

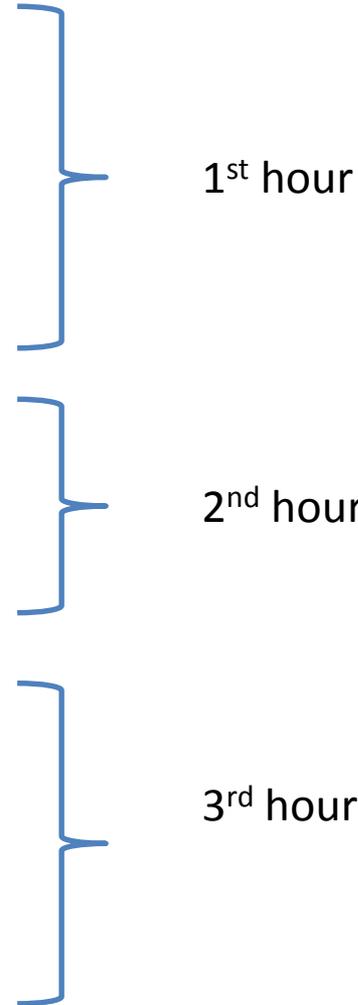
TRANSFER LINES – PAPER STUDIES

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CAS, Erice, March 2017

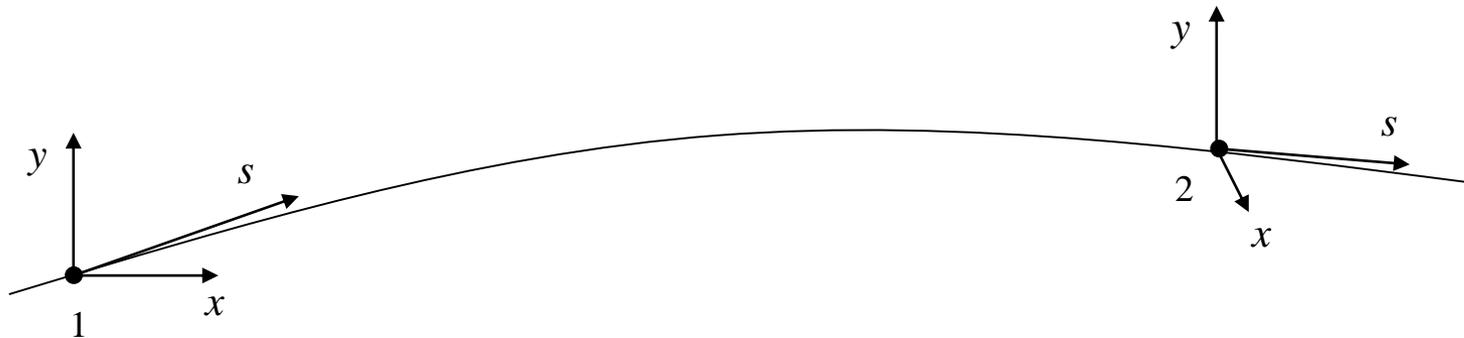
Outline

- Introduction
 - What is a transfer line?
- Paper studies
 - Geometry – estimate of bend angles
 - Optics – estimate of quadrupole gradients and apertures
 - Error estimates and tolerances on fields
- Examples of using MADX for
 - Optics and survey matching
 - Achromats
 - Final focus matching
 - Error and correction studies
- Special cases in transfer lines
 - Particle tracking
 - Plane exchange
 - Tilt on a slope
 - Gantry matching
 - Dilution
 - Stray fields



General beam transport

...moving from s_1 to s_2 through n elements, each with transfer matrix M_i



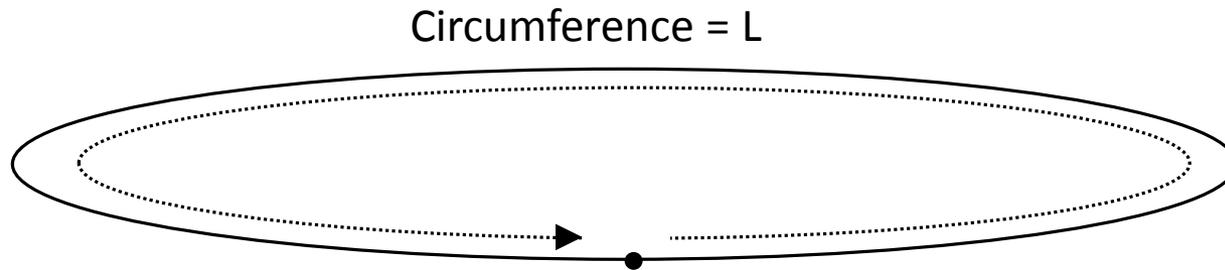
$$\begin{bmatrix} x_2 \\ x_2' \end{bmatrix} = \mathbf{M}_{1 \rightarrow 2} \cdot \begin{bmatrix} x \\ x' \end{bmatrix} = \begin{bmatrix} C & S \\ C' & S' \end{bmatrix} \cdot \begin{bmatrix} x \\ x' \end{bmatrix}$$

$$\mathbf{M}_{1 \rightarrow 2} = \prod_{i=1}^n \mathbf{M}_n$$

Twiss
parameterisation

$$\mathbf{M}_{1 \rightarrow 2} = \begin{bmatrix} \sqrt{\beta_2/\beta_1} (\cos \Delta\mu + \alpha_1 \sin \Delta\mu) & \sqrt{\beta_1\beta_2} \sin \Delta\mu \\ \sqrt{1/\beta_1\beta_2} [(\alpha_1 - \alpha_2) \cos \Delta\mu - (1 + \alpha_1\alpha_2) \sin \Delta\mu] & \sqrt{\beta_1/\beta_2} (\cos \Delta\mu - \alpha_2 \sin \Delta\mu) \end{bmatrix}$$

Circular Machine

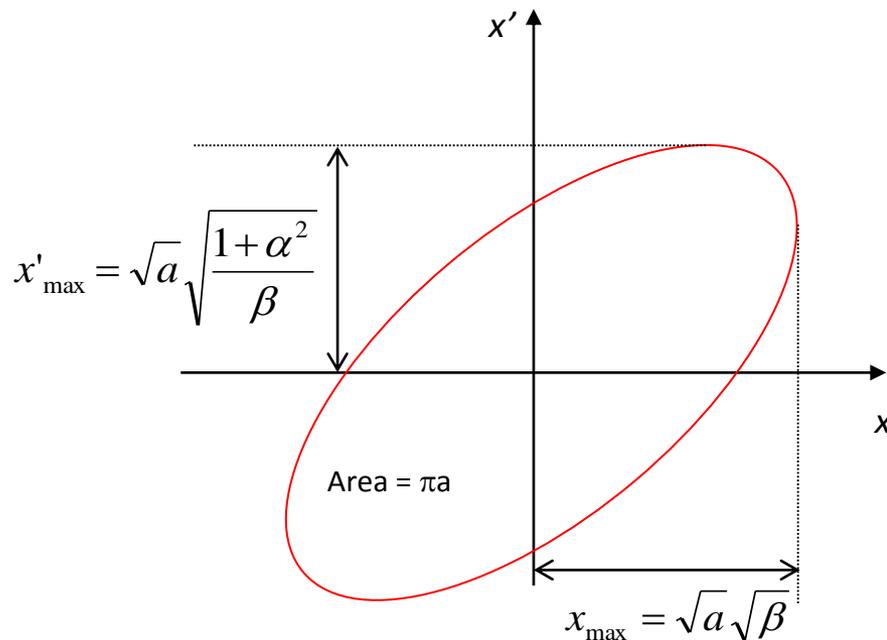


One turn $\mathbf{M}_{1 \rightarrow 2} = \mathbf{M}_{0 \rightarrow L} = \begin{bmatrix} \cos 2\pi Q + \alpha \sin 2\pi Q & \beta \sin 2\pi Q \\ -\frac{1}{\beta} (1 + \alpha^2) \sin 2\pi Q & \cos 2\pi Q - \alpha \sin 2\pi Q \end{bmatrix}$

- The solution is *periodic*
- Periodicity condition for one turn (closed ring) imposes $\alpha_1 = \alpha_2$, $\beta_1 = \beta_2$, $D_1 = D_2$
- This condition *uniquely* determines $\alpha(s)$, $\beta(s)$, $\mu(s)$, $D(s)$ around the whole ring

Circular Machine

- Periodicity of the structure leads to regular motion
 - Map single particle coordinates on each turn at any location
 - Describes an ellipse in phase space, defined by one set of a and b values \Rightarrow Matched Ellipse (for this location)

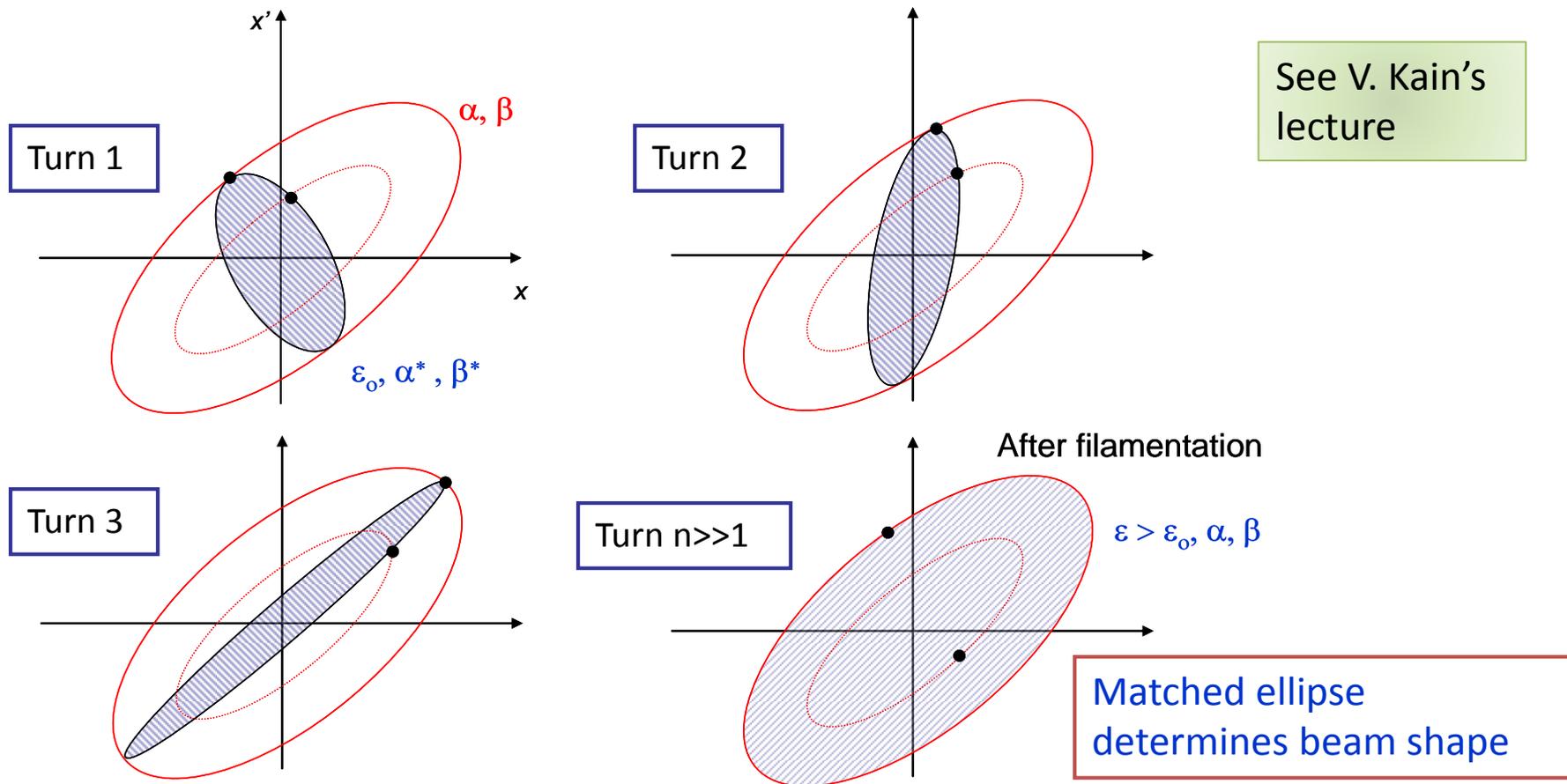


$$a = \gamma \cdot x^2 + 2\alpha \cdot x \cdot x' + \beta \cdot x'^2$$

$$\gamma = \frac{1 + \alpha^2}{\beta}$$

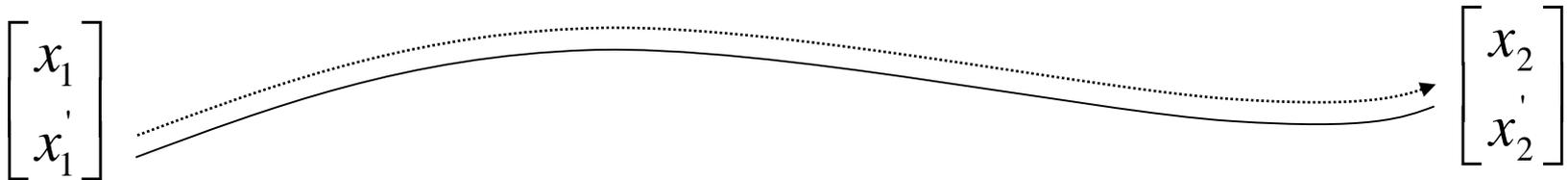
Circular Machine

- For a location with matched ellipse (a, b), an injected beam of emittance ϵ , characterised by a different ellipse (a^*, b^*) generates (via filamentation) a large ellipse with the original a, b , but larger ϵ



Transfer line

Single pass:
$$\begin{bmatrix} x_2 \\ x_2' \end{bmatrix} = \mathbf{M}_{1 \rightarrow 2} \cdot \begin{bmatrix} x \\ x' \end{bmatrix}$$

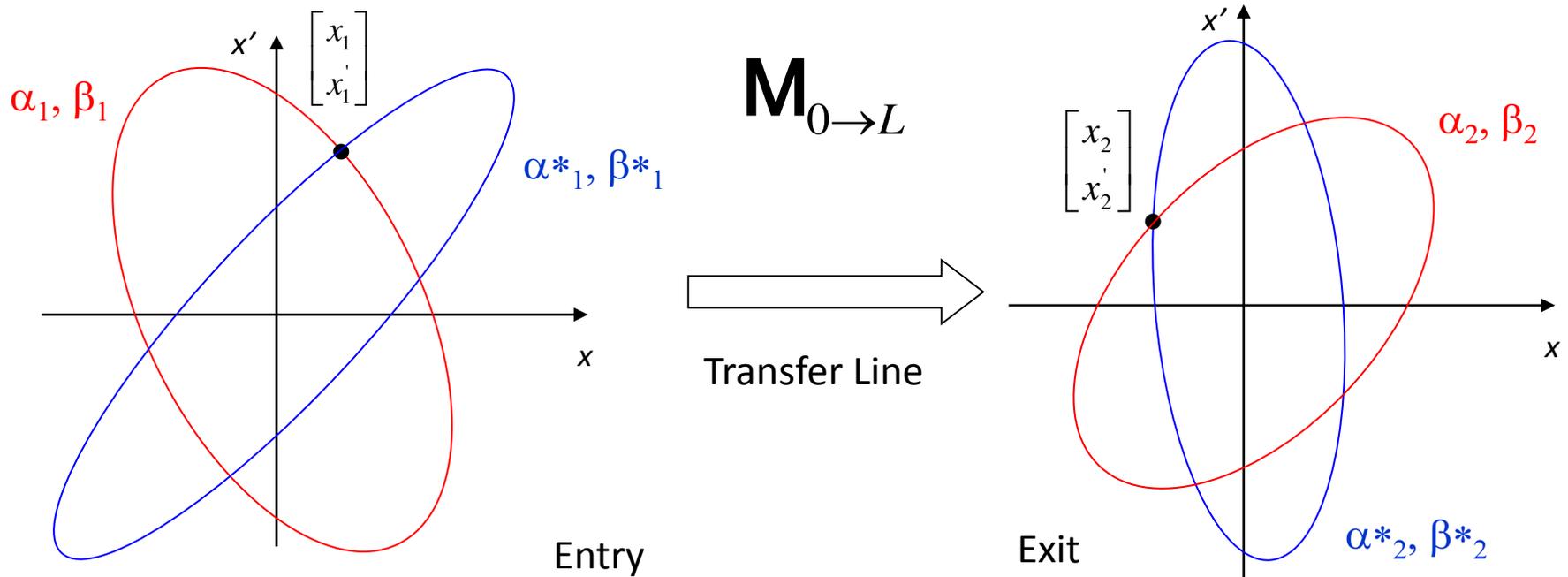


$$\mathbf{M}_{1 \rightarrow 2} = \begin{bmatrix} \sqrt{\beta_2/\beta_1} (\cos \Delta\mu + \alpha_1 \sin \Delta\mu) & \sqrt{\beta_1\beta_2} \sin \Delta\mu \\ \sqrt{1/\beta_1\beta_2} [(\alpha_1 - \alpha_2) \cos \Delta\mu - (1 + \alpha_1\alpha_2) \sin \Delta\mu] & \sqrt{\beta_1/\beta_2} (\cos \Delta\mu - \alpha_2 \sin \Delta\mu) \end{bmatrix}$$

- No periodic condition exists
- The Twiss parameters are simply propagated from beginning to end of line
- At any point in line, $\alpha(s) \beta(s)$ are functions of $\alpha_1 \beta_1$

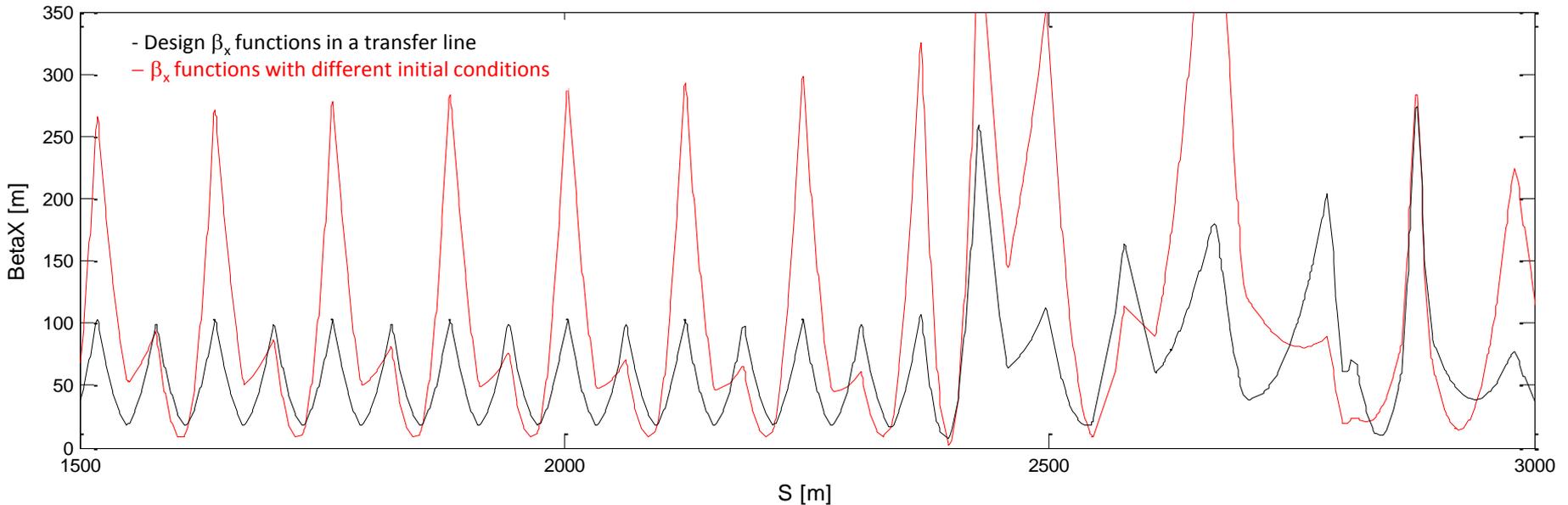
Transfer line

- On a single pass there is no regular motion
 - Map single particle coordinates at entrance and exit.
 - Infinite number of equally valid possible starting ellipses for single particle
.....transported to infinite number of final ellipses...



Transfer Line

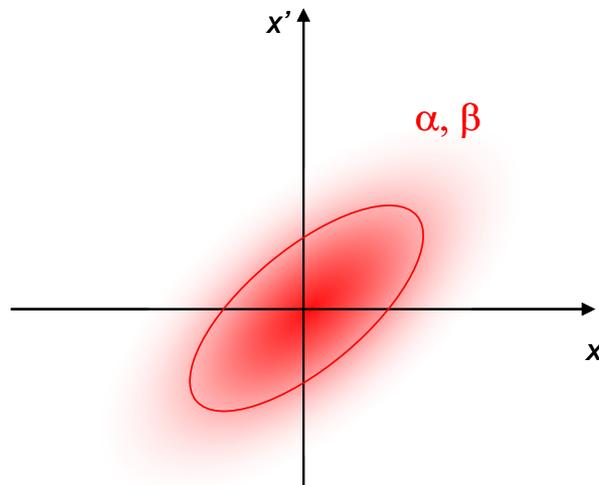
- The optics functions in the line depend on the initial values



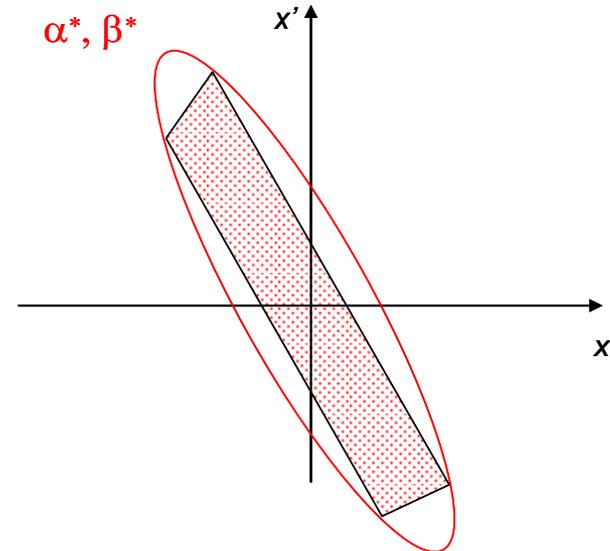
- Same considerations are true for Dispersion function:
 - Dispersion in ring defined by periodic solution \rightarrow ring elements
 - Dispersion in line defined by initial D and D' and line elements

Transfer Line

- Initial a, b defined for transfer line by beam shape at entrance



Gaussian beam

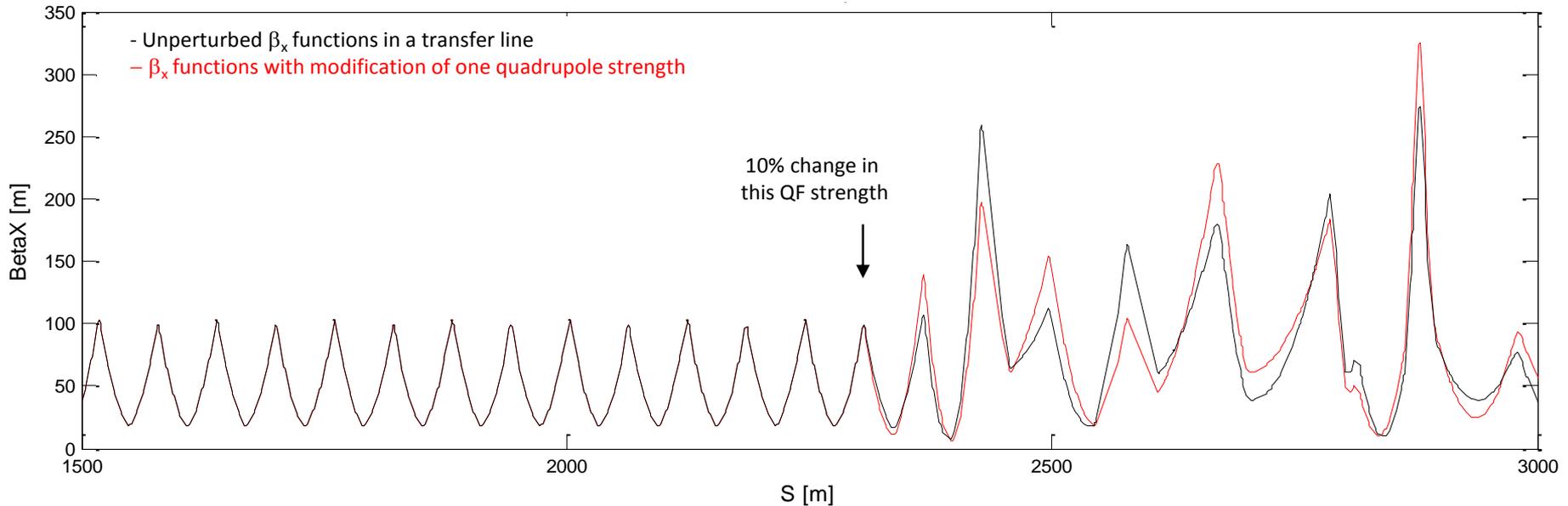


Non-Gaussian beam
(e.g. slow extracted)

- Propagation of this beam ellipse depends on line elements
- [A transfer line optics is different for different input beams](#)

Transfer Line

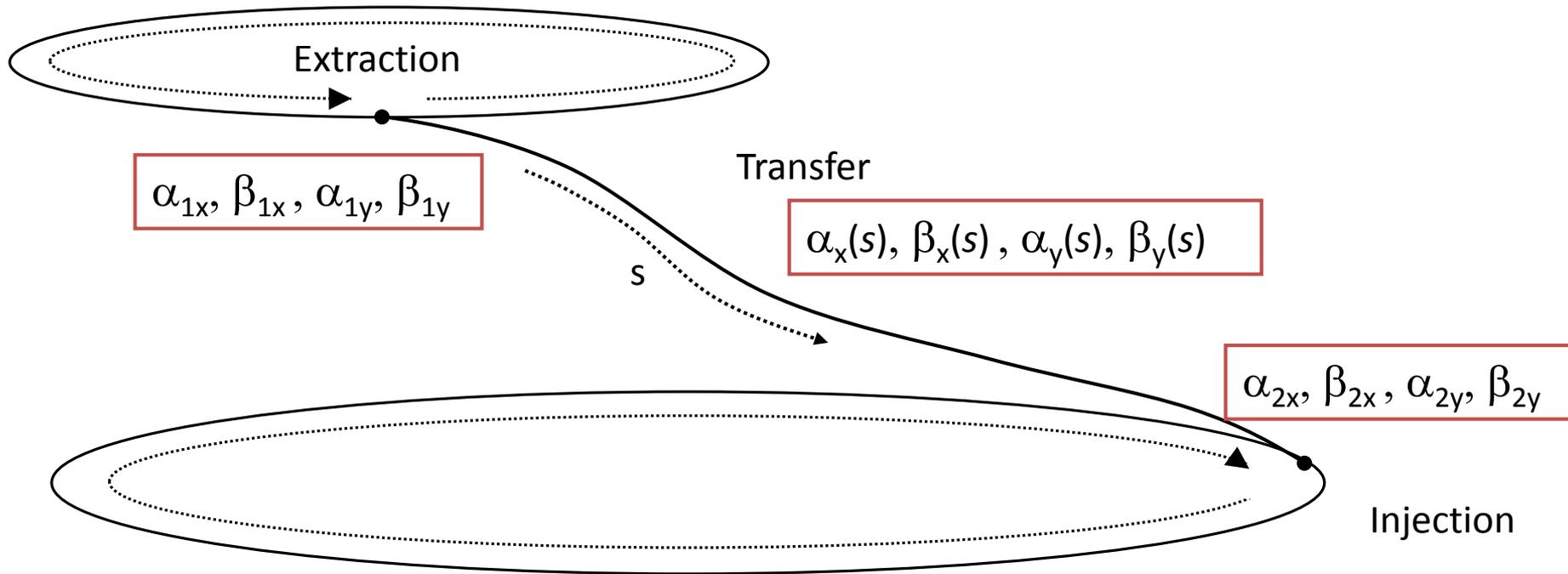
- Another difference....unlike a circular ring, a change of an element in a line affects *only* the downstream Twiss values (including dispersion)



Linking Machines

- Beams have to be transported from extraction of one machine to injection of next machine
 - Trajectories must be matched, ideally in all 6 geometric degrees of freedom (x,y,z,theta,phi,psi)
- Other important constraints can include
 - Minimum bend radius, maximum quadrupole gradient, magnet aperture, cost, geology

Linking Machines



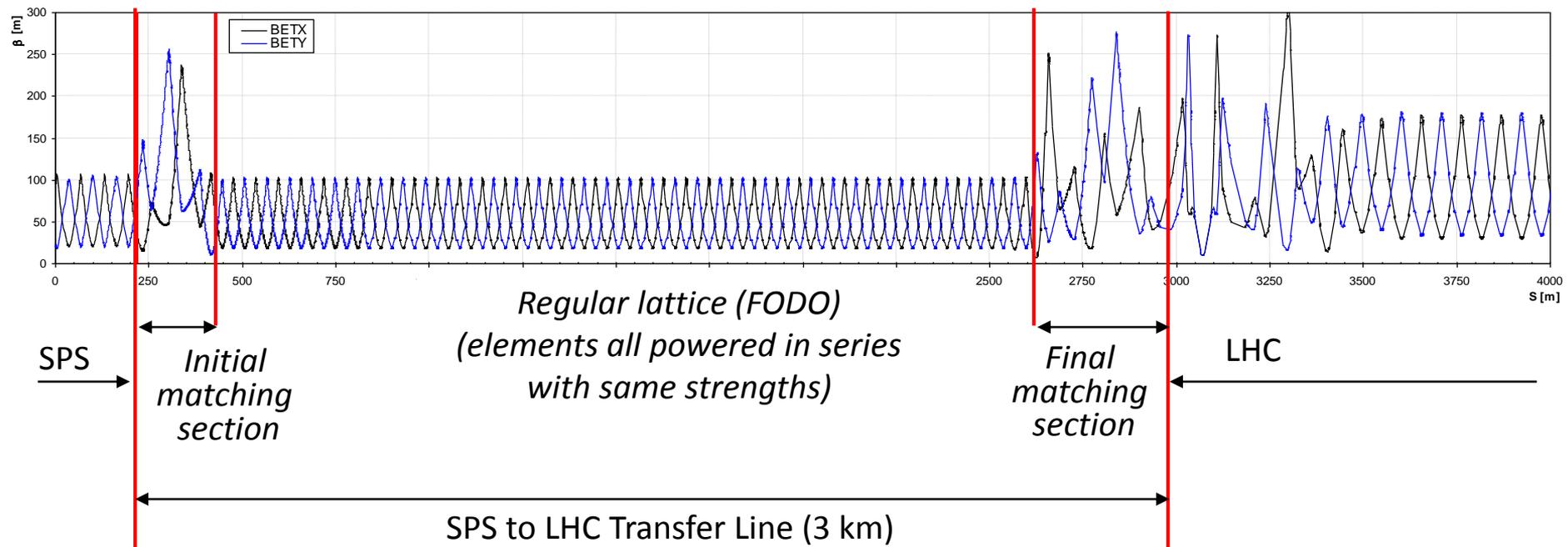
The Twiss parameters can be propagated when the transfer matrix M is known

$$\begin{bmatrix} x_2 \\ x_2' \end{bmatrix} = \mathbf{M}_{1 \rightarrow 2} \cdot \begin{bmatrix} x \\ x' \end{bmatrix} = \begin{bmatrix} C & S \\ C' & S' \end{bmatrix} \cdot \begin{bmatrix} x \\ x' \end{bmatrix}$$

$$\begin{bmatrix} \beta_2 \\ \alpha_2 \\ \gamma_2 \end{bmatrix} = \begin{bmatrix} C^2 & -2CS & S^2 \\ -CC' & CS' + SC' & -SS' \\ C'^2 & -2C'S' & S'^2 \end{bmatrix} \cdot \begin{bmatrix} \beta_1 \\ \alpha_1 \\ \gamma_1 \end{bmatrix}$$

Linking Machines

- For long transfer lines we can simplify the problem by designing the line in separate sections
 - Regular central section – e.g. FODO or doublet, with quads at regular spacing, (+ bending dipoles), with magnets powered in series
 - Initial and final matching sections – independently powered quadrupoles, with sometimes irregular spacing.



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1st hour



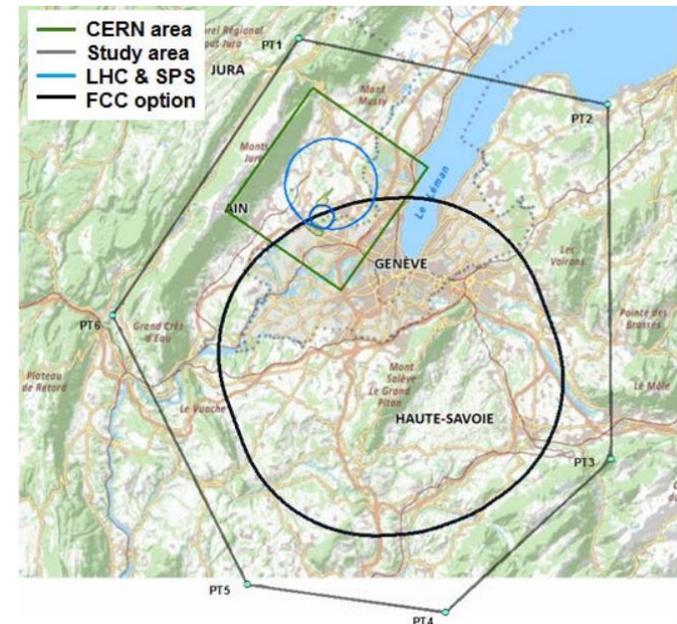
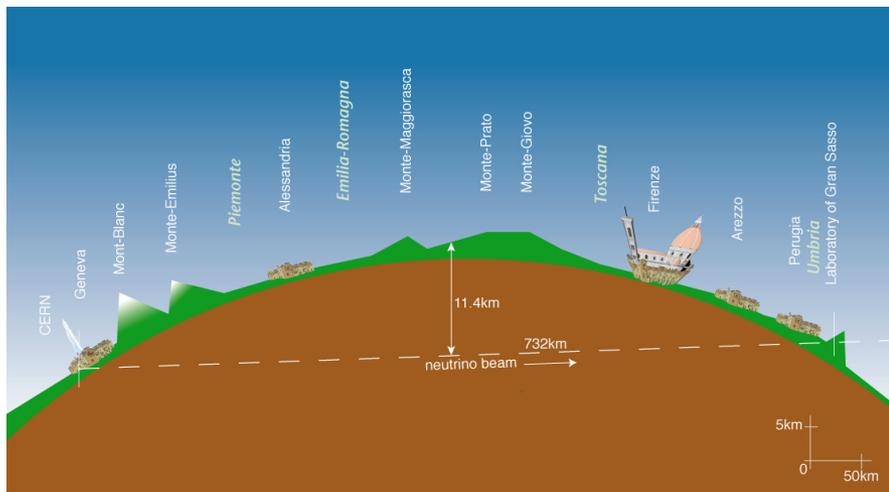
2nd hour



3rd hour

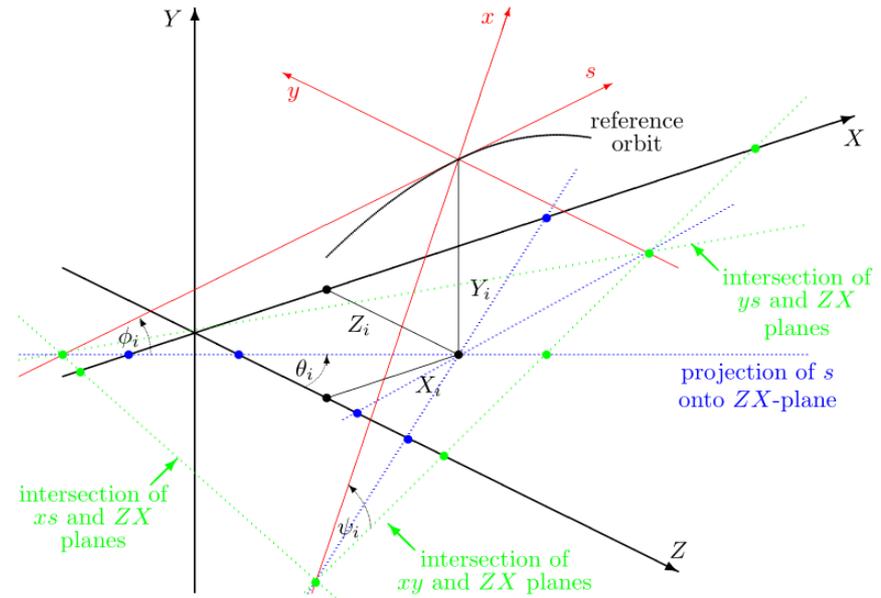
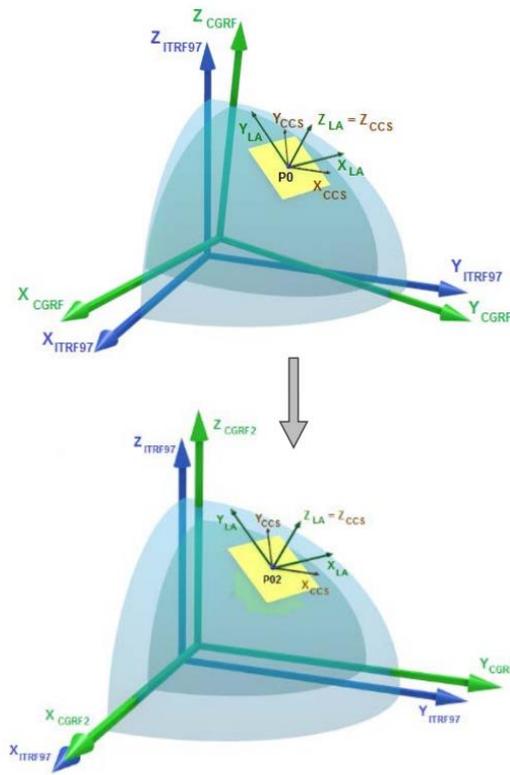
Survey

- Need coordinates and angles of points to be linked in a common coordinate system
- Linking CNGS to Gran Sasso in Italy the CERN reference frame had to be connected to the global systems of Switzerland and Italy – small rotations seen but negligible
- FCC study covers an area ten times bigger than existing installations



Survey vs MADX

- Clear definition of coordinate system with survey colleagues is essential!



MADX reference system

Bending fields

- Magnetic and electric rigidity:

$$B\rho = \frac{p}{q} \quad \Rightarrow \quad B\rho [Tm] = 3.3356 \frac{A}{n} p [GeV/c]$$

A ... atomic mass number
n ... charge state
p ... average momentum per nucleon

$$E\rho = \frac{pv}{q} \quad \Rightarrow \quad E\rho [kV] = \frac{\gamma+1}{\gamma} \frac{A}{n} T [keV]$$

T ... average kinetic energy per nucleon

- Deflection angle: $\theta = \frac{Bdl}{B\rho}$ or $\frac{Edl}{E\rho}$

Where is the limit between electric and magnetic?

- Electric devices are limited by the applied voltage – one can assume several 10s of kV as limit for reasonable accelerator apertures
- Magnets are limited by the field quality at low fields
 - Strong dependence on material properties
 - Remnant fields become important
 - Measuring the field becomes a challenge
- Example
 - 100 keV antiprotons
 - Electrostatic quadrupoles with 60 mm diameter require applied voltages of below 10 kV
 - Electrostatic bends of up to 30 kV

If you are in the energy grey zone...how to choose between magnetic and electric?

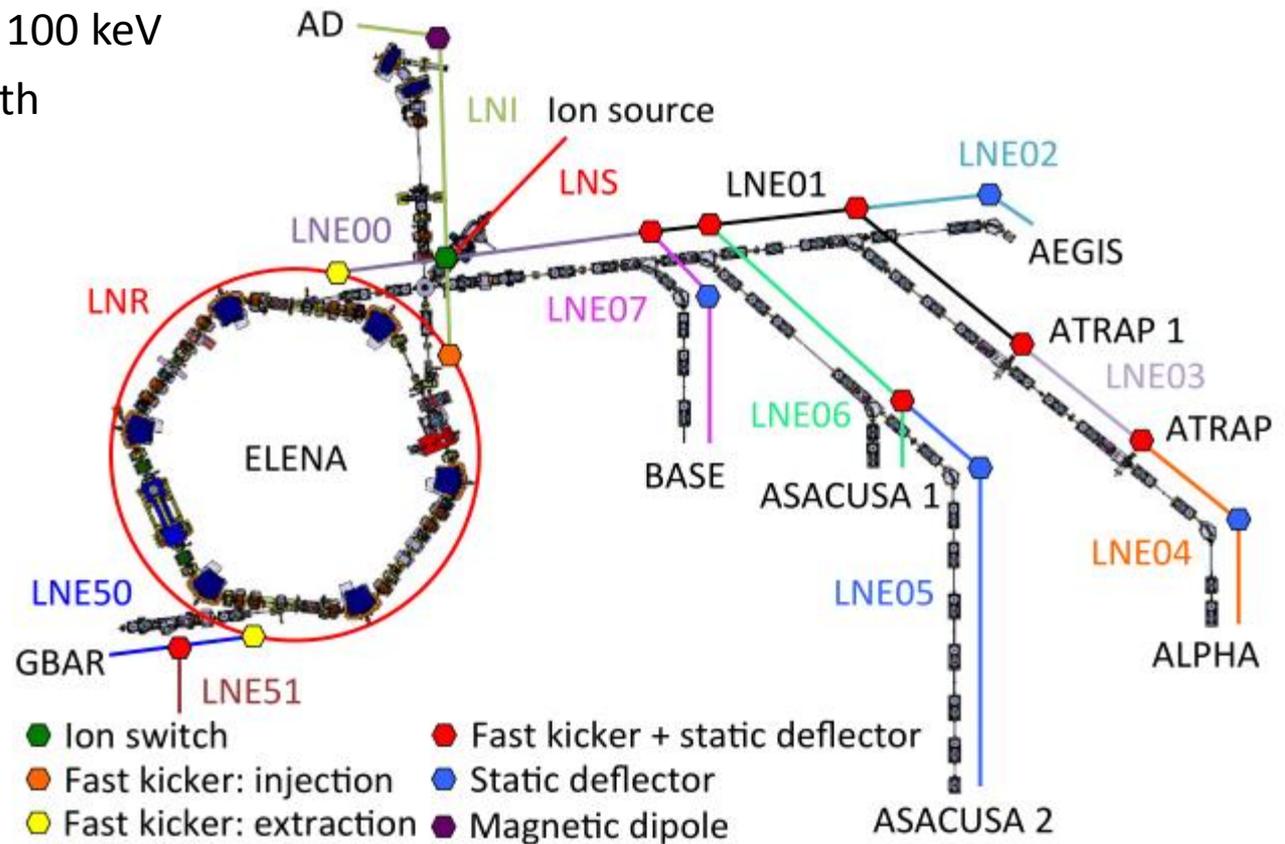
Pros and cons of electrostatic beam lines:

- Cheap element production
- Cheap power supplies and cabling
- Mass-independent
- No hysteresis effects (easy operation)
- No power consumption – no cooling
- Transverse field shape easy to optimize

- Difficult to measure field shape – effective length
- Diagnosis of bad connections
 - Inside vacuum
 - Large outgassing surface area
 - Vulnerable to dirt inside vacuum
 - Requires vacuum interlock for sparking and safety
 - Repair requires opening the vacuum
 - Limited choice of vacuum and bake-out compatible insulators

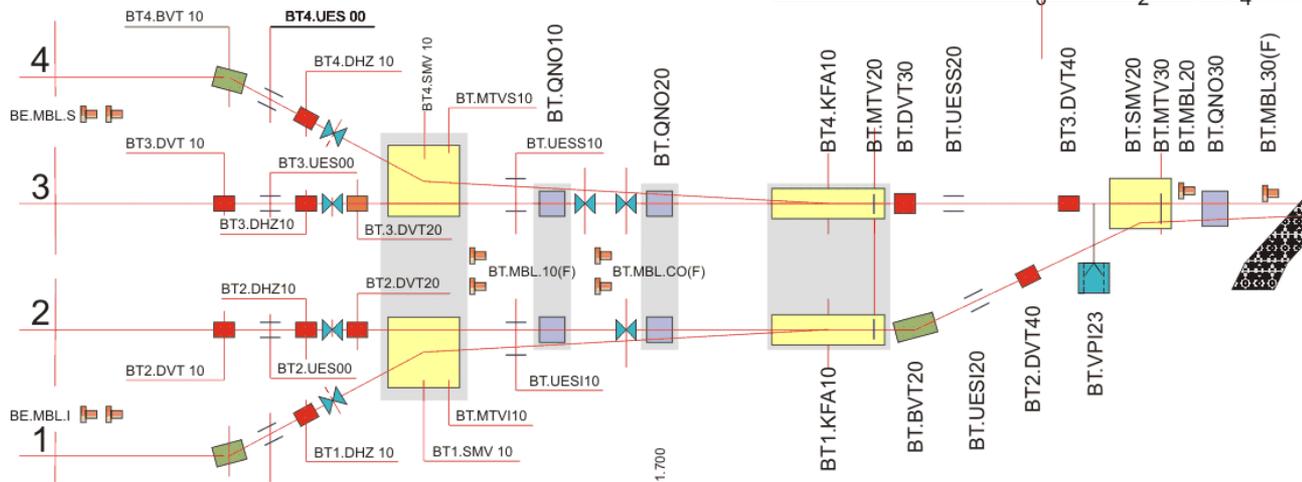
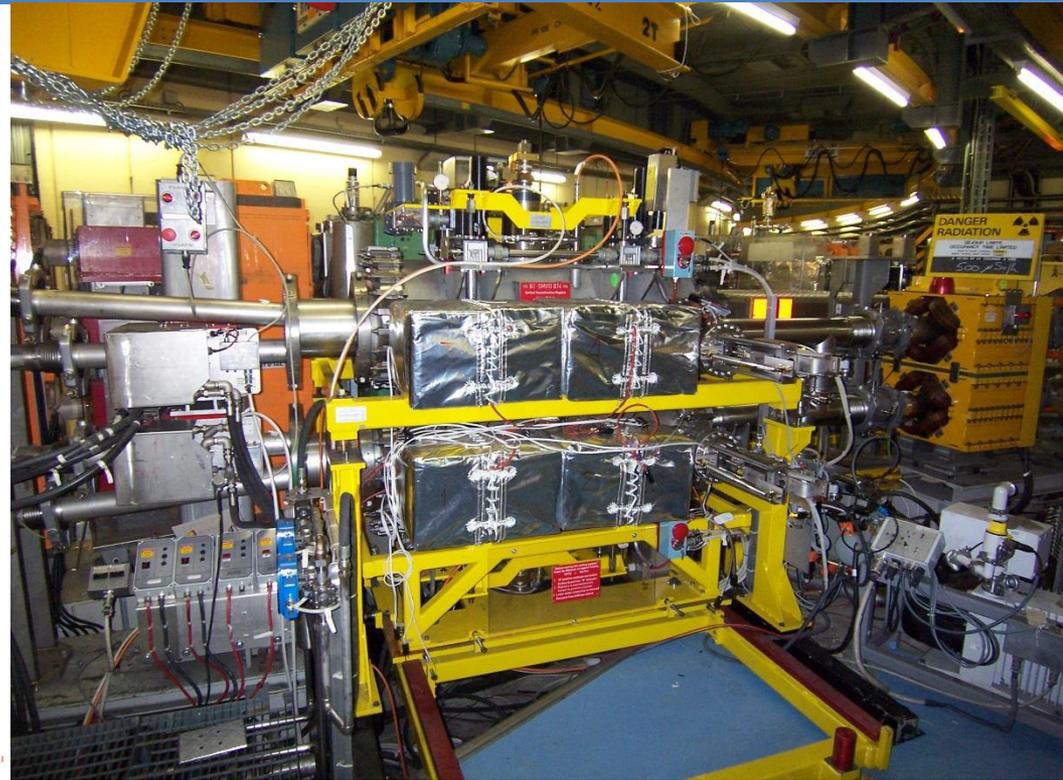
2D geometry

- Very low energy of 100 keV
- Short bending length



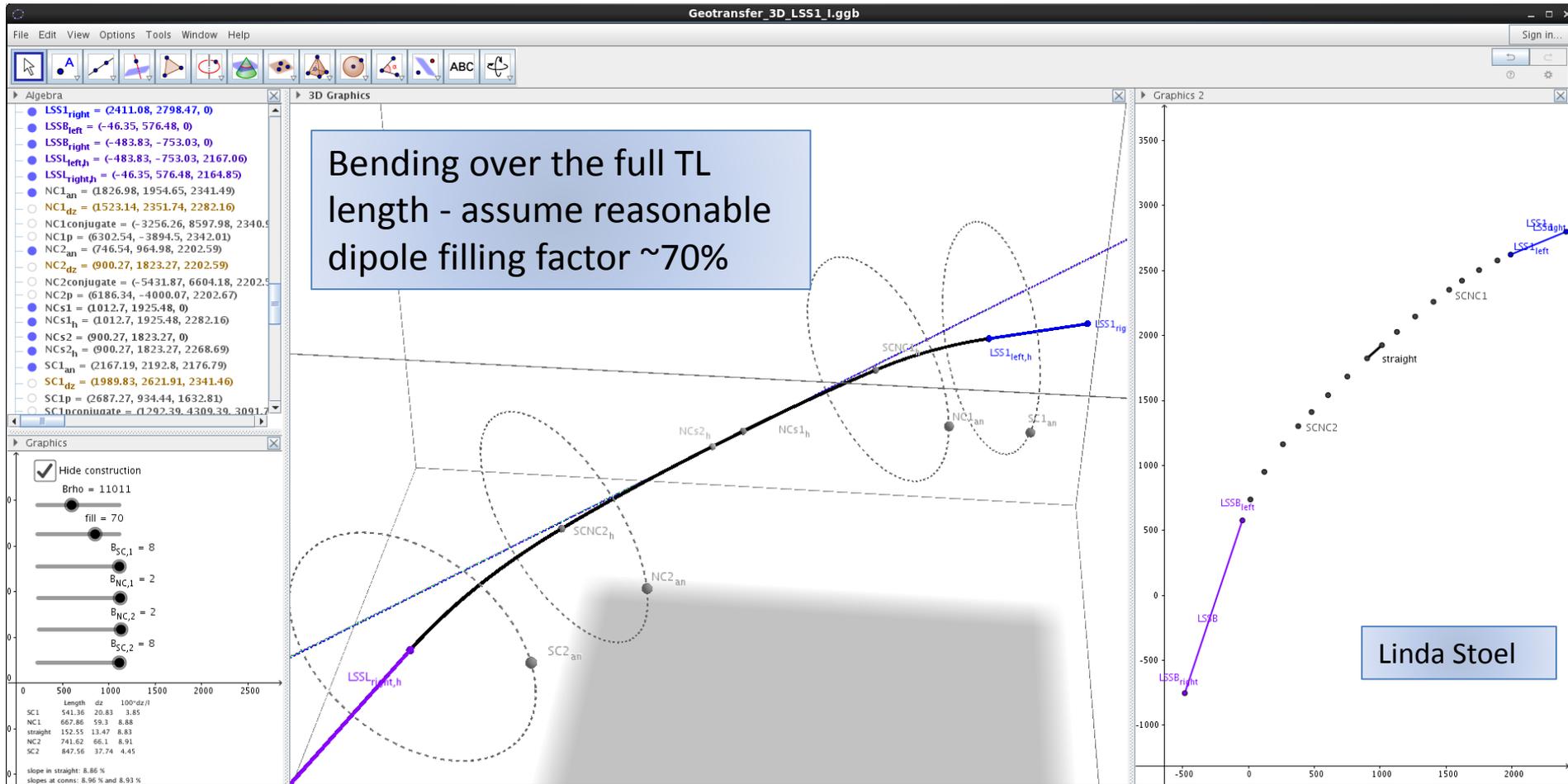
2D geometry

- 36 cm vertical height difference over several m
- 1-2 GeV



More complex 3D geometry

Several 100 m vertical, several km length, 3.3 TeV...distributed bending



Bending field limits

- So far we considered the bending fields in transfer lines limited solely by hardware
- A few 10 kV on electrostatic devices to avoid sparking
- 2 T for normal conducting magnets
- Something like existing LHC dipole reach 9-9.5 T for superconducting magnets

- But is there anything else which might limit the bending field?

Lorentz Stripping

- Transfer of H⁻ ions
- Extra electron binding energy is 0.755 eV
- A moving ion sees the magnetic field in its rest frame – Lorentz transform gives electric field as

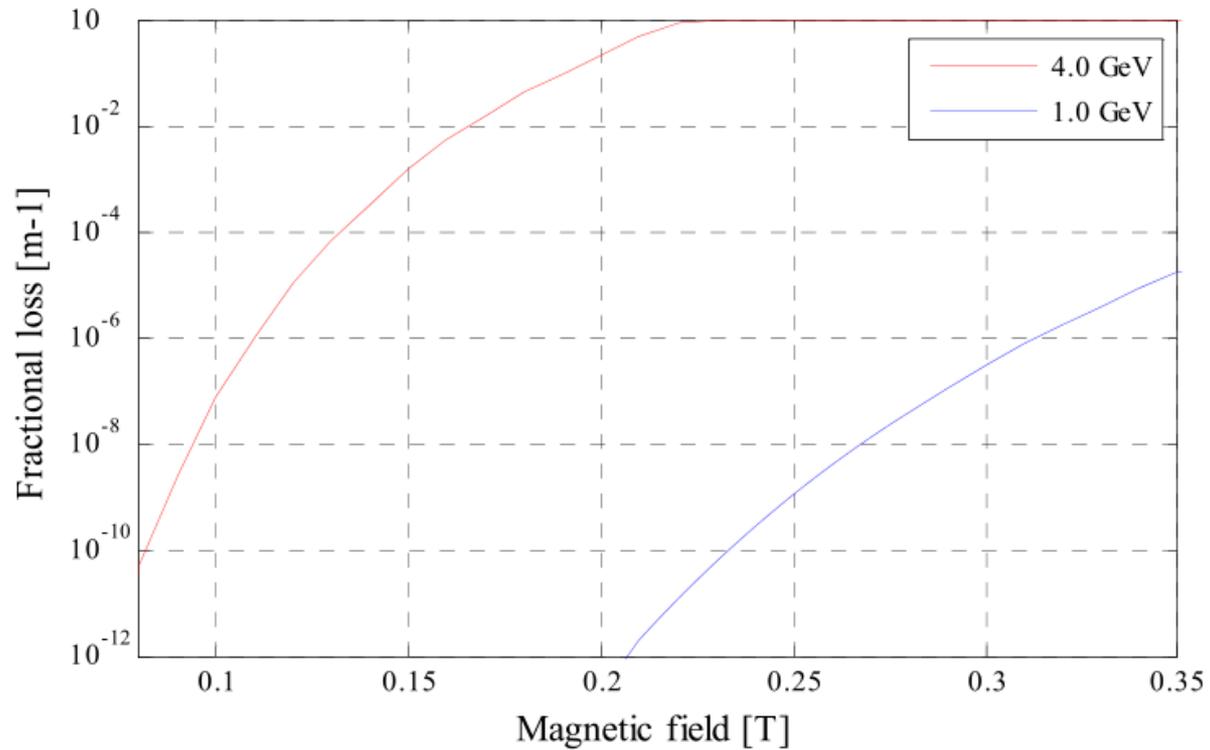
$$E \left[\frac{MV}{cm} \right] = 3.197 \cdot p \left[\frac{GeV}{c} \right] \cdot B [T]$$

- Lifetime

$$\tau = \frac{A}{E} e^{\left(\frac{C}{E}\right)} \quad A = 7.96 \cdot 10^{-14} \text{ s MV/cm}, C = 42.56 \text{ MV/cm}$$

Example of PS2

- 4 GeV injection
- Fractional loss below 10^{-4} limits magnetic field to 0.13 T

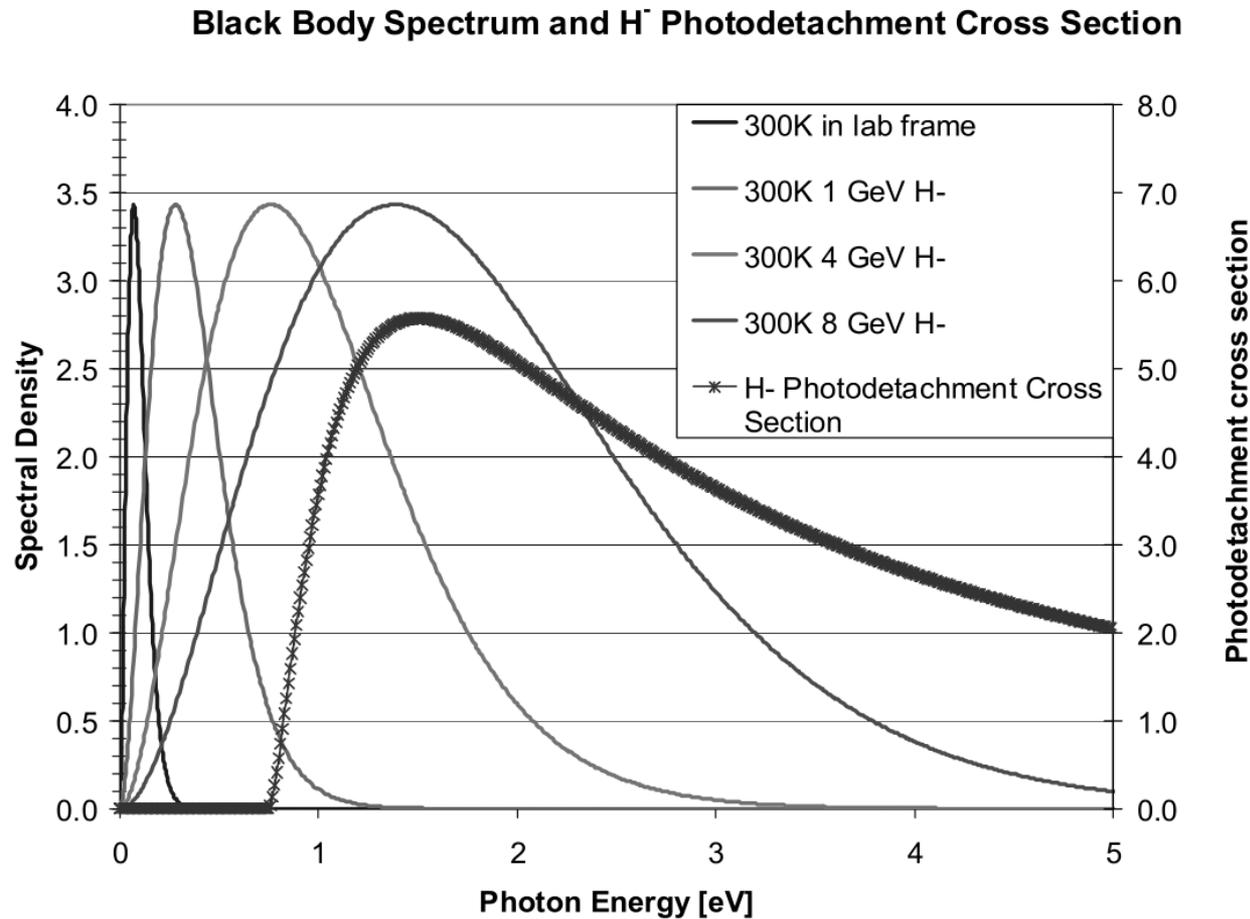


Example of Fermilab Project-X

- Was considered as proton source with **8 GeV H-** into recycler ring for neutrino program
- Usual power loss limit in lines of $\sim 1\text{W/m}$
- Activation was found to be not acceptable for 8 GeV ions
- Reduction to 0.05 W/m power loss to meet radioprotection requirements
- In this regime also other loss processes become relevant...

[4] D. Johnson. Challenges Associated with 8 GeV H- Transport and Injection for FERMILAB PROJECT-X *. (Proceedings of Hadron Beam 2008), 2008.

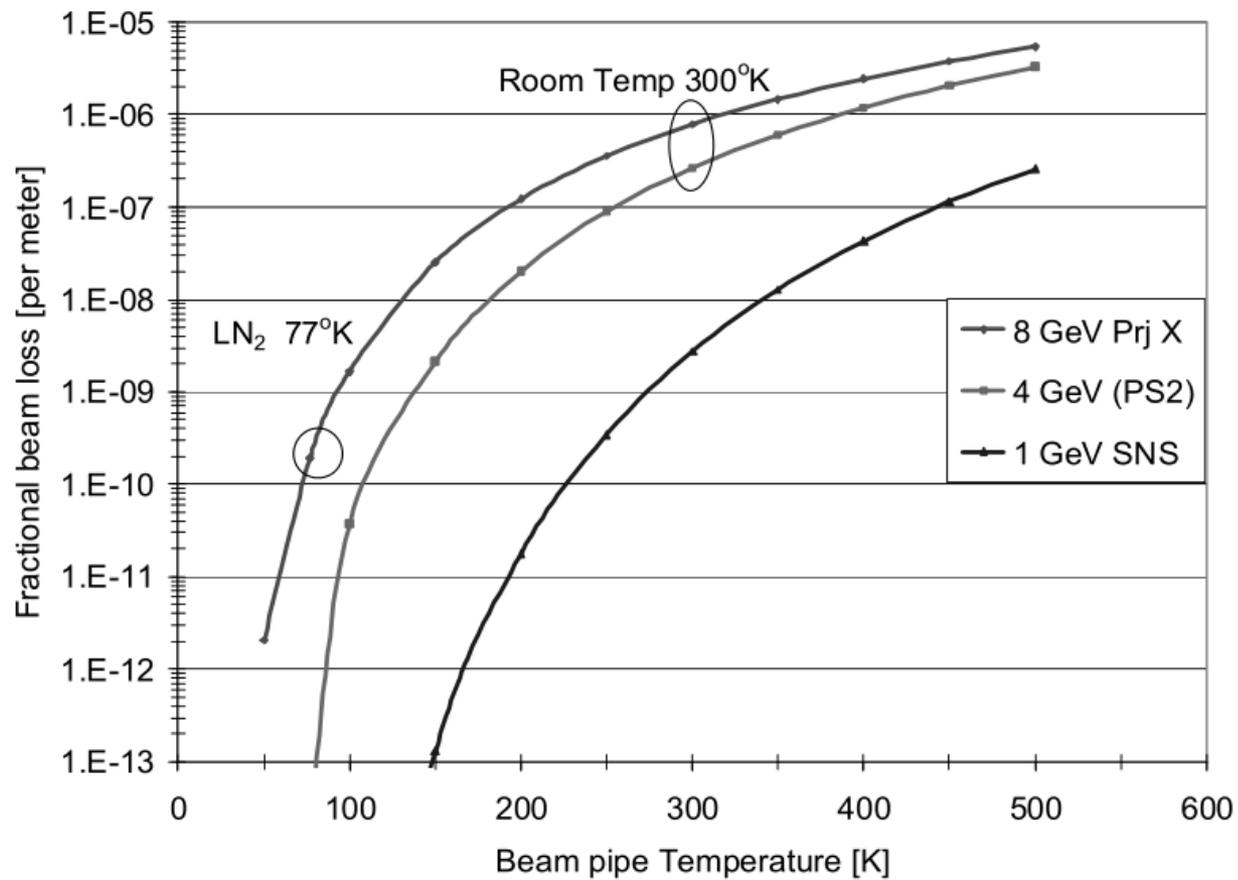
Black body radiation



[6] D. Johnson. Challenges Associated with 8 GeV H⁻ Transport and Injection for FERMILAB PROJECT-X *. (Proceedings of Hadron Beam 2008), 2008.

Black body radiation

Loss Rate vs Beam Pipe Temperature



Rest gas stripping

- Power loss due to stripping on rest gas per length l

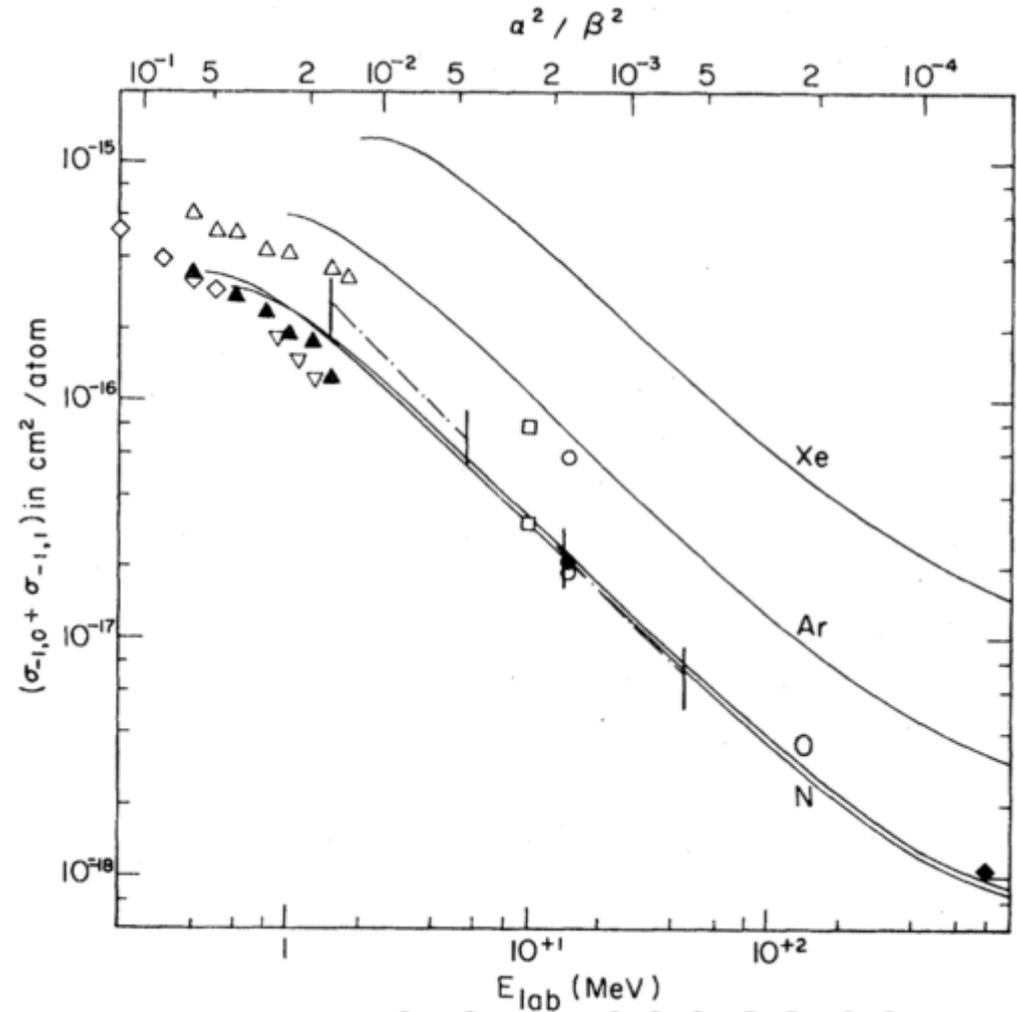
$$P = E I \frac{d\sigma}{d\Omega} \rho l$$

Beam energy
and intensity

Gas density
...fct of T, p

$$\frac{d\sigma}{d\Omega} = \frac{7 \cdot 10^{-19}}{\beta^2} \text{cm}^2 \text{ per atom of nitrogen or oxygen}$$

$$\frac{d\sigma}{d\Omega} = \frac{1 \cdot 10^{-19}}{\beta^2} \text{cm}^2 \text{ per atom of hydrogen}$$



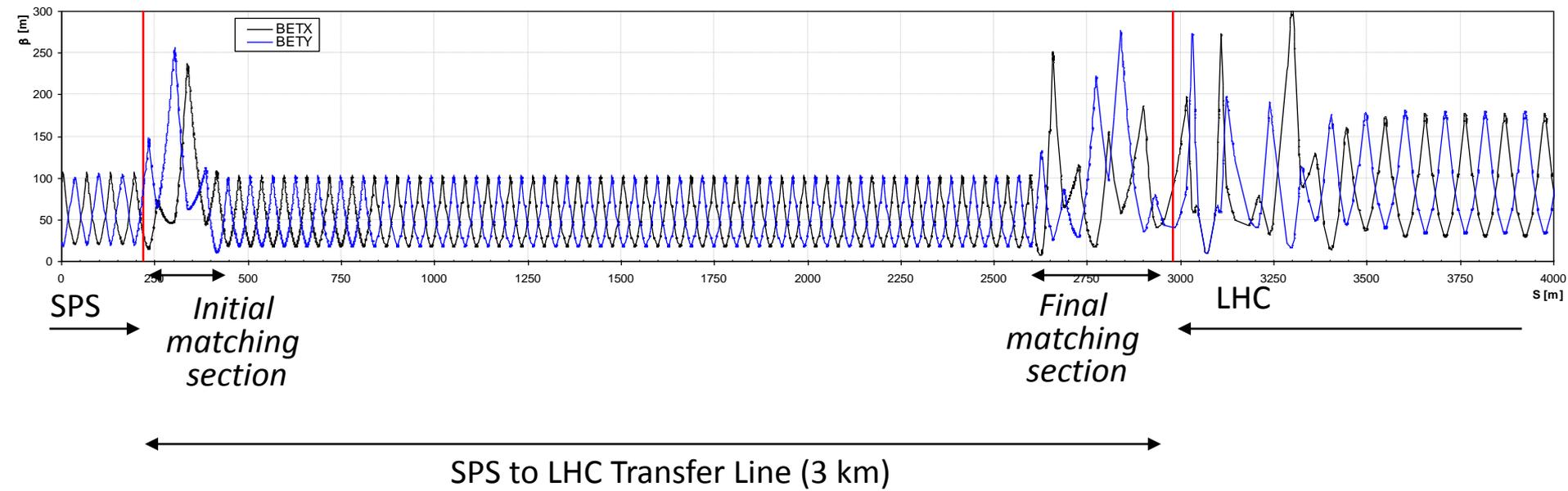
$$\text{Electron loss} \propto \frac{1}{\beta^2}$$

Example Fermilab Project X

- Loss rate from black body radiation at 300 K not acceptable
- Installing a cool beam screen (77 K)
 - Reduces black body radiation by factor ~ 16
 - Improves vacuum pumping (better than $1e-8$ Torr)
- Lorentz stripping limits dipole fields to 0.05 T

Focussing structure

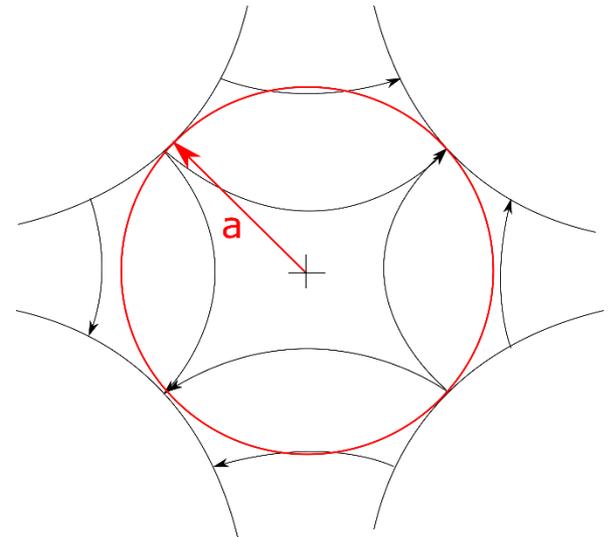
- Cell length optimised for dipole filling at extraction energy
- Can assume this as a good starting point for our transfer line
- For transfer lines often a 90 deg FODO structure is chosen
 - Good ratio of max/min in beta function
 - Same aperture properties
 - Provides good locations for trajectory correctors and instrumentation
 - Good phase advance for injection/extraction and protection equipment



Quadrupole field

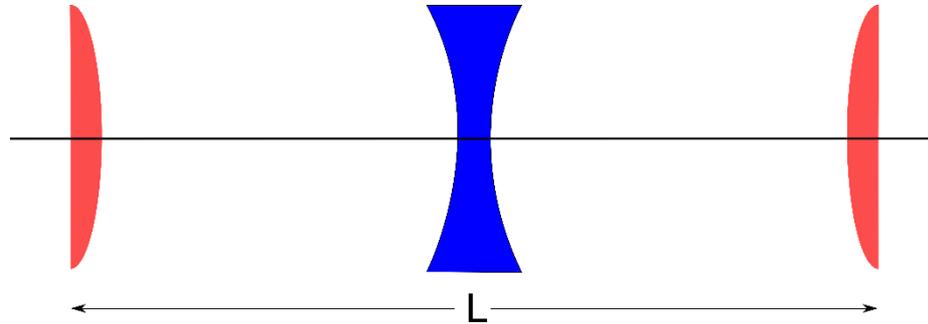
- What is needed to specify the quadrupole pole tip field:

$$B = g \cdot x \quad B_{poletip} = g \cdot a$$



- Need to define quadrupole gradient g [T/m] and pole radius a [m]

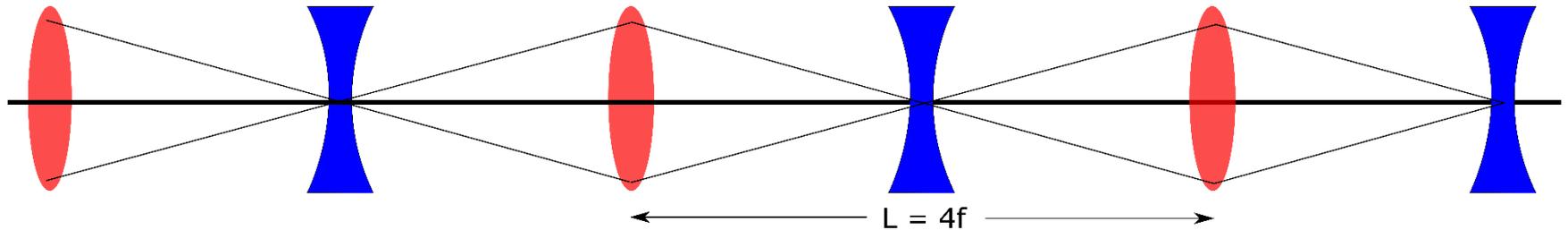
FODO cell



$$\begin{aligned} R &= \begin{pmatrix} 1 & 0 \\ 1/2f & 1 \end{pmatrix} \begin{pmatrix} 1 & L/2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -1/f & 1 \end{pmatrix} \begin{pmatrix} 1 & L/2 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 1/2f & 1 \end{pmatrix} \\ &= \begin{pmatrix} 1 - L^2/8f^2 & L - L^2/4f \\ -L/4f^2 - L^2/16f^3 & 1 - L^2/8f^2 \end{pmatrix} \\ &= \begin{pmatrix} \cos \mu + \alpha \sin \mu & \beta \sin \mu \\ -\gamma \sin \mu & \cos \mu - \alpha \sin \mu \end{pmatrix} \end{aligned}$$

FODO stability

Stability for: $f > \frac{L}{4}$

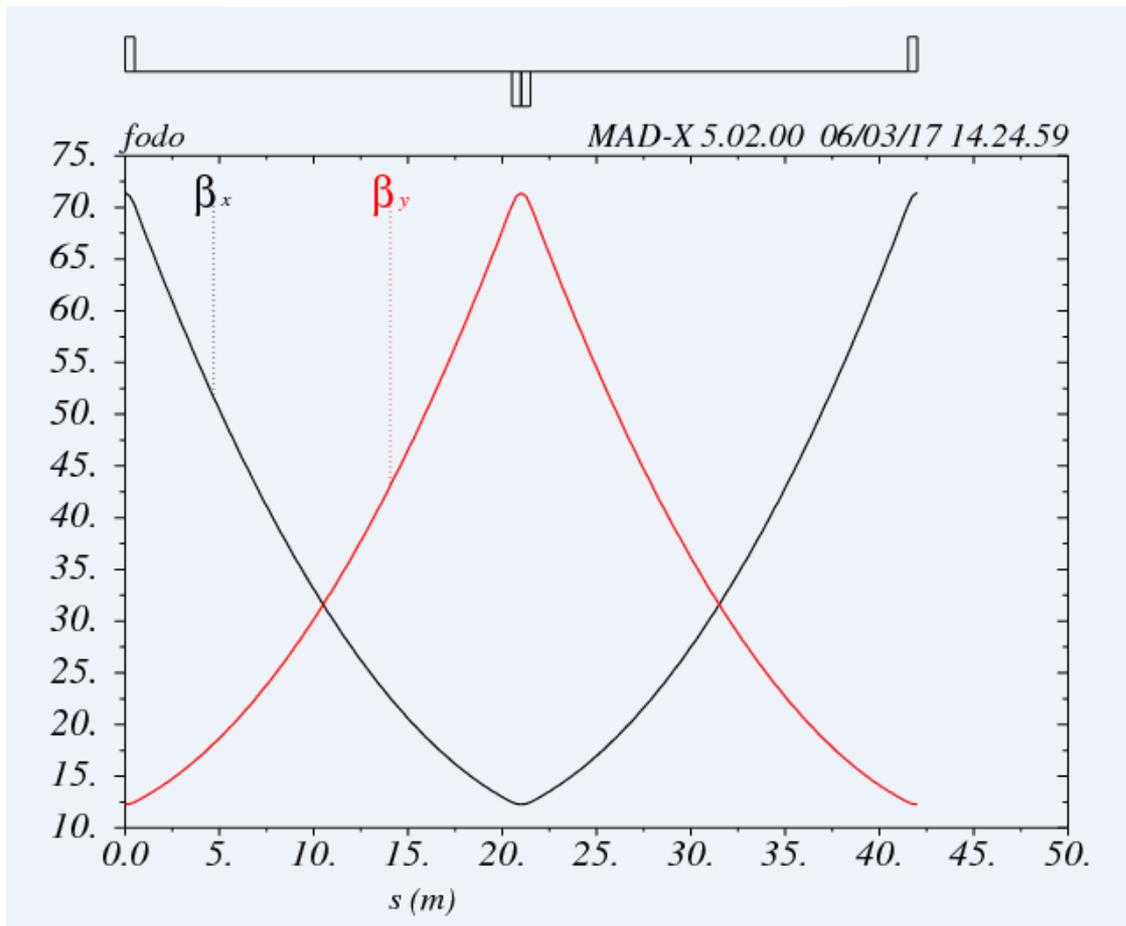


$$\frac{L}{f} = 4 \sin \frac{\mu}{2}$$

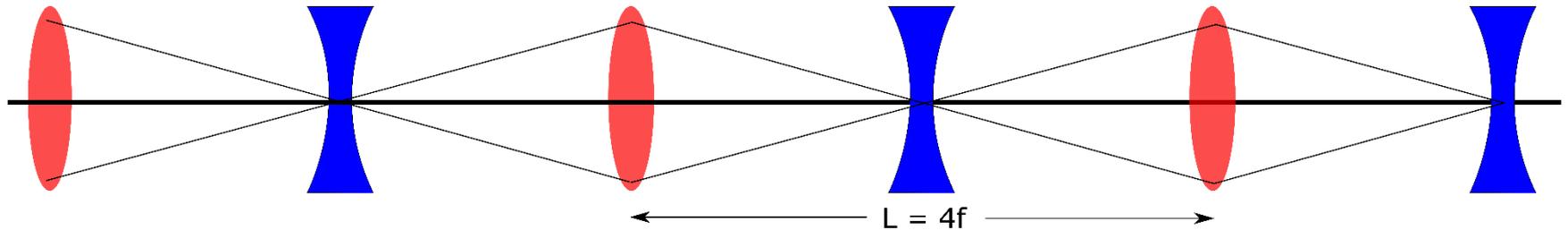
Estimate required gradient of quadrupoles

FODO optics

$$\mu = 90^\circ \Rightarrow f = \frac{L}{2\sqrt{2}}$$



FODO stability



$$\frac{L}{f} = 4 \sin \frac{\mu}{2}$$

Estimate required gradient of quadrupoles

$$\beta = \left(L + \frac{L^2}{4f} \right) / \sin \mu$$

Use maximum betatron function to estimate beam size and pole tip field of quadrupoles

Apertures

- Arbitrary choice and depends on beam energy (destructive?)
- While in a collider have O(10 sig + few mm)
- Less in TLs

$$A_{x,y} = \pm n_{sig} \cdot \sqrt{k_{\beta} \cdot \beta_{x,y} \cdot \frac{\epsilon_{x,y}}{\beta\gamma}} \pm D_{x,y} \cdot k_{\beta} \cdot \frac{\Delta p}{p} \pm CO \cdot \sqrt{\frac{\beta_{x,y}}{\beta_{xmax,ymax}}} \pm \text{align}$$

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Optics uncertainty in
TLs vs rings

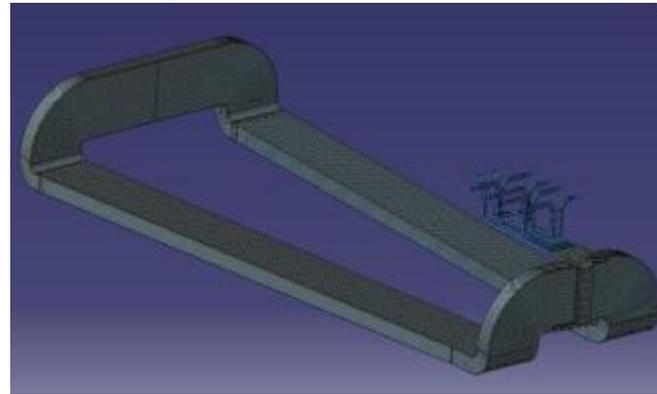
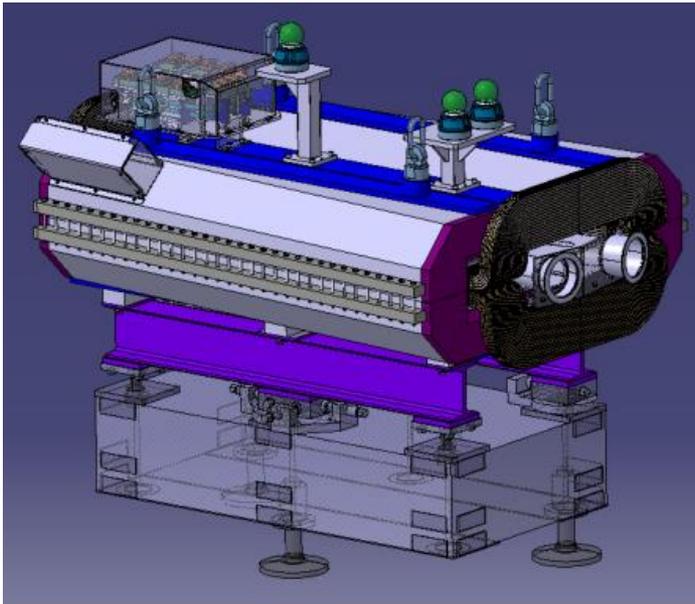
Conservative approach
But be aware when
you specify minimum
beam sizes

$$x = \theta \cdot \sqrt{\beta_1 \beta_2} \cdot \sin(\mu_2 - \mu_1)$$

Apertures

- Arbitrary choice and depends on beam energy (destructive?)
- While in a collider have $O(10 \text{ sig} + \text{few mm})$
- Less in TLs

$$A_{x,y} = \pm n_{sig} \cdot \sqrt{k_{\beta} \cdot \beta_{x,y} \cdot \frac{\epsilon_{x,y}}{\beta\gamma}} \pm D_{x,y} \cdot k_{\beta} \cdot \frac{\Delta p}{p} \pm CO \cdot \sqrt{\frac{\beta_{x,y}}{\beta_{xmax,ymax}}} \pm \text{align}$$



Aperture calculation examples

- Low energy transfer lines (100 keV, ELENA)

- 10 mm trajectory variation
- 4 mm alignment

$$\Sigma_{x,y}(95\%) = \sqrt{(\beta_{x,y} \kappa_{x,y} \epsilon_{x,y}(95\%))} + (D_{x,y} \kappa_{x,y} \frac{\delta p}{p}(95\%))$$

$$\kappa_{x,y} = 1.2$$

- Medium energy (1-2 GeV, PS Booster)

- $A_{x,y} = \pm n_{sig} \cdot \sqrt{k_{\beta} \cdot \beta_{x,y} \cdot \frac{\epsilon_{x,y}}{\beta\gamma}} \pm D_{x,y} \cdot k_{\beta} \cdot \frac{\Delta p}{p} \pm CO \cdot \sqrt{\frac{\beta_{x,y}}{\beta_{xmax,ymax}}} \pm \text{align}$
- $n_{sig} = 3, k_{\beta} = 1.2, CO = 3 \text{ mm}, \text{align} = 0$ (usually $\sim 2 \text{ mm}$)

- High energy transfer lines (0.45 – 3 TeV, LHC, FCC)

- $n_{sig} = 6, k_{\beta} = 1.0, CO = 1.5 \text{ mm}$

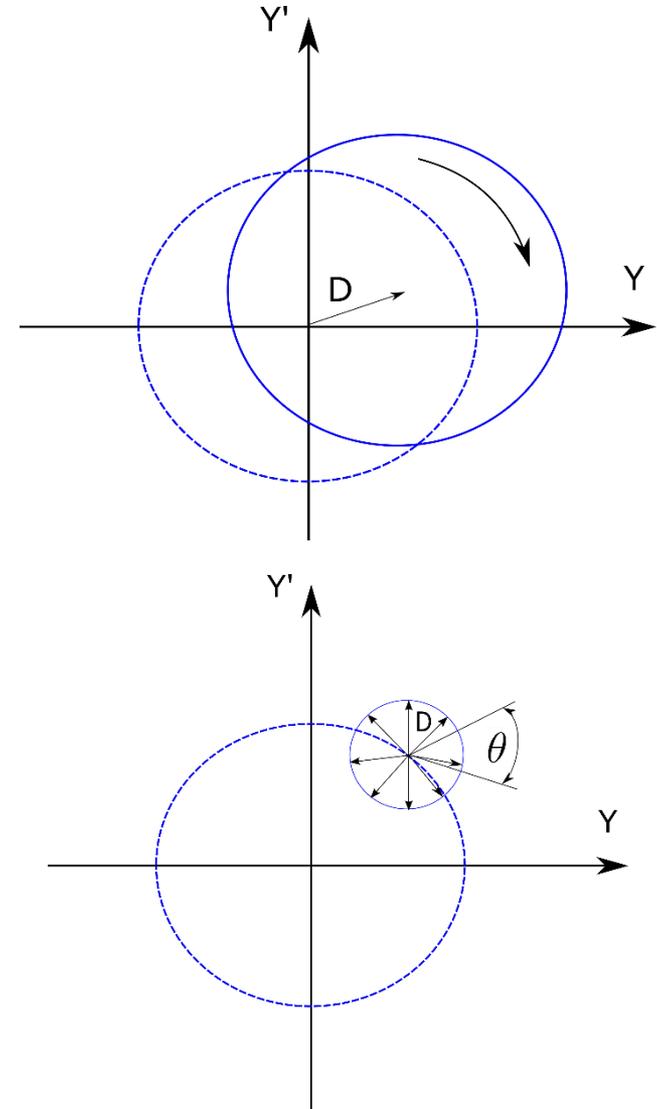
First estimate of field error specification

- Impact of field errors on aperture requirements should be negligible
- Beam quality is the constraint – emittance growth

$$Y_2 = Y_1 + D \cos \theta$$

$$Y_2^2 = Y_1^2 + Y_1 D \cos \theta + D^2 \cos^2 \theta$$

$$\langle Y_2^2 \rangle = \langle Y_1^2 \rangle + 2 \langle Y_1 D \cos \theta \rangle + \langle D^2 \cos^2 \theta \rangle$$



First estimate of field error specification

$$2 \langle Y_1 D \cos \theta \rangle = 2 \langle Y_1 \rangle \langle D \cos \theta \rangle$$

D and Y_1 are uncorrelated

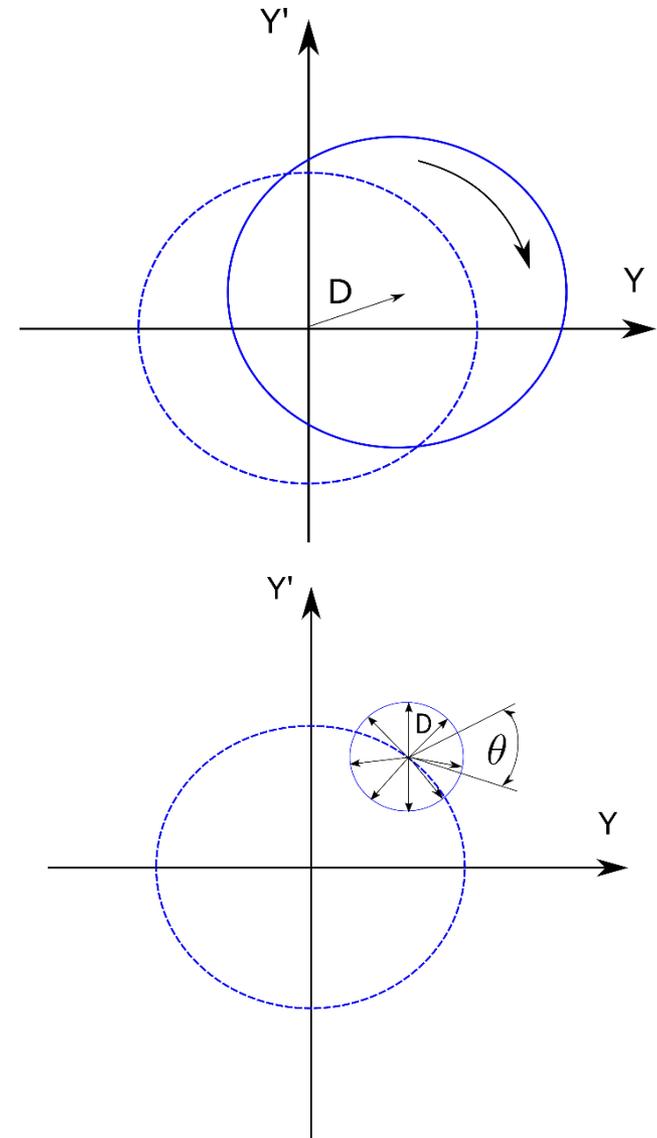
And averaging over the constant D gives 0

$$\langle Y_2^2 \rangle = \langle Y_1^2 \rangle + 2 \langle Y_1 D \cos \theta \rangle + \langle D^2 \cos^2 \theta \rangle$$

$$\langle Y_2^2 \rangle = \langle Y_1^2 \rangle + \frac{1}{2} D^2$$

This is valid for any point P on any circle...

$$\epsilon_2 = \epsilon_1 + \frac{1}{2} D^2$$



First estimate of field error specification

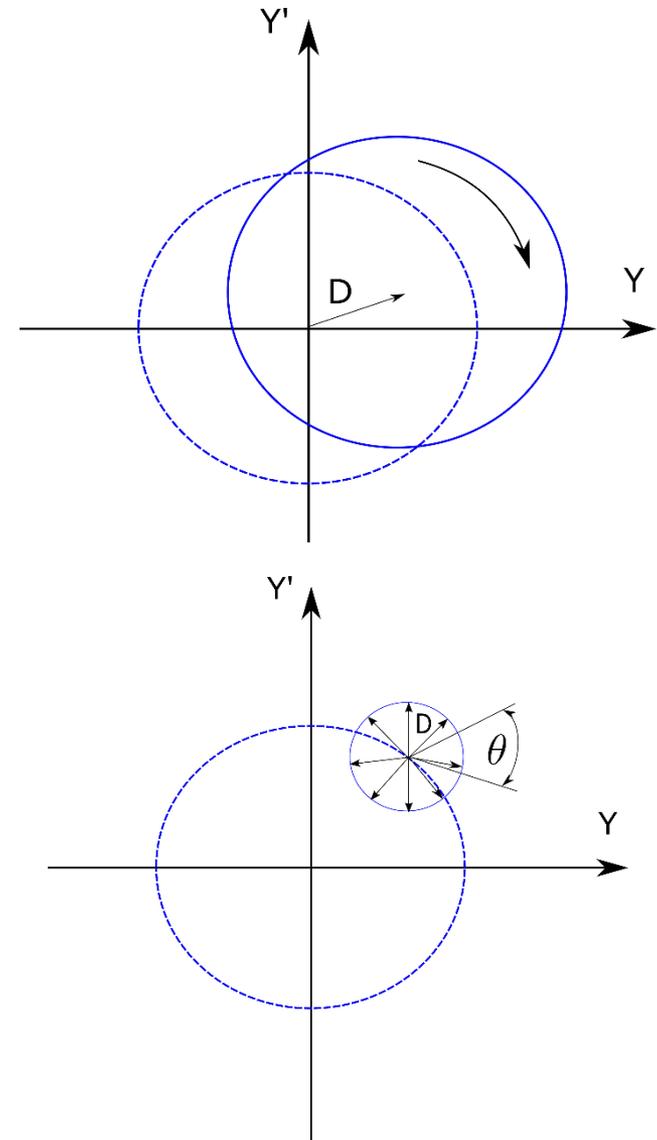
$$D^2 = (\Delta Y)^2 + (\Delta Y'^2) = (\Delta y)^2 \frac{1+\alpha^2}{\beta} + (\Delta y')^2 \beta$$

$$\epsilon_2 = \epsilon_1 + \frac{1}{2} \left((\Delta y)^2 \frac{1+\alpha^2}{\beta} + (\Delta y')^2 \beta \right)$$

Magnet
misalignment

Dipole field error

$$\Delta y' = \frac{\Delta B l}{B \rho}$$

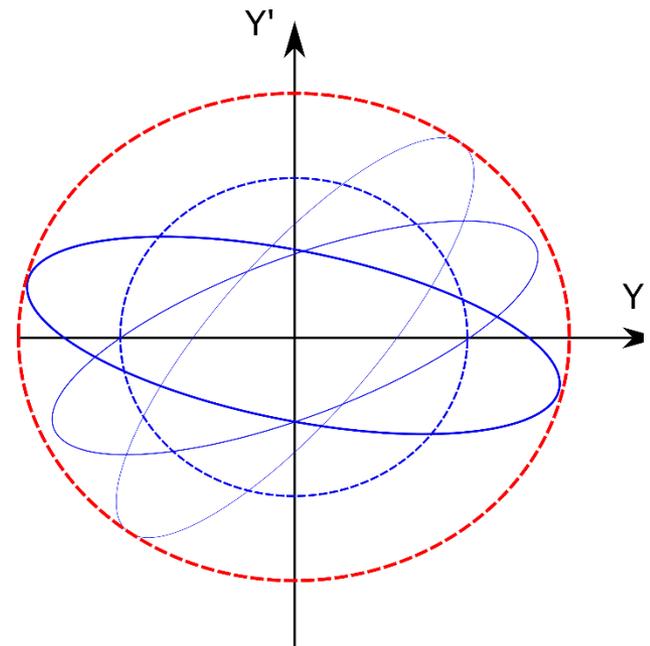
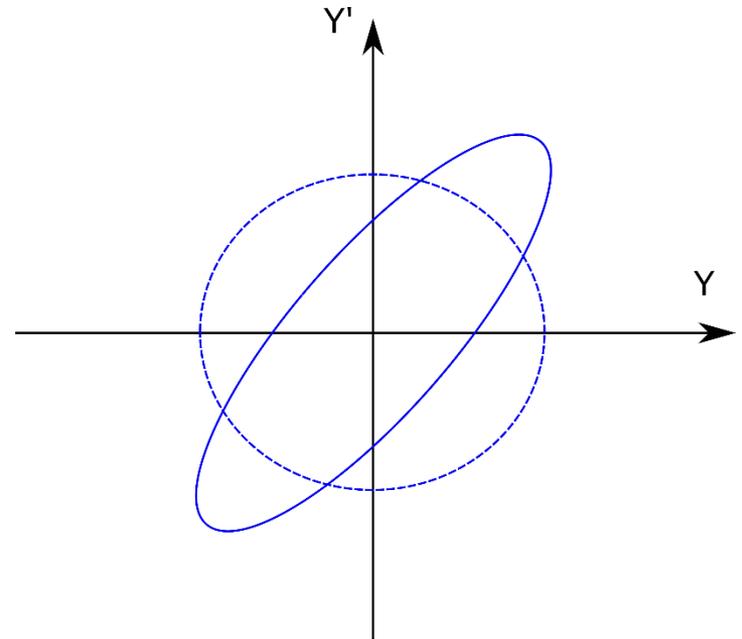


Gradient errors

$$k = -\frac{\Delta Gl}{B\rho}$$



$$\epsilon_2 = \frac{1}{2} (k^2 \beta^2 + 2) \epsilon_1$$



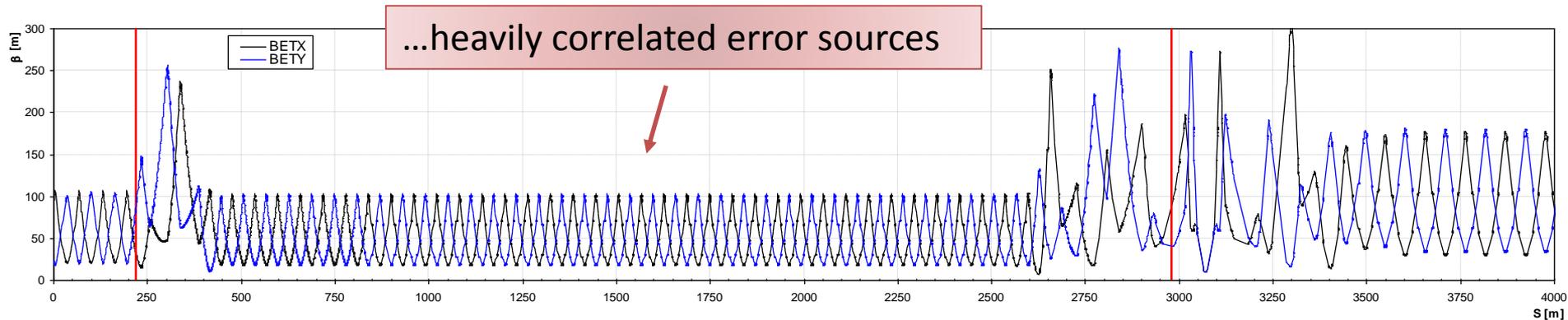
Combining errors

- Averaging over a distribution of uncorrelated errors

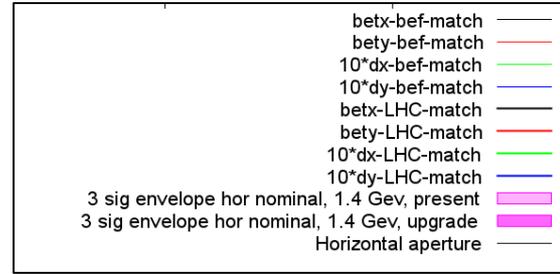
$$\langle y^2 \rangle = \frac{\beta(s)}{2} \sum_n \beta_n \langle \delta_n^2 \rangle \quad \delta_n = \frac{\Delta B l}{B_0 \rho_0} \quad \text{and} \quad \delta_n = -lk \Delta y$$

$$\langle y^2 \rangle = \frac{\beta_{aver} n \delta_{rms}^2}{2}$$

- But here we have to be very careful...

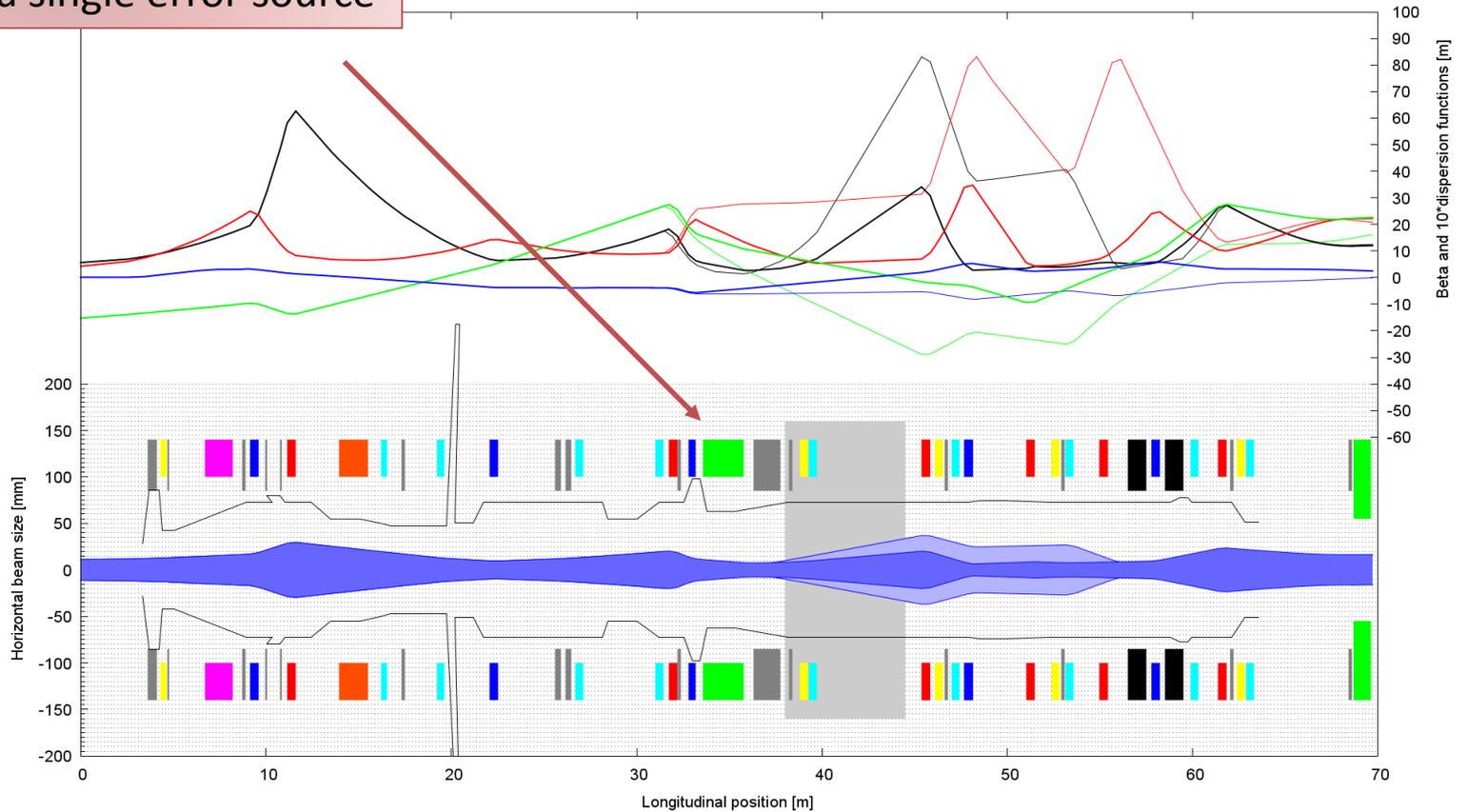


Combining errors example



Not correlated at all but dominated by a single error source

BT-BTP4: from PSB ej to PS inj, optics in [m] and horizontal beam envelope in [mm]



Typical specifications from correction studies

- Number of monitors and required resolution
 - Every $\frac{1}{4}$ betatron wavelength
 - Grid resolution: ~ 3 wires/sigma
- Number of correctors and strength
 - Every $\frac{1}{2}$ betatron wavelength H - same for V
 - Displace beam by few betatron sigma per cell
- Dipole and quadrupole field errors
 - Integral main field known to better than $1-10e-4$
 - Higher order field errors $< 1-10e-4$ of the main field
- Dynamic errors from power converter stability
 - $1-10e-5$
- Alignment tolerances
 - 0.1-0.5 mm
 - 0.1-0.5 mrad

Take values with caution!

They can strongly vary depending on energy, intensity, machine purpose, etc.

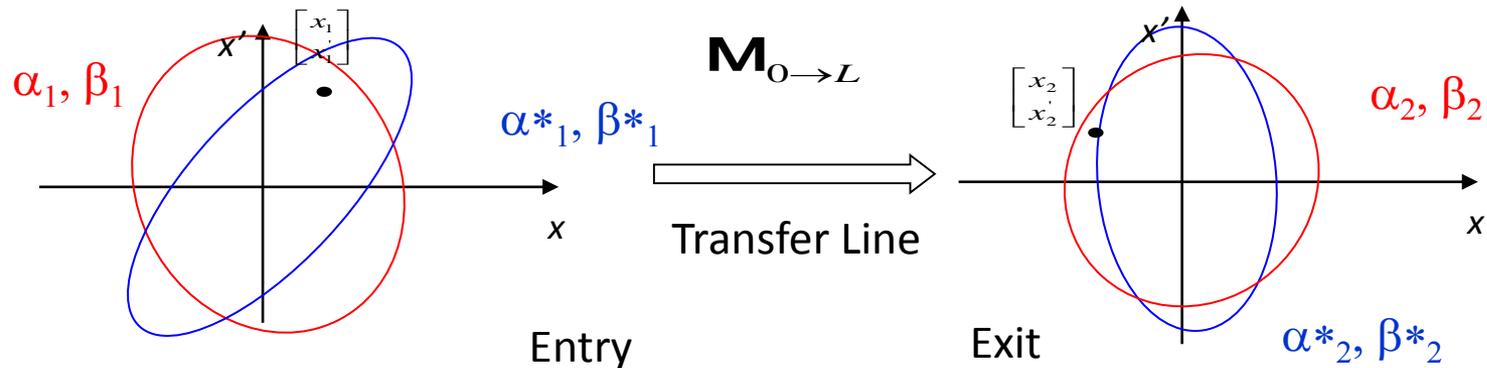
Summary

Before switching on a computer we can define for a transfer line:

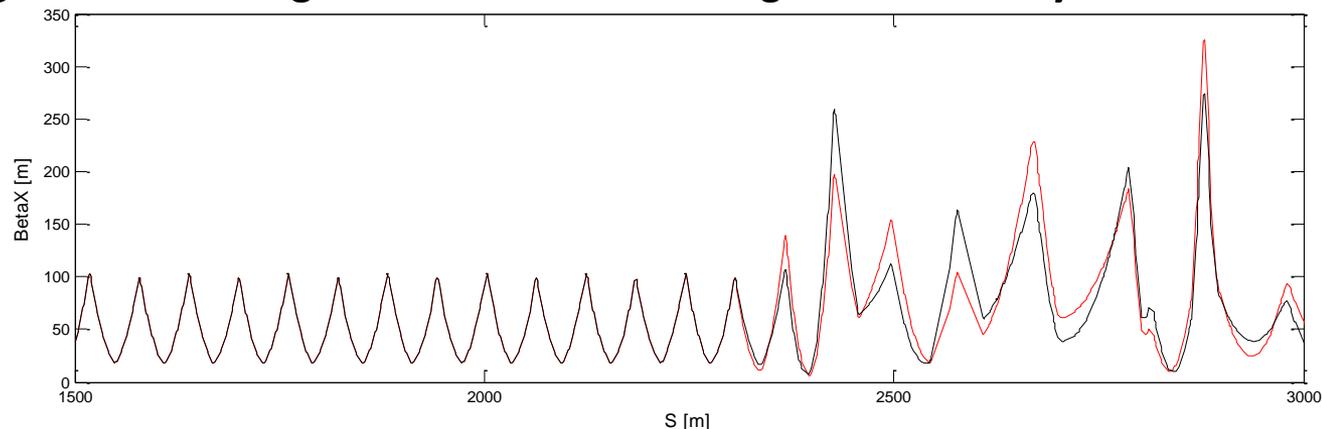
- Number of dipoles and quadrupoles, correctors and monitors
- Dipole field and quadrupole pole tip field
- Aperture of magnets and beam instrumentation
- Rough estimate of required field quality and alignment accuracy

Wrap up

- Optics in a ring is defined by ring elements and **periodicity** – optics in a transfer line is dependent line elements and **initial conditions**



- Changes** of the strength of a transfer line magnet **affect only downstream optics**



Wrap up

- Geometry calculations require a set of coordinates in a common reference frame
- **Bending fields** are defined by geometry and the **magnetic or electric rigidity**:

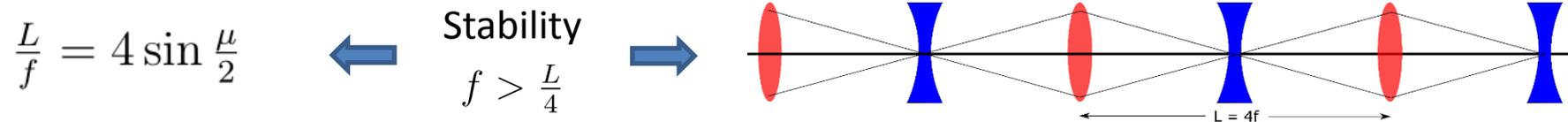
$$B\rho [Tm] = 3.3356 \frac{A}{n} p [GeV/c] \quad E\rho [kV] = \frac{\gamma+1}{\gamma} \frac{A}{n} T [keV]$$

$$\theta = \frac{Bdl}{B\rho} \text{ or } \frac{Edl}{E\rho}$$

- The choice between magnetic and electric depends mainly on the beam energy
- If you are in the grey zone, consider: field design and measurement, power consumption, vacuum, interlocking
- For the estimates of bending radii in lines remember to take into account the filling factor (~70%) and Lorentz-Stripping in case of H⁻ ions

Wrap up

- **Quadrupole gradients and apertures** can be estimated in case of simple focussing structure like FODO cells



$\beta = \left(L + \frac{L^2}{4f} \right) / \sin \mu$
→
Defines beam size and quadrupole pole tip field

- **Aperture specifications** require safety factors for the optics and constant contributions for trajectory variations and alignment errors

- $A_{x,y} = \pm n_{sig} \cdot \sqrt{k_{\beta} \cdot \beta_{x,y} \cdot \frac{\epsilon_{x,y}}{\beta_{\gamma}}} \pm D_{x,y} \cdot k_{\beta} \cdot \frac{\Delta p}{p} \pm CO \cdot \sqrt{\frac{\beta_{x,y}}{\beta_{xmax,ymax}}} \pm \text{alignment}$

Wrap up

- **Estimating tolerances** from emittance growth:

- Dipole field and alignment:
$$\epsilon_2 = \epsilon_1 + \frac{1}{2} \left((\Delta y)^2 \frac{1+\alpha^2}{\beta} + (\Delta y')^2 \beta \right)$$

The diagram illustrates the components of the emittance growth equation. A box labeled "Magnet misalignment" has an arrow pointing to the $(\Delta y)^2$ term in the equation. Another box labeled "Dipole field error" has an arrow pointing to the $(\Delta y')^2$ term. Below the "Dipole field error" box is a smaller box containing the equation $\Delta y' = \frac{\Delta B l}{B \rho}$.

- Gradient errors:

$$\epsilon_2 = \frac{1}{2} (k^2 \beta^2 + 2) \epsilon_1 \quad k = -\frac{\Delta G l}{B \rho}$$

Thank you for your attention

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