

Plasma Diagnostics of Ion Sources

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Diagnostics – The Window to the Knowledge



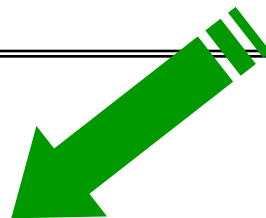
- ▶ Langmuir probes: ϕ_{pl}, n_e, T_e
- ▶ Absorption techniques: $n_s \rightarrow Cs, H^-$
- ▶ Emission spectroscopy: $n_s, T_s \rightarrow e, H, H_2, H^-$

Monitoring and Quantification – Spatial and Temporal Resolution



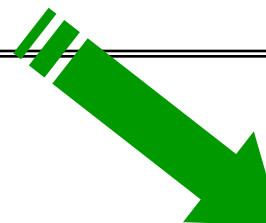
What do I want to know?

- ▶ identify the quantity (and the reason for it)
- ▶ define the required precision
- ▶ temporal behaviour and required time resolution
- ▶ necessity for spatial resolution
- ▶ ...



Adequate diagnostic technique?

- ▶ extensive or simple setup
- ▶ data acquisition and evaluation
- ▶ reliability
- ▶ costs and time (manpower) needed
- ▶ ...



Accessibility of the ion source?

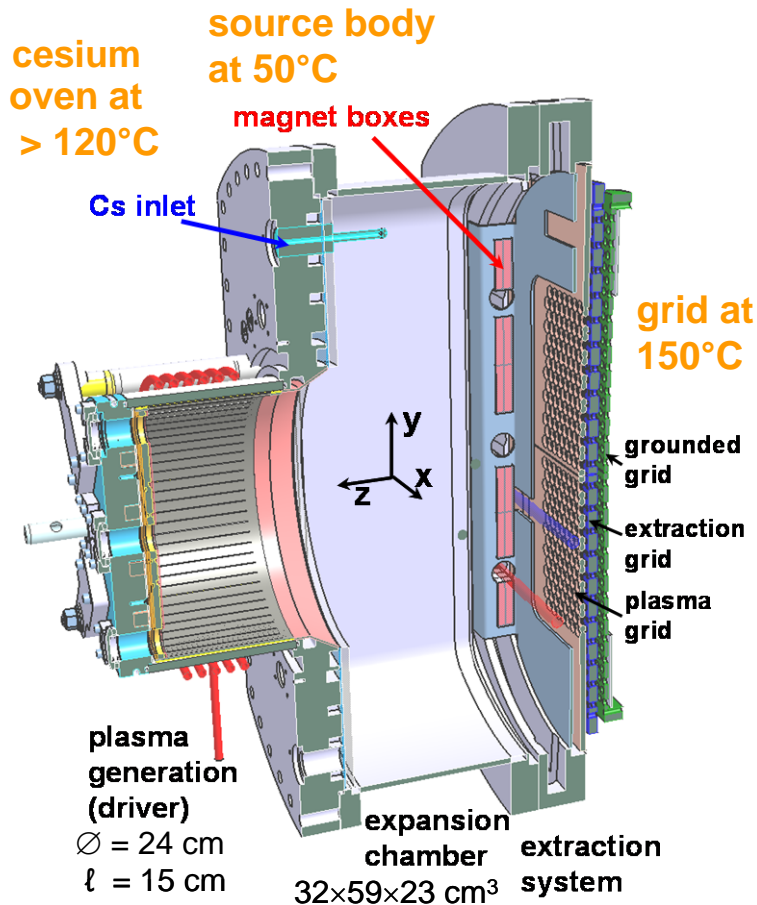
- ▶ diagnostic ports
- ▶ test stand or continuous operation
- ▶ risks and feasibility
- ▶ reliability
- ▶ ...

invasive – non-invasive ; active – passive ; basic – specific parameters

Method	Standard	Sophisticated	Extras
Langmuir probe	single probe (cylindrical or planar)	double or triple probe emissive probe	special method: Boyd-Twiddy
emission spectroscopy	optical wavelength range with fibre optics & survey spectrometer	extended wavelength range VUV, UV or IF	sophisticated system spectral resolution, type of detector
absorption spectroscopy	white light absorption technique	tuneable laser absorption	cavity-ringdown spectroscopy
Laser methods	laser induced fluorescence	TDLAS	
mass spectrometry	residual gas analyser	energy resolved mass spectrometry	

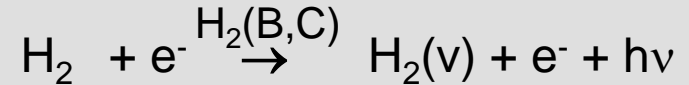
Introduction into techniques and applications for example in [1] – [7]

ICP: $f = 1$ MHz, $P = 70$ kW, $p = 0.3$ Pa
 6s plasma (4s beam), every 3 min: BATMAN
 cw, up to 1 hour, every 3 min: MANITU

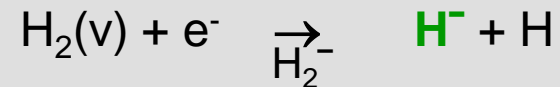


H^- formation and losses ...

volume process



followed by



dissociative attachment

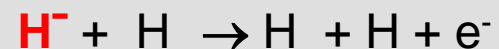
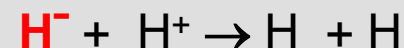
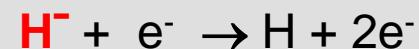
surface process



low work function

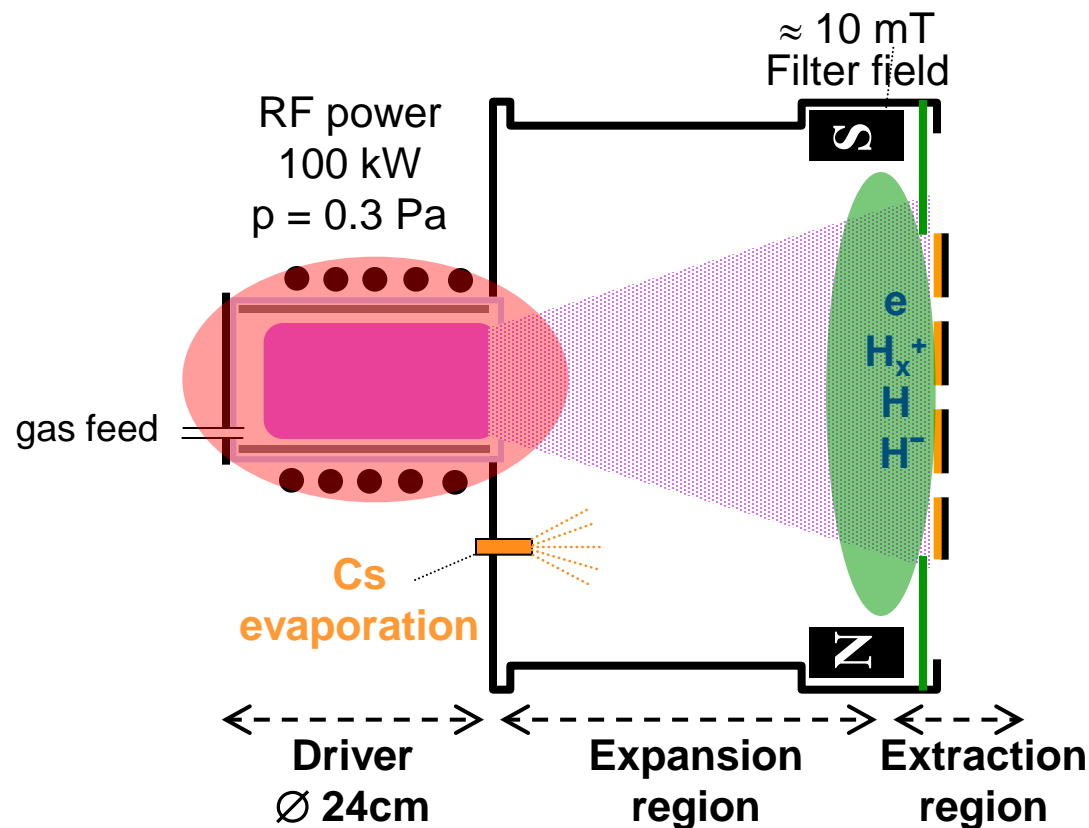
caesium layer

destruction processes



... determine source optimisation

Ion sources for negative hydrogen ions: ionising – recombining plasma



$H, H_x^+ + \text{surface } e^- \rightarrow H^-$
Cs evaporation \rightarrow low work function

Plasma generation ionising plasma

- ▶ ionisation: $\alpha \approx 0.1$
- ▶ dissociation: $\delta \approx 0.3$
- ▶ $T_e \approx 10 \text{ eV}$, $n_e \approx 5 \times 10^{18} \text{ m}^{-3}$

H^- generation recombining plasma

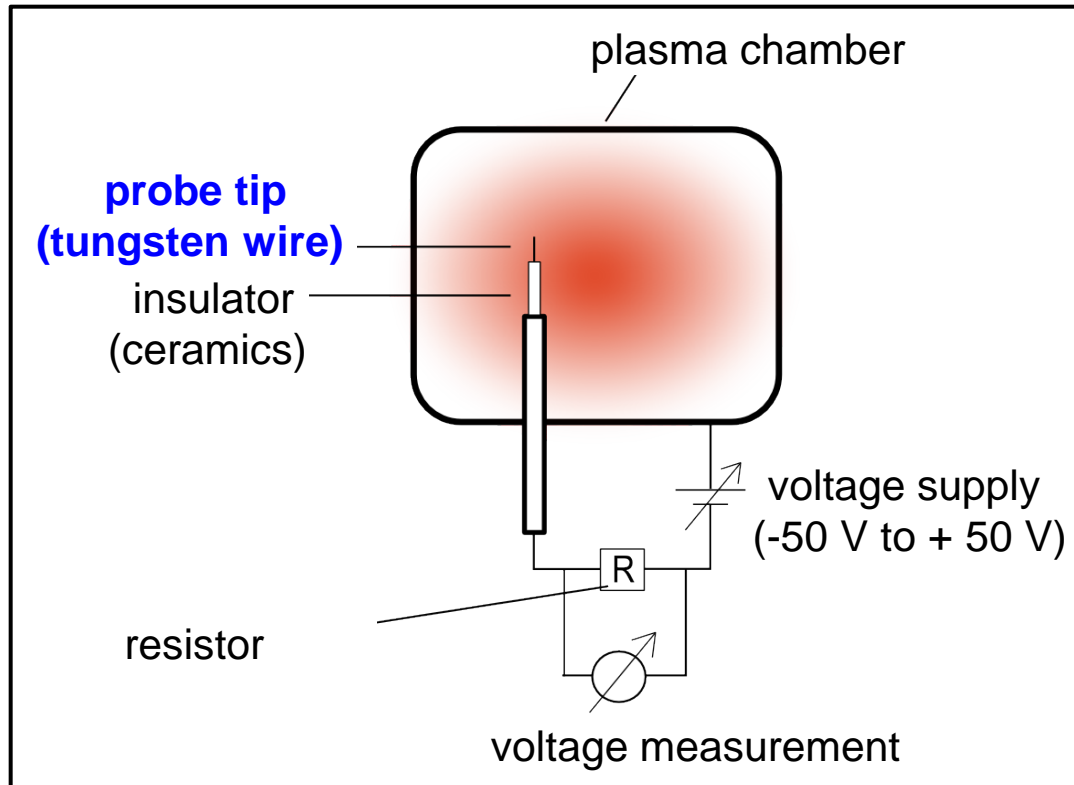
- ▶ $T_e \approx 1 \text{ eV}$, $n_e \approx 5 \times 10^{17} \text{ m}^{-3}$
- ▶ $H^-/n_e \approx 0.1 - 5$
- ▶ $Cs^+/n_e \approx 0.01 - 0.1$

Main issues

- ▶ Production and destruction of negative ions
- ▶ Extraction of negative ions
- ▶ Reduction of co-extracted electrons

focus of diagnostics
(and modelling)
on plasma close to the grid

Main principle



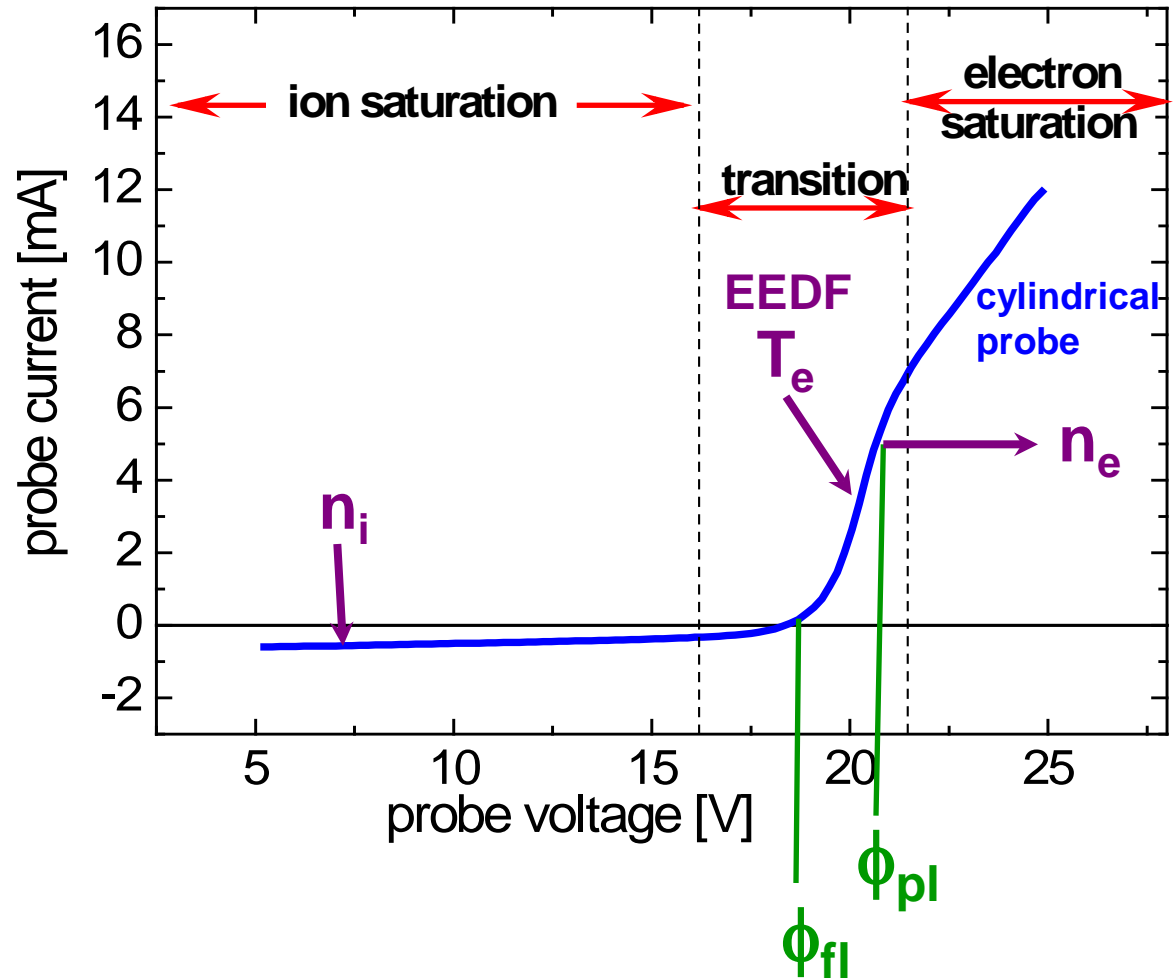
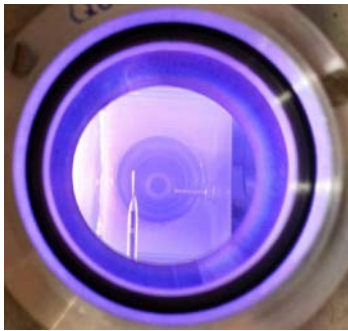
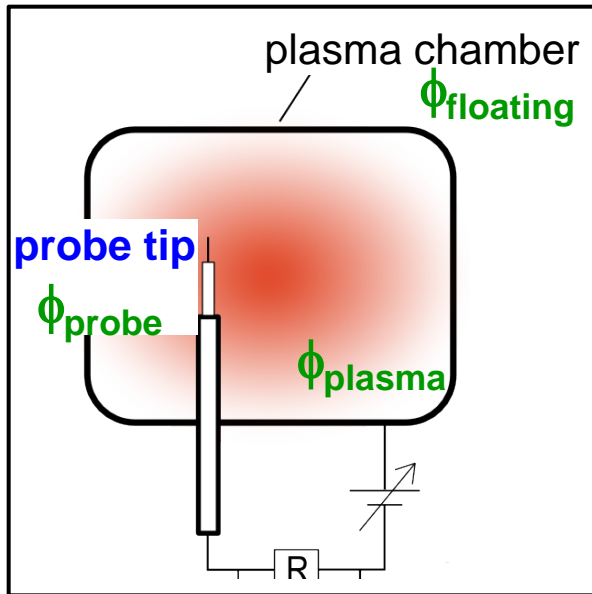
- ▶ stick a wire into the plasma
tungsten, $\varnothing \approx 100 \mu\text{m}$, $l = 1 \text{ cm}$
- ▶ choose a reference electrode
potential of plasma chamber
- ▶ apply a variable voltage
typ. from - 50 V to + 50 V



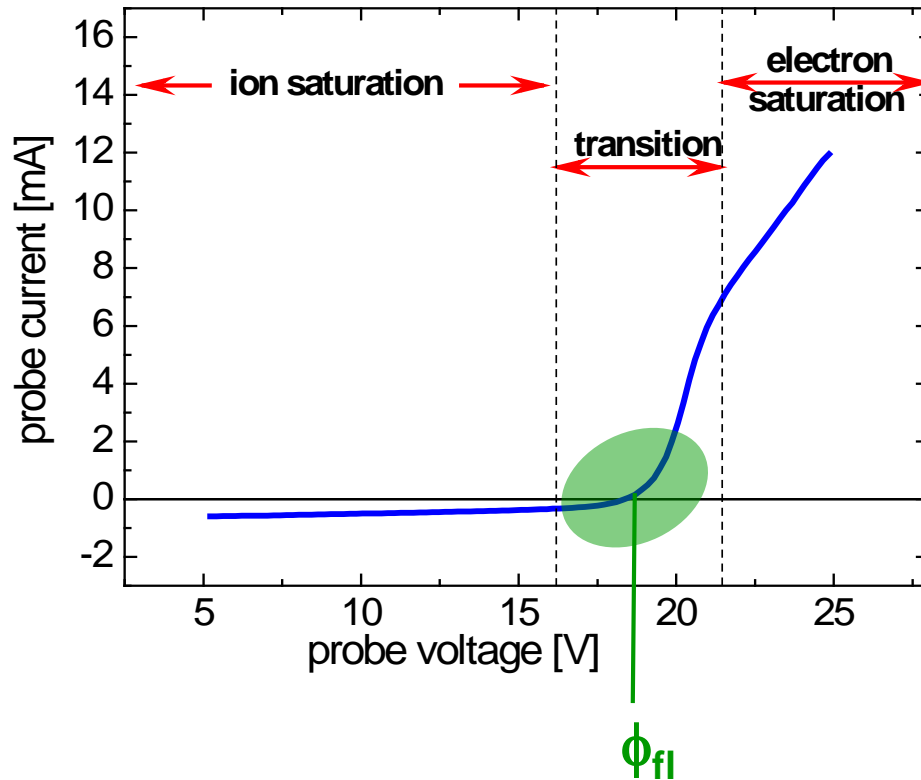
**I - V
characteristics**

Recommended text books [1], [2], [5], [6] & publications and programs by F.F. Chen [8]

Main principle



Floating potential



- ▶ same fluxes for ions and electrons

$$\Gamma_{\text{ions}} = \Gamma_{\text{electrons}}$$

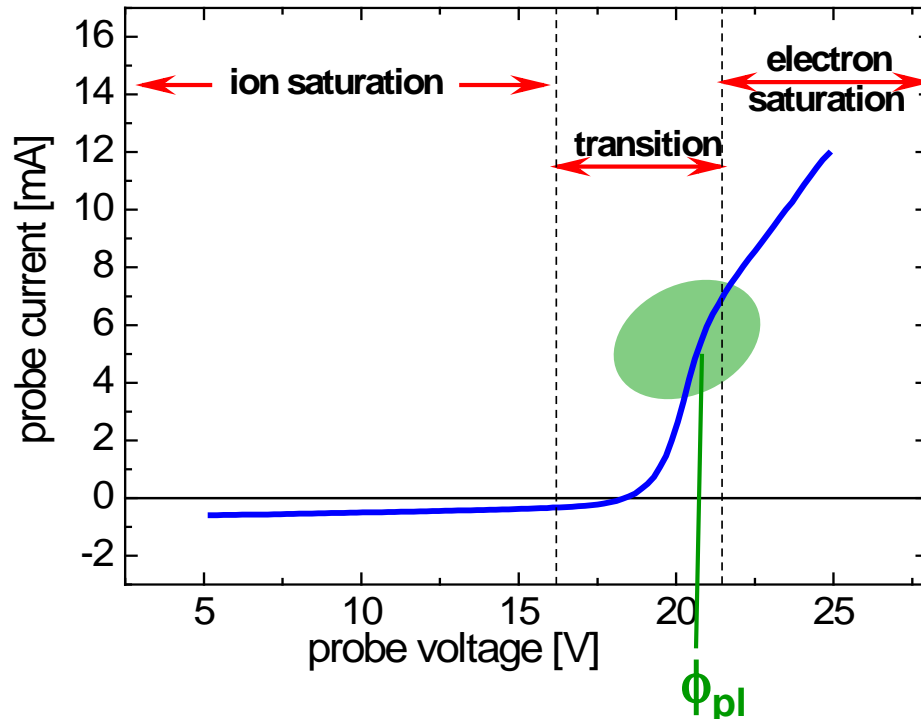
same currents

- ▶ no probe current

↓

$$I_{\text{probe}} = 0 \rightarrow \phi_{fl}$$

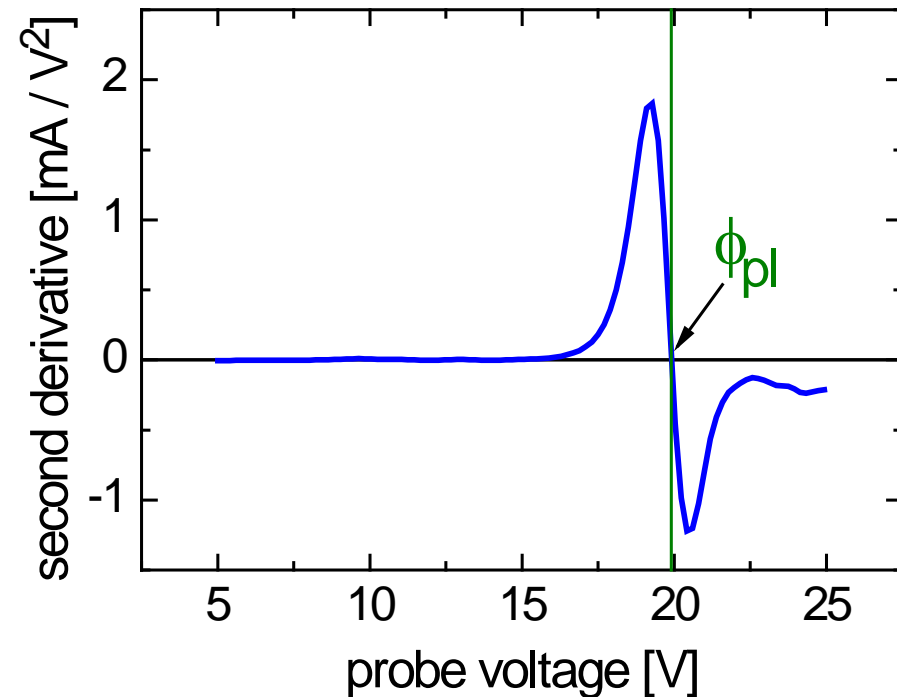
Plasma potential



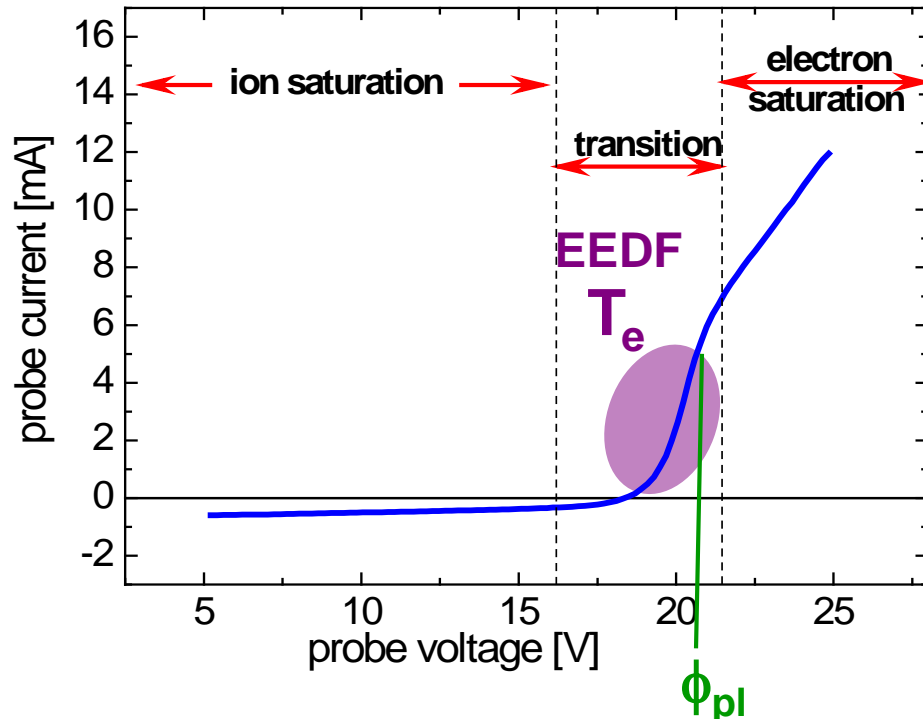
- ▶ for curves with high noise level crossing of linear fits to electrons saturation current and transition close to turning point

- ▶ determined by ambipolar diffusion
- ▶ turning point of I-V characteristics
- ▶ zero-crossing of second derivative

$$d^2 I / d \phi_{pr}^2 \rightarrow \phi_{pl}$$



Electron energy distribution function (EEDF) and electron temperature

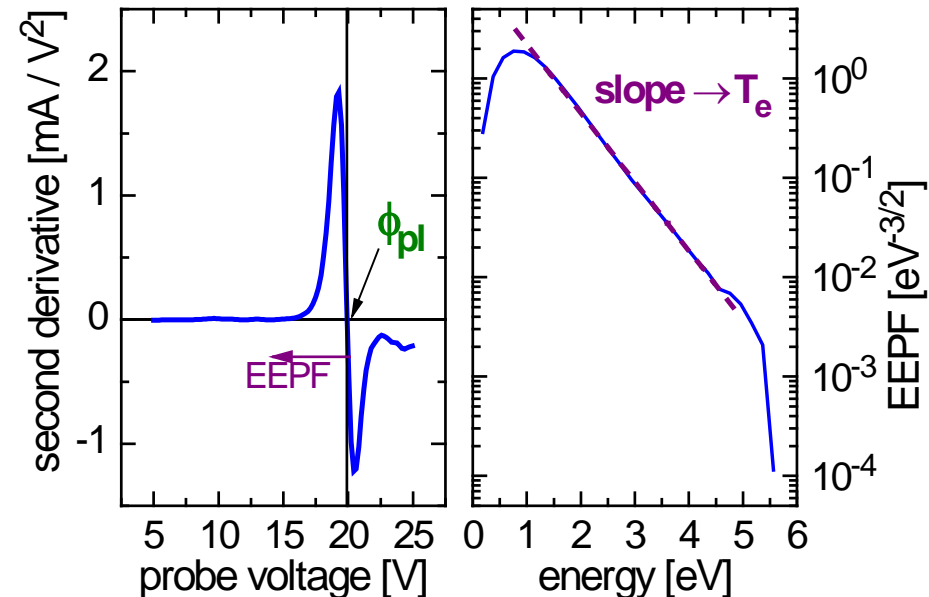


- ▶ T_e also from potential difference

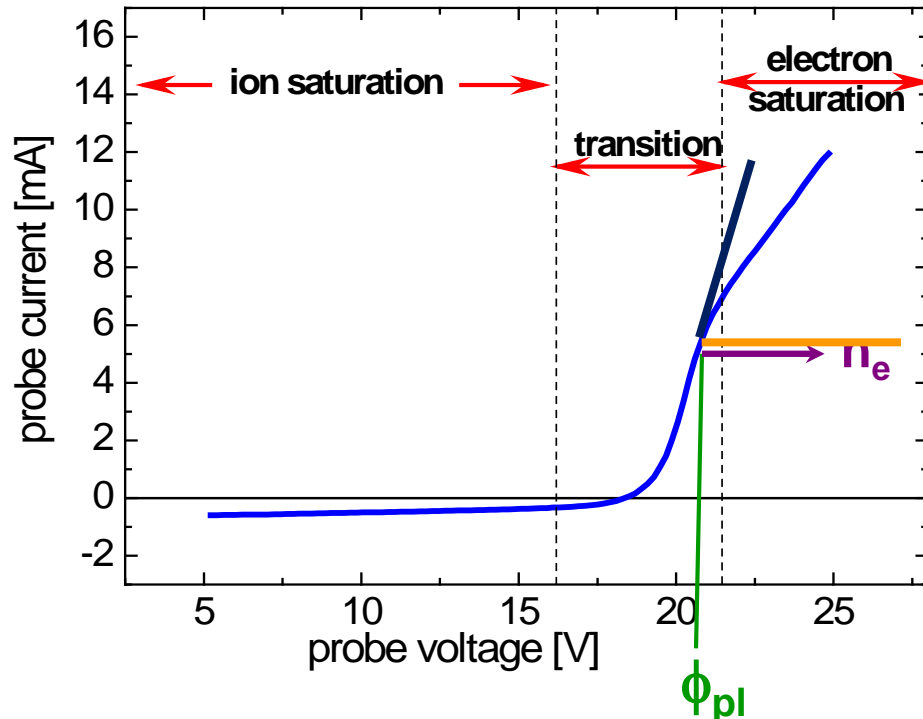
$$\phi_{pl} - \phi_{fl} = k_b T_e \times \ln(\sqrt{m_i / (2\pi m_e)})$$

$$\approx 2-3 \times T_e$$

- ▶ distinguish: **EEDF = sqrt(E) × EEPF**
e.g.: Maxwell function probability function
for both: normalisation
- ▶ plot $\log(I_{pr})$ versus $E = \phi_{pl} - \phi_{pr}$
- ▶ slope yields T_e for Maxwell EEDF
 $d(\ln I) / dE = e / (k_b T_e)$



The electron density



- ▶ electron saturation current

$$I_{e,sat} = \frac{1}{4} n_e e v_e A_{eff}$$

- ▶ problem: effective probe area due to increase of plasma sheath



- ▶ take current at plasma potential

$$\rightarrow A_{eff} = A_{probe}$$

$$n_e = \frac{I(\phi_{pl})}{r_{pr} l_{pr} e} \sqrt{\frac{m_e}{2\pi k_B T_e}}$$

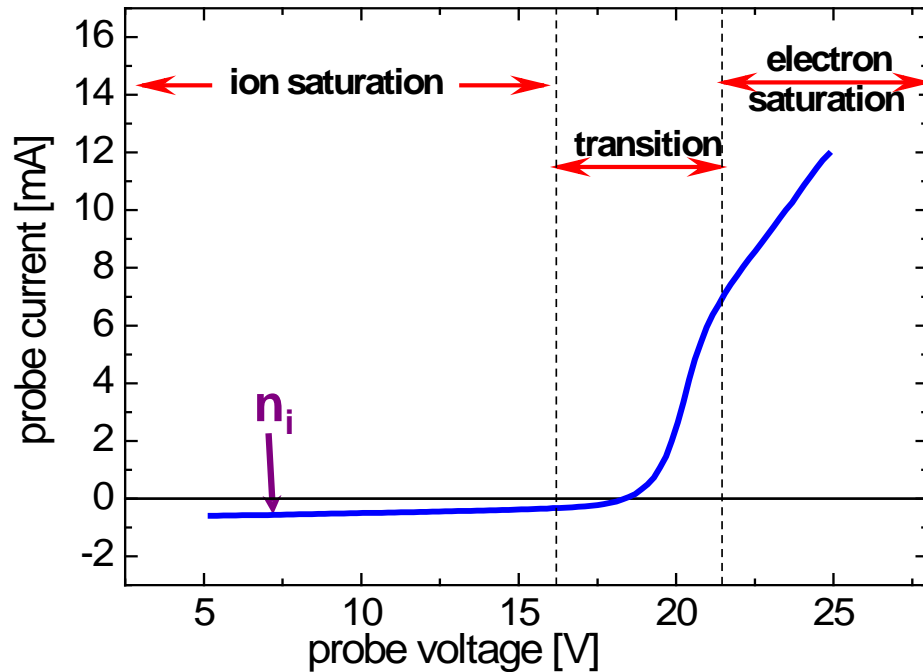
needs T_e

Probe geometry

influences shape of electron saturation current

- ▶ cylindrical probe (standard case)
- ▶ planar probe $\rightarrow A_{pr} \gg$ sheath
- ▶ spherical probe

Ion density (positive ions)



- ▶ choose proper ϕ_{pr}
guide line: $\phi = \phi_{pl} - 10 \times kT_e$
- ▶ check if $n_e = n_i$ is fulfilled
- ▶ $I_{sat,i}$ at fixed ϕ_{pr} : useful as monitor signal

- ▶ ion saturation current
- ▶ basically **three** theories available
OML: Orbital-Motion-Limited
ABR: Allen-Boyd-Reynolds
BRL: Bernstein-Rabinowitz-Langmuir

all of them assuming collision-less plasma sheath, i.e. $\lambda(\text{ions}) > r(\text{sheath})$

- ▶ simplest case: **OML** ($r_{pr}/r_{sheath} < 3$)

$$I_i = n_i e A_{pr} \sqrt{\frac{k_B T_e}{m_i}} \leftarrow \text{needs } T_e$$


needs m_i

often unclear, e.g. hydrogen

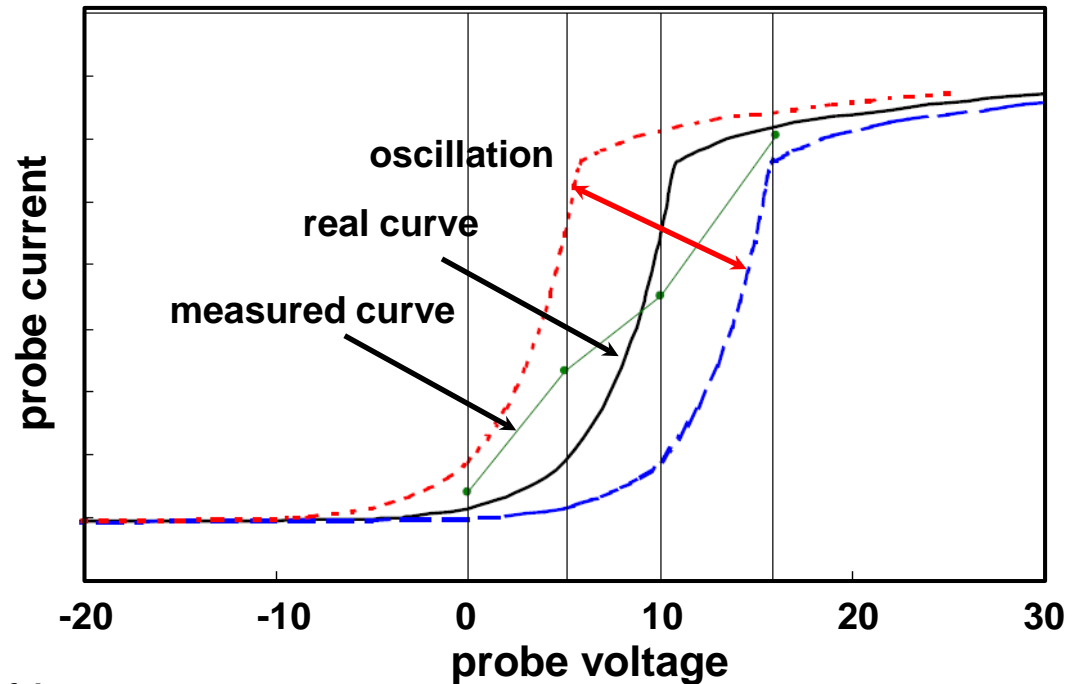
→ H^+ , H_2^+ , H_3^+

Specific features – to keep in mind

- ▶ **invasive method** → probe size versus plasma volume
- ▶ **level of noise** → EEPF for typically three orders of magnitude
average, smoothing, filtering
- ▶ **RF field** → oscillating ϕ_{pl}
measured curve \neq real curve

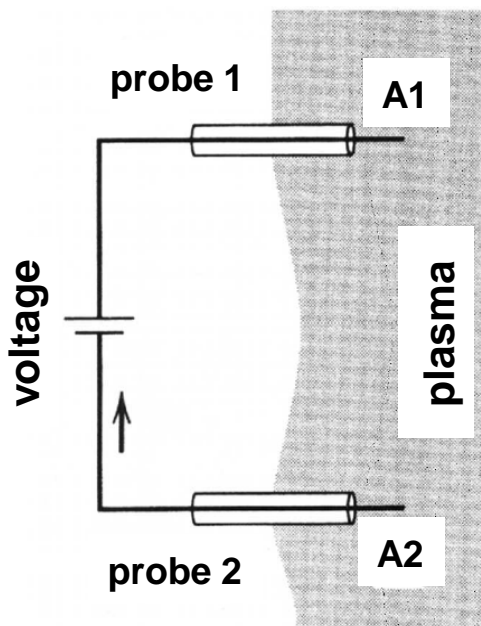


RF compensation
active or passive
- ▶ **magnetic field** → gyro motion
use I_{ion} instead of I_e
- ▶ **negative ions** → $I_{neg.ion}$ instead of I_e
for same mass $I_{sat,ion} = I_{sat,neg.ion}$ → symmetric curves



For monitoring or quantification with spatial and temporal resolution !

Double probe



- ▶ two probe tips of same size
distance > 2 Debye length
- ▶ voltage between the probes
both probes are floating
no reference potential needed

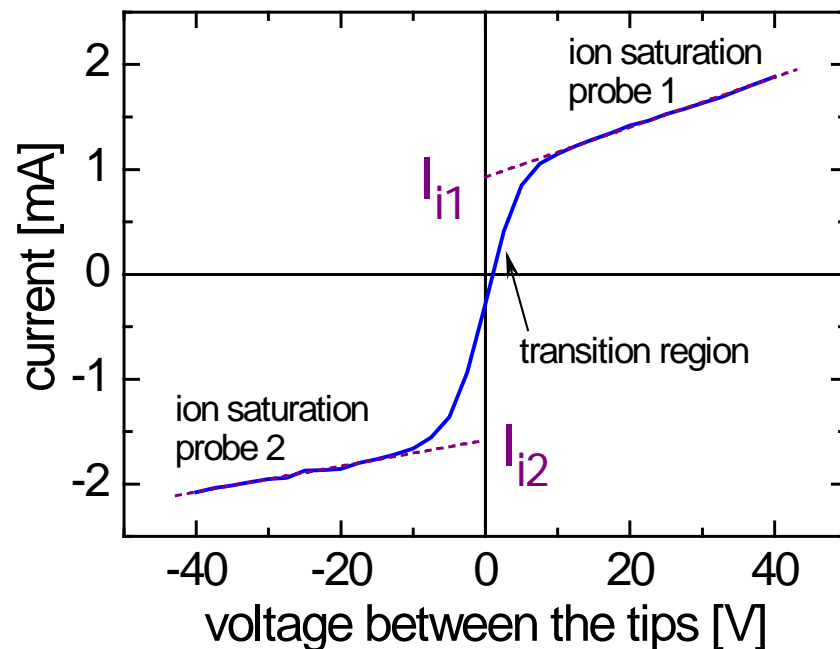


compatible with
quartz or ceramic
chamber
and RF field

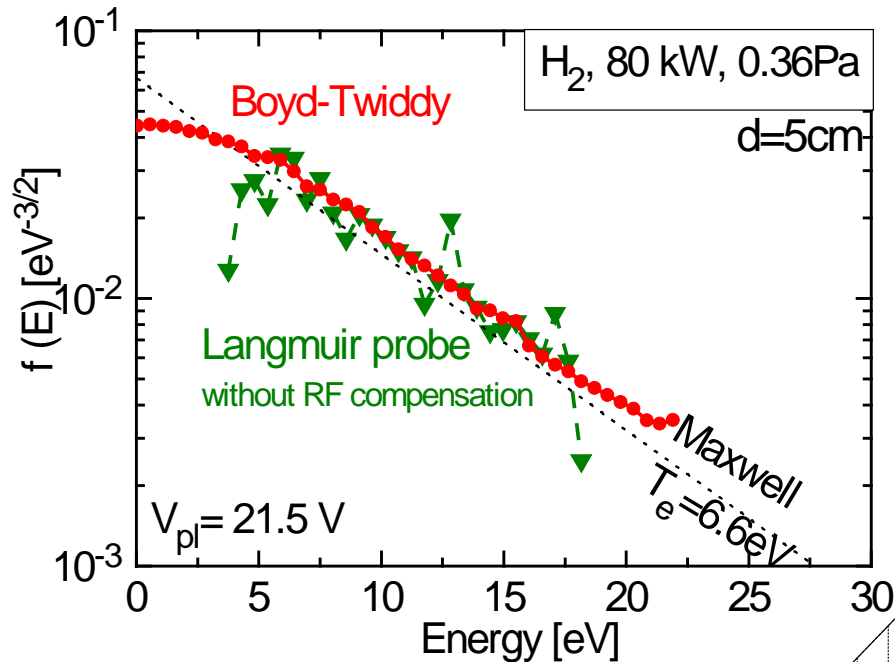
- ▶ symmetric curve: ion saturation current

$$I_{i,sat} = \frac{I_{isat,+} + I_{isat,-}}{2} \quad n_i = \frac{I_{i,sat}}{eA \sqrt{\frac{k_B T_e}{m_i}}}$$

- ▶ T_e from transition region



Comparison of EEDF



Sophisticated Langmuir probe system

for RF ion sources

McNeely et al.

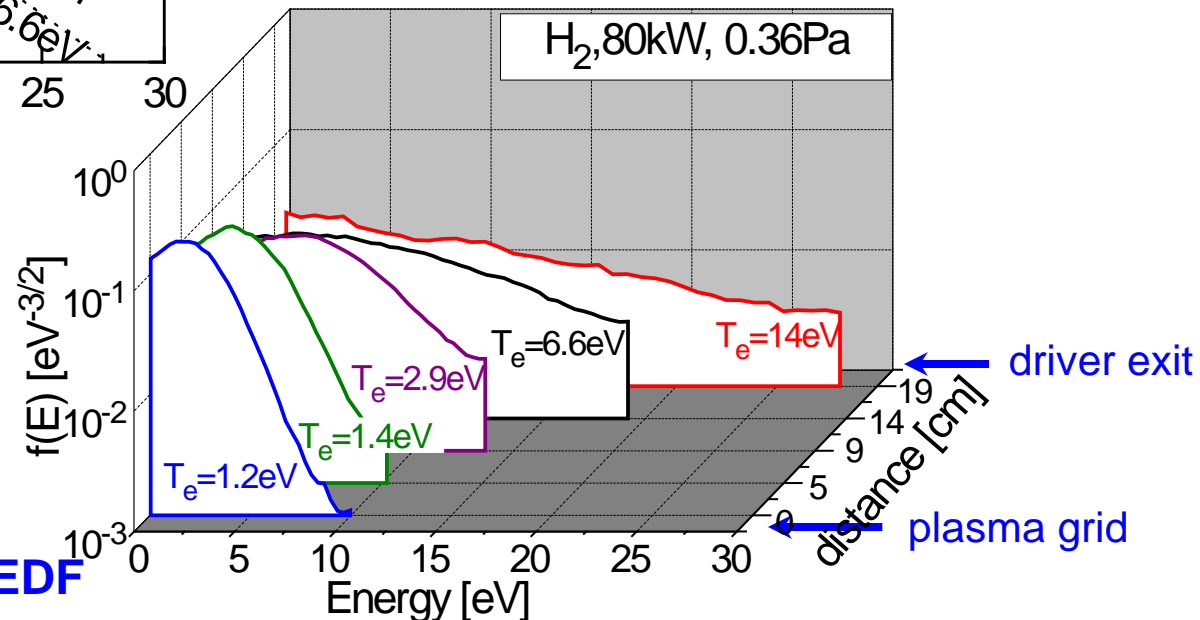
PSST 18 (2009) 014011

Maxwellian EEDF

Boyd-Twiddy method: direct measurement of EEDF

- ▶ voltage ramp is superimposed by AC modulated signal
- ▶ measure frequency spectrum of probe current

B. Crowley, S. Dietrich PSST 18 (2009) 014010



Radiation of low temperature plasmas

Colourful plasmas !



He
pink



Ne
red



N₂, air
orange

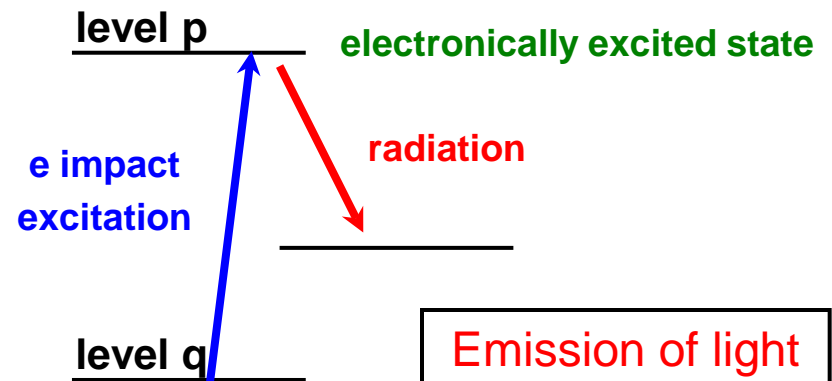


H₂, H
purple

- ▶ **Neutrals** atoms and molecules
- ▶ **Ions** single charged
- ▶ **Electrons** $n_e \ll n_n$



collisions and spontaneous emission



Emission of light
in the UV-VIS-IR

Emission versus absorption spectroscopy

Emission
passive

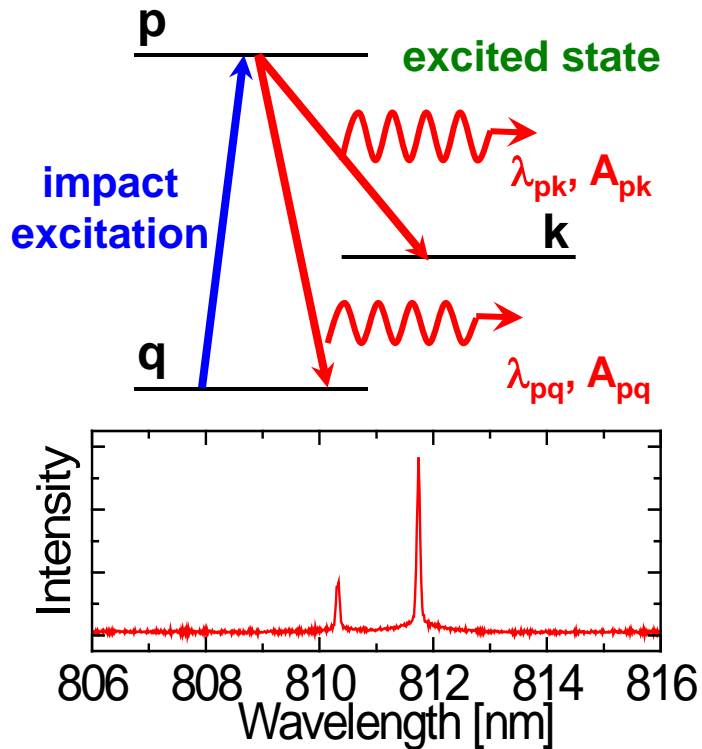
- ▶ photon energy
- ▶ wavelength
- ▶ Einstein coefficients

$$E = h\nu$$

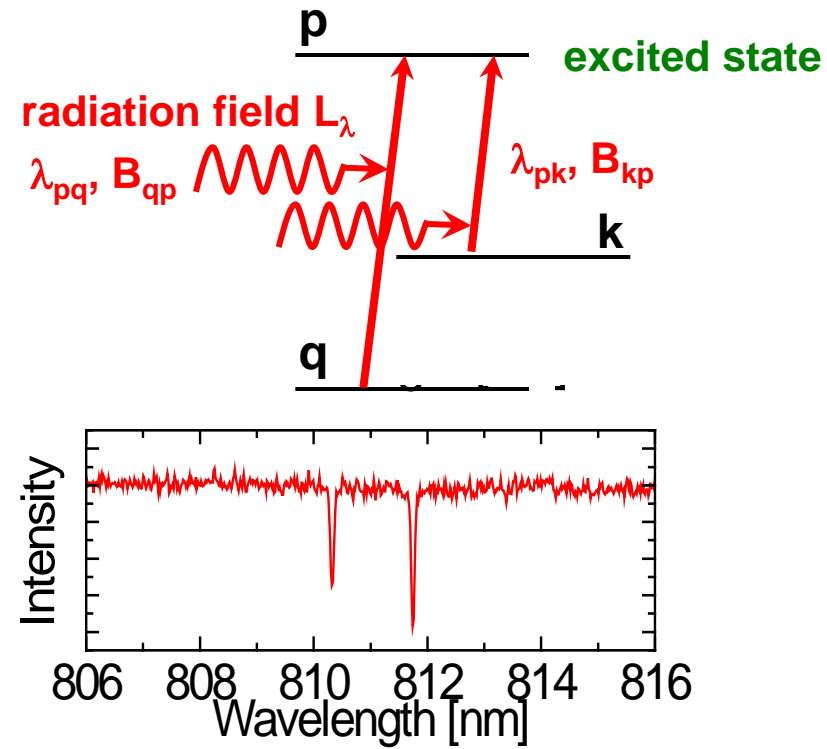
$$\lambda = (E_p - E_k) / hc$$

$$A_{pk}, B_{kp}$$

Absorption
active

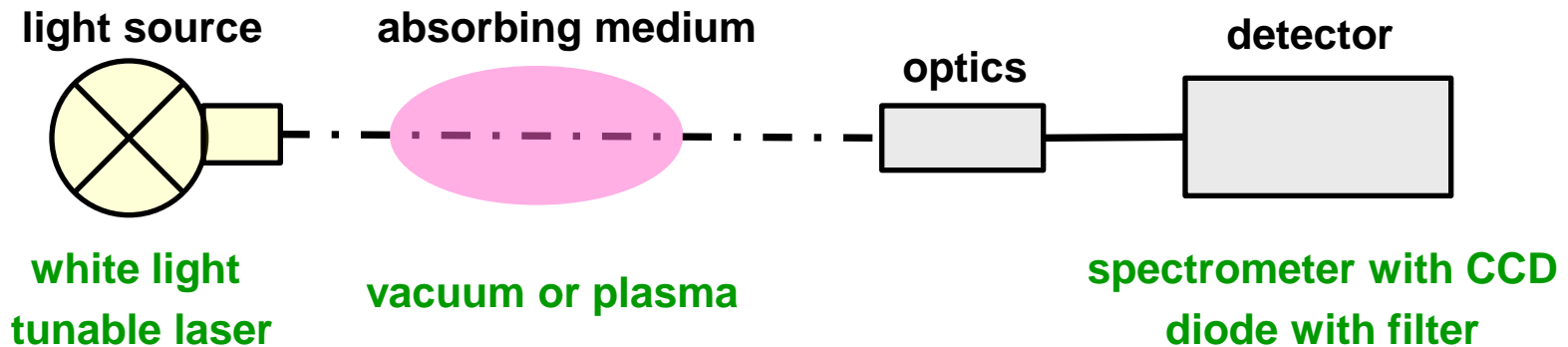


- ▶ VIS: simple equipment
- ▶ information of upper level p



- ▶ expensive equipment
- ▶ information of lower level q

Main principle of line absorption



Non-invasive and line of sight integrated method !

Absorption in a medium
with path length l

$$I(\lambda, l) = I(\lambda, 0) \exp[-\kappa(\lambda)l]$$

absorption coefficient $\kappa(\lambda)$
statistical weights g_i, g_k

$$k(\lambda) = \frac{1}{l} \ln \left(\frac{I(\lambda, 0)}{I(\lambda, l)} \right) \quad \text{with Ladenburg relation} \quad \int_{\text{line}} \kappa_{ki}(\lambda) d\lambda = n_k \frac{g_i}{g_k} \frac{\lambda_0^4}{c} \frac{A_{ik}}{8\pi}$$

Density of lower state
ground state, metastables

$$n_k = \frac{8\pi c}{\lambda_0^4} \frac{g_k}{g_i} \frac{1}{A_{ik} l} \int_{\text{line}} \ln \left(\frac{I(\lambda, 0)}{I(\lambda, l)} \right) d\lambda$$

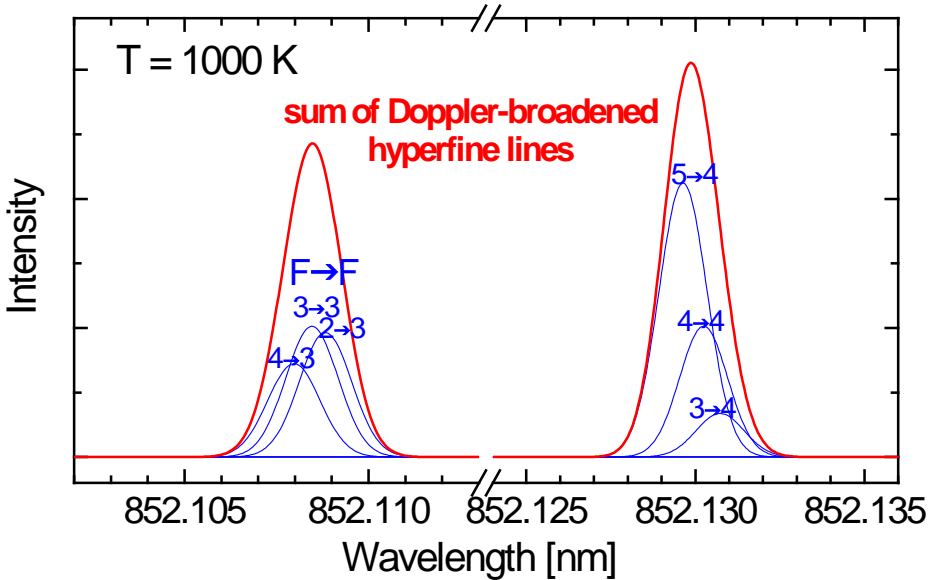
Example: Cs line at 852.1 nm

resonance line $6^2S_{1/2} - 6^2P_{3/2}$ with hyperfine structure

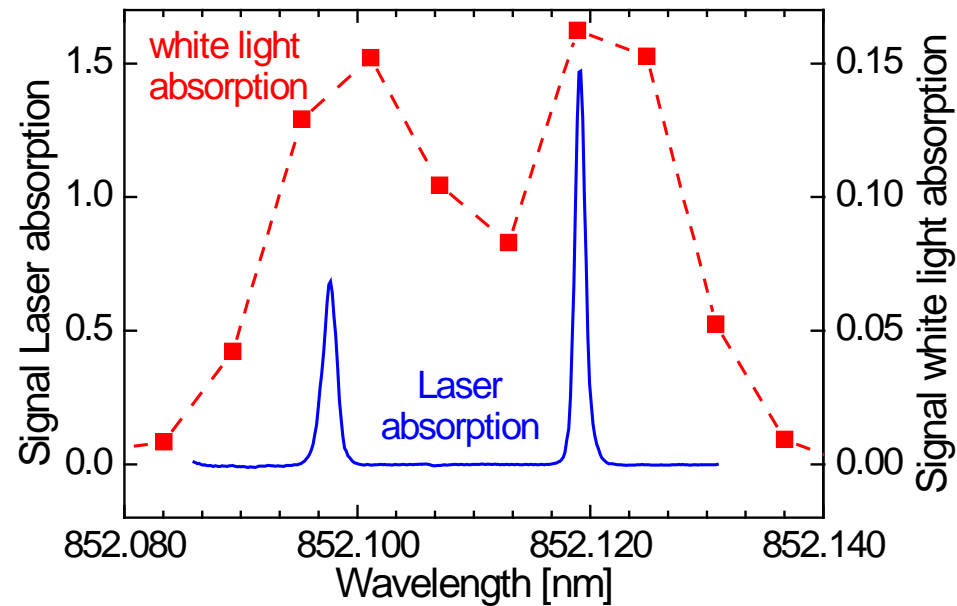
take line broadening into account

Doppler profile

$$\Delta\lambda_{D,FWHM} = \frac{\lambda}{c} \sqrt{\frac{8k_B T \ln 2}{m}}$$



appearance depends on technique



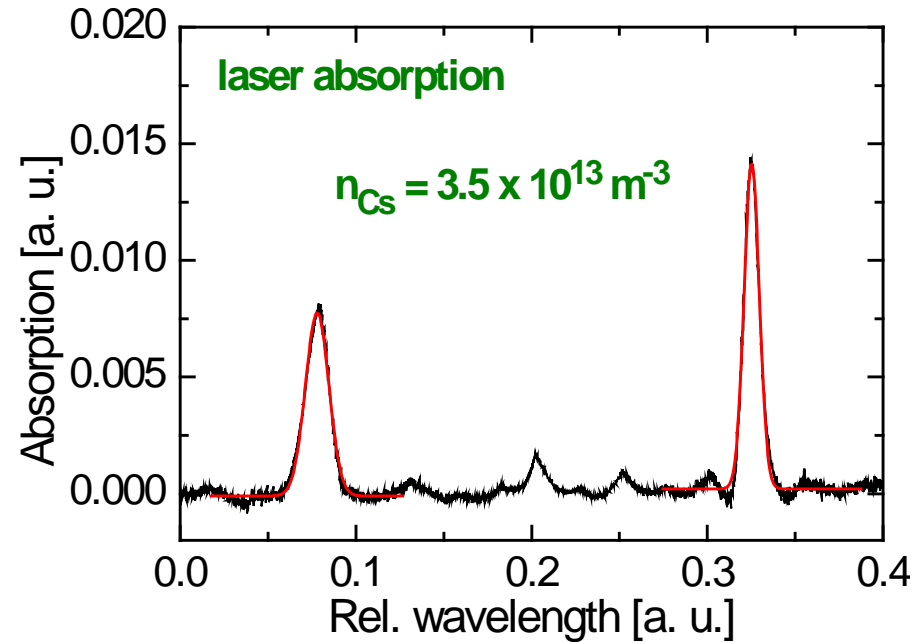
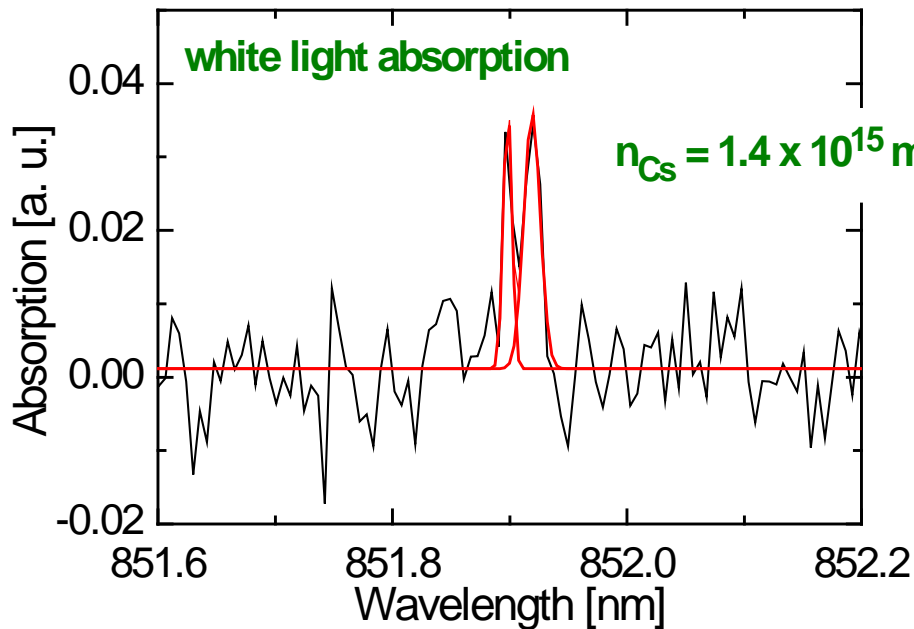
white light absorption \leftrightarrow laser absorption

apparatus profile (spectrometer) Doppler profile

six hyperfine lines overlap to two peaks: $\Delta\lambda \approx 21$ pm

Example: Cs line at 852.1 nm

white light absorption versus laser absorption

**Improved detection limit for laser absorption: factor > 40 !**

sensitivity range: $3 \times 10^{13} \text{ m}^{-3}$ - 10^{17} m^{-3} (path length = 15 cm)
being perfectly in the range required for the ion sources

Straightforward analysis

- ▶ in vacuum
- ▶ with plasma \rightarrow subtract emission

$$n_k = \frac{8\pi c}{\lambda_0^4} \frac{g_k}{g_i} \frac{1}{A_{ik} l} \int_{\text{line}} \ln \left(\frac{I(\lambda, 0)}{I(\lambda, l)} \right) d\lambda$$

... but ...

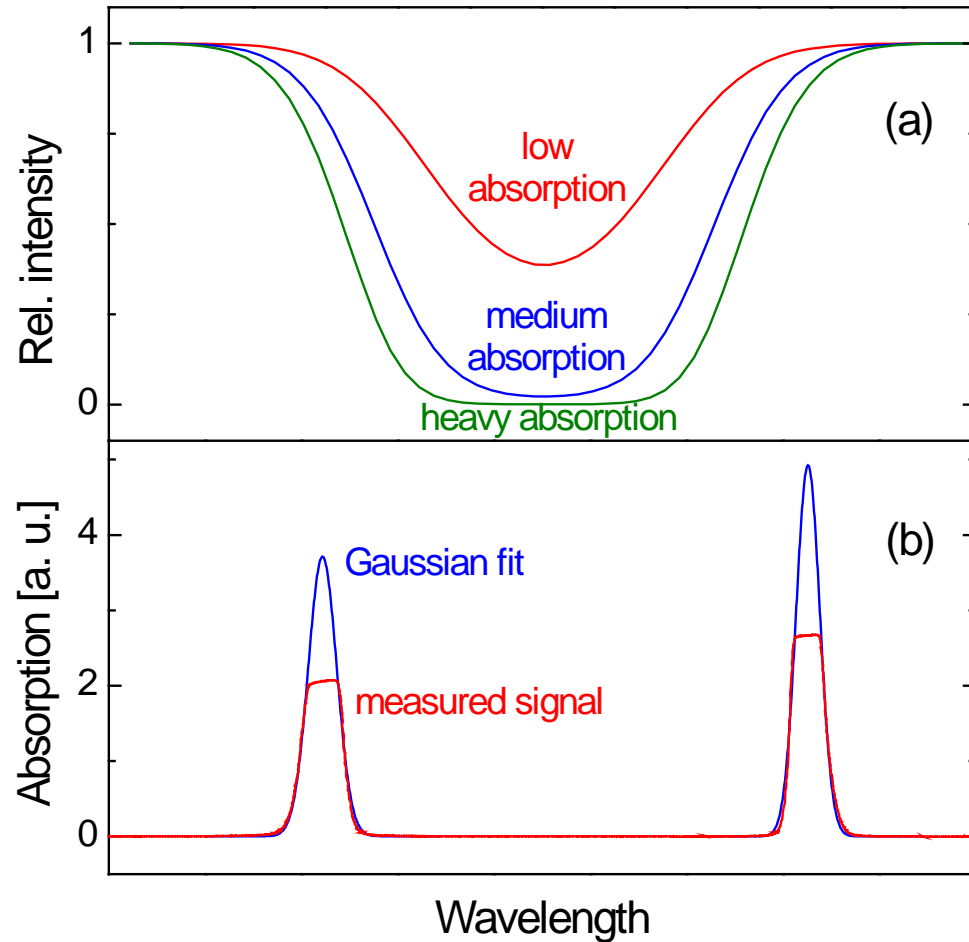
Line saturation

- ▶ strong absorption: $n_k \times l$
- ▶ correction factors by profile fitting

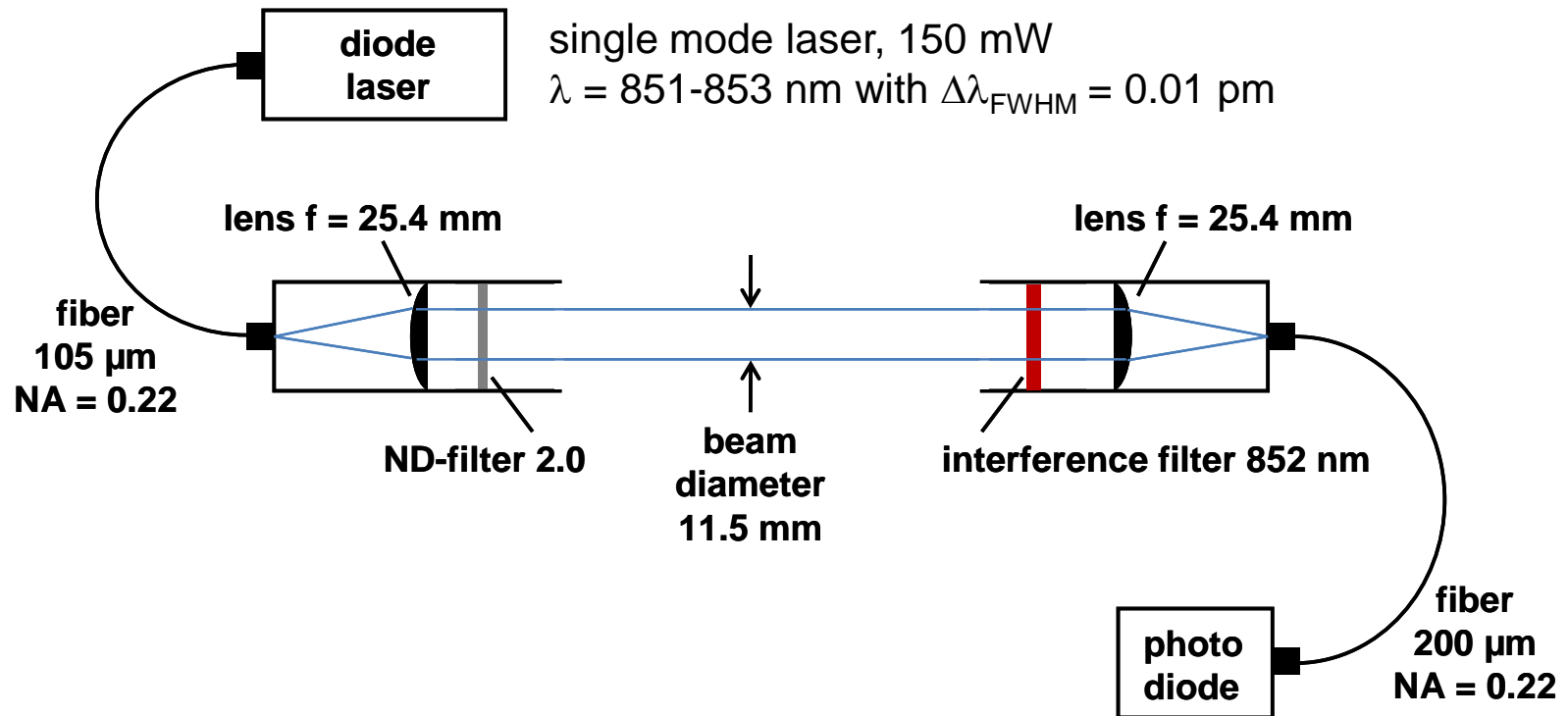
... and ...

Depopulation effect

- ▶ strong intensity
- ▶ attenuation of laser to $\approx 1\%$ \rightarrow trade-off with temporal resolution



Tunable diode laser – Fibre optics – Photo diode with interference filter

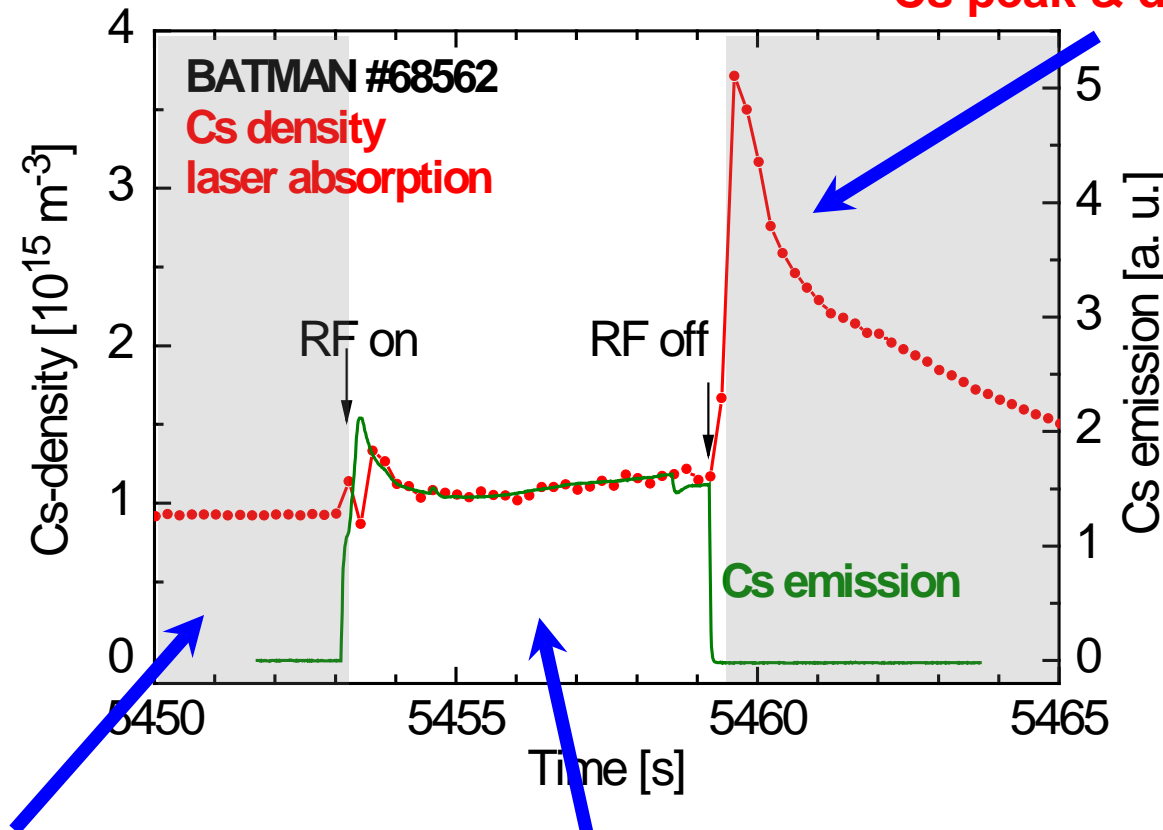


Simple and robust setup for application to ion sources !

On-line monitoring: vacuum – plasma (6s with 4s beam) - vacuum

③ 30 s after pulse:

Cs peak & decay time

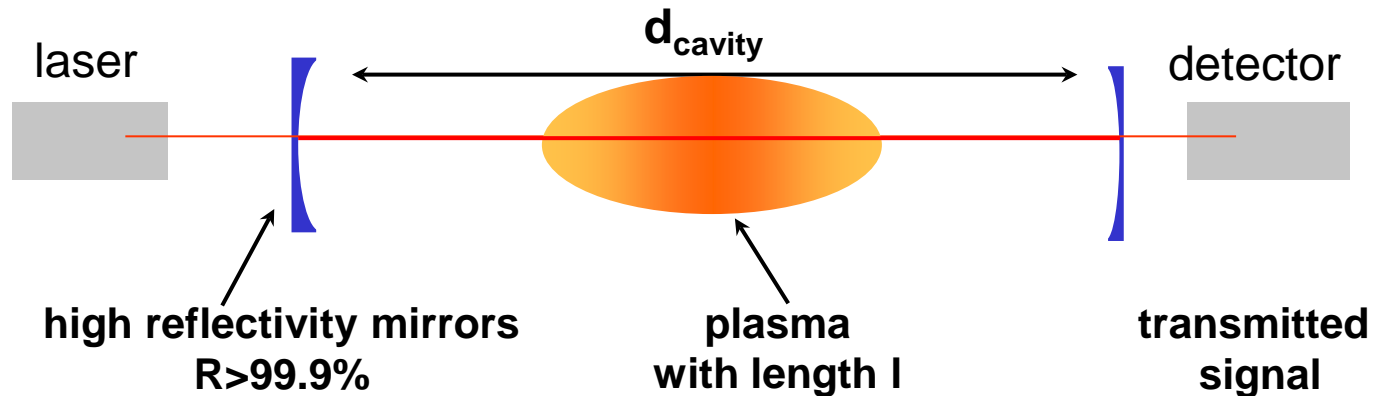


① 3 s before pulse: vacuum

② 6 s plasma with 4 s beam: plasma
90% of Cs is ionized

Cavity – Ringdown – Absorption – Spectroscopy → CRDS

pumping of an optical cavity by a (tunable) laser source
 measurement of signal decay after laser source is switched off

**Measurement of laser light attenuation trapped in a high-finesse optical cavity**

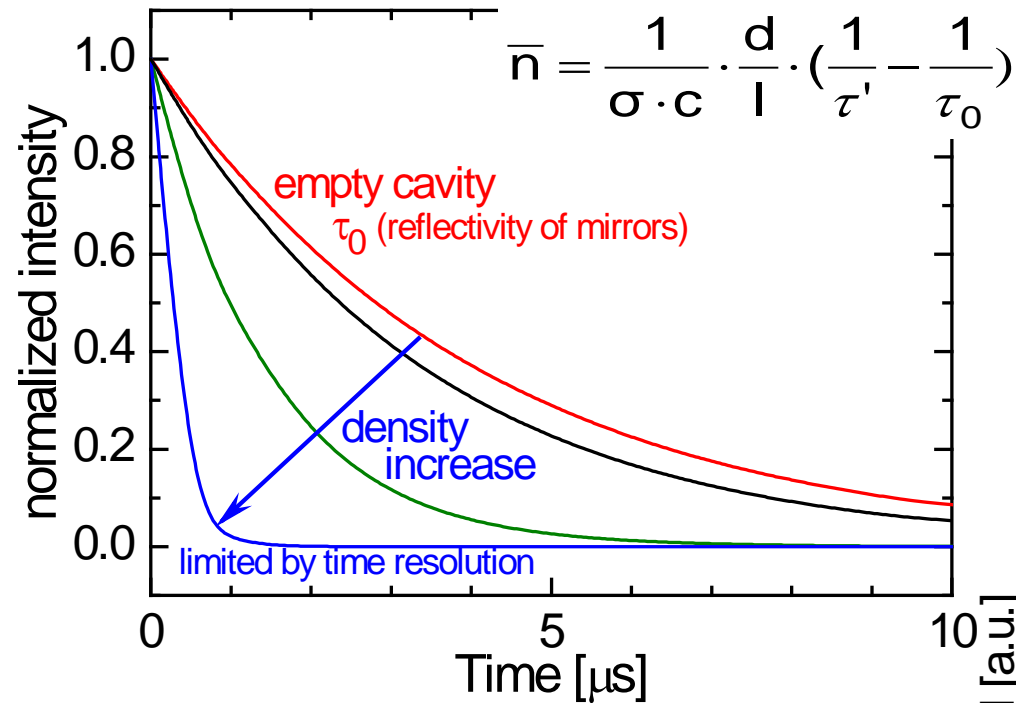
transfers absorption signal from wavelength into time dependence

empty cavity with decay time τ_0

$$I(t) = I_0 \cdot e^{-t/\tau_0} ; \quad \tau_0 = \frac{d}{c(1-R)}$$

additional absorption $\tau_0 \rightarrow \tau'$

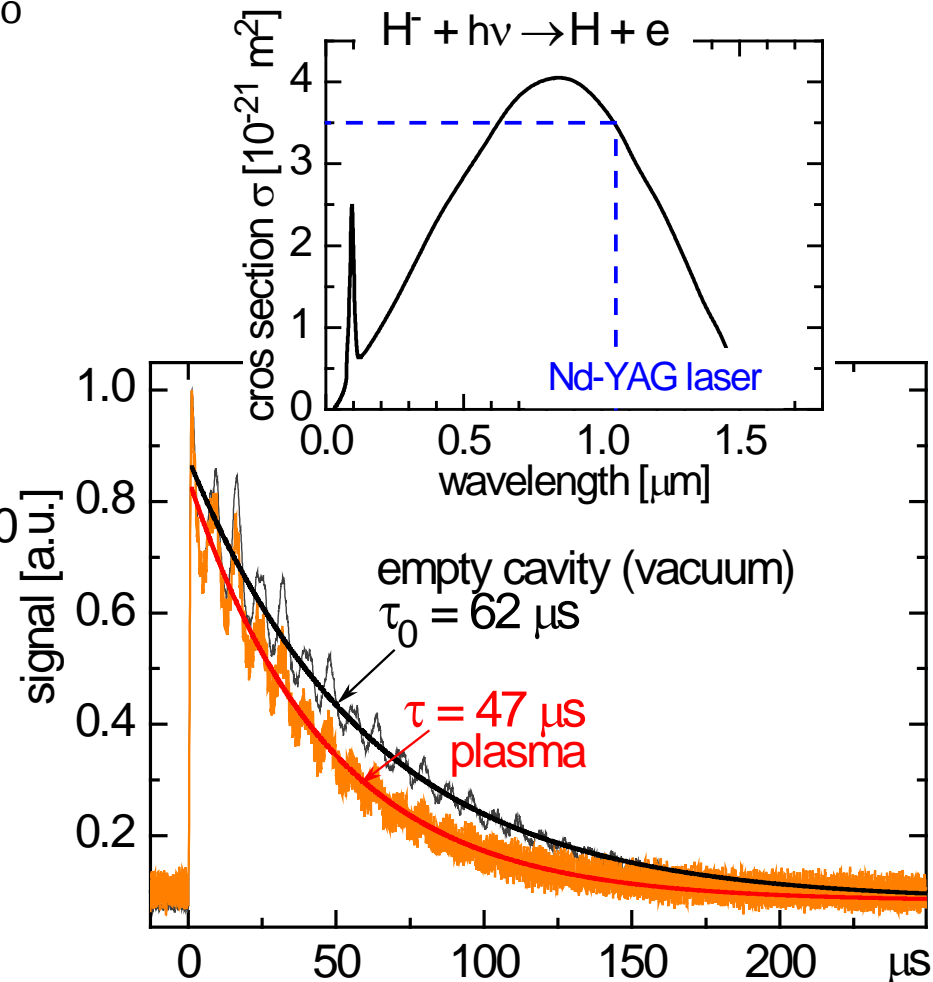
$$\bar{n} = \frac{1}{\sigma \cdot c} \cdot \frac{d}{l} \cdot \left(\frac{1}{\tau'} - \frac{1}{\tau_0} \right)$$

Cavity – Ringdown – Absorption – Spectroscopy \rightarrow CRDS H^- density: line of sight averaged

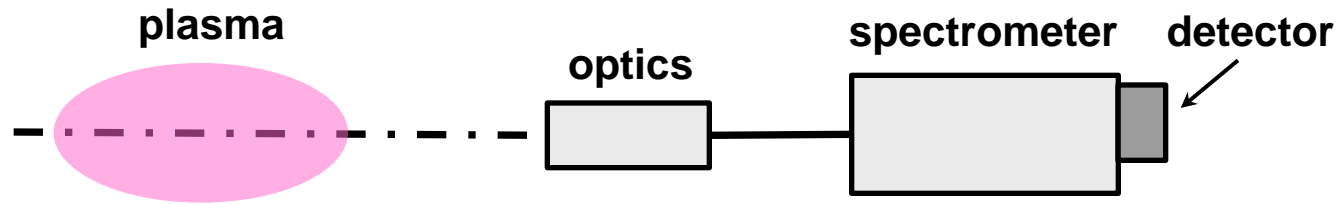
- ▶ detection limit $\approx 10^{15} \text{ m}^{-3}$
- ▶ $\text{H}^- = 5 \times 10^{17} \text{ m}^{-3} \rightarrow t = 8 \mu\text{s}$

Example H^-

cross section: photodetachment



The main principle



UV and VUV: vacuum system
IF: detector

Non-invasive and line of sight integrated method !

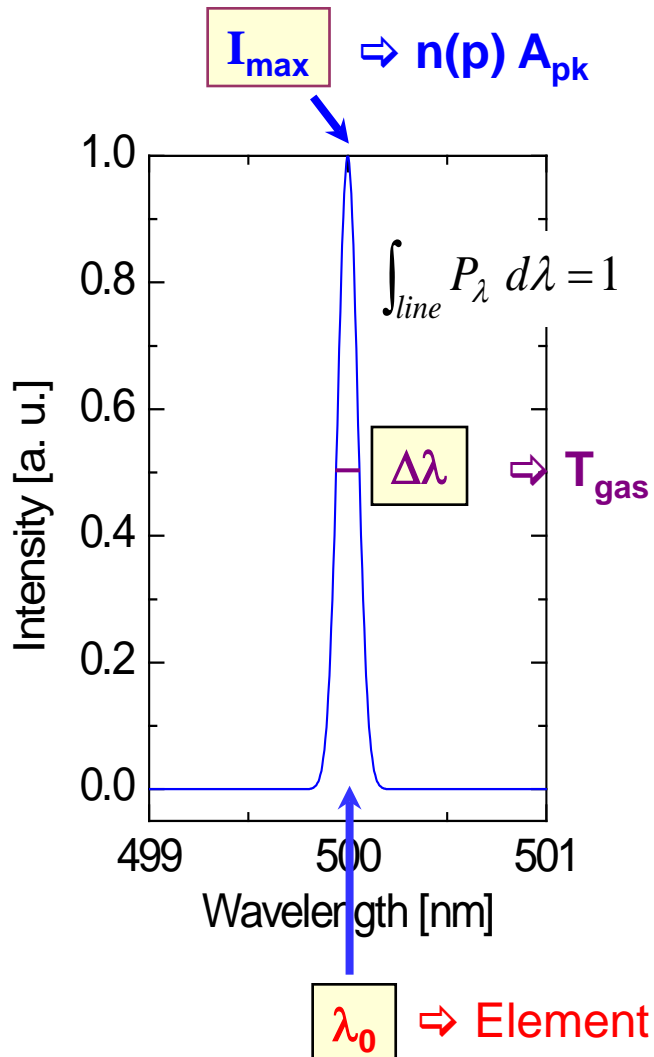
Measures density of excited state ...

$$\varepsilon_{pk} = n(p) A_{pk} \frac{hc / \lambda}{4\pi}$$

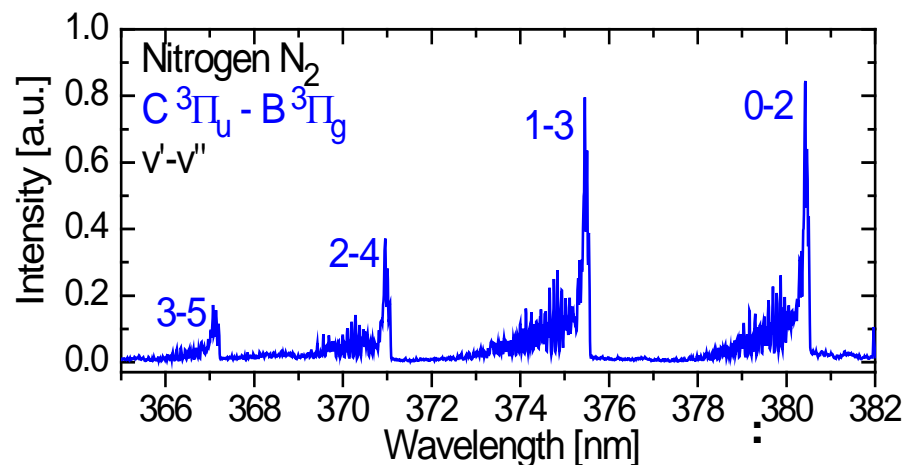
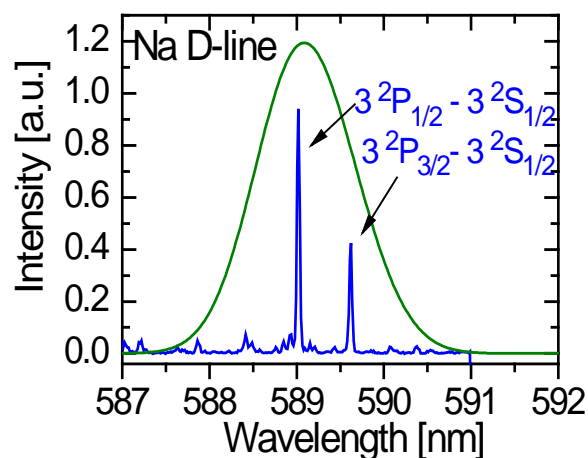
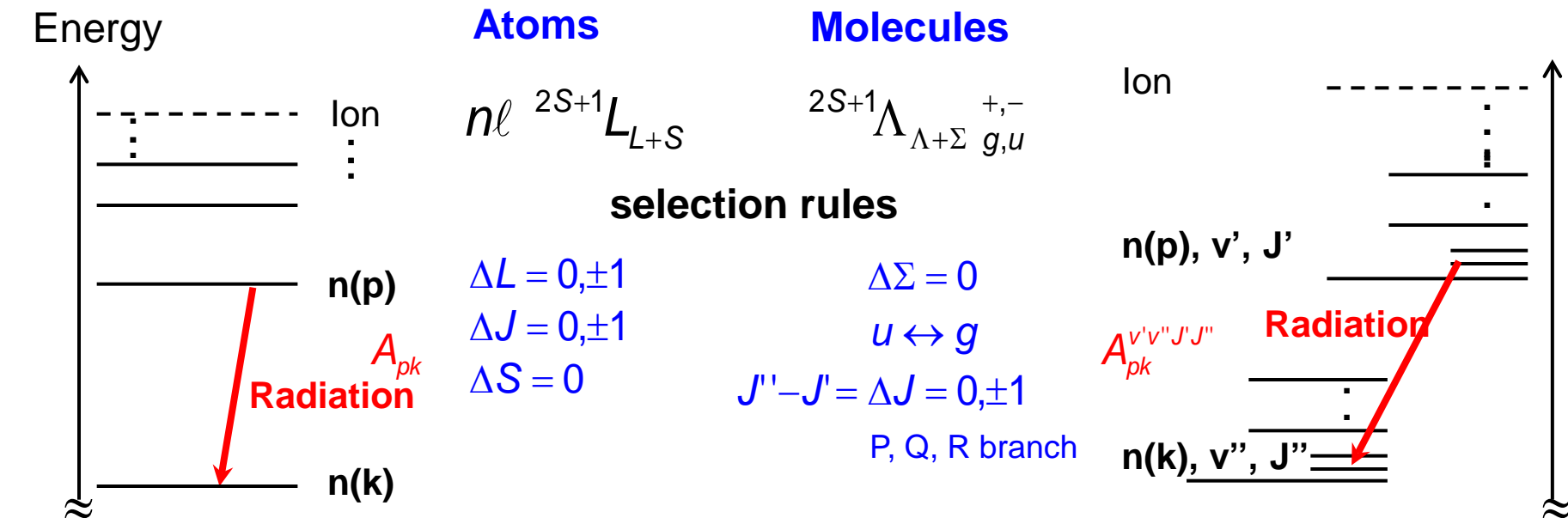
... which depends on plasma parameters !

Recommended text books [1], [4], [7], [9], [10]

What information can be obtained from the line emission ?

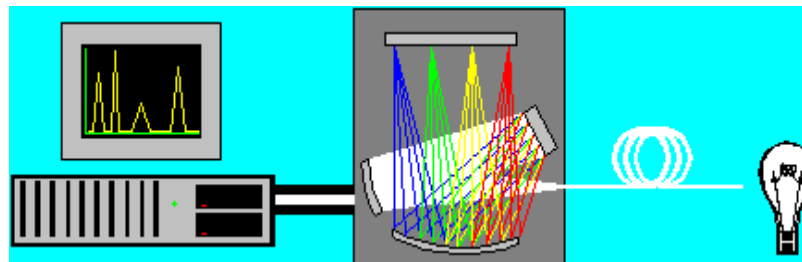


- ▶ **Intensity:** plasma parameters
density and temperature of neutrals, ions, electrons
insight in plasma processes
- ▶ **Line profile:** broadening mechanism
- ▶ **Doppler broadening:** particle temperature
- ▶ **Wavelength:** species
- ▶ **Wavelength shift:** particle velocity



Appearance depends on spectral resolution !

Spectroscopic system



Detector

- ▶ **photomultiplier**
 λ scan
 $\Delta\lambda, \Delta t$
- ▶ **diode array**
 λ range
- ▶ **CCD, ICCD**
pixel size - $\Delta\lambda$
intensity

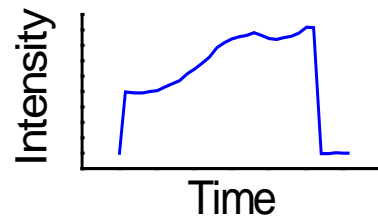
Spectrometer

- ▶ **focus length**
spectral resolution $\Delta\lambda$
- ▶ **grating**
spectral resolution
Blaze - intensity
- ▶ **slits**
entrance slit $\Delta\lambda$
exit slit - detector

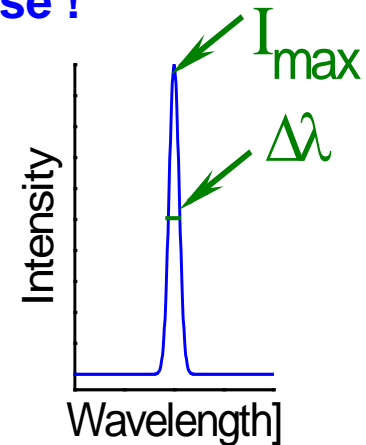
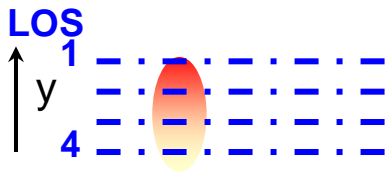
Optics

- ▶ **fibre**
very flexible
VIS: glass, quartz,
UV enhanced
- ▶ **lens and aperture**
Imaging optics
solid angle

The spectroscopic system is determined by the purpose !



- ▶ **time resolution** detector
- ▶ **spatial resolution** detector, lines of sight
- ▶ **intensity** detector, spec., optics
- ▶ **spectral resolution** detector, spec., optics



survey spectrometer	pocket size	$\Delta\lambda \approx 1-2 \text{ nm}$
1m spectrometer	good optics	$\Delta\lambda \approx 20 \text{ pm}$
Echelle spectrometer	high resolution	$\Delta\lambda \approx 1-2 \text{ pm}$

→ line profile
line shift

line monitoring

very simple

Δt , poor $\Delta\lambda$

less information

relative intensities

common technique

poor Δt , $\Delta\lambda$, Δx , flexible

moderate information

absolute intensities

expensive technique

poor Δt , $\Delta\lambda$, Δx , flexible

powerful tool

Calibration of the spectroscopic system

Wavelength: pixel \rightarrow nm
spectral lamps, plasma, λ tables

Radiance - Intensity

counts \rightarrow W/m²/sr, ph/m²/s

Ulbricht sphere



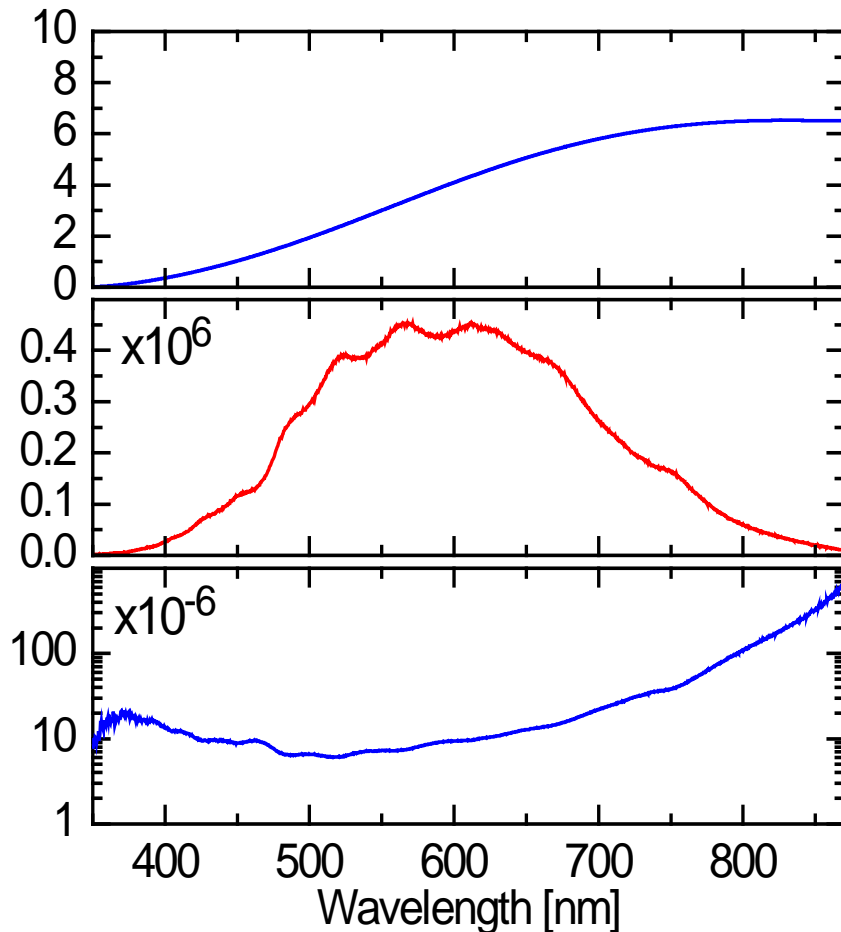
Calibrated spectrum
spectral radiance
[W/m²/sr/nm]

Measurement
intensity [counts/s]

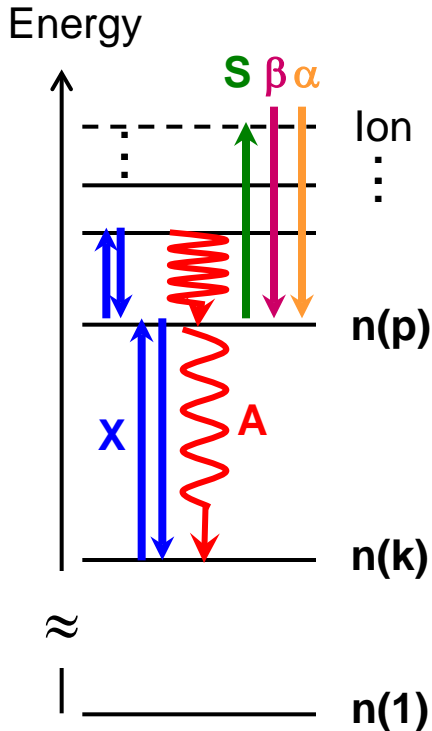
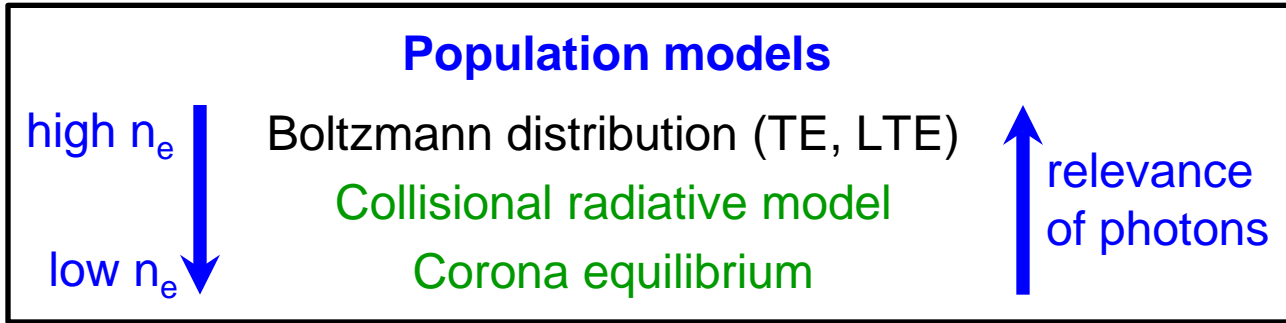
Conversion factor
spectral sensitivity

$$\left[\frac{\text{W}}{\text{m}^2 \text{sr nm (counts/s)}} \right] \times \frac{4\pi \lambda}{hc} = \left[\frac{\text{photons}}{\text{m}^2 \text{sr nm (counts/s)}} \right]$$

exposure time \nearrow



From intensity $I_{pk} = n(p) A_{pk}$ to plasma parameter



Rate equations for excitation and de-excitation processes

- ▶ electron impact excitation and de-excitation
- ▶ absorption and emission, heavy particle collisions, ...

Corona model $\downarrow n_1 \cong n_0$

$$\frac{dn(p)}{dt} = n_1 n_e X_{1p}(T_e) - n(p) \sum_k A_{pk} = 0$$

rate coefficient
emission rate coefficient

$I_{pk} = n_0 n_e X_{pk}(T_e)$ with $X_{pk} = X_{1p}(T_e) A_{pk} / \sum_k A_{pk}$

CR model $I_{pk} = n_0 n_e X_{pk}^{eff}(T_e, n_e, \dots)$

Electron temperature from absolute line emission

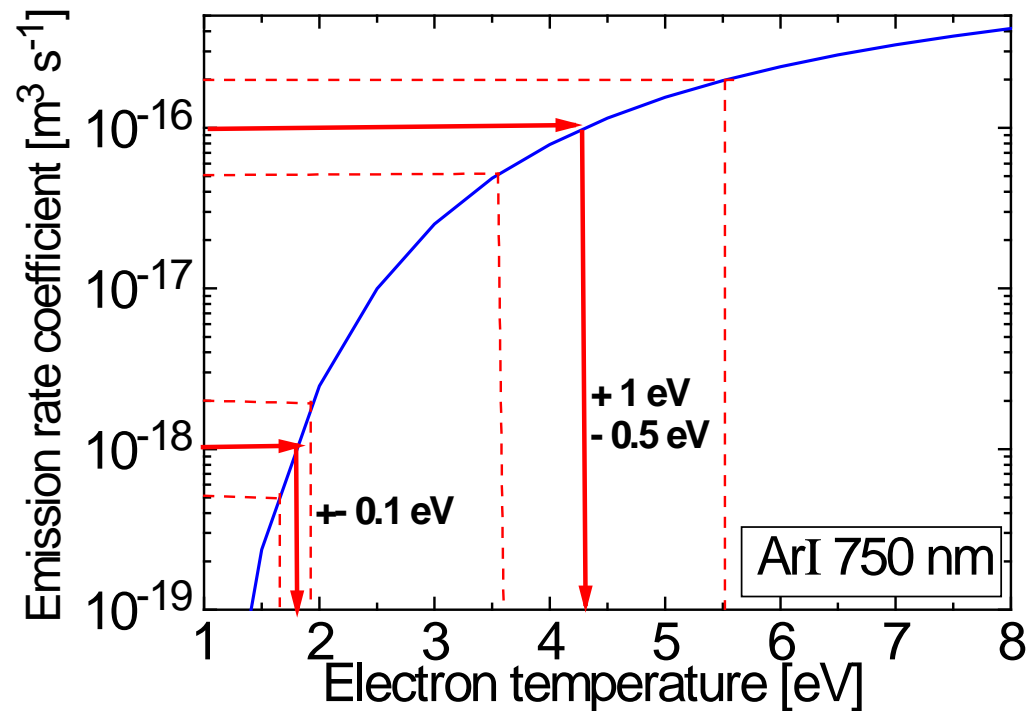
$$I_{pk} = n_0 n_e X_{pk}^{eff}(T_e, n_e, \dots)$$

$$n_0, n_e \text{ known} \Rightarrow X_{pk}^{eff}(T_e, n_e, \dots) = \frac{I_{pk}}{n_0 n_e}$$

Find suitable gases and diagnostic lines

- ▶ admixture of small amount of diagnostic gas
- ▶ prominent example: Ar

Very sensitive for low T_e !

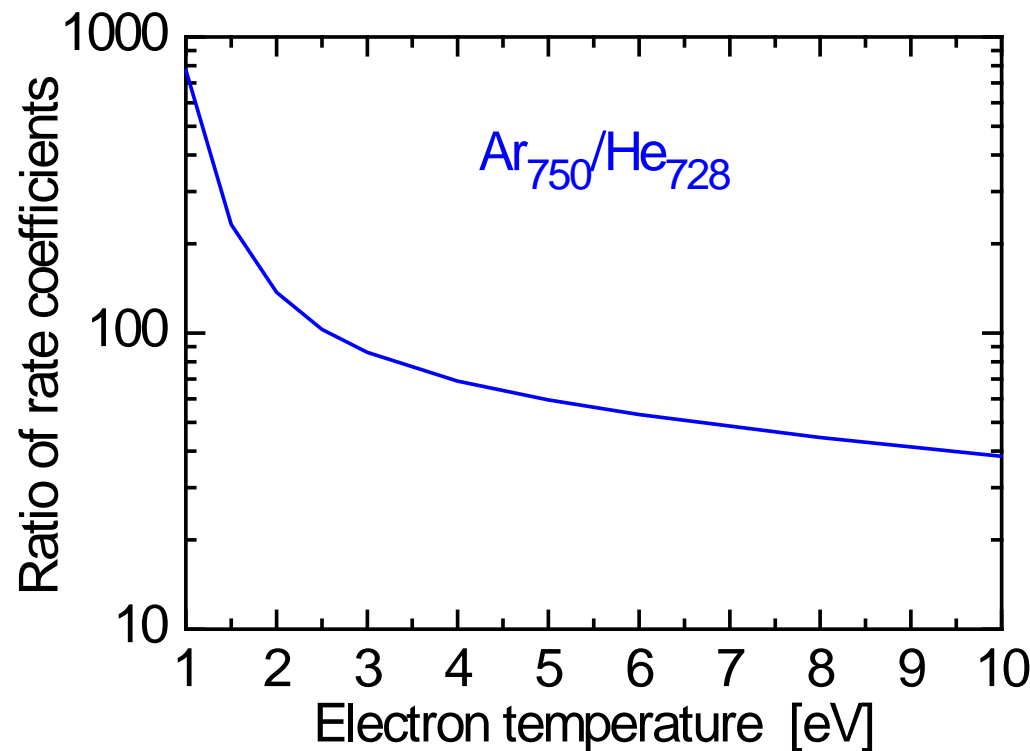


Electron temperature from line ratio (relative calibration)

$$\frac{I_{pk}^1}{I_{pk}^2} = \frac{n_1 \cancel{n_e} X_{pk}^1(T_e)}{n_2 \cancel{n_e} X_{pk}^2(T_e)} \rightarrow \text{ratio of rate coefficients for known densities}$$

Find suitable gases and diagnostic lines

- ▶ n_1, n_2 inert gases or $n_1 = n_2$
- ▶ I_{pk} undisturbed lines
- ▶ ground state excitation
- ▶ X_{pk} ratio depends on T_e



U. Fantz et al., Nucl. Fusion 49 (2009) 125007

Actinometry: density ratio from line ratio (relative calibration)

$$\frac{I_{pk}^1}{I_{pk}^2} = \frac{n_1 \cancel{n_e} X_{pk}^1(T_e)}{n_2 \cancel{n_e} X_{pk}^2(T_e)} \rightarrow \text{ratio of densities for known rate coefficients}$$

▶ density ratio (n_H/n_{H_2})

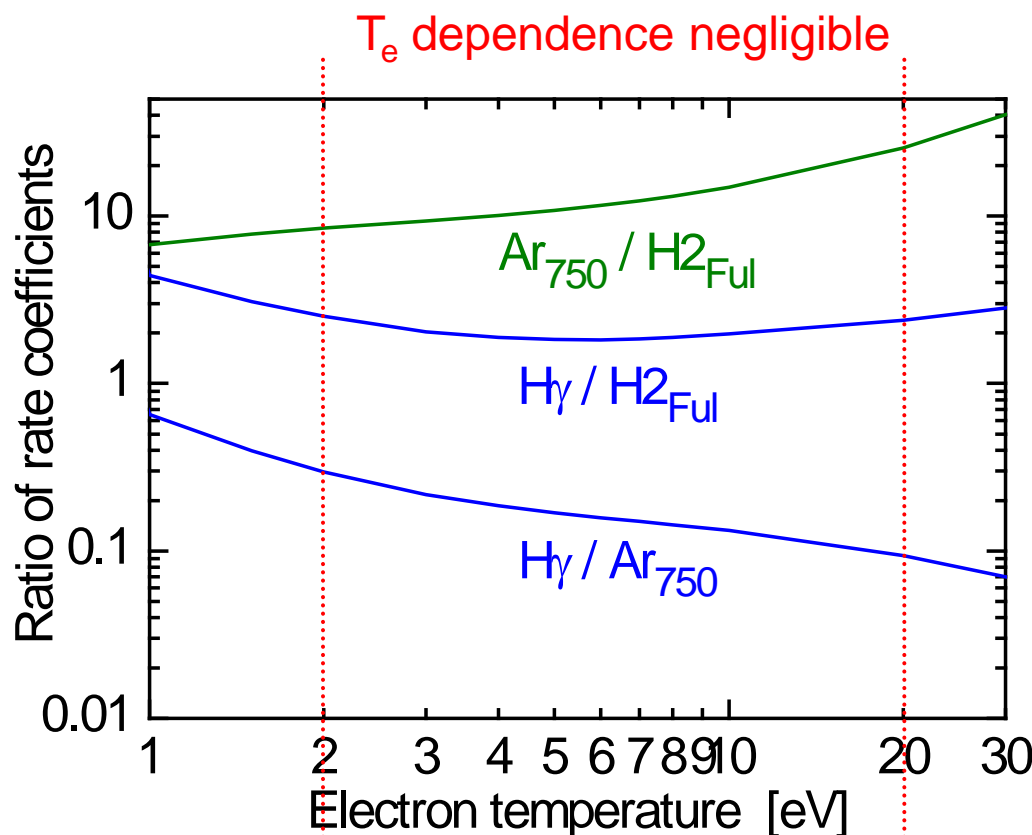
$$\frac{I_{434}^H}{I_{Ful}^{H_2}} = \frac{n_H \cancel{n_e} X_{434}^H(T_e)}{n_{H_2} \cancel{n_e} X_{Ful}^{H_2}(T_e)} \quad \text{independent on } n_e$$

▶ density ratio (n_{Ar}/n_{H_2})

$$\frac{I_{750}^{Ar}}{I_{Ful}^{H_2}} = \frac{n_{Ar} \cancel{n_e} X_{750}^{Ar}(T_e)}{n_{H_2} \cancel{n_e} X_{Ful}^{H_2}(T_e)}$$

▶ density ratio (n_{He}/n_{H_2})

dependence on T_e (factor of 10)
→needs iteration



Particle density from absolute line emission

$$I_{pk} = n_0 n_e X_{pk}^{eff}(T_e, n_e, \dots)$$

$$n_e, T_e \text{ known} \Rightarrow n_0 = \frac{I_{pk}}{n_e X_{pk}^{eff}(T_e, n_e, \dots)}$$

Knowledge of dominant excitation mechanism is essential !

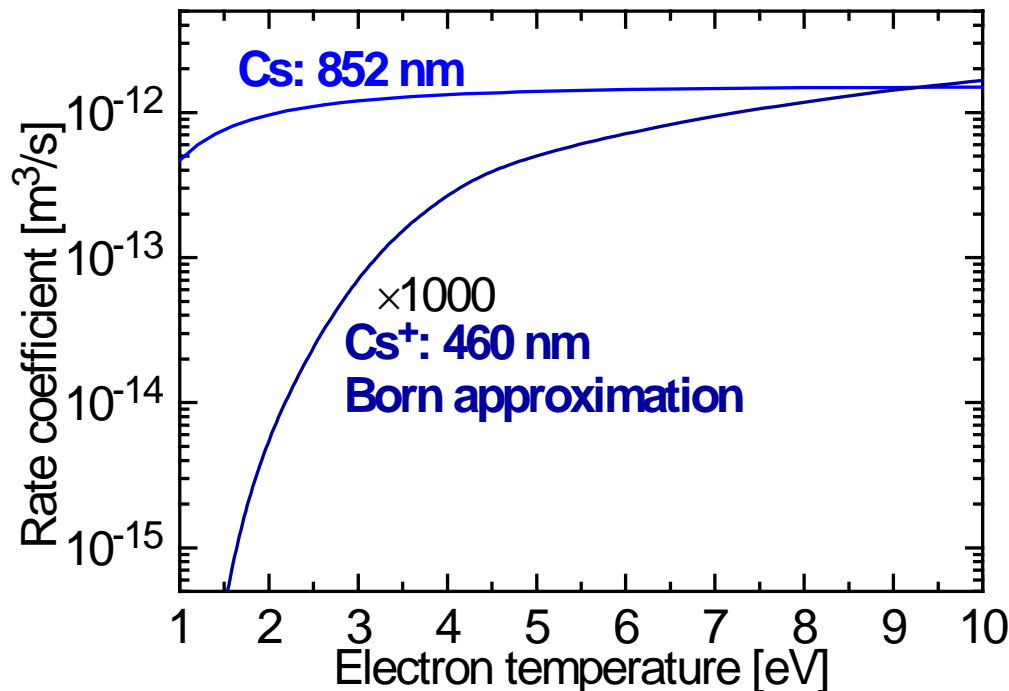
Example: Cs and Cs⁺ lines

$$\text{Cs: } I_{852}^{Cs} = n_{Cs} n_e X_{852}^{Cs}(T_e)$$

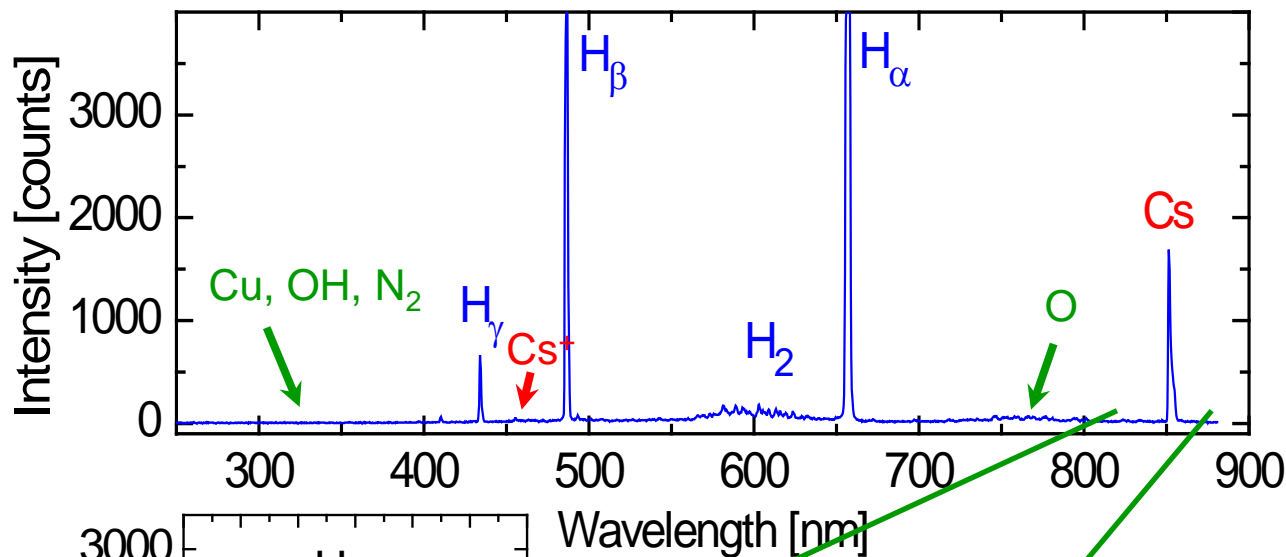
needs n_e , almost independent of T_e

$$\text{Cs}^+: I_{460}^{Cs^+} = n_{Cs^+} n_e X_{460}^{Cs^+}(T_e)$$

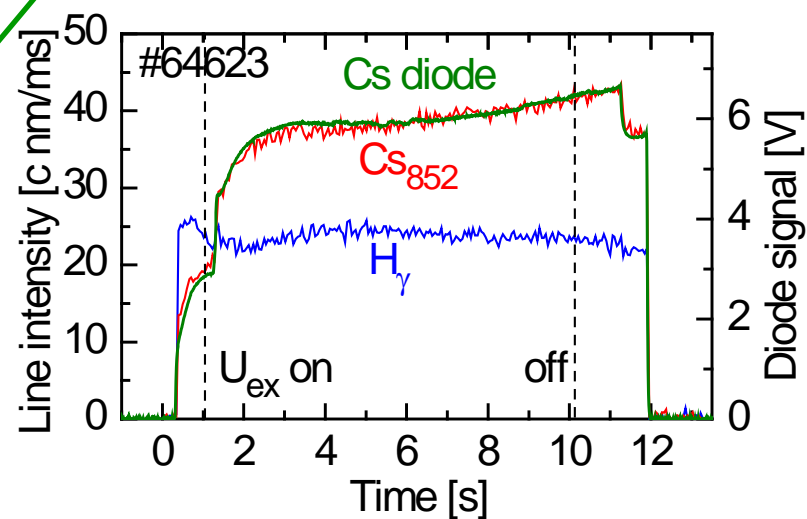
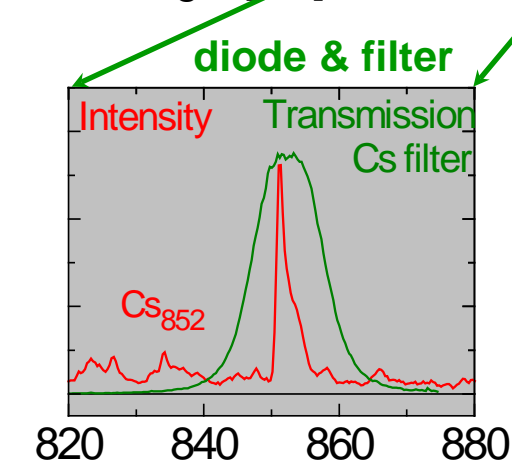
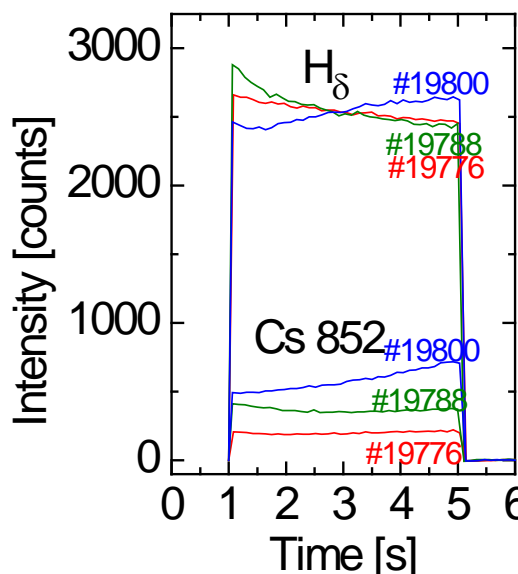
needs n_e , strong dependence on T_e



Survey spectrometer and on-line monitoring

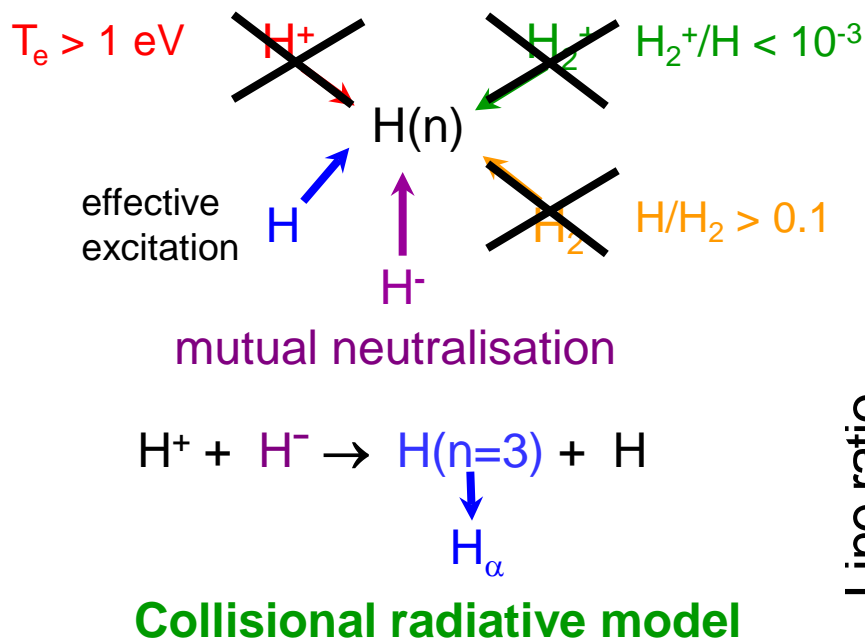


- ▶ impurities
- ▶ Cs, Cs⁺
- ▶ H from H_γ
- ▶ n_e from H_β/H_γ
- ▶ n_e, T_e diagnostic gas Ar
- ▶ H^- from H_α/H_β



A novel diagnostic technique for H^- volume density NJP 8 (2006) 301

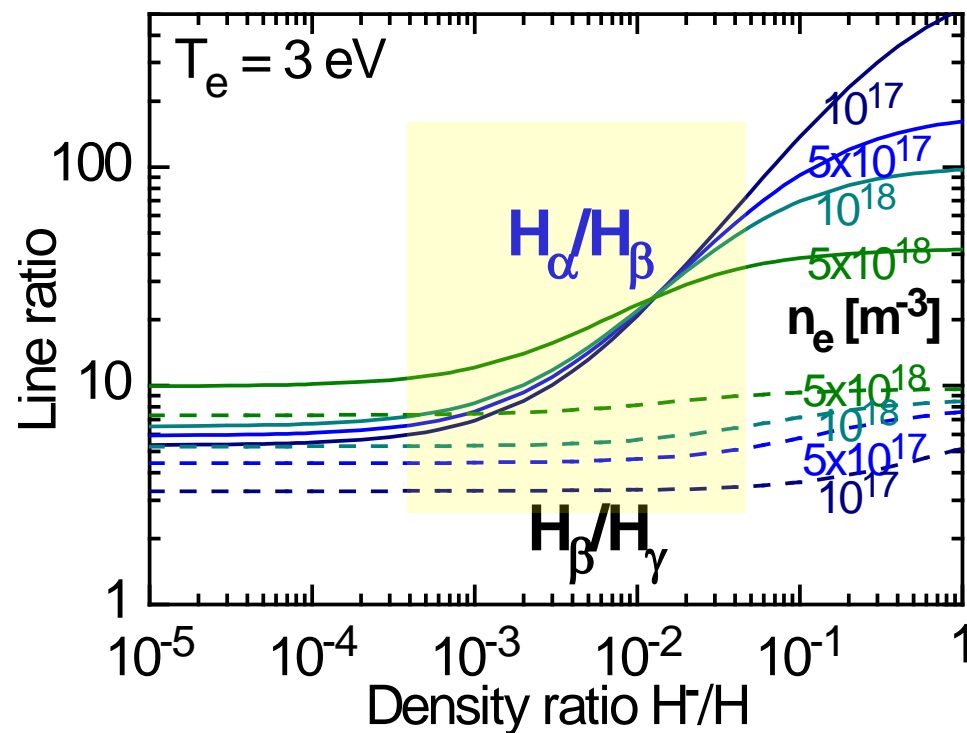
Population mechanisms for H



Measurement of Balmer line ratios

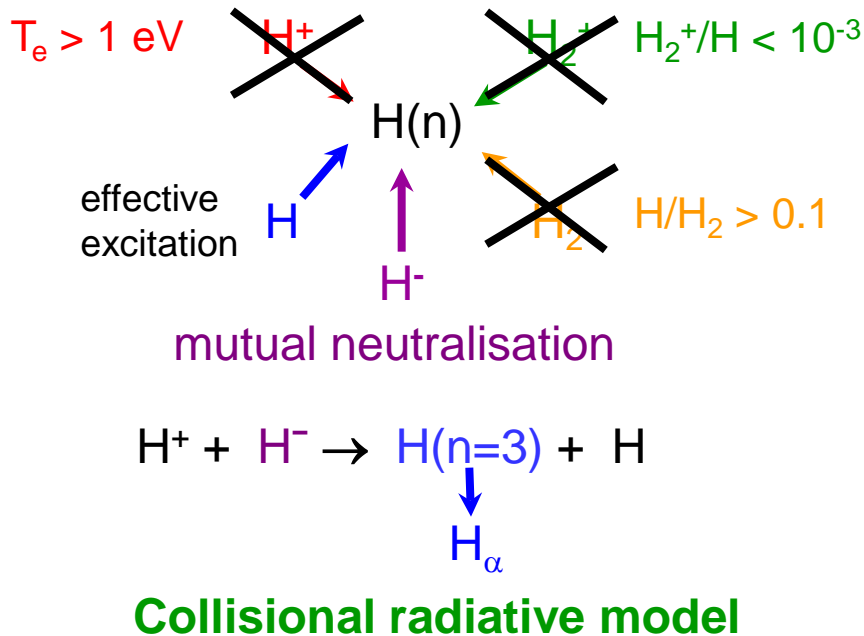
- ▶ H_α/H_β depends on H^-/H
- ▶ H_β/H_γ reflects n_e and T_e

Line ratios depend on n_e , T_e and H^-/H



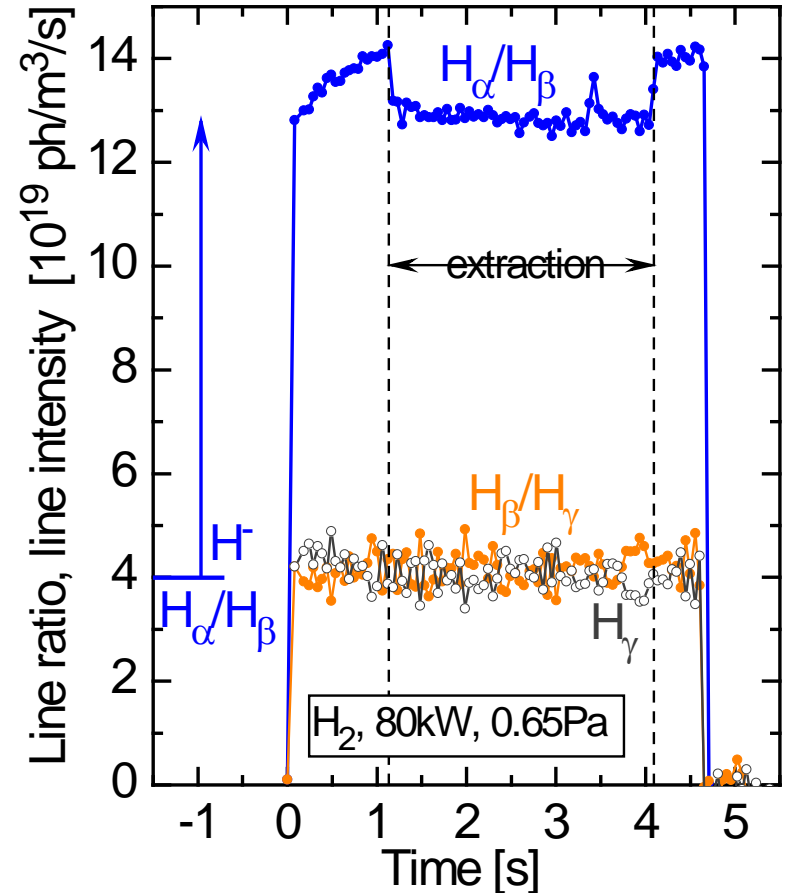
A novel diagnostic technique for H^- volume density NJP 8 (2006) 301

Population mechanisms for H



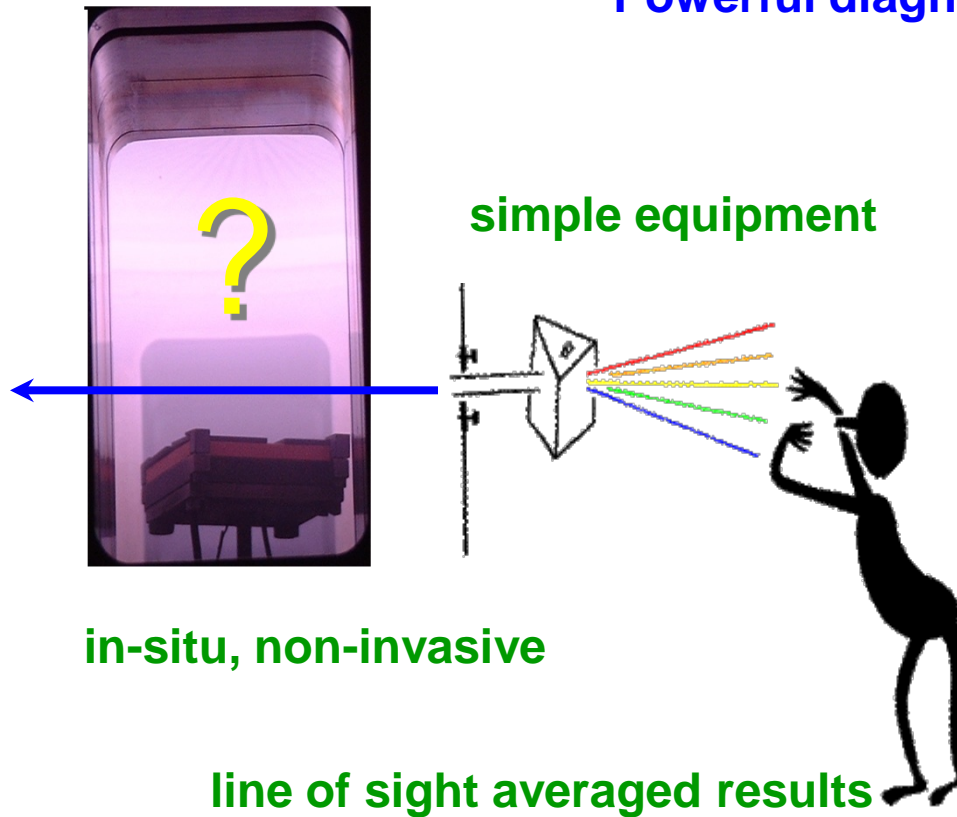
Measurement of Balmer line ratios

- ▶ H_α/H_β depends on H^-/H
- ▶ H_β/H_γ reflects n_e and T_e



- ▶ high H_α/H_β ratio: $H^- = 1 \times 10^{17} \text{ m}^{-3}$
- ▶ stable H_β/H_γ ratio, i.e. stable n_e and T_e

Powerful diagnostic tool



- ▶ identification of species
- ▶ particle densities
- ▶ particle temperatures
- ▶ on-line monitoring
- ▶ insight in plasma processes
- ▶ spatial resolution by several lines of sight

Analysis

based on atomic and molecular physics

simple

quite complex

supported by collisional radiative models

Plasma Diagnostics of Ion Sources

The three W's

- ▶ What do I want to know ? → **and why?**
- ▶ What is the adequate technique ? → **effort versus gain!**
- ▶ What is the accessibility of the source ? → **feasibility !**

The three examples

- ▶ Langmuir probes → ϕ_{pl}, n_e, T_e , (EEDF)
- ▶ Absorption techniques → $n_{\text{species}} \rightarrow \text{Cs}, \text{H}^-$
- ▶ Emission spectroscopy → $n_s, T_s \rightarrow \text{e}, \text{H}, \text{H}_2, \text{H}^-$

The three “keep-in-mind’s”

- ▶ Monitoring versus quantification → **trends or full information**
- ▶ Spatial resolution → **averaged or x-resolved (step width!)**
- ▶ Temporal resolution → **averaged or t-resolved (time scale!)**

Diagnostics – The Window to the Knowledge !

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