{\sigma_s} \tan \frac{\phi}{2}\right)^2}} \frac{1}{\sqrt{1 + \left(\frac{\sigma_s}{\sigma_x} \tan \frac{\phi}{2}\right)^2}}$
- For small crossing angles and $\sigma_s \gg \sigma_{x,y}$
- $\Rightarrow S = \frac{1}{\sqrt{1 + \left(\frac{\sigma_s}{\sigma_x} \tan \frac{\phi}{2}\right)^2}} \approx \frac{1}{\sqrt{1 + \left(\frac{\sigma_s \phi}{\sigma_x 2}\right)^2}}$

Example LHC:

$\Phi = 285 \mu\text{rad}$, $\sigma_s = 7.5 \text{ cm}$, $S = 0.81$

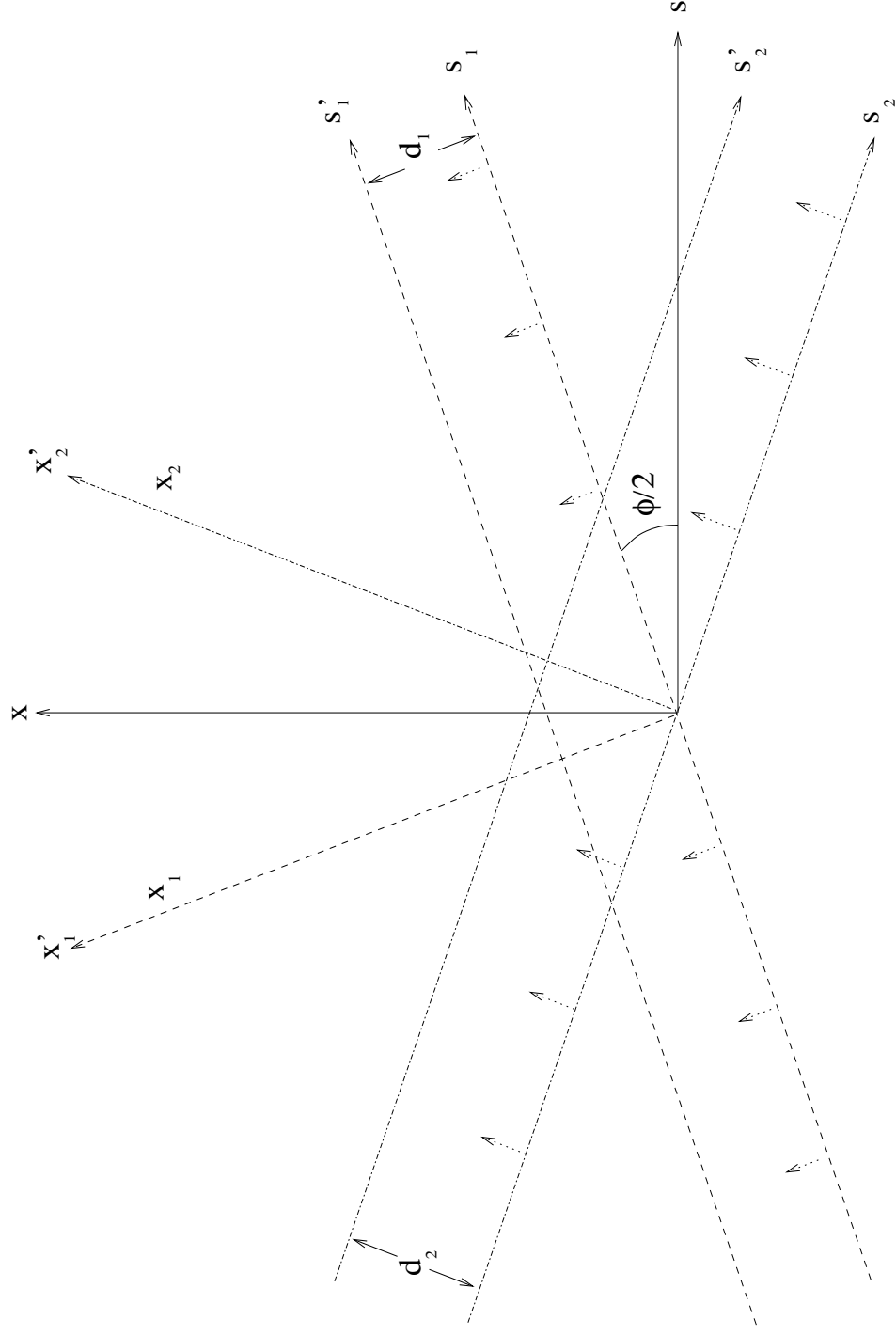
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Example LHC:

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Offset and crossing angle



Offset and crossing angle

- Offsets in crossing plane: d_1 and d_2

- Transformations slightly changed:

$$\begin{cases} x_1 = d_1 + x \cos \frac{\phi}{2} - s \sin \frac{\phi}{2}, & s_1 = s \cos \frac{\phi}{2} + x \sin \frac{\phi}{2}, \\ x_2 = d_2 + x \cos \frac{\phi}{2} + s \sin \frac{\phi}{2}, & s_2 = s \cos \frac{\phi}{2} - x \sin \frac{\phi}{2} \end{cases}$$

- Gives after integration over y and s_0 :

$$\mathcal{L} = \frac{N_1 N_2 f N_b}{8\pi^2 \sigma_s \sigma_x^2 \sigma_y} \int \int e^{-\frac{x^2 \cos^2(\phi/2) + s^2 \sin^2(\phi/2)}{\sigma_x^2}} e^{-\frac{x^2 \sin^2(\phi/2) + s^2 \cos^2(\phi/2)}{\sigma_s^2}} \times e^{-\frac{d_1^2 + d_2^2 + 2(d_1 + d_2)x \cos(\phi/2) - 2(d_2 - d_1)s \sin(\phi/2)}{2\sigma_x^2}} dx ds.$$

Offset and crossing angle

- After integration over x :

$$\mathcal{L} = \frac{N_1 N_2 f N_b}{8\pi^{\frac{3}{2}} \sigma_s} 2 \cos \frac{\phi}{2} \int_{-\infty}^{+\infty} W \frac{e^{-(As^2+2Bs)}}{\sigma_x \sigma_y} ds$$

- with:

$$A = \frac{\sin^2 \frac{\phi}{2}}{\sigma_x^2} + \frac{\cos^2 \frac{\phi}{2}}{\sigma_s^2} \quad B = \frac{(d_2 - d_1) \sin(\phi/2)}{2\sigma_x^2}$$

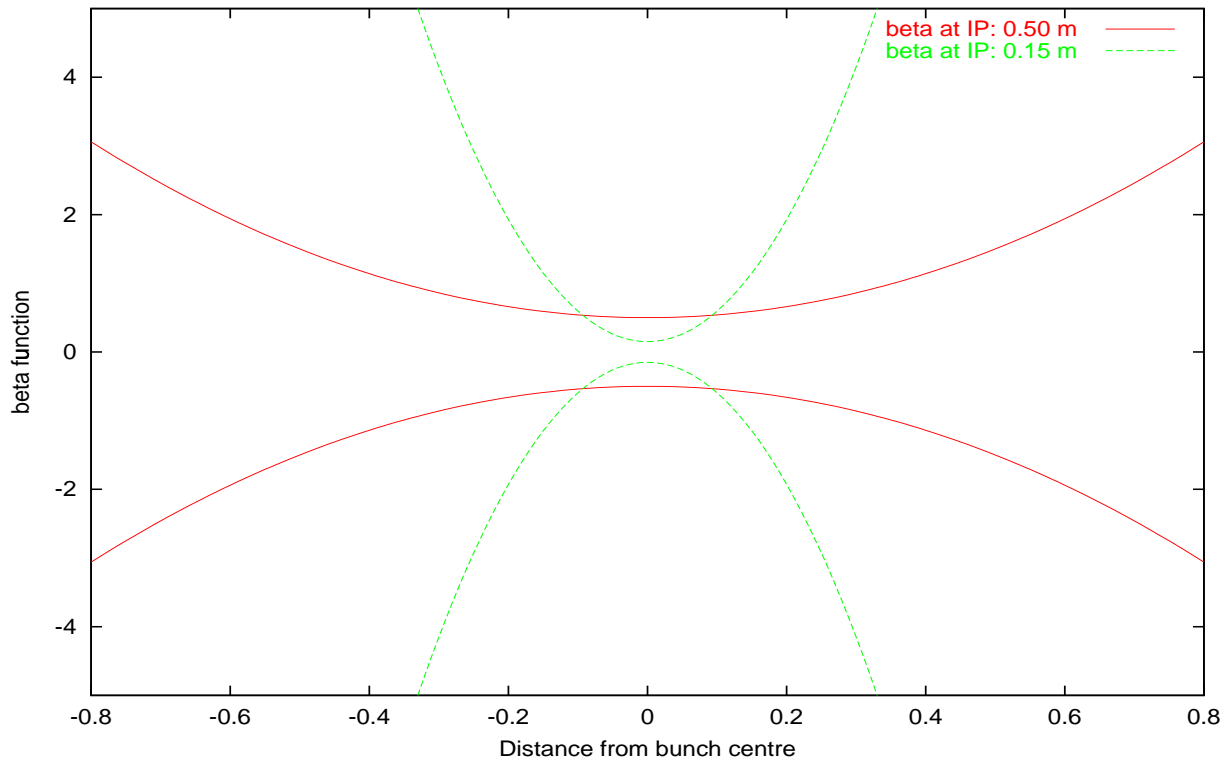
- and

$$W = e^{-\frac{1}{4\sigma_x^2}(d_2-d_1)^2}$$

⇒ Luminosity with correction factors:

$$\mathcal{L} = \frac{N_1 N_2 f N_b}{4\pi \sigma_x \sigma_y} \cdot W \cdot e^{\frac{B^2}{A}} \cdot S$$

Hour glass effect



- β -functions depends on position s
- Usually: $\beta(s) = \beta^* \left(1 + \left(\frac{s}{\beta^*}\right)^2\right)$
i.e. $\sigma \implies \sigma(s) \neq \text{const.}$
- Important when β^* comparable to bunch length σ_s

Hour glass effect

- Replace σ by $\sigma(s)$ in formulae
- $$\mathcal{L} = \left(\frac{N_1 N_2 f B}{8\pi\sigma_x^* \sigma_y^*} \right) \frac{2 \cos \frac{\phi}{2}}{\sqrt{\pi}\sigma_s} \int_{-\infty}^{+\infty} \frac{e^{-s^2 A}}{1 + \left(\frac{s}{\beta^*}\right)^2} ds$$
- $$A = \frac{\sin^2 \frac{\phi}{2}}{(\sigma_x^*)^2 \left[1 + \left(\frac{s}{\beta^*}\right)^2\right]} + \frac{\cos^2 \frac{\phi}{2}}{\sigma_s^2}$$
- \rightarrow Numerical Integration

Calculations for the LHC

- $N_1 = N_2 = 1.15 \times 10^{11}$ particles/bunch
- 2808 bunches/beam
- $f = 11.2455$ kHz, $\phi = 285$ μ rad
- $\beta_x^* = \beta_y^* = 0.55$ m
- $\sigma_x^* = \sigma_y^* = 16.6$ μ m, $\sigma_s = 7.7$ cm

- **Simplest case (Head on collision):**
$$\mathcal{L} = 1.200 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$
- **Effect of crossing angle:**
$$\mathcal{L} = 0.973 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$
- **Effect of crossing angle & Hourglass:**
$$\mathcal{L} = 0.969 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$$

If the beams are not Gaussian ??

Exercise:

- Assume flat distributions:
(normalized to 1) $\rho_1 = \rho_2 = \frac{1}{2a}$,
for $[-a \leq z \leq a]$, $z = x, y$
- Calculate r.m.s. in x and y:
 $\langle (x, y)^2 \rangle = \int_{-\infty}^{+\infty} (x, y)^2 \cdot \rho(x, y) dx dy$
- and $\mathcal{L} = \int_{-\infty}^{+\infty} \rho_1(x, y) \rho_2(x, y) dx dy$
- Compute: $\mathcal{L} \sqrt{\langle x^2 \rangle \langle y^2 \rangle}$
- Repeat for Gaussian distribution and compare
- Try other distributions (parabolic, triangular, cosinelike etc.)

Integrated luminosity

- $\mathcal{L}_{\text{int}}(s) = \int_0^T \int_{-s}^{+s} \mathcal{L}(s', t) ds' dt$
- The real figure of merit:
 $\mathcal{L}_{\text{int}}(s) \cdot \sigma_p = \text{number of events}$
- During data taking: continuous recording of \mathcal{L}
- For studies: assume some life time behaviour. E.g.
 $\mathcal{L}(t) \longrightarrow \mathcal{L}_0 \exp\left(-\frac{t}{\tau}\right)$
- Contributions to life time from:
intensity decay, emittance growth, bunch length increase etc.

Integrated luminosity

- Knowledge of preparation time allows optimization of \mathcal{L}_{int}



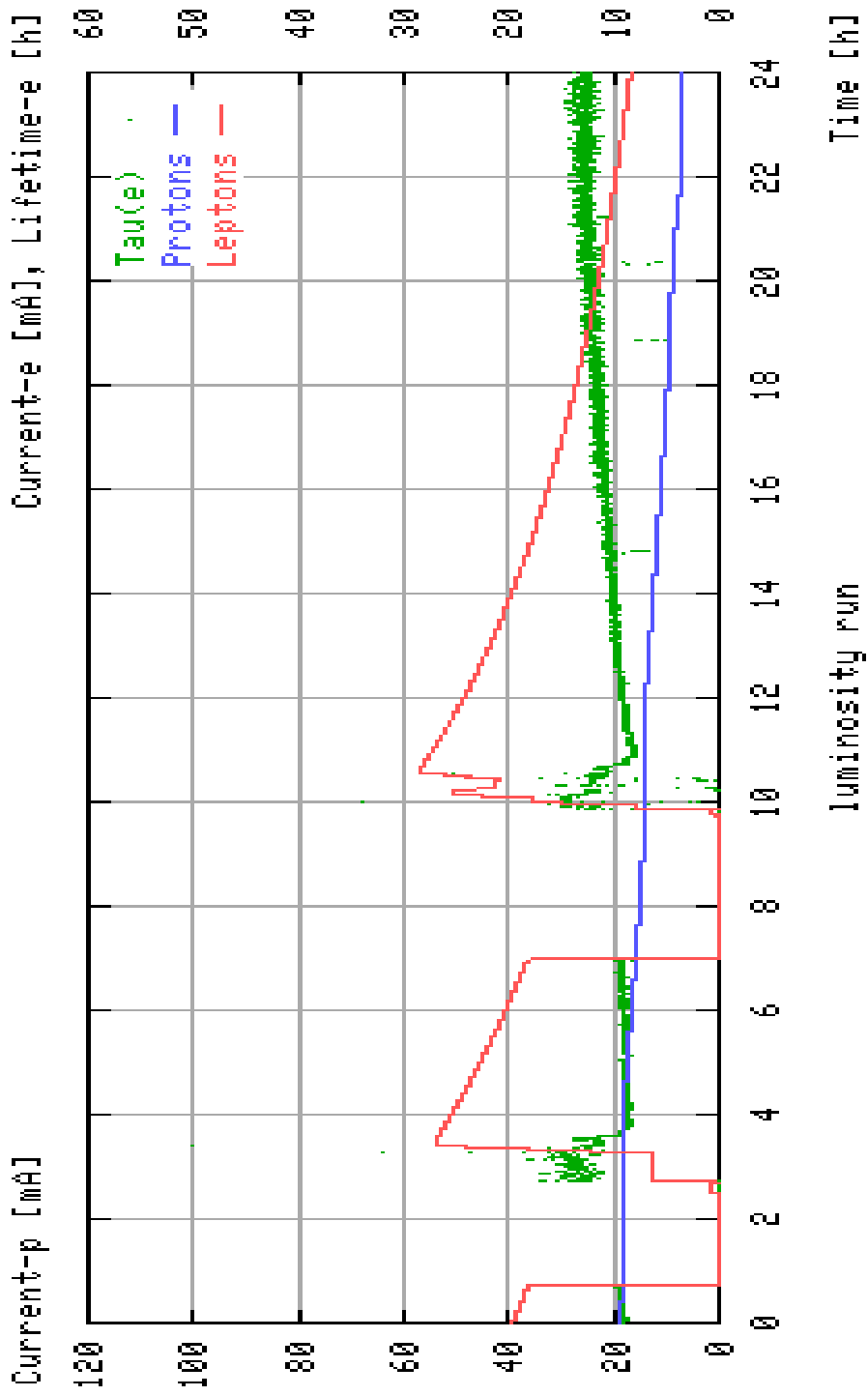
- Typical run times LEP:
 $t_r \approx 8 - 10$ hours
- For LHC long preparation time
 t_p expected

Maximising Integrated luminosity

- Assume exponential decay of luminosity $\mathcal{L}(t) = \mathcal{L}_0 \cdot e^{t/\tau}$
- Average luminosity $\langle \mathcal{L} \rangle$
$$\langle \mathcal{L} \rangle = \frac{\int_0^{t_r} dt \mathcal{L}(t)}{t_r + t_p} = \mathcal{L}_0 \cdot \tau \cdot \frac{1 - e^{-t_r/\tau}}{t_r + t_p}$$
- t_r is a "free" parameter, i.e. can be chosen
- (Theoretical) maximum for:
$$t_r \approx \tau \cdot \ln(1 + \sqrt{2t_p/\tau} + t_p/\tau)$$
- Example LHC (OB): $t_p \approx 10\text{h}$,
 $\tau \approx 15\text{h}$, $\Rightarrow t_r \approx 15\text{h}$
- Exercise: Would you try to improve τ or t_p ???

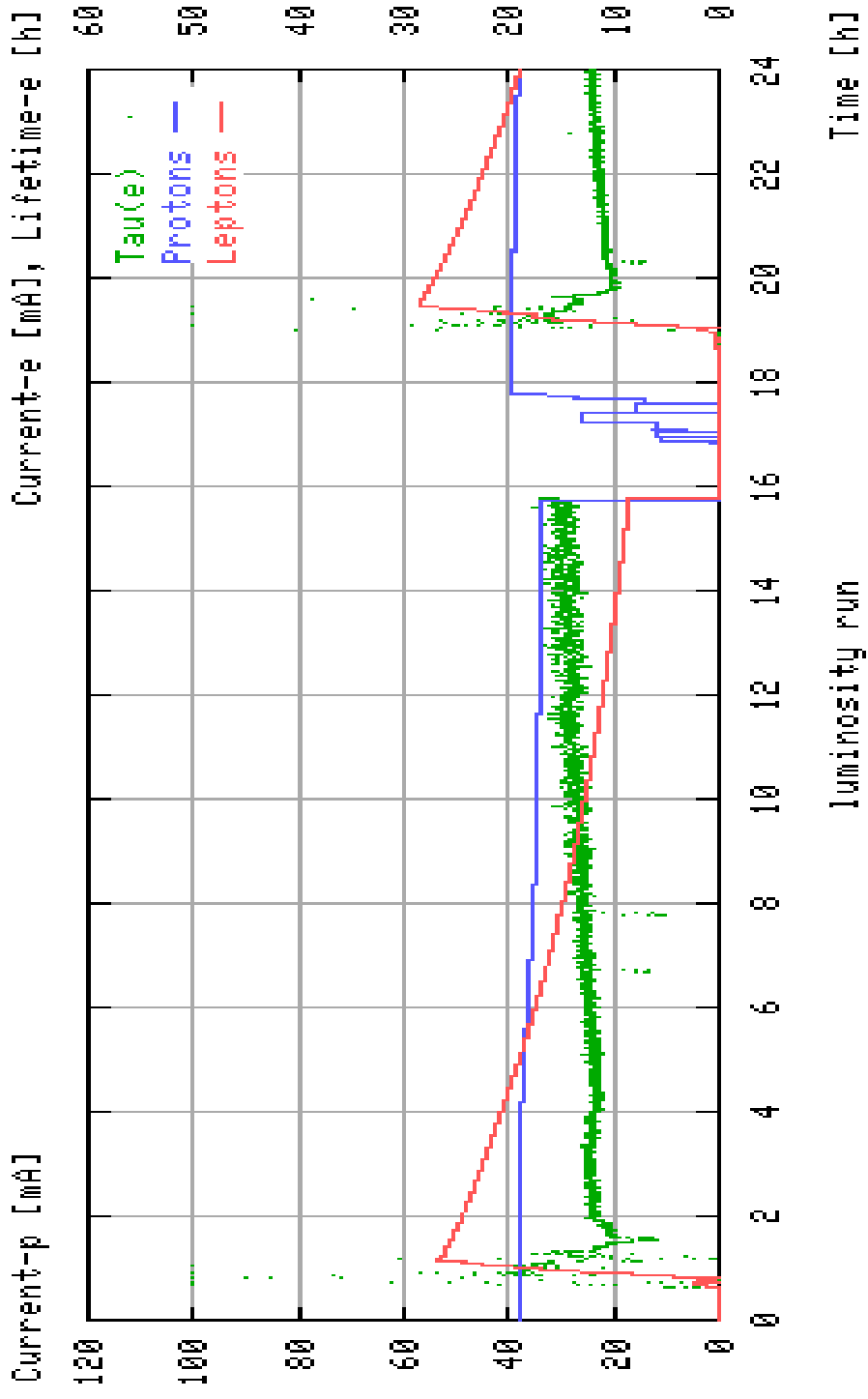
HERA on Saturday November 23 2002

p:7.2[mA] 73.9[h] 920[GeV] e:8.4[mA] 13.9[h] 27.6[GeV]



HERA on Thursday December 12 2002

p:38.1[mA] 123.9[h] 920[GeV] e+:18.9[mA] 12.1[h] 27.6[GeV]



Luminous region

- Density distribution of interaction vertices
- OR: Which fraction of collisions occur $\pm s$ from the IP ?
- Important for experiments !
- Depends on σ_x , σ_y , and σ_s
- But also on crossing angle ϕ

$$\mathcal{L}_0 = \int_{-\infty}^{+\infty} \mathcal{L}(s') ds' \longrightarrow \mathcal{L}(s) = \int_{-s}^{+s} \mathcal{L}(s') ds'$$

$$\mathcal{L}(s) = \left(\frac{N_1 N_2 f B}{8\pi\sigma_x^* \sigma_y^*} \right) \frac{2 \cos \frac{\phi}{2}}{\sqrt{\pi}\sigma_s} \sqrt{\frac{\pi}{A}} \operatorname{erf} \left(\sqrt{A} s \right)$$

Some results for LHC

- $\sigma_s = 7.7$ cm, $\beta^* = 0.55$ m, $\phi = 285$ μ rad:
- 100% lumi $\rightarrow s = \pm 12$ cm $\rightarrow s = \pm 12$ cm
- 95% lumi $\rightarrow s = \pm 8$ cm $\rightarrow s = \pm 9$ cm
- 90% lumi $\rightarrow s = \pm 7$ cm $\rightarrow s = \pm 8$ cm
- 85% lumi $\rightarrow s = \pm 6$ cm $\rightarrow s = \pm 6.5$ cm
- 80% lumi $\rightarrow s = \pm 5.5$ cm $\rightarrow s = \pm 6$ cm

Interactions per crossing

- Luminosity $\propto B N_1 N_2$
- In LHC: crossing every 25 ns
- Per crossing approximately 20 interactions
- May be undesirable (pile up in detector)
- \implies more bunches B , or smaller N ??

Luminosity measurement

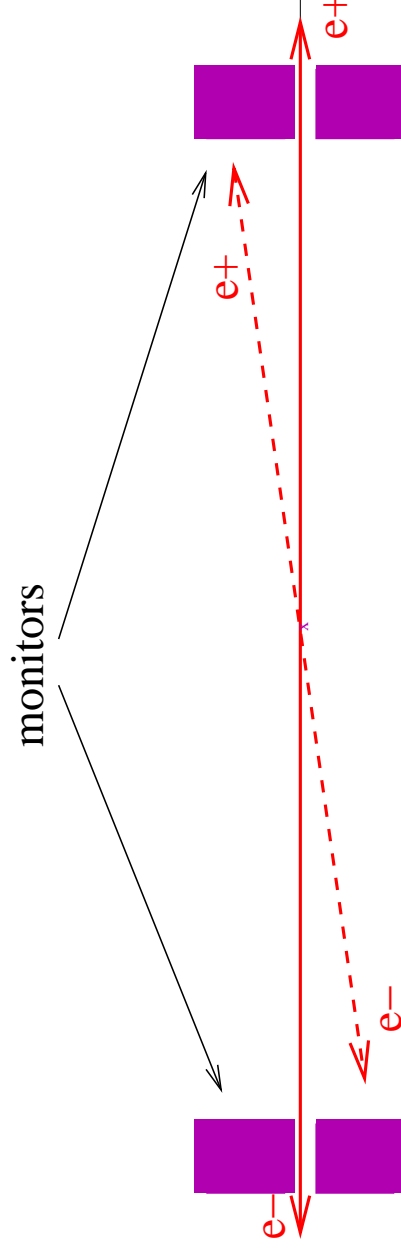
- Provide a signal proportional to interaction rate
- Large dynamic range:
 $10^{27} \text{ cm}^{-2}\text{s}^{-1}$ to $10^{34} \text{ cm}^{-2}\text{s}^{-1}$
- Very fast, if possible for individual bunches
- Used for optimization
- For absolute luminosity need calibration method

Luminosity calibration

$$(e^+e^-)$$

- Use well known and calculable process
- $e^+e^- \rightarrow e^+e^-$ elastic scattering (Bhabha scattering)
- Have to go to small angles
($\sigma_{el} \propto \Theta^{-3}$)
- Small rates at high energy
($\sigma_{el} \propto \frac{1}{E^2}$)

Luminosity calibration



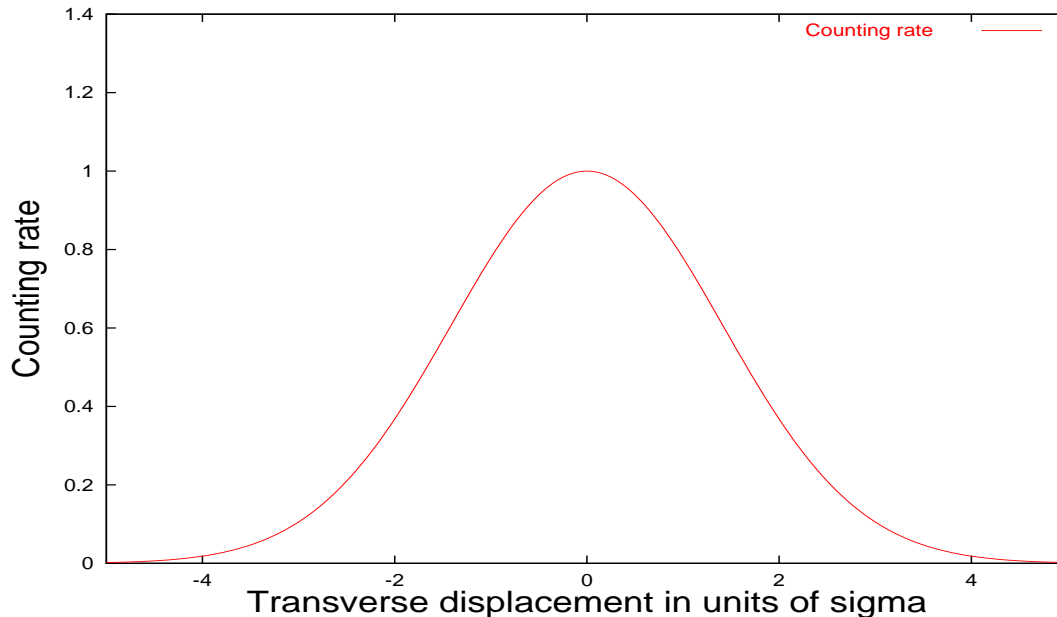
- Measure coincidence at small angles
- Low counting rates (LEP: 25 Hz for $\mathcal{L} = 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$)
- Background may be problematic

Luminosity calibration

(hadrons, e.g. pp or $p\bar{p}$)

- Have to measure beam current and beam sizes
- Beam size measurement:
 - wire scanner or synchrotron light monitors
 - measurement with beam ... → remember luminosity with offset
 - move the two beams against each other in transverse planes (van der Meer scan)
 - already used at Intersecting Storage Rings (ISR)

Luminosity optimization



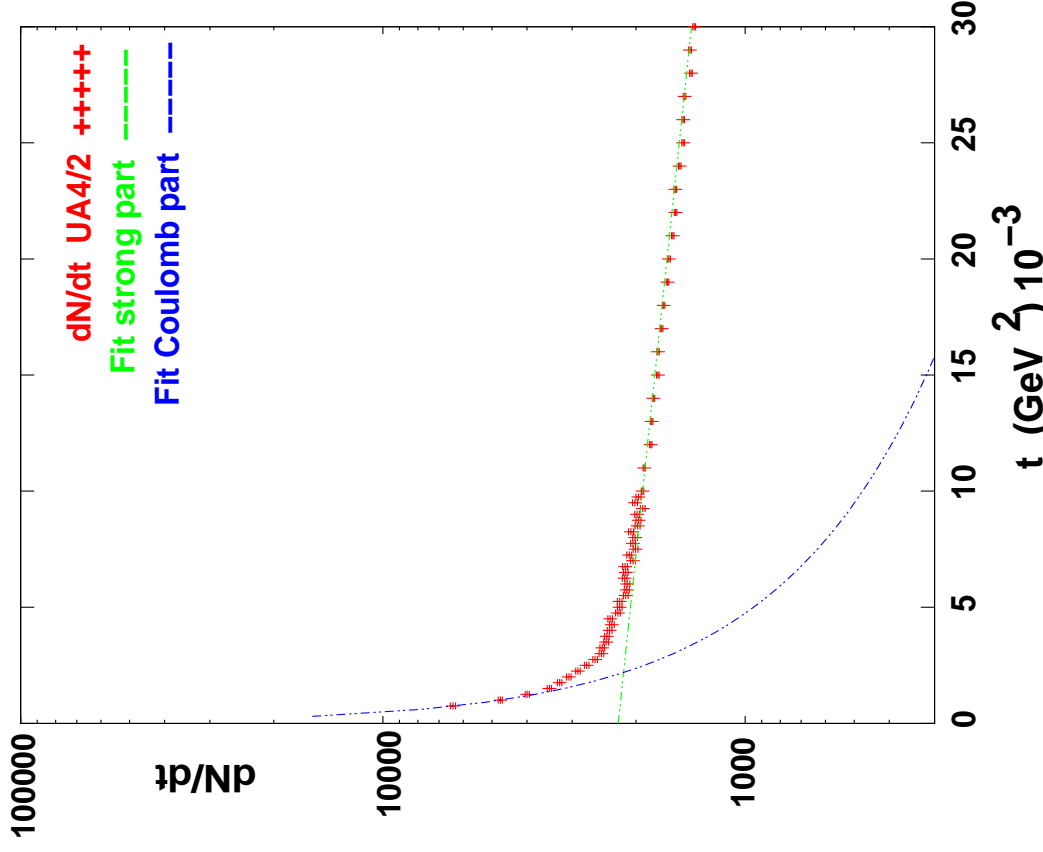
- Remember: $W = e^{-\frac{1}{4\sigma^2}(d_2 - d_1)^2}$
- \rightarrow ratio of luminosity $\mathcal{L}(d)/\mathcal{L}_0$
- Determines σ
- ... and centres the beams !
- ”beam-beam deflection scans ...”



Absolute value of \mathcal{L} (pp or $p\bar{p}$)

- By total rate and optical theorem (luminosity independent determination of σ_{tot}):
 - $\sigma_{tot} \cdot \mathcal{L} = N_{inel} + N_{el}$ (Total counting rate)
 - $\lim_{t \rightarrow 0} \frac{d\sigma_{el}}{dt} = (1 + \rho^2) \frac{\sigma_{tot}^2}{16\pi} = \frac{1}{\mathcal{L}} \frac{dN_{el}}{dt} \Big|_{t=0}$
 - $\mathcal{L} = \frac{(1+\rho^2)}{16\pi} \frac{(N_{inel} + N_{el})^2}{(dN_{el}/dt)_{t=0}}$
- Luminosity determined from experimental rates

Differential elastic cross section



Measure dN/dt at small t ($0.01 < \text{GeV}^2$) and extrapolate to $t = 0.0$

Needs special optics to allow measurement at very small t

Measure total counting rate $N_{el} + N_{inel}$
Needs good detector coverage

Often use slightly modified method, precision 1 – 2 %

Absolute value of \mathcal{L} (pp or $p\bar{p}$)

- By Coulomb normalization:
- Coulomb amplitude exactly calculable:
 - $\lim_{t \rightarrow 0} \frac{d\sigma_{el}}{dt} = \frac{1}{\mathcal{L}} \left. \frac{dN_{el}}{dt} \right|_{t=0} = \pi |f_c + f_N|^2$
 - $\simeq \pi \left| \frac{2\alpha_{em}}{-t} + \frac{\sigma_{tot}}{4\pi} (\rho + i) e^{B \frac{t}{2}} \right|^2 \simeq \frac{4\pi\alpha_{em}^2}{t^2} \Big|_{|t| \rightarrow 0}$
 - Fit gives: σ_{tot}, ρ, B and \mathcal{L}
- Can be done measuring **only** elastic scattering
(No N_{inel} needed !)
- \implies Roman Pots

Not treated :

- Coasting beams (e.g. ISR)
- Asymmetric colliders (e.g. PEP, HERA)
- Linear colliders (SLC, TESLA etc.)

How to cook high Luminosity ?

- Get high intensity
- Get small beam sizes (small ϵ and β^*)
- Get short bunches
- Get many bunches
- Get exact head-on collisions