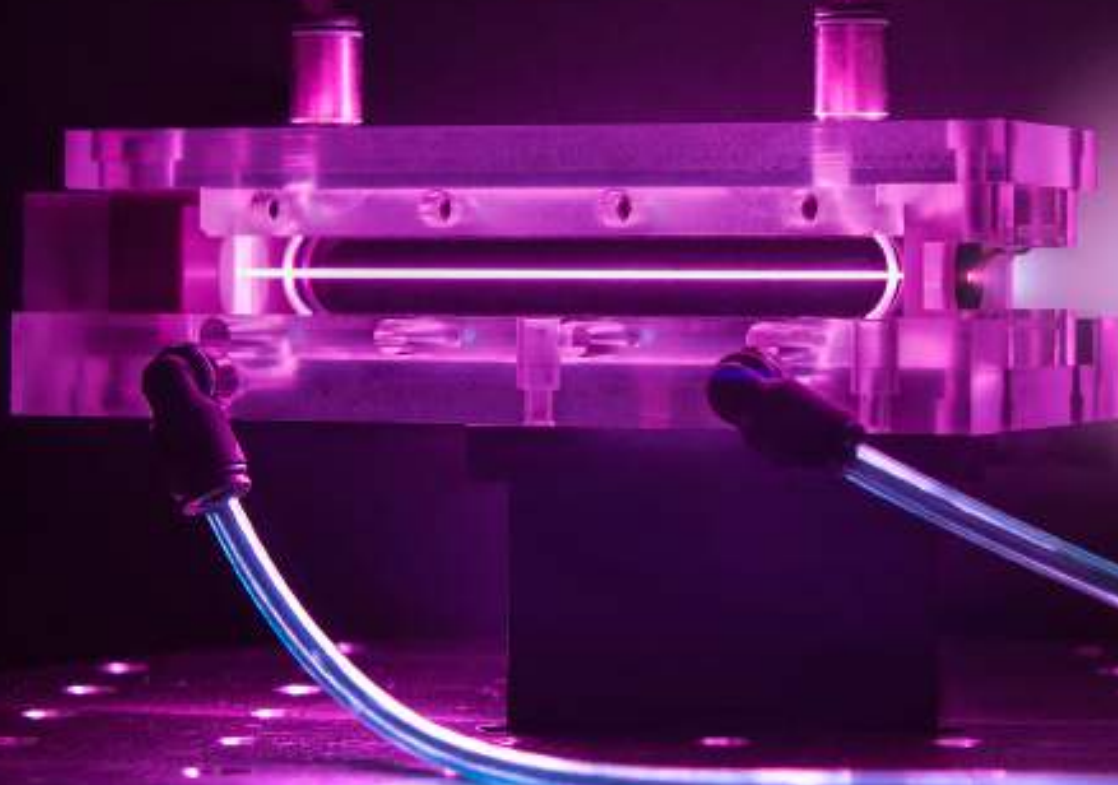


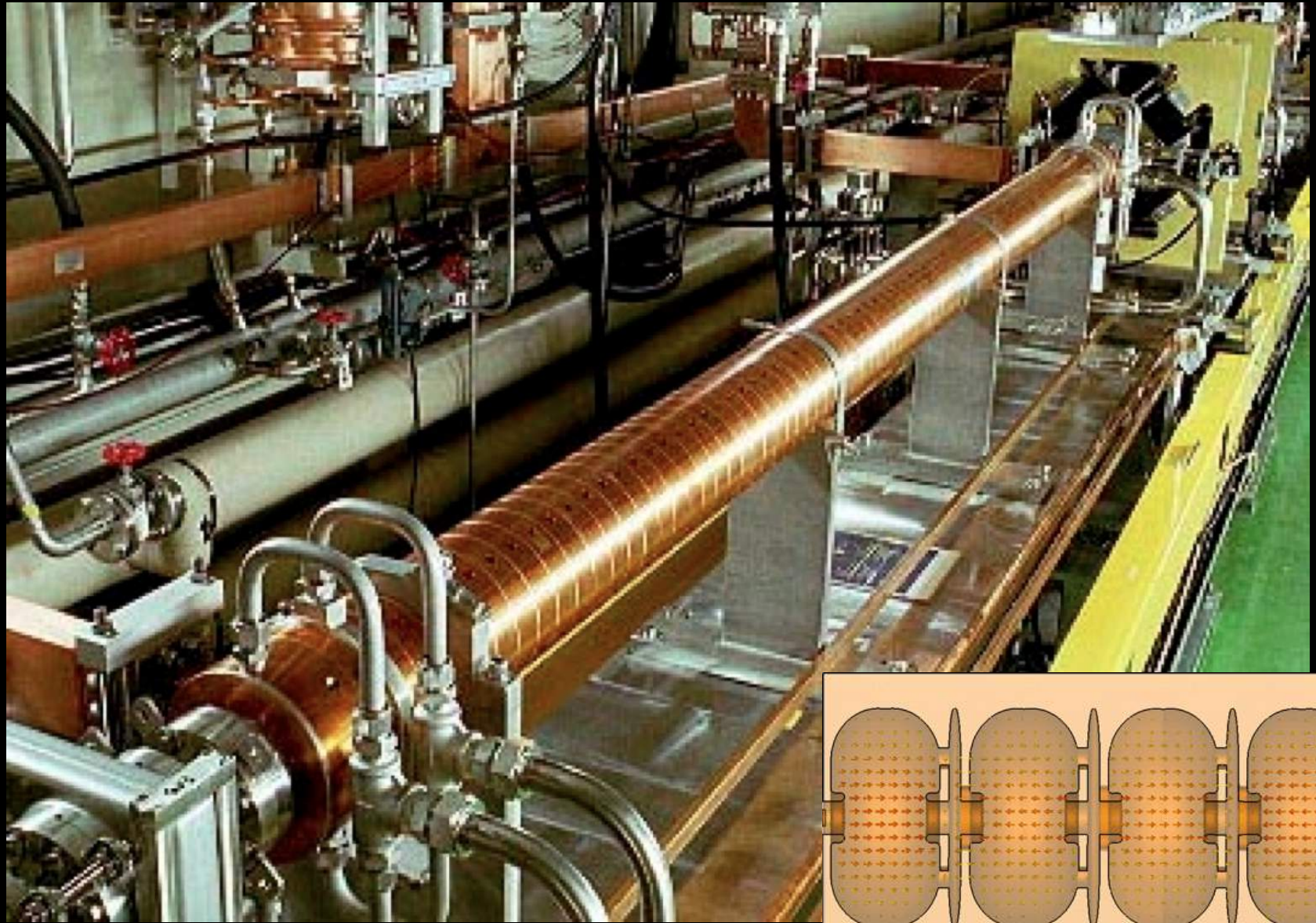
# Beam Driven Plasma Acceleration

Massimo.Ferrario@LNF.INFN.IT



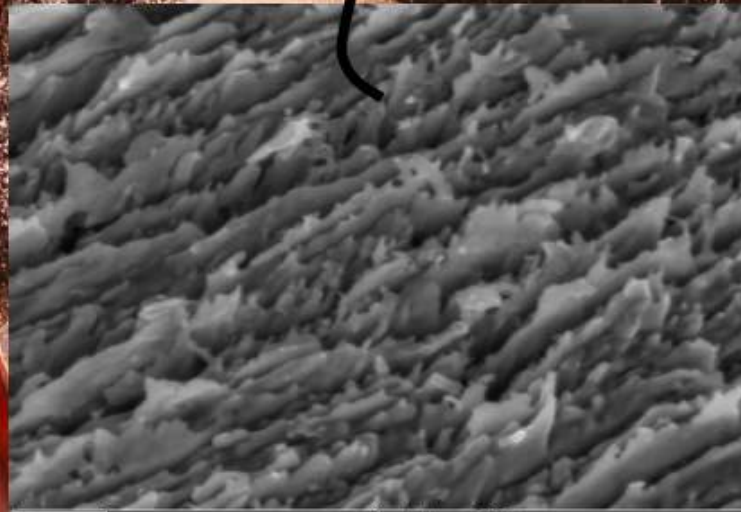
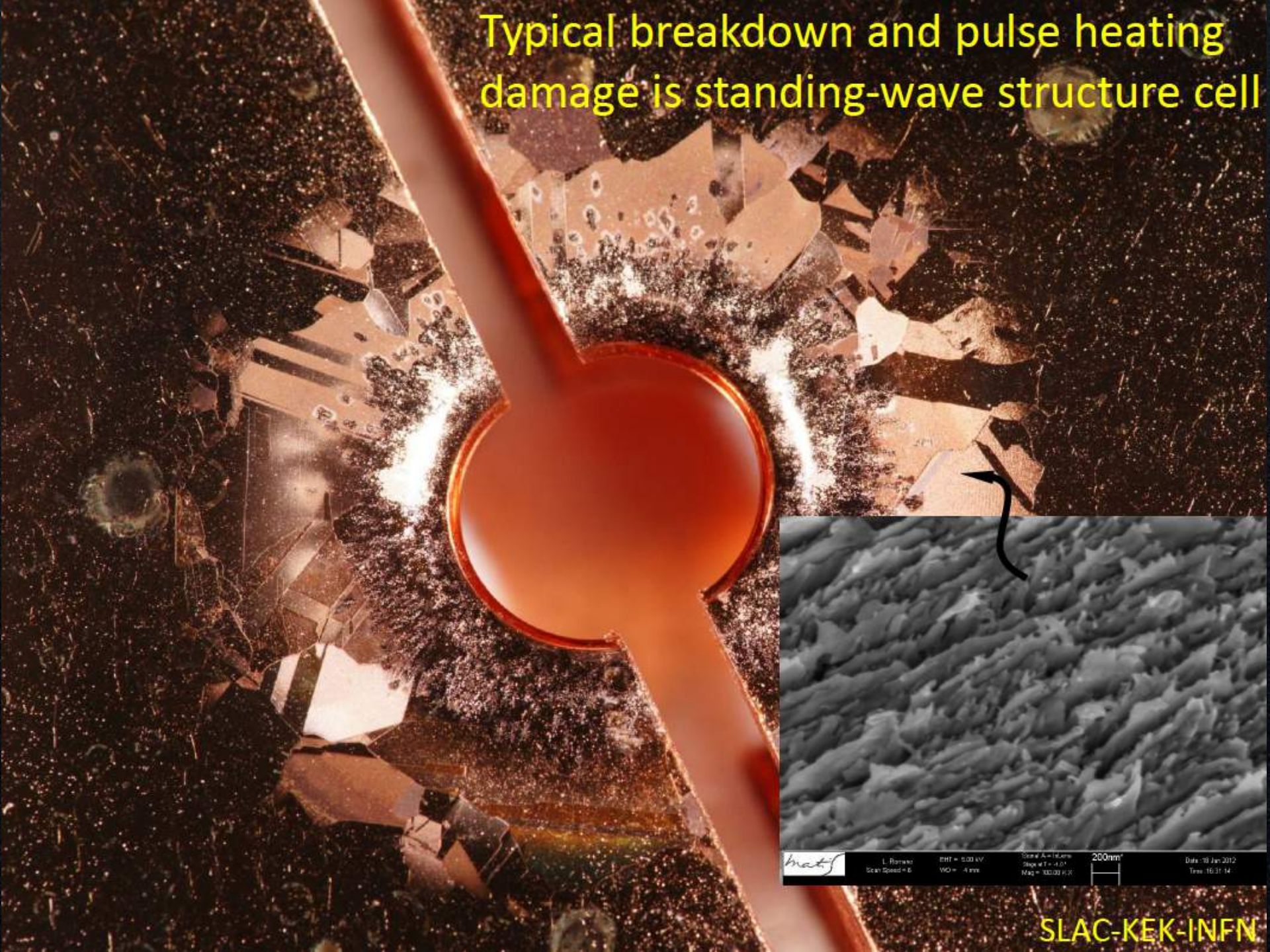
Egham - September 9<sup>th</sup> 2017

# Conventional RF accelerating structures





Typical breakdown and pulse heating damage is standing-wave structure cell

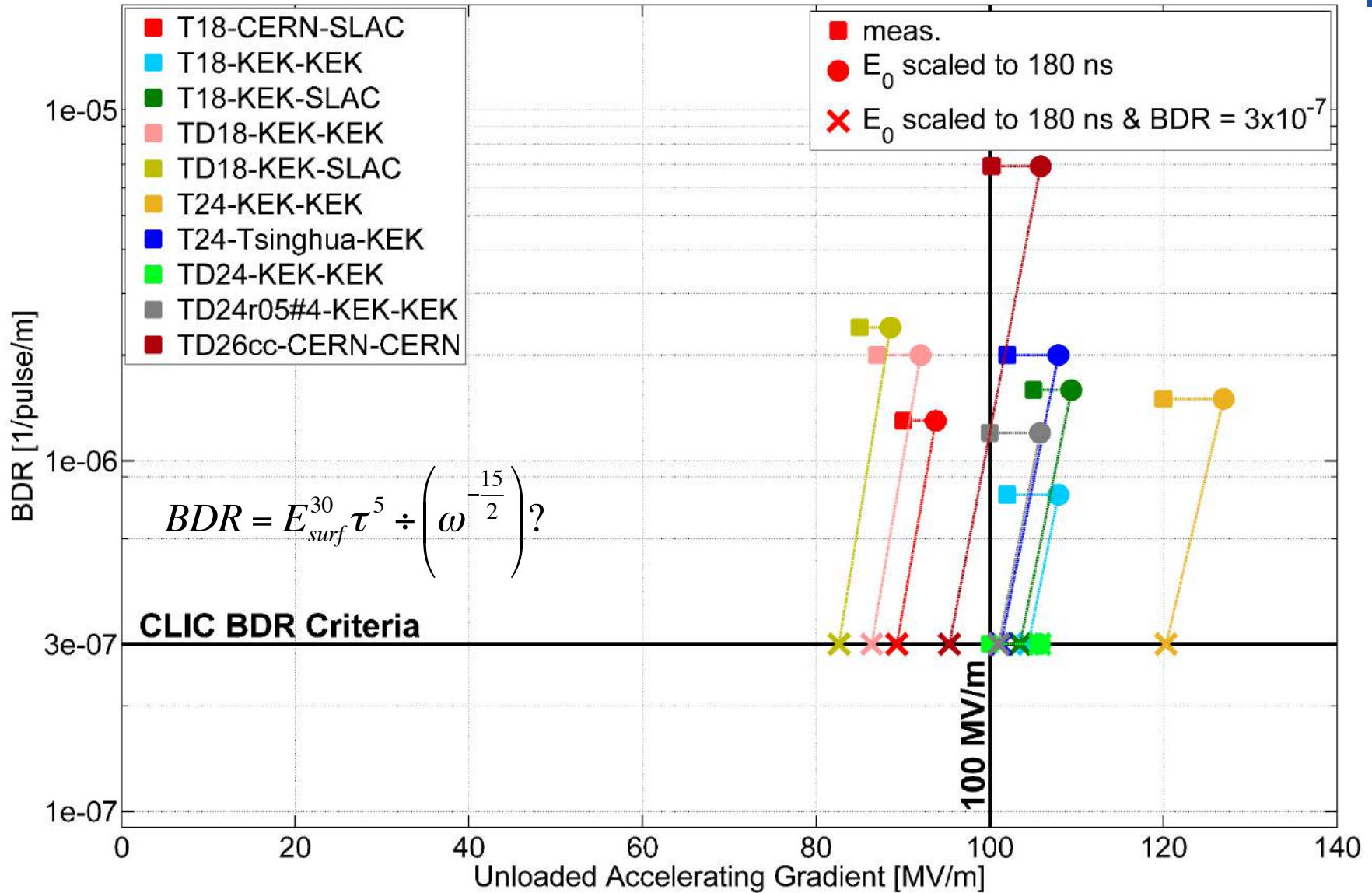


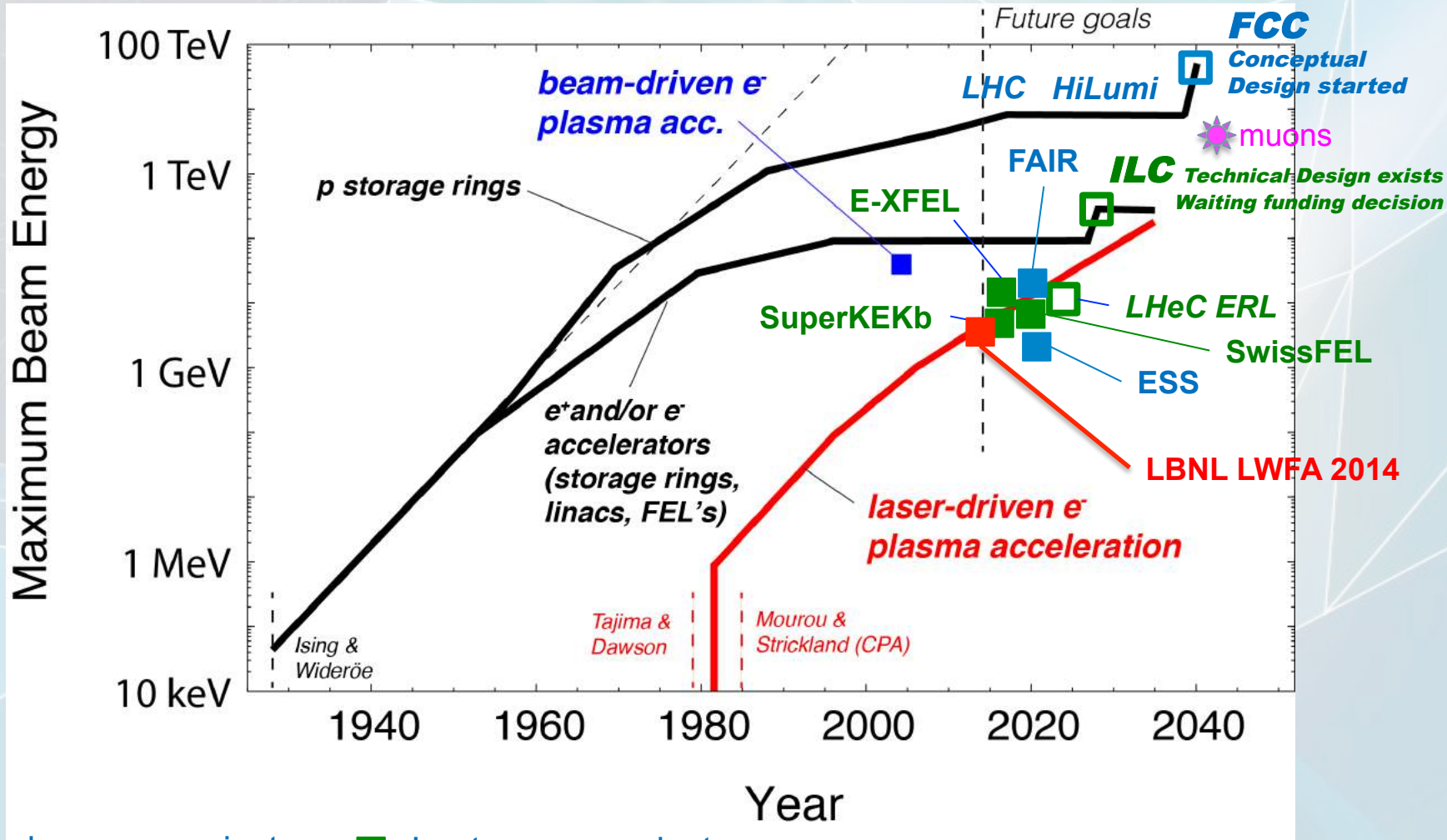
met  
L. Borrero  
Scan Speed = 8  
EHT = 5.00 kV  
WD = 4 mm  
Stral. A. = 15.0 um  
Spot #1 = 4.01  
Mag = 330.00 KX  
200um  
Date: 10 Jun 2012  
Time: 16:31:14





# Performance summary at CLIC specifications





- Hadron acc. project
- Lepton acc. project
- Hadron acc. proposal
- Lepton acc. proposal

## Laser Electron Accelerator

T. Tajima and J. M. Dawson

*Department of Physics, University of California, Los Angeles, California 90024*

(Received 9 March 1979)

An intense electromagnetic pulse can create a weak of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density  $10^{18}\text{W/cm}^2$  shone on plasmas of densities  $10^{18}\text{cm}^{-3}$  can yield giga-electronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.

## Acceleration of Electrons by the Interaction of a Bunched Electron Beam with a Plasma

Pisin Chen<sup>(a)</sup>

*Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305*

and

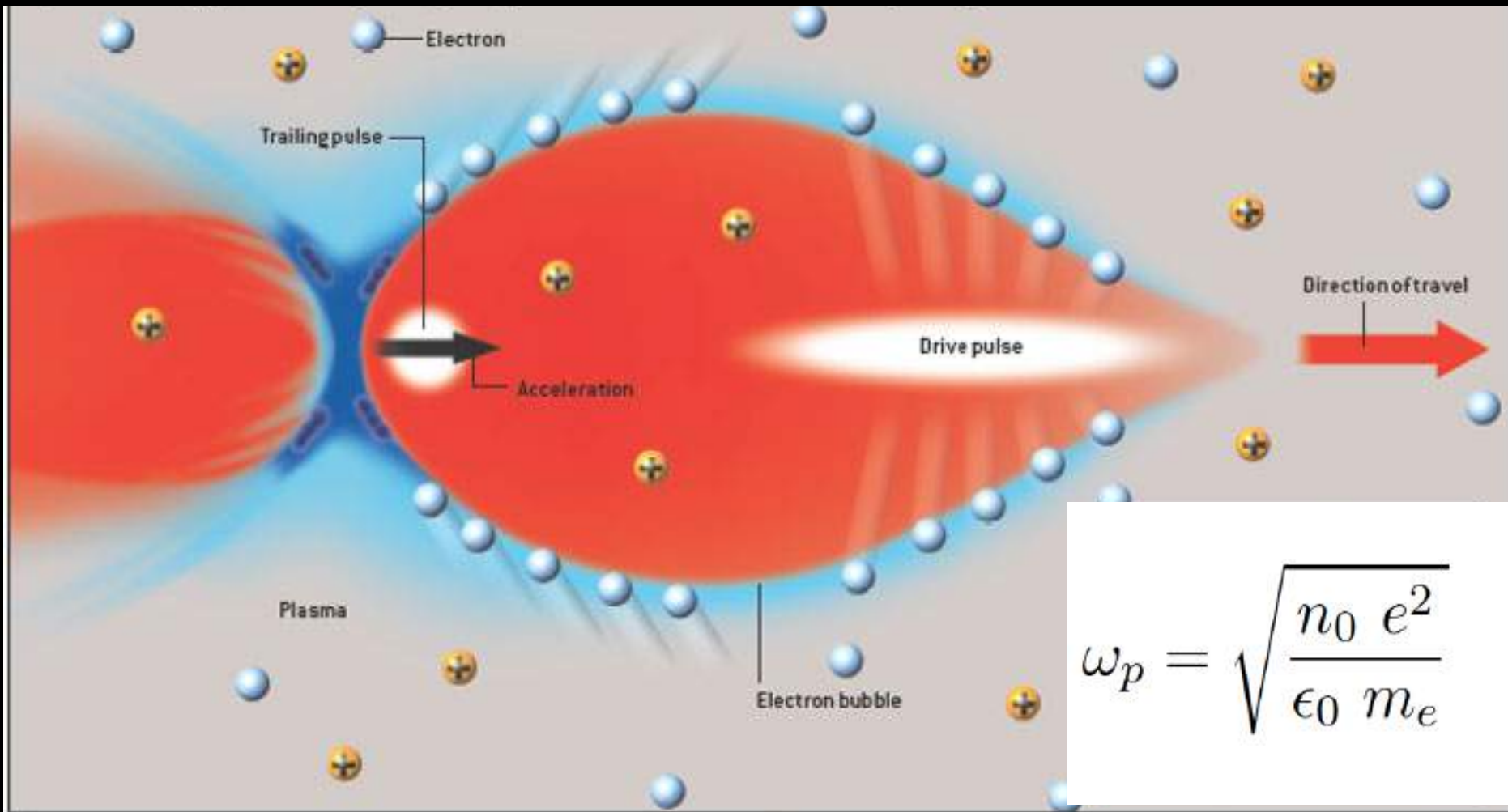
J. M. Dawson, Robert W. Huff, and T. Katsouleas

*Department of Physics, University of California, Los Angeles, California 90024*

(Received 20 December 1984)

A new scheme for accelerating electrons, employing a bunched relativistic electron beam in a cold plasma, is analyzed. We show that energy gradients can exceed  $1\text{ GeV/m}$  and that the driven electrons can be accelerated from  $\gamma_0 mc^2$  to  $3\gamma_0 mc^2$  before the driving beam slows down enough to degrade the plasma wave. If the driving electrons are removed before they cause the collapse of the plasma wave, energies up to  $4\gamma_0 mc^2$  are possible. A noncollinear injection scheme is suggested in order that the driving electrons can be removed.

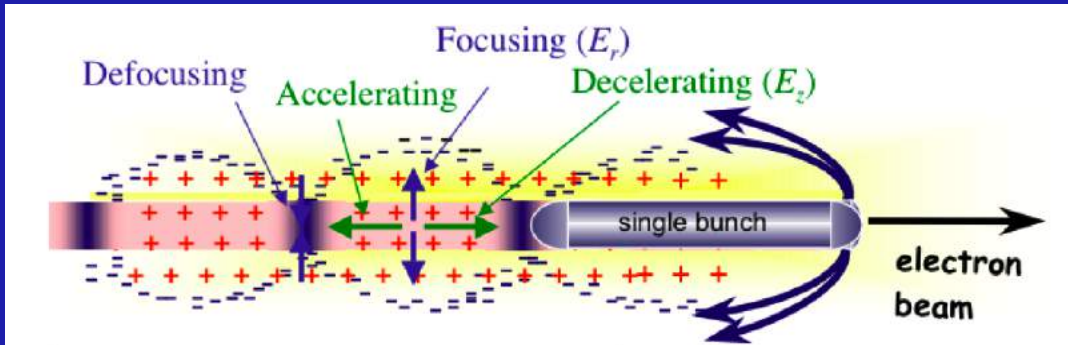




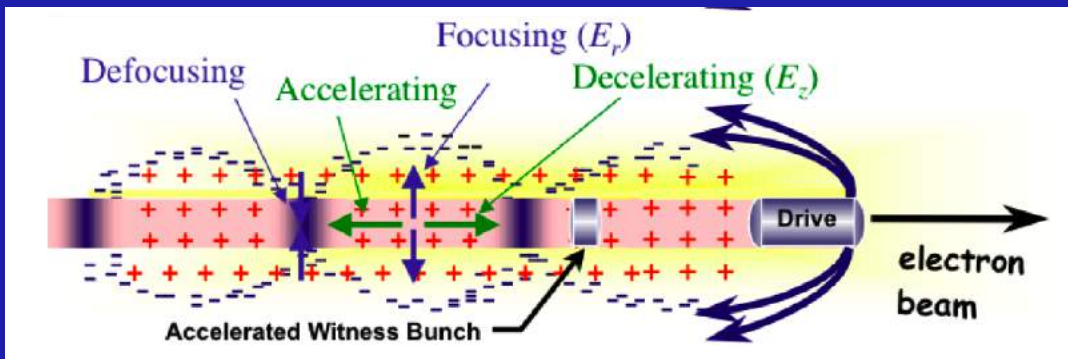
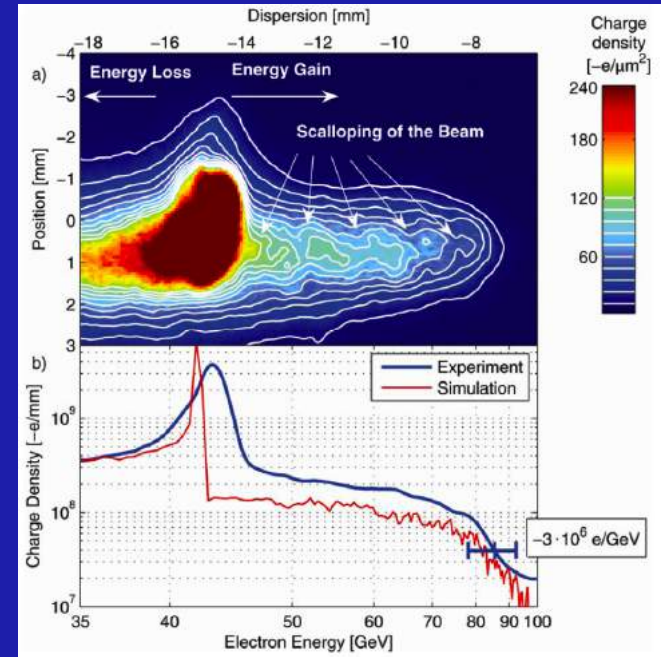
$$\omega_p = \sqrt{\frac{n_0 e^2}{\epsilon_0 m_e}}$$

$$\lambda_p \approx 1\text{mm} \cdot \sqrt{\frac{10^{15}\text{cm}^{-3}}{n_0}}$$

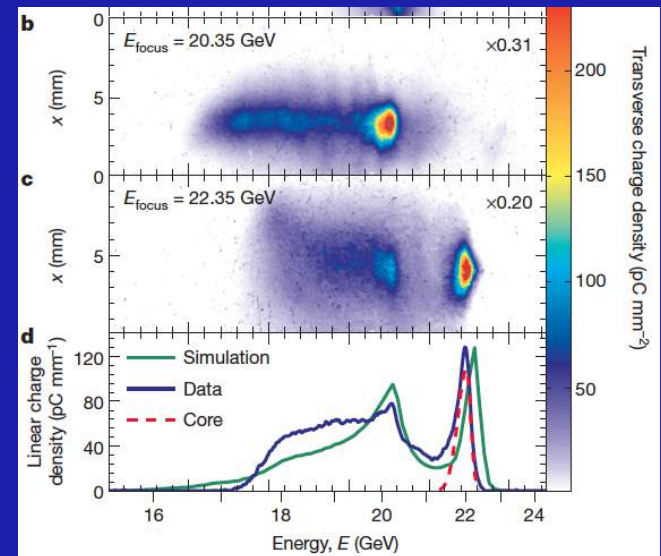
**0.3 mm for  $n_0 = 10^{16}\text{cm}^{-3}$**



Blumenfeld, I. et al. *Energy doubling of 42 GeV electrons in a metre-scale plasma wakefield accelerator.* *Nature* 445, 741–744 (2007).



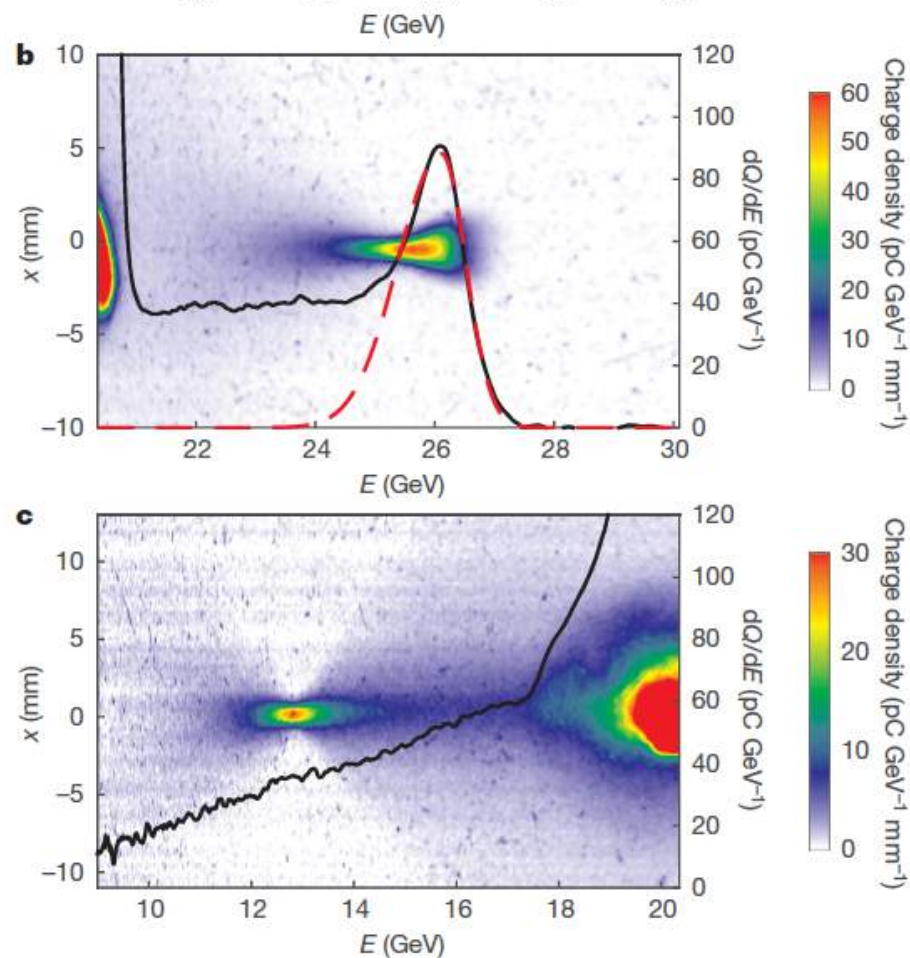
Litos, M. et al. *High-efficiency acceleration of an electron beam in a plasma wakefield accelerator.* *Nature* 515, 92–95 (2014).





# Multi-gigaelectronvolt acceleration of positrons in a self-loaded plasma wakefield

S. Corde<sup>1,2</sup>, E. Adli<sup>1,3</sup>, J. M. Allen<sup>1</sup>, W. An<sup>4,5</sup>, C. I. Clarke<sup>1</sup>, C. E. Clayton<sup>4</sup>, J. P. Delahaye<sup>1</sup>, J. Frederico<sup>1</sup>, S. Gessner<sup>1</sup>, S. Z. Green<sup>1</sup>, M. J. Hogan<sup>1</sup>, C. Joshi<sup>4</sup>, N. Lipkowitz<sup>1</sup>, M. Litos<sup>1</sup>, W. Lu<sup>6</sup>, K. A. Marsh<sup>4</sup>, W. B. Mori<sup>4,5</sup>, M. Schmeltz<sup>1</sup>, N. Vafaei-Najafabadi<sup>4</sup>, D. Walz<sup>1</sup>, V. Yakimenko<sup>1</sup> & G. Yocky<sup>1</sup>



# CONCEPTUAL DESIGN OF THE DRIVE BEAM FOR A PWFA-LC\*

S. Pei<sup>#</sup>, M. J. Hogan, T. O. Raubenheimer, A. Seryi, SLAC, CA 94025, U.S.A.  
H. H. Braun, R. Corsini, J. P. Delahaye, CERN, Geneva

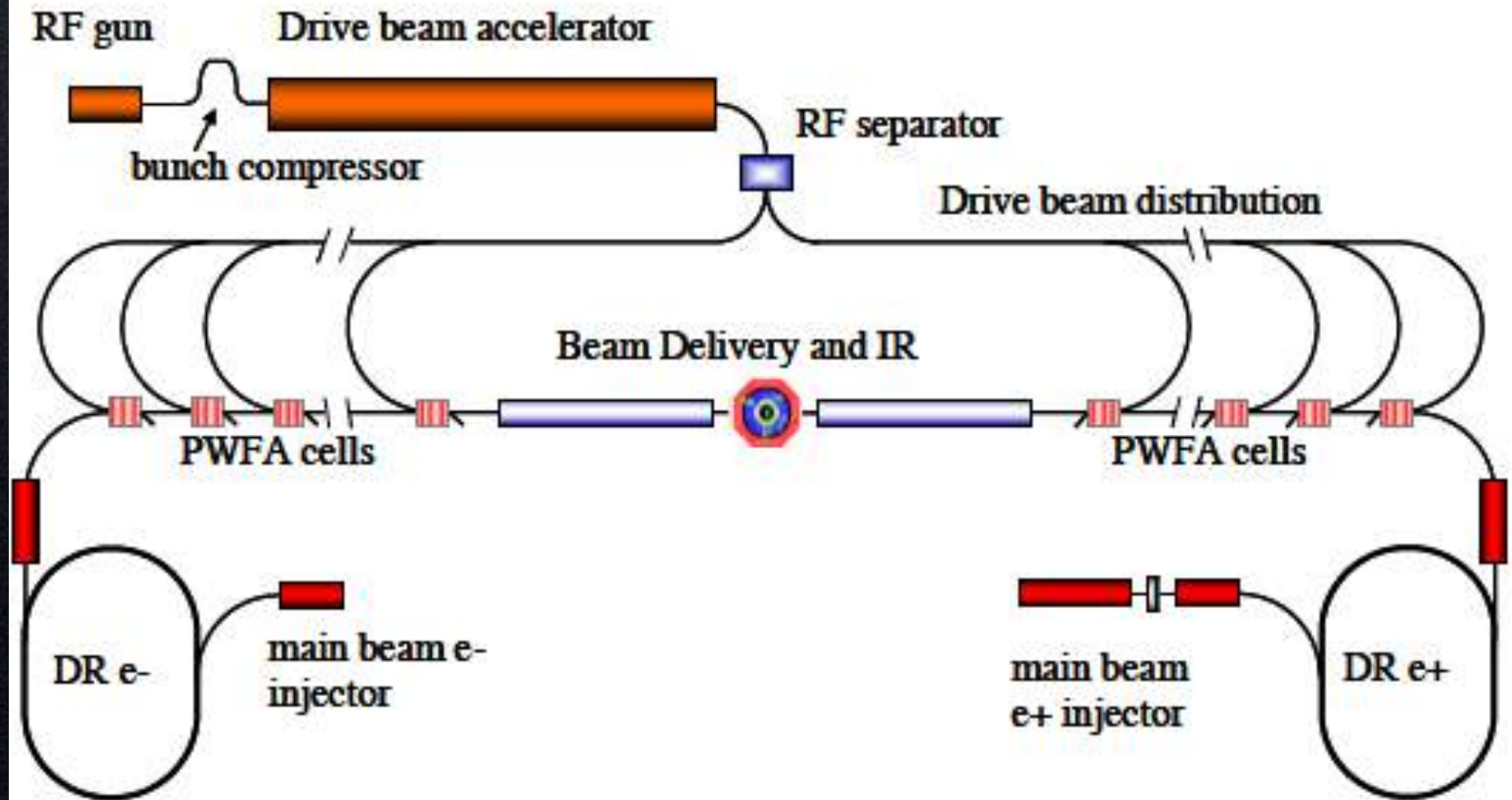


Fig. 1: Concept for a multi-stage PWFA Linear Collider.



Table 1: Key Parameters of the Conceptual Multi-Stage PWFA-based Linear Collider

Main beam: bunch population, bunches per train, rate	$1 \times 10^{10}$ , 125, 100 Hz
Total power of two main beams	20 MW
Drive beam: energy, peak current and active pulse length	25 GeV, 2.3 A, 10 $\mu$ s
Average power of the drive beam	58 MW
Plasma density, accelerating gradient and plasma cell length	$1 \times 10^{17}$ cm <sup>-3</sup> , 25 GV/m, 1 m
Power transfer efficiency drive beam=>plasma =>main beam	35%
Efficiency: Wall plug=>RF=>drive beam	50% $\times$ 90% = 45%
Overall efficiency and wall plug power for acceleration	15.7%, 127 MW
Site power estimate (with 40MW for other subsystems)	170 MW
Main beam emittances, x, y	2, 0.05 mm-mrad
Main beam sizes at Interaction Point, x, y, z	0.14, 0.0032, 10 $\mu$ m
Luminosity	$3.5 \times 10^{34}$ cm <sup>-2</sup> s <sup>-1</sup>
Luminosity in 1% of energy	$1.3 \times 10^{34}$ cm <sup>-2</sup> s <sup>-1</sup>

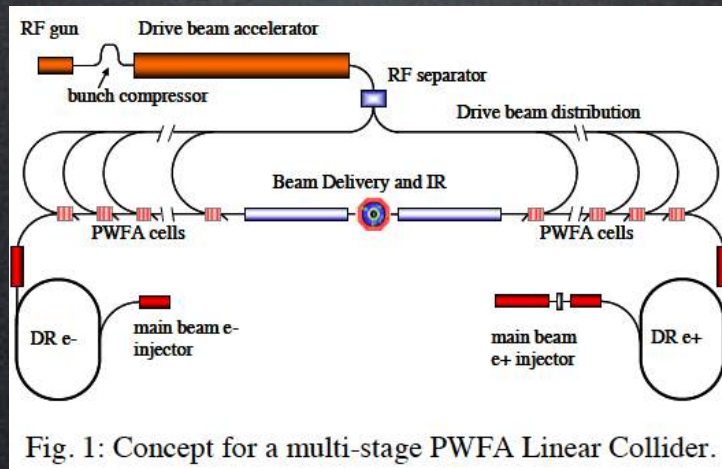


Fig. 1: Concept for a multi-stage PWFA Linear Collider.



ATWAKE

P. Muggli, 06/04/2013, EAAC 2103

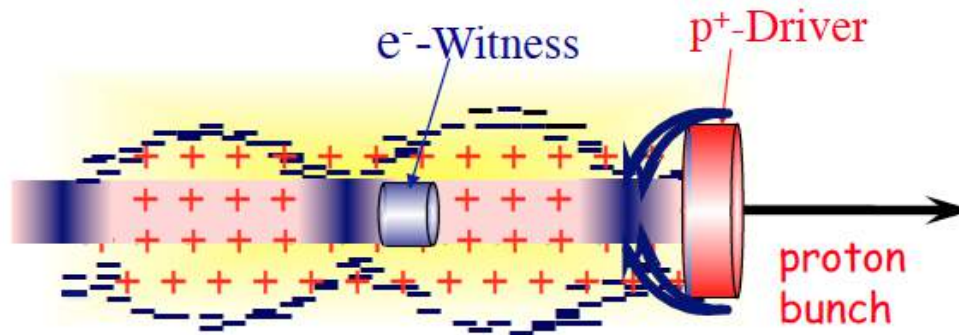
**Proton-driven  
Plasma Wakefield Acceleration  
Collaboration:  
Accelerating  $e^-$  on the wake of a  $p^+$  bunch**







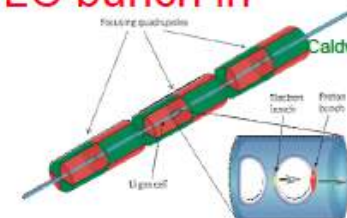
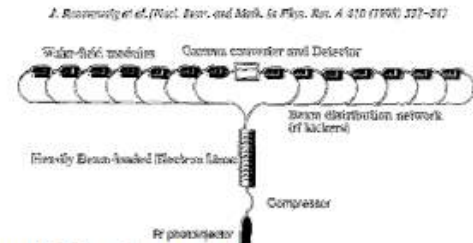
# WHY p<sup>+</sup>-DRIVEN PWFA?



- ✦ ILC, 0.5TeV bunch with  $2 \times 10^{10} e^-$  ~1.6kJ
- ✦ SLAC, 20GeV bunch with  $2 \times 10^{10} e^-$  ~60J
- ✦ SLAC-like driver for staging (FACET= 1 stage, collider 10<sup>+</sup> stages)
- ✦ SPS, 400GeV bunch with  $10^{11} p^+$  ~6.4kJ
- ✦ LHC, 7TeV bunch with  $10^{11} p^+$  ~112kJ

✦ A single SPS or LHC bunch could produce an ILC bunch in a single PWFA stage!

✦ Large average gradient! ( $\geq 1 \text{ GeV/m}$ , 100's m)



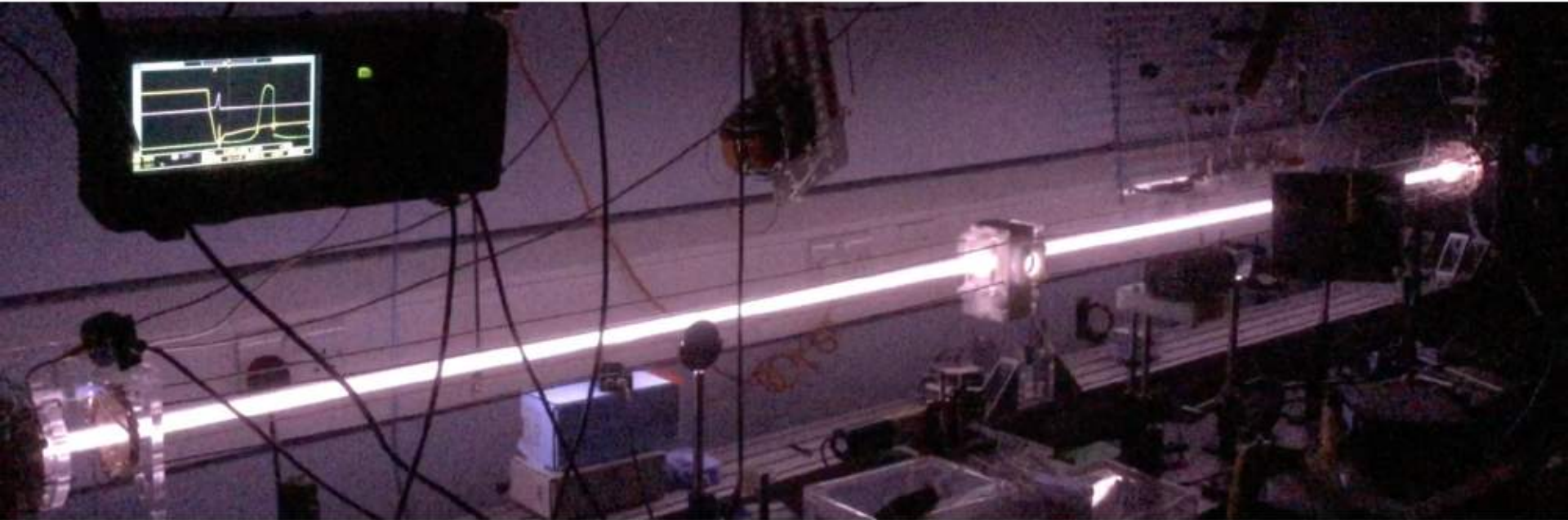
Caldwell, Nat. Phys. 5, 363, (2009)



P. Muggli, 06/04/2013, EAAC 2103

# Discharge configuration II

preliminary tests with the AWAKE 3 meter test tube at IC - 2016



very promising results

... reliable, low jitter plasma formation

scalability of electric circuit for plasmas  $> 10$  m seem achievable...





# EUROPEAN PLASMA RESEARCH ACCELERATOR WITH EXCELLENCE IN APPLICATIONS



This project has received funding from the European Union's Horizon 2020 research and innovation programme under grant agreement No 653782.

## PRESENT EXPERIMENTS

Demonstrating **100 GV/m** routinely

Demonstrating **GeV** electron beams

Demonstrating basic **quality**

## EuPRAXIA INFRASTRUCTURE

Engineering a high quality, compact plasma accelerator  
**5 GeV electron beam for the 2020's**

Demonstrating user readiness

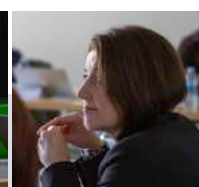
Pilot users from FEL, HEP, medicine, ...

## PRODUCTION FACILITIES

Plasma-based **linear collider** in **2040's**

Plasma-based **FEL** in **2030's**

**Medical, industrial** applications soon

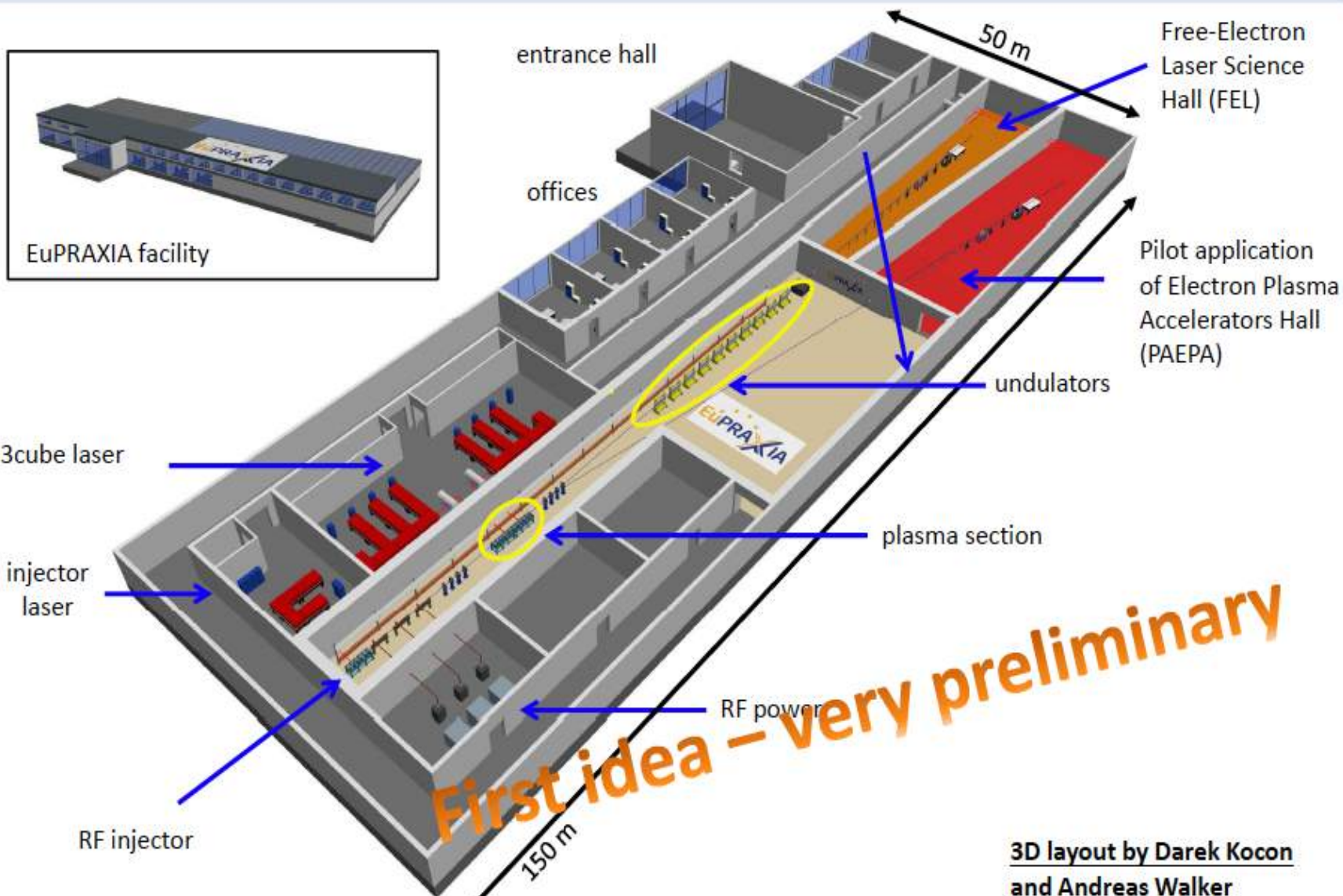




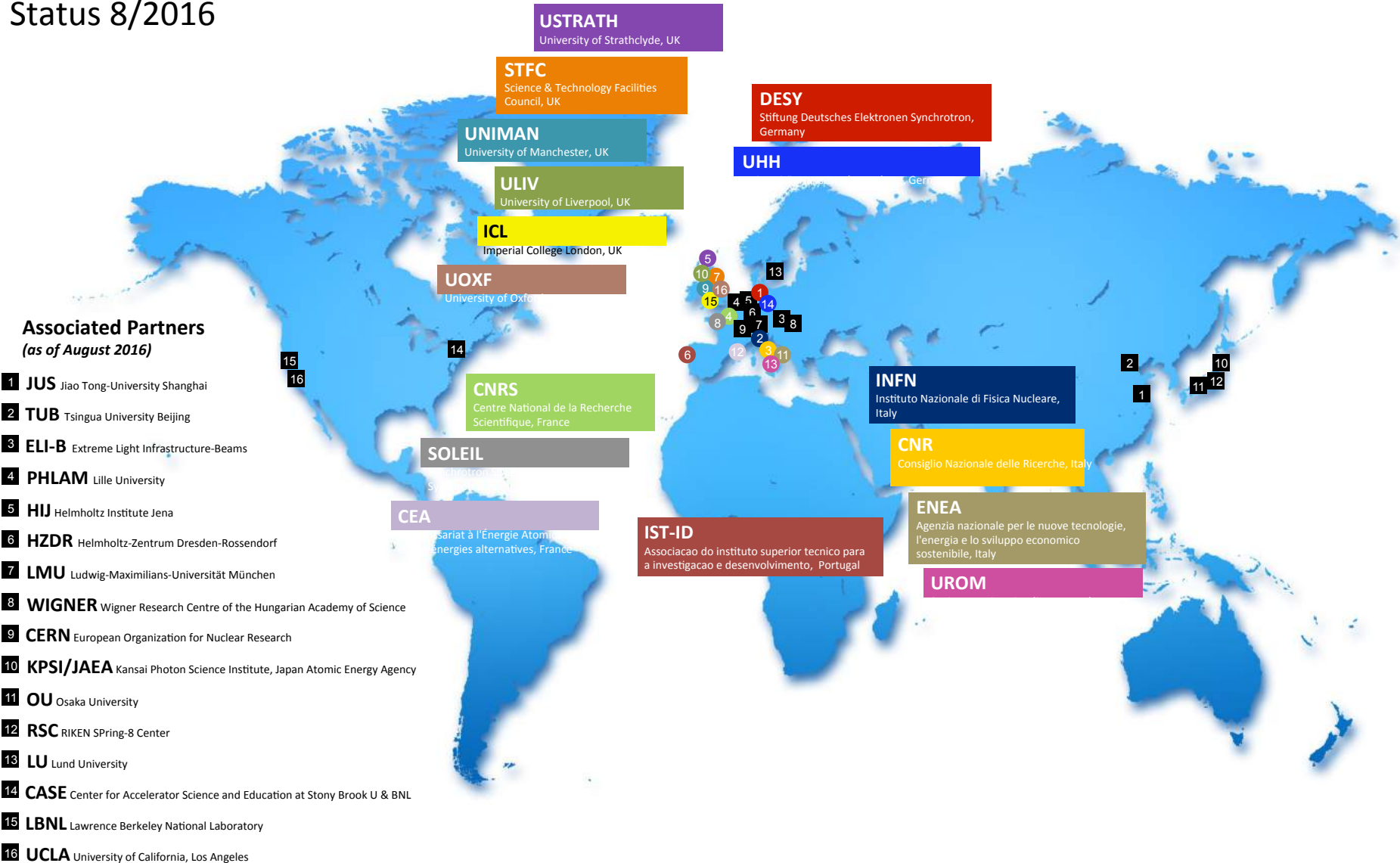
# EuPRAXIA Design Study

Approved as HORIZON 2020 INFRADEV, 4 years, 3 M€

Coordinator: Ralph Assmann (DESY)



Status 8/2016



### Associated Partners (as of August 2016)

- 1 **JUS** Jiao Tong-University Shanghai
- 2 **TUB** Tsinghua University Beijing
- 3 **ELI-B** Extreme Light Infrastructure-Beams
- 4 **PHLAM** Lille University
- 5 **HIJ** Helmholtz Institute Jena
- 6 **HZDR** Helmholtz-Zentrum Dresden-Rossendorf
- 7 **LMU** Ludwig-Maximilians-Universität München
- 8 **WIGNER** Wigner Research Centre of the Hungarian Academy of Science
- 9 **CERN** European Organization for Nuclear Research
- 10 **KPSI/JAEA** Kansai Photon Science Institute, Japan Atomic Energy Agency
- 11 **OU** Osaka University
- 12 **RSC** RIKEN SPring-8 Center
- 13 **LU** Lund University
- 14 **CASE** Center for Accelerator Science and Education at Stony Brook U & BNL
- 15 **LBNL** Lawrence Berkeley National Laboratory
- 16 **UCLA** University of California, Los Angeles



# PWFA accelerating field

$$E_z(r, \xi) \approx (\alpha)(k_p^2 \sigma_z) e^{-k_p^2 \sigma_z^2 / 2} \cos k_p \xi R(r)$$

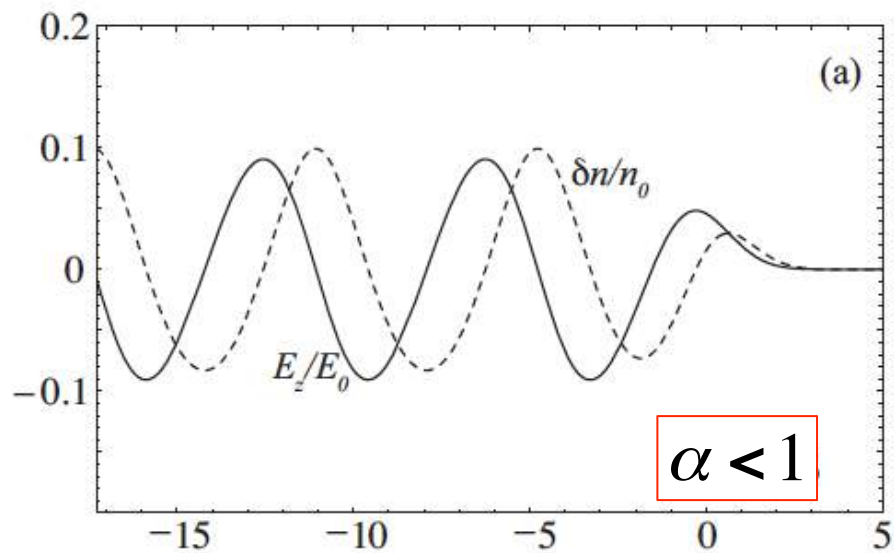
$$\alpha = \frac{n_b}{n_0}$$

$$\sigma_z \approx \frac{\lambda_p}{\pi \sqrt{2}}$$

$$R(0) = - \left( \frac{k_p^2 \sigma_r^2}{2} \right) \left( e^{k_p^2 \sigma_r^2 / 2} \right) \Gamma \left( 0, \frac{k_p^2 \sigma_r^2}{2} \right)$$

$$E_z \propto \frac{n_b}{n_0} \sigma_z \sigma_r^2 \propto Q$$

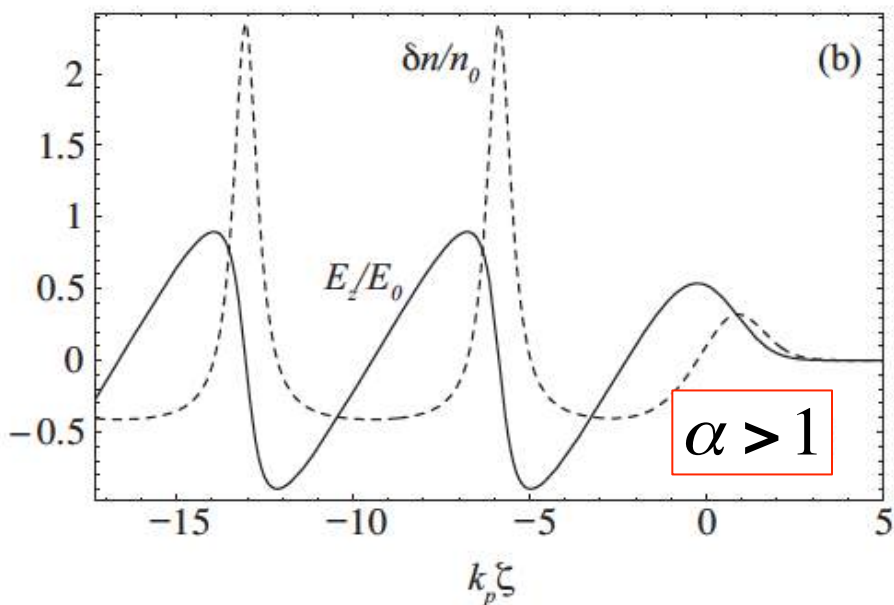
# Regimes: Linear & Non-Linear



**Linear**



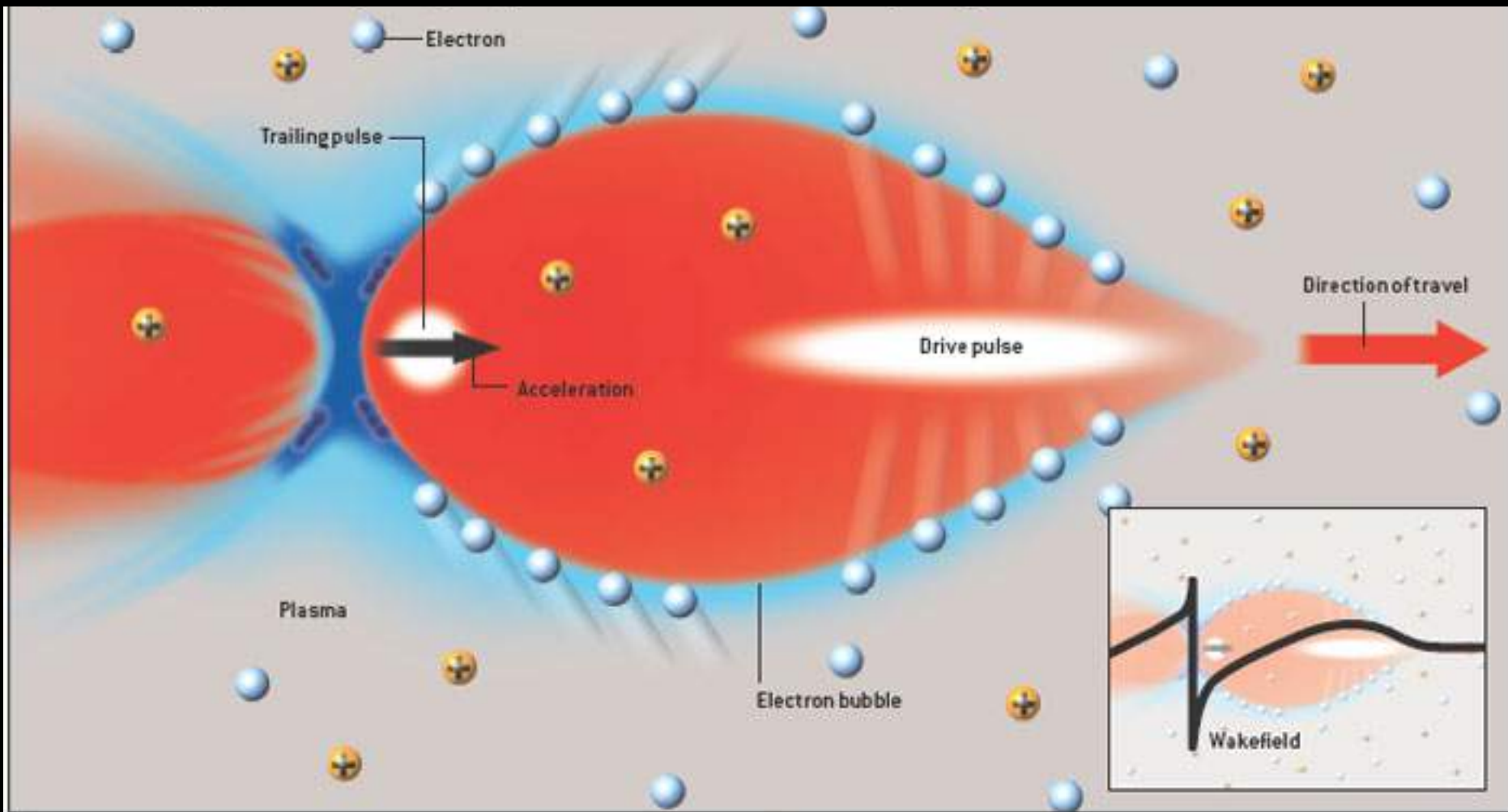
FIG. 8. Time-averaged density variation  $\delta n/n_0$  (dashed curve) and axial electric field  $E_z/E_0$  (solid curve) in an LWFA driven by a Gaussian laser pulse (pulse is moving to the right, centered at  $k_p \zeta = 0$  with rms intensity length  $L_{\text{rms}} = k_p^{-1}$ ) for (a)  $a_0 = 0.5$  and (b)  $a_0 = 2.0$ .



**Non-Linear**



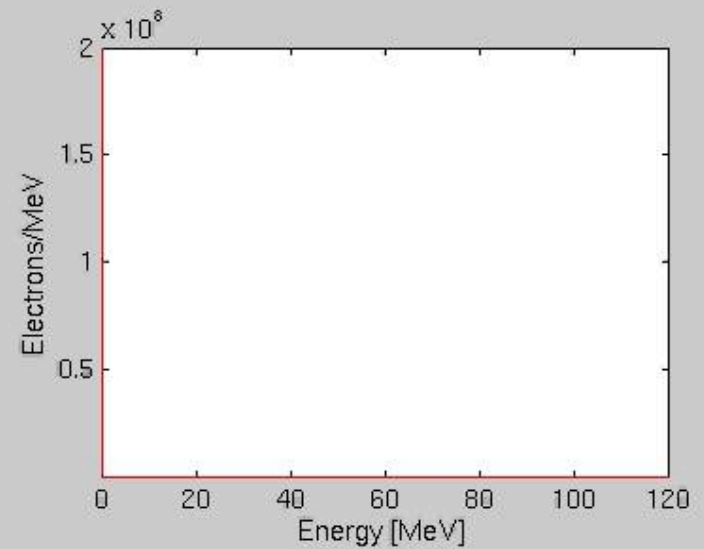
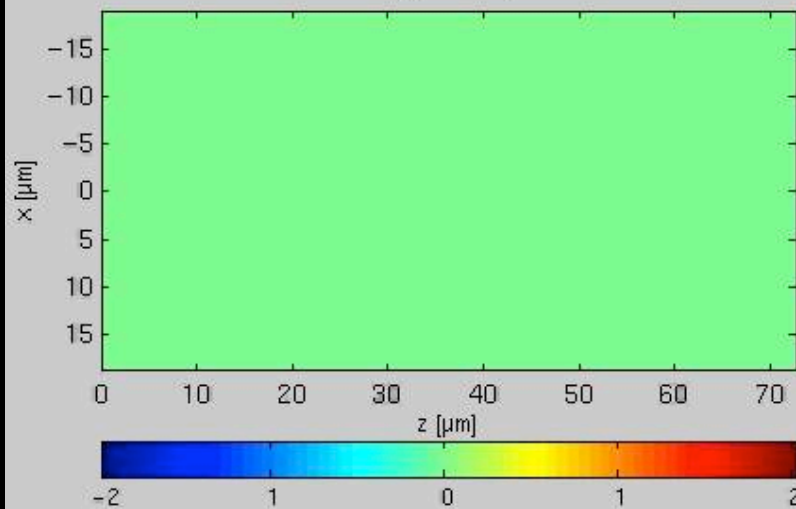
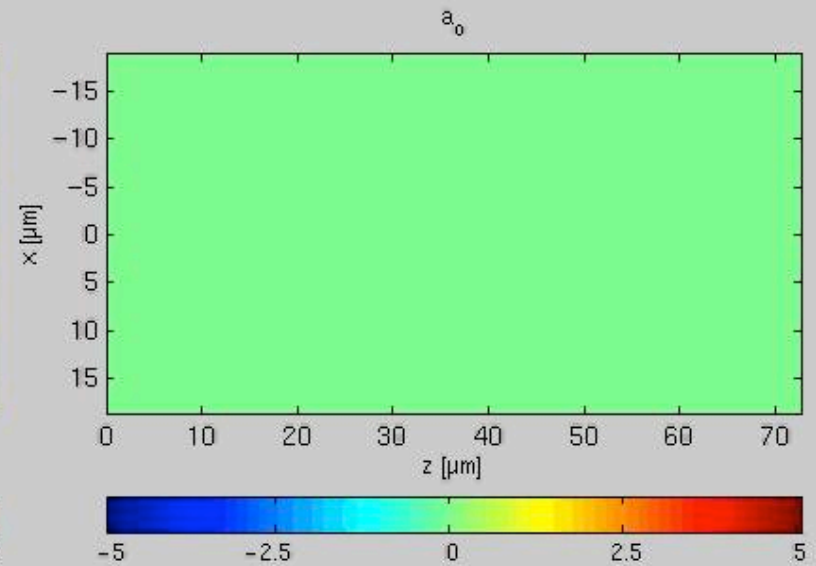
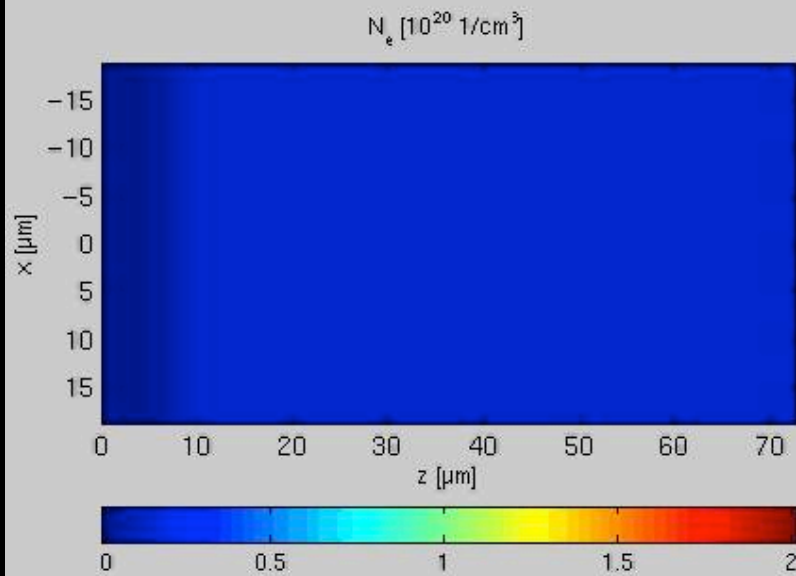




## Breakdown limit?

$$E_0 = \frac{m_e c \omega_p}{e} \approx 100 \left[ \frac{\text{GeV}}{m} \right] \cdot \sqrt{n_0 [10^{18} \text{ cm}^{-3}]}$$

# Self-injection

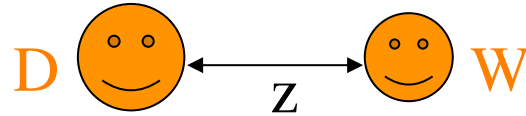




# External-injection



# Wilson theorem for collinear wake field acceleration



Energy lost by an infinitely short bunch:  $U_d = -q_d^2 w_{||}(0)$

Energy change of the second bunch  $U_w = -q_w^2 w_{||}(0) + q_d q_w w_{||}(z_w)$

The sum of the energy exchange by the two bunches must be smaller or equal to zero:

$$U_d + U_w = -q_d^2 w_{||}(0) - q_w^2 w_{||}(0) + q_d q_w w_{||}(z_w) \leq 0$$

$$w_{||}(z_w) \leq \frac{q_d^2 + q_w^2}{q_d q_w} w_{||}(0)$$

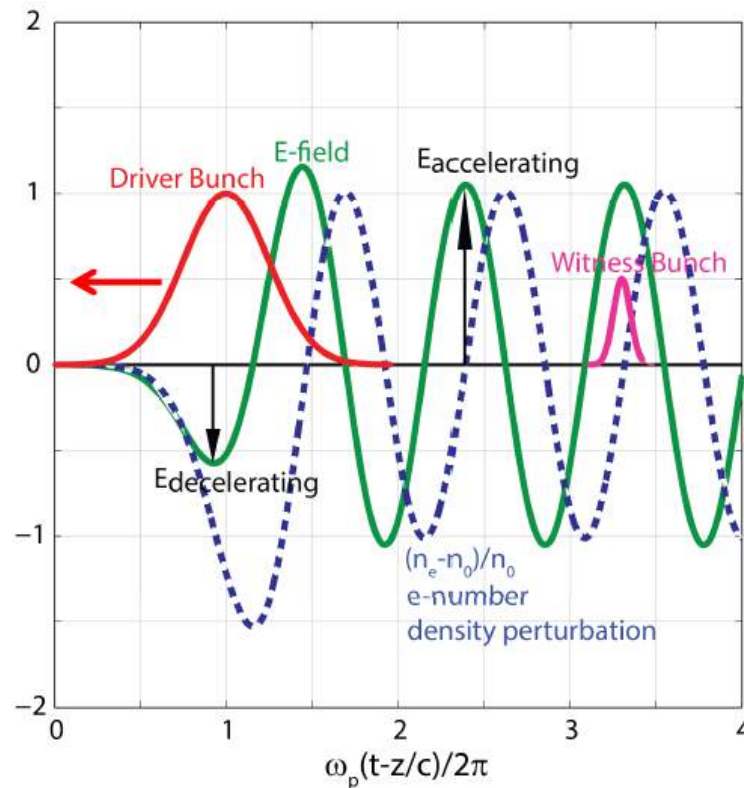
This relation has to be true for all values of charges (the w-potential doesn't depend on q), also when  $q_d = q_w$  which minimize the r.h.s, we obtain:

$$w_{||}(z_w) \leq 2w_{||}(0)$$



The **transformer ratio**  $R_T$  is a figure of merit of the accelerating field.

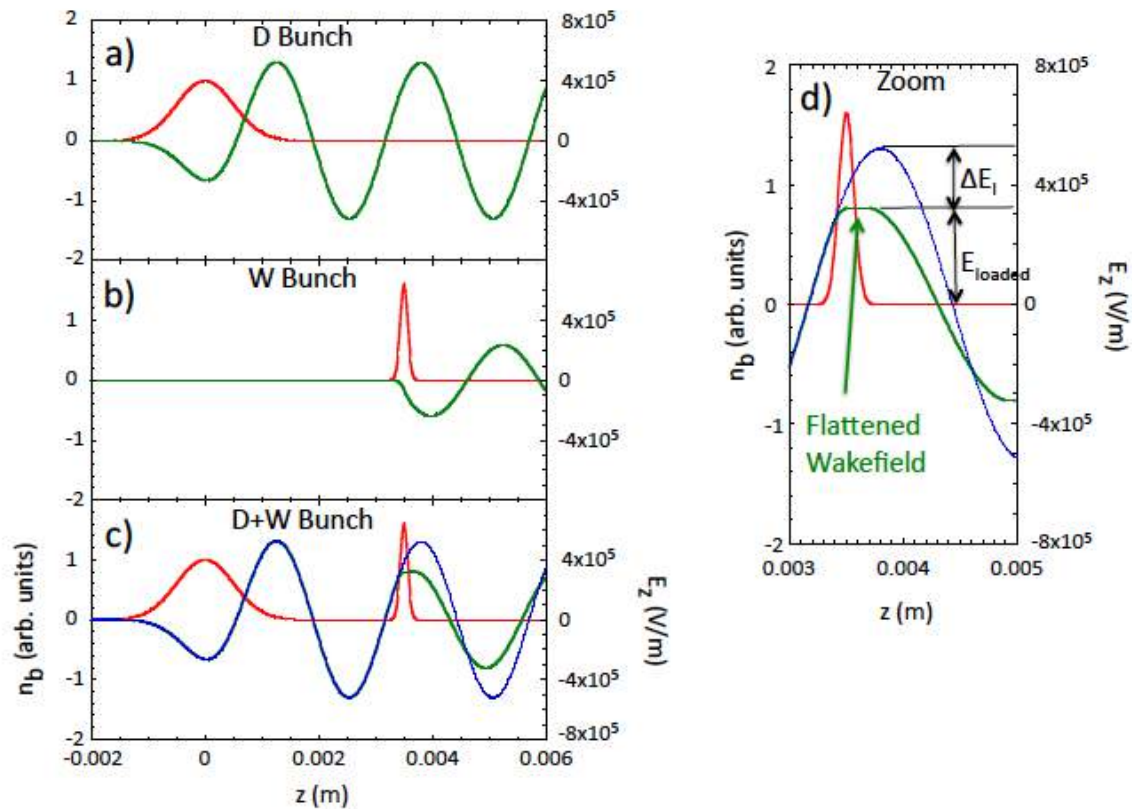
According to previous discussion is defined as the ratio between the peak accelerating field behind drive bunch and the peak decelerating field within drive bunch:



$$R_T = \frac{E_+}{E_-}$$

$$R_T \leq 2$$

# Beam loading



**Fig. 5:** Linear beam loading example: (a) drive bunch density profile (red line) and longitudinal wakefield  $E_z$  (green line), (b) same for the witness bunch, (c) same for the drive and witness bunches together. The field of the drive bunch only is shown as the blue line in panel (c). A zoom around the witness bunch is shown in panel (d). The bunches move to the left.

# Energy gain and efficiency

Assuming  $R_T=2$  and that the Driver will lose all its energy in the plasma, the **active length** can be defined as:

$$\Delta T_d = eE_- L_{act} \qquad L_{act} = \frac{\Delta T_d}{eE_-} = \frac{2\Delta T_d}{eE_+}$$

The **Witness energy gain** is:

$$\Delta T_w = eL_{act}E_+ - eL_{act}E_-^w$$

Where the second term on the r.h.s is the energy lost in the plasma by the witness itself.

Substituting the definition of  $L_{act}$ , and using the field scaling  $E_z \propto Q$  w.r.t the charge it results:

$$\Delta T_w = 2 \left( 1 - \frac{E_+^w}{2E_+} \right) \Delta T = \left( 2 - \frac{q_w}{Q_d} \right) \Delta T_d$$

And **the efficiency** of the process is given by:

$$\eta = \frac{q_w \Delta T_w}{Q_d \Delta T_d} = \left( 2 - \frac{q_w}{Q_d} \right) \frac{q_w}{Q_d}$$



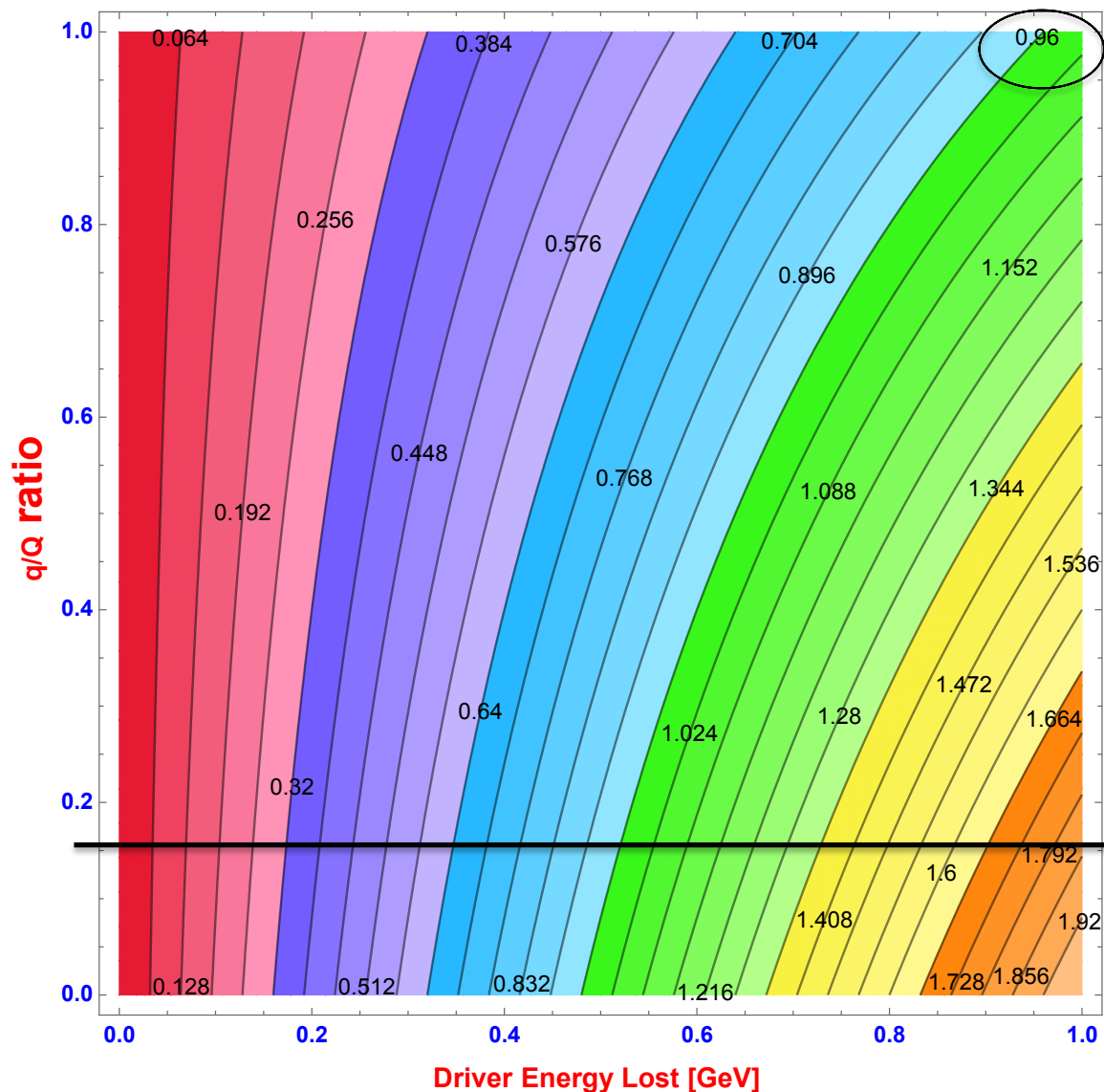
$$\Delta T_w \propto \left(2 - \frac{q}{Q}\right) \Delta T_d$$

$q=30 \text{ pC}$

$Q=200 \text{ pC}$

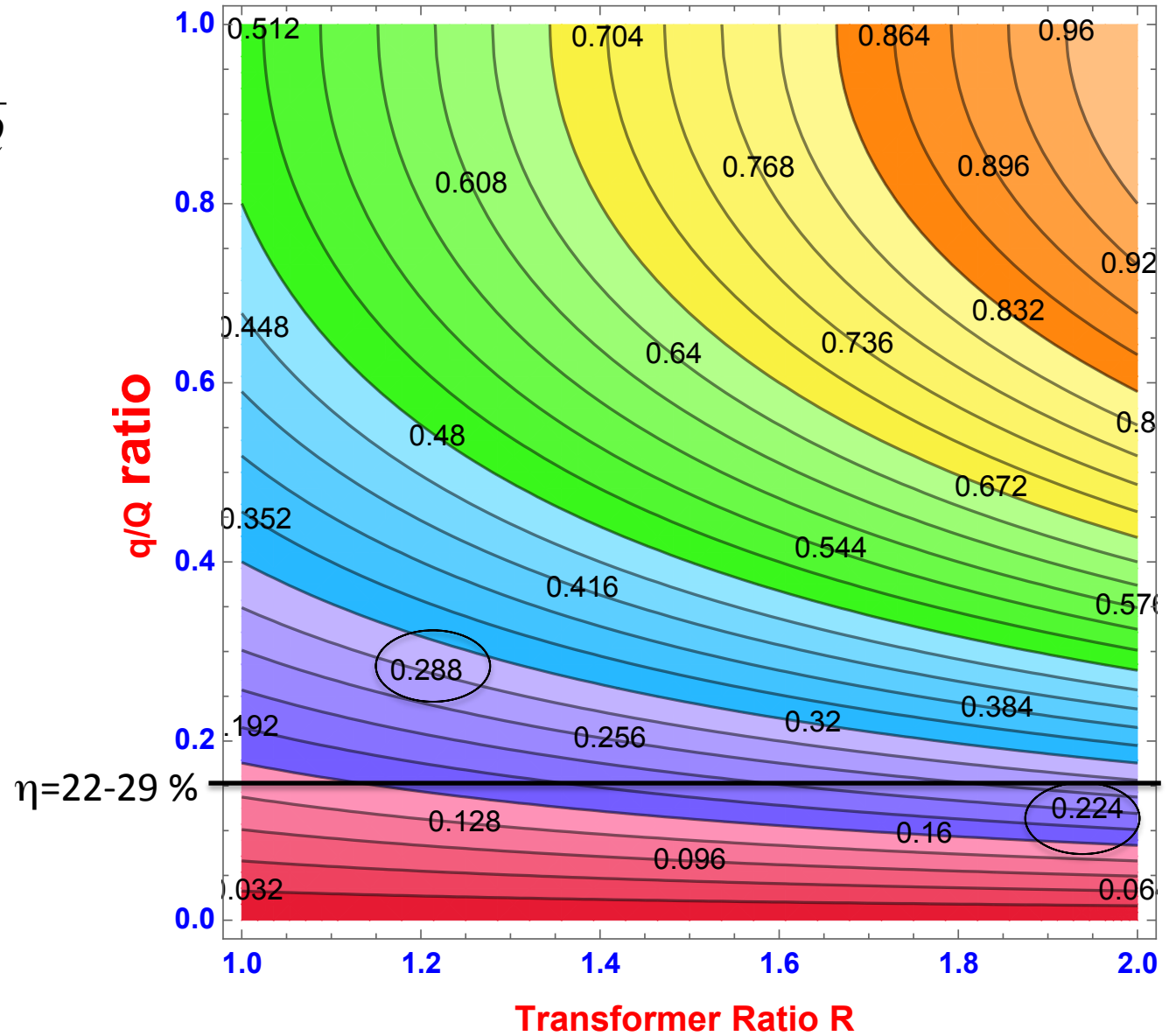
$q/Q=0.15$

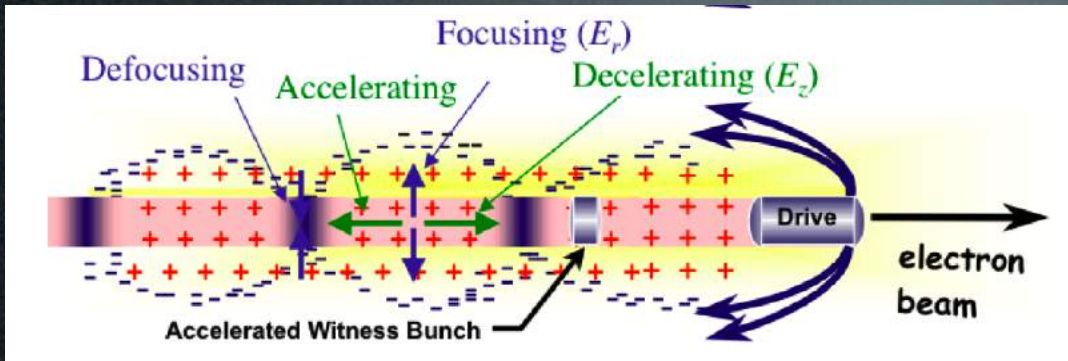
### Witness Energy Gain [GeV]



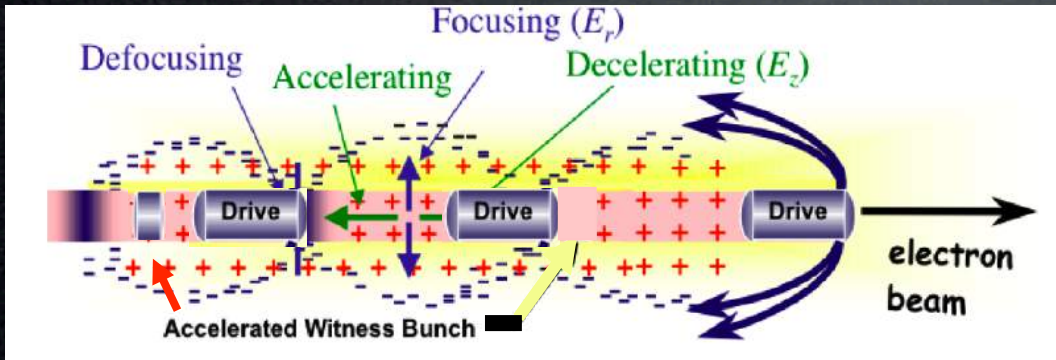
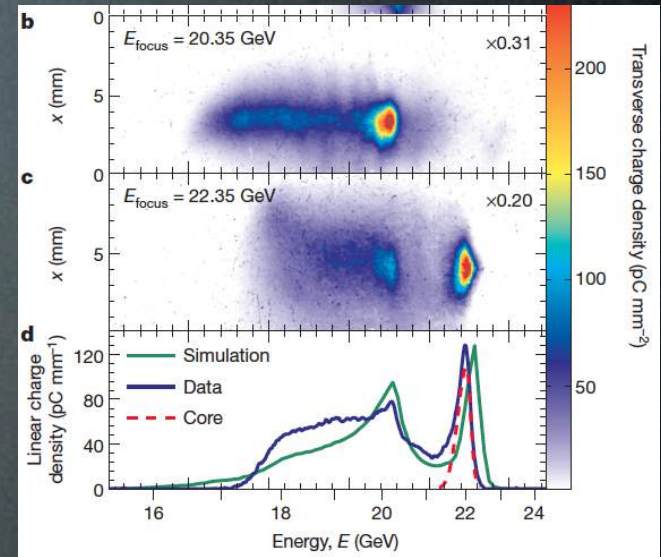
# Energy Transfer Efficiency

$$\eta = \frac{\Delta U_w}{\Delta U_d} = R \left( 1 - \frac{q}{2Q} \right) \frac{q}{Q}$$

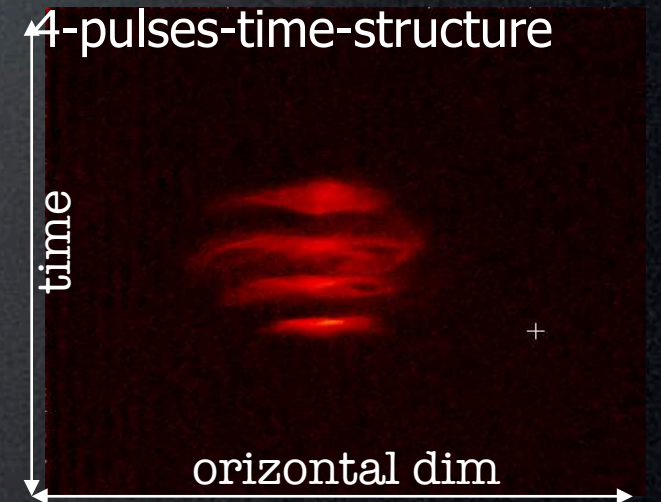




Litos, M. et al. *High-efficiency acceleration of an electron beam in a plasma wakefield accelerator.* **Nature** 515, 92–95 (2014).



In progress at SPARC\_LAB







MAX-PLANCK-GESELLSCHAFT

# MULTIBUNCH PWFA



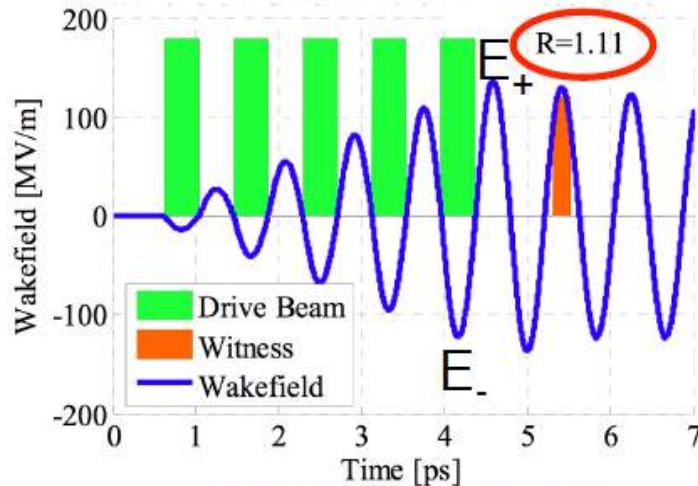
Transformer Ratio:  $R = E_+ / E_-$       Energy Gain:  $\leq RE_0$

$\sigma_r = 125 \mu\text{m}$ ,  $n_e = 1.8 \times 10^{16} \text{ cm}^{-3}$ ,  $\lambda_p = 250 \mu\text{m}$

$E_0$ : incoming energy

$Q = 30 \text{ pC/bunch}$ ,  $\Delta z = 250 \mu\text{m} \approx \lambda_p$

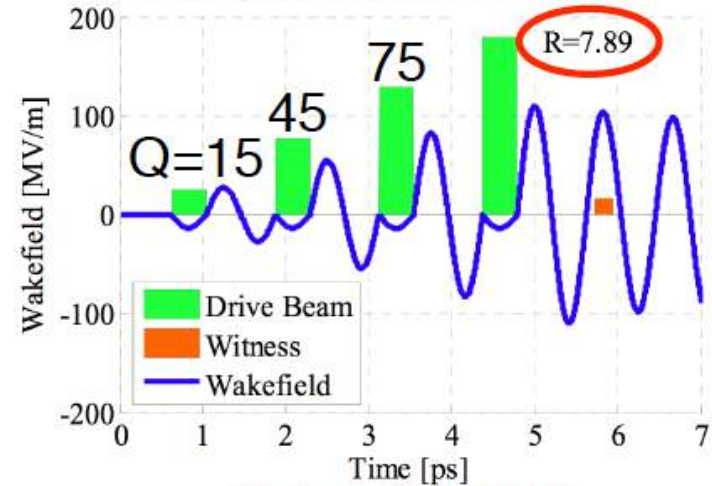
Bunch Train



Kallos, PAC'07 Proceedings

$\Delta z = 375 \mu\text{m} \approx 1.5 \lambda_p$

Ramped Bunch Train\*



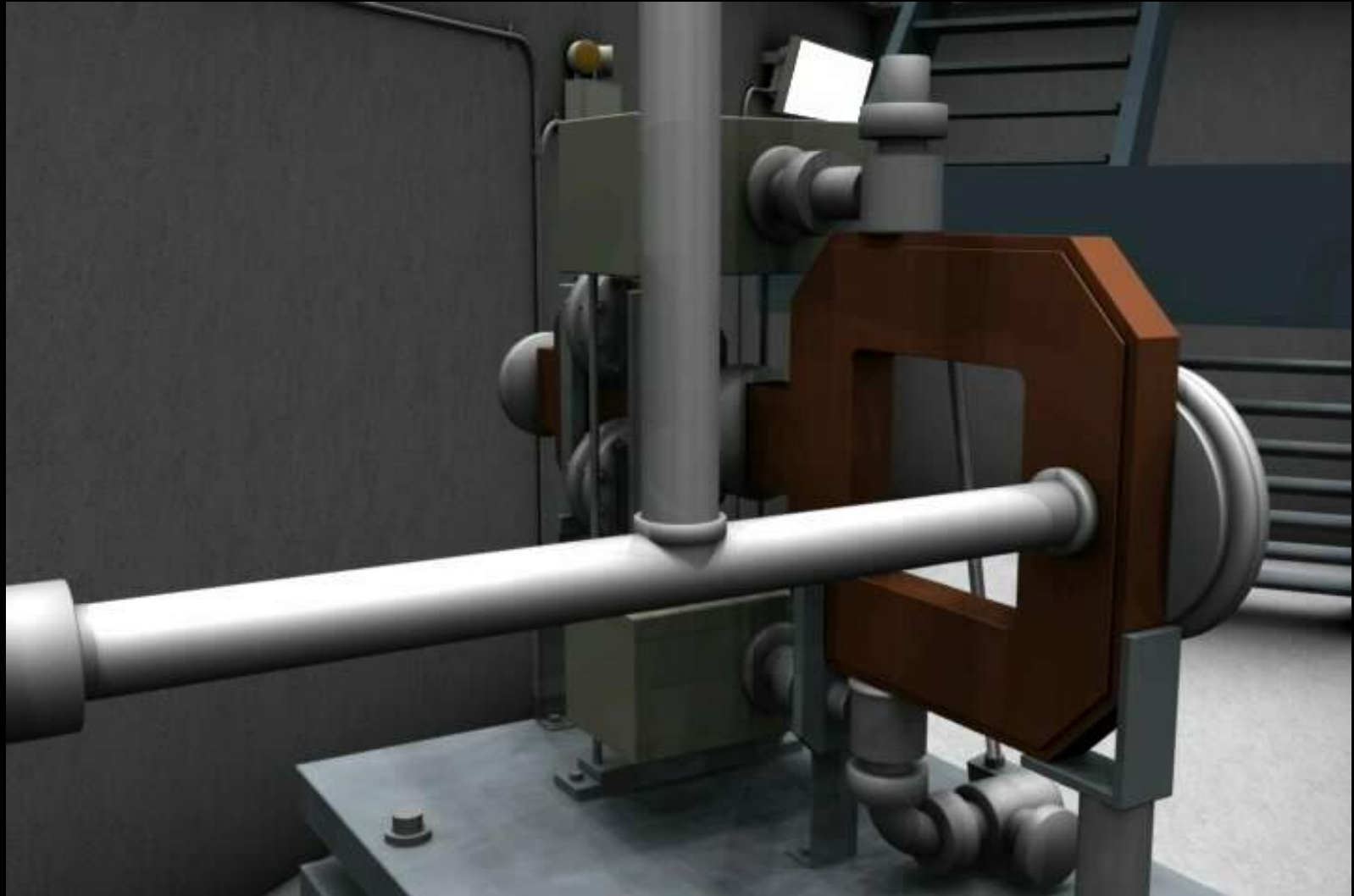
\*Tsakanov, NIMA, 1999

➔ Linear (2D) theory for  $n_b \ll n_e$ !

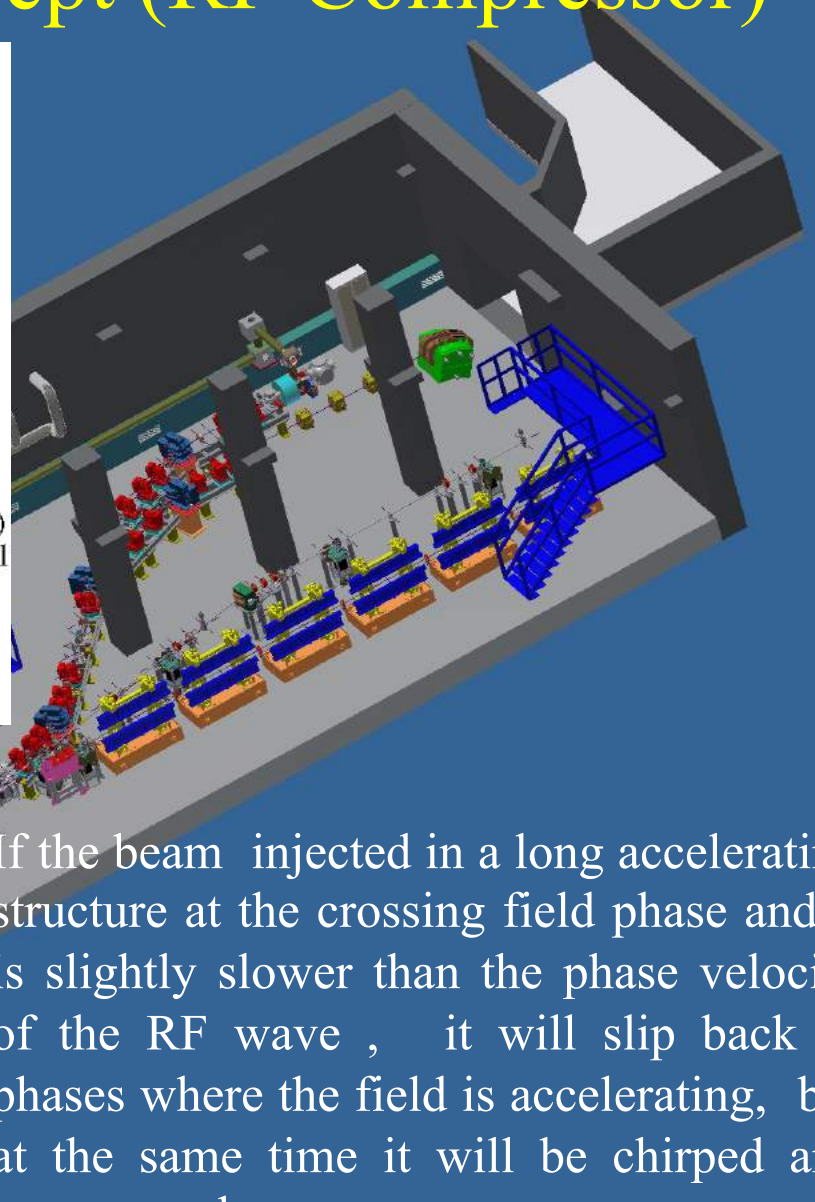
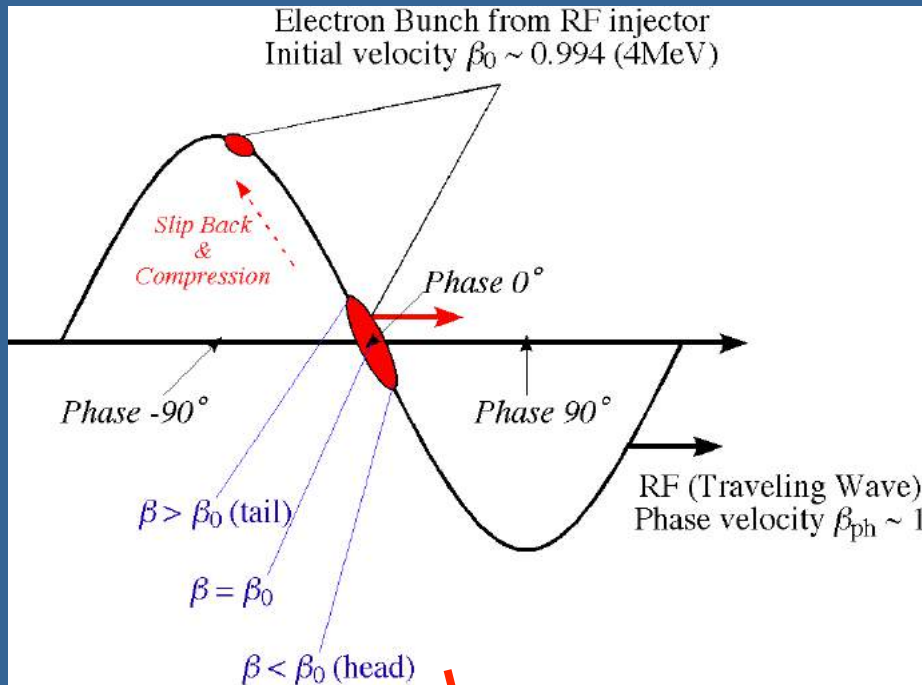
➔  $R = 7.9 \Rightarrow$  multiply energy by  $\sim 8$  in a single PWFA stage!



# High Brightness Photo-Injector



# Velocity bunching concept (RF Compressor)



If the beam injected in a long accelerating structure at the crossing field phase and it is slightly slower than the phase velocity of the RF wave, it will slip back to phases where the field is accelerating, but at the same time it will be chirped and compressed.



# Laser Comb technique: generation of a train of short bunches

(Parmela code)

Charge vs. Time

Energy vs. Time

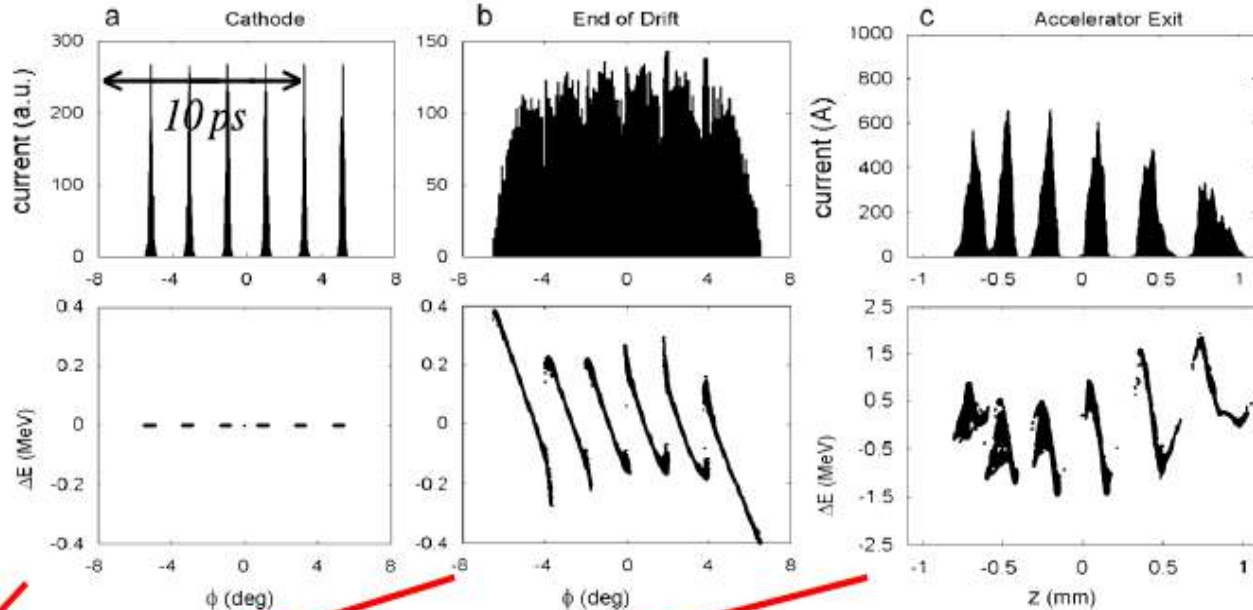
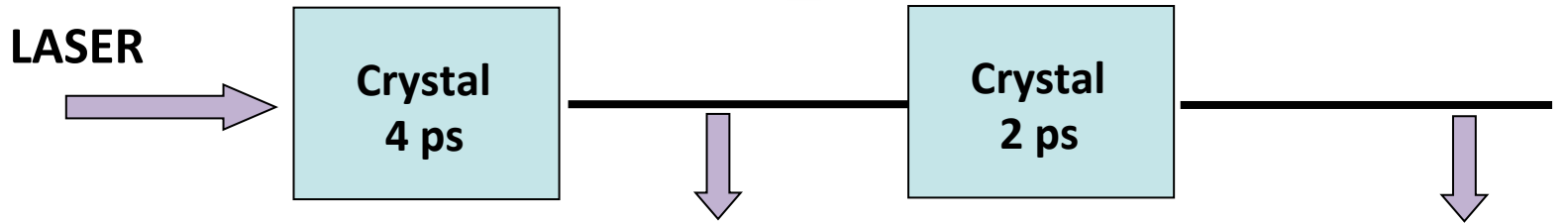
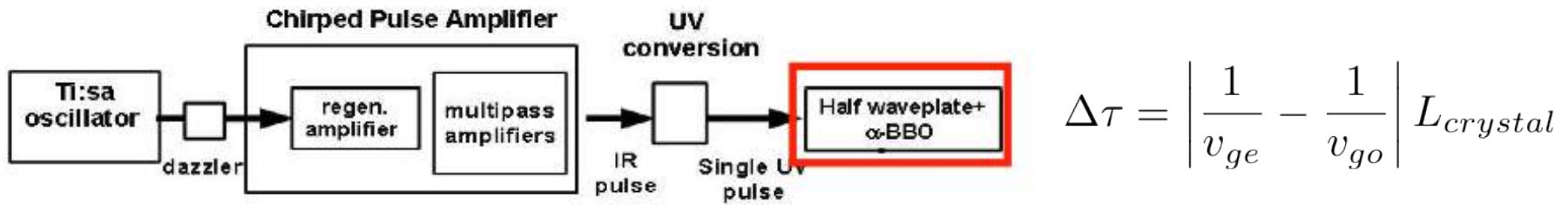


Fig. 1. Evolution of a six bunches electron beam train: the columns from left refer respectively, to (a) the cathode, (b) the end of the drift at 150 cm and (c) the end of linac at 12 m far from cathode. The rows from top refer, respectively, to longitudinal profile and to energy modulation  $\Delta E$  (MeV).

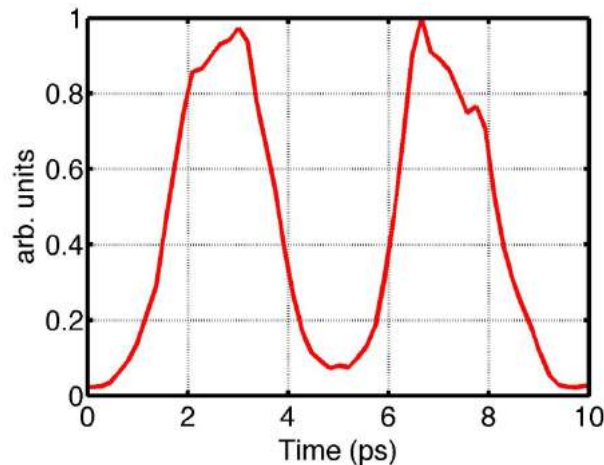


- P.O. Shea et al., Proc. of 2001 IEEE PAC, Chicago, USA (2001) p.704. (Low charge regime only)
- M. Ferrario, M. Boscolo et al., Int. J. of Mod. Phys. B, 2006 (High charge, Beam Echo)

# Laser Pulse Train Generation

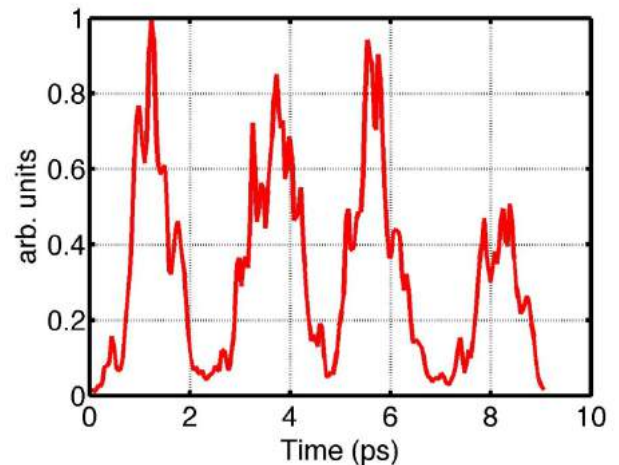


UV pulses



Streak camera

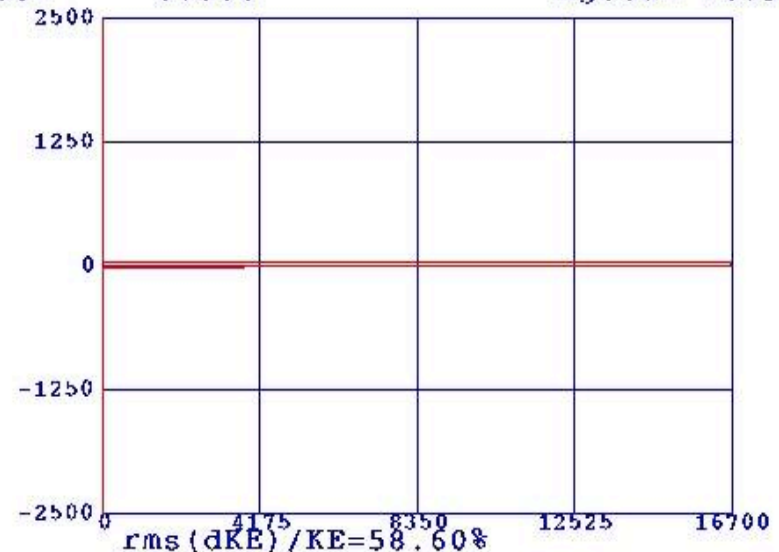
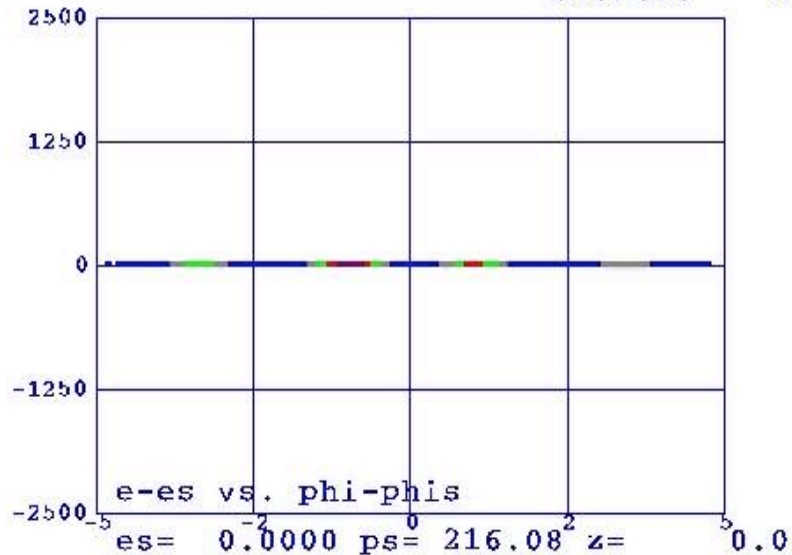
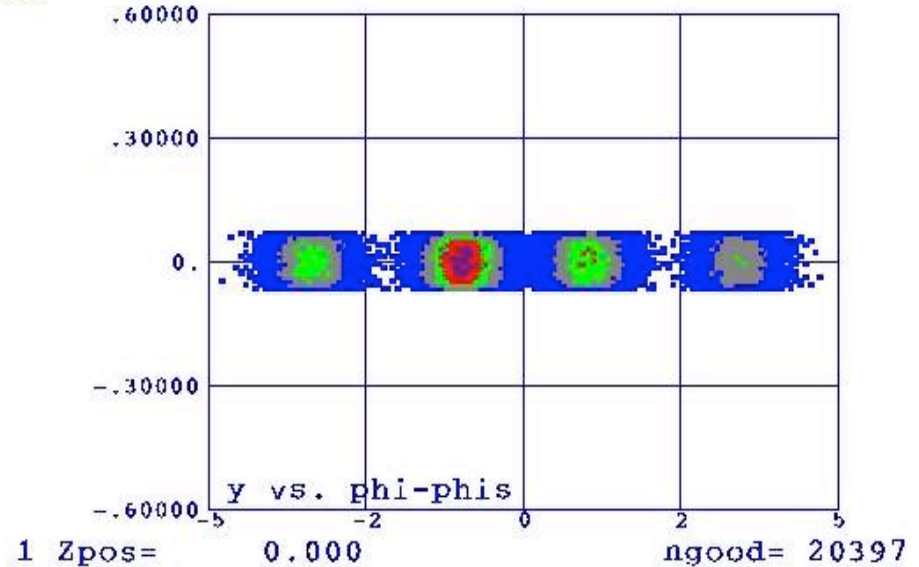
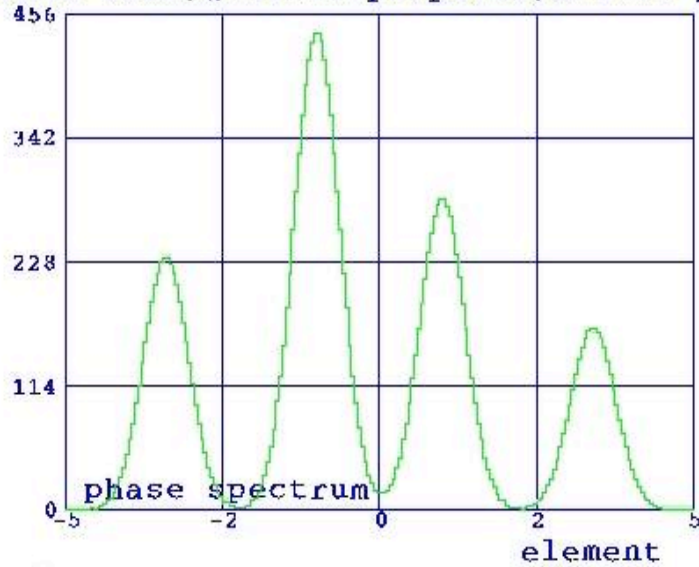
UV pulses



Streak camera

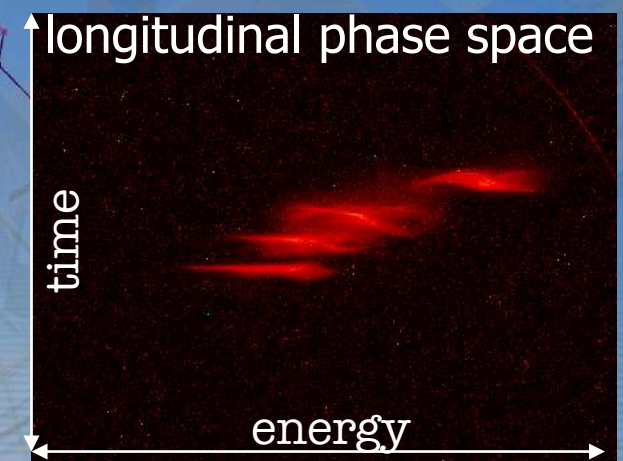
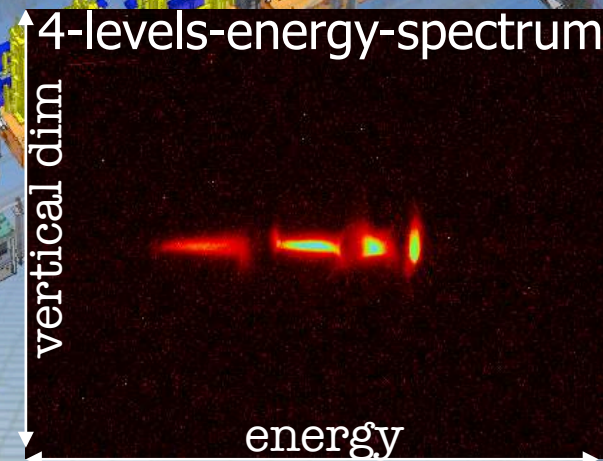
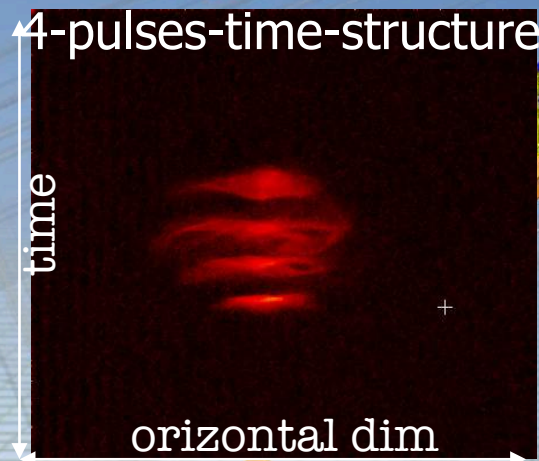
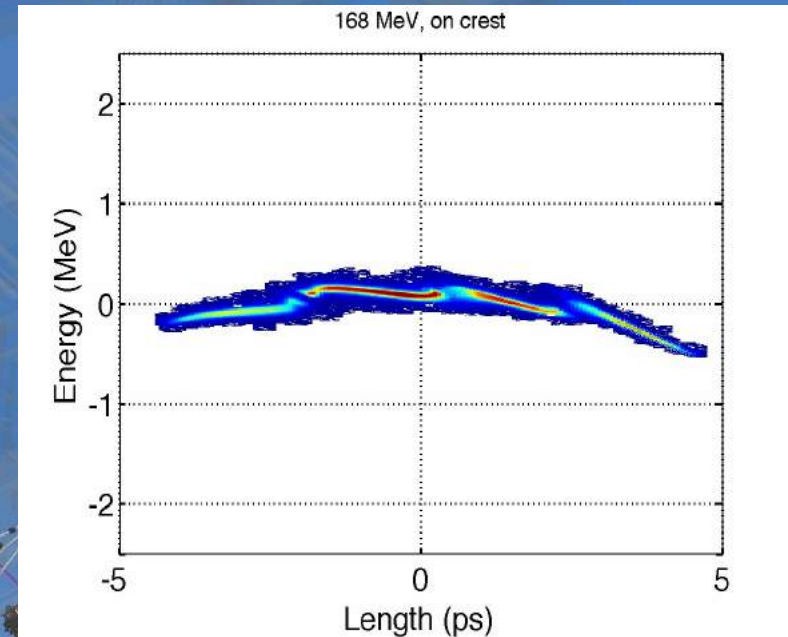
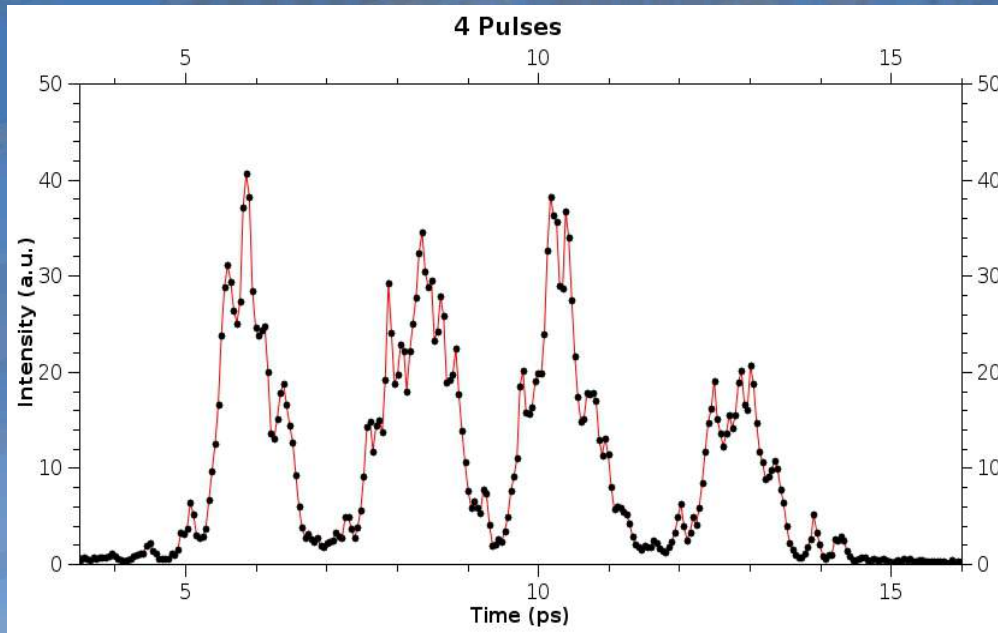
# Overcompression

SPARC COMB,  $q_{tot}=220\text{pC/pulse}$ ,  $d=4.27\text{ psec}$



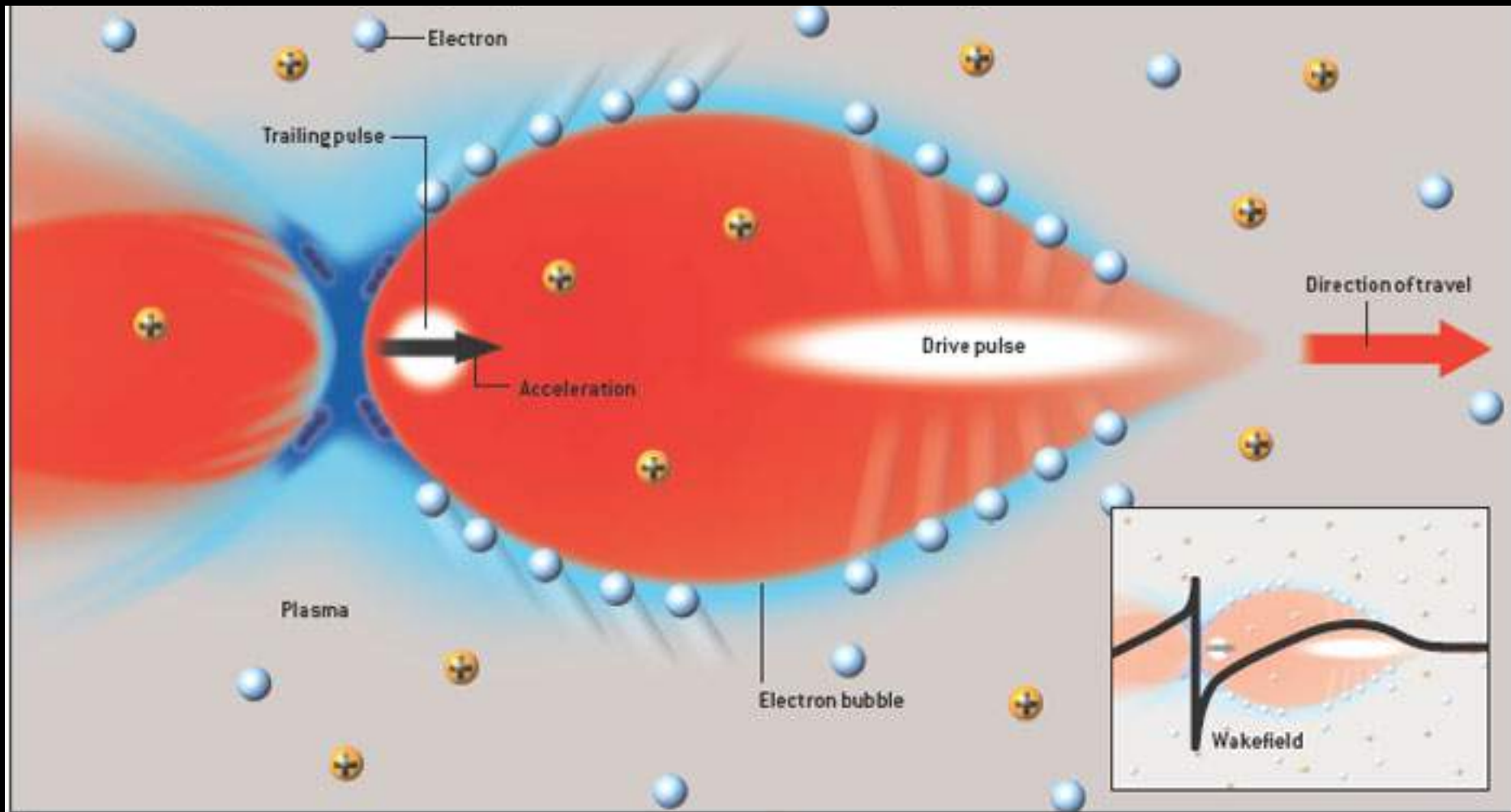


# Laser COMB: experimental results

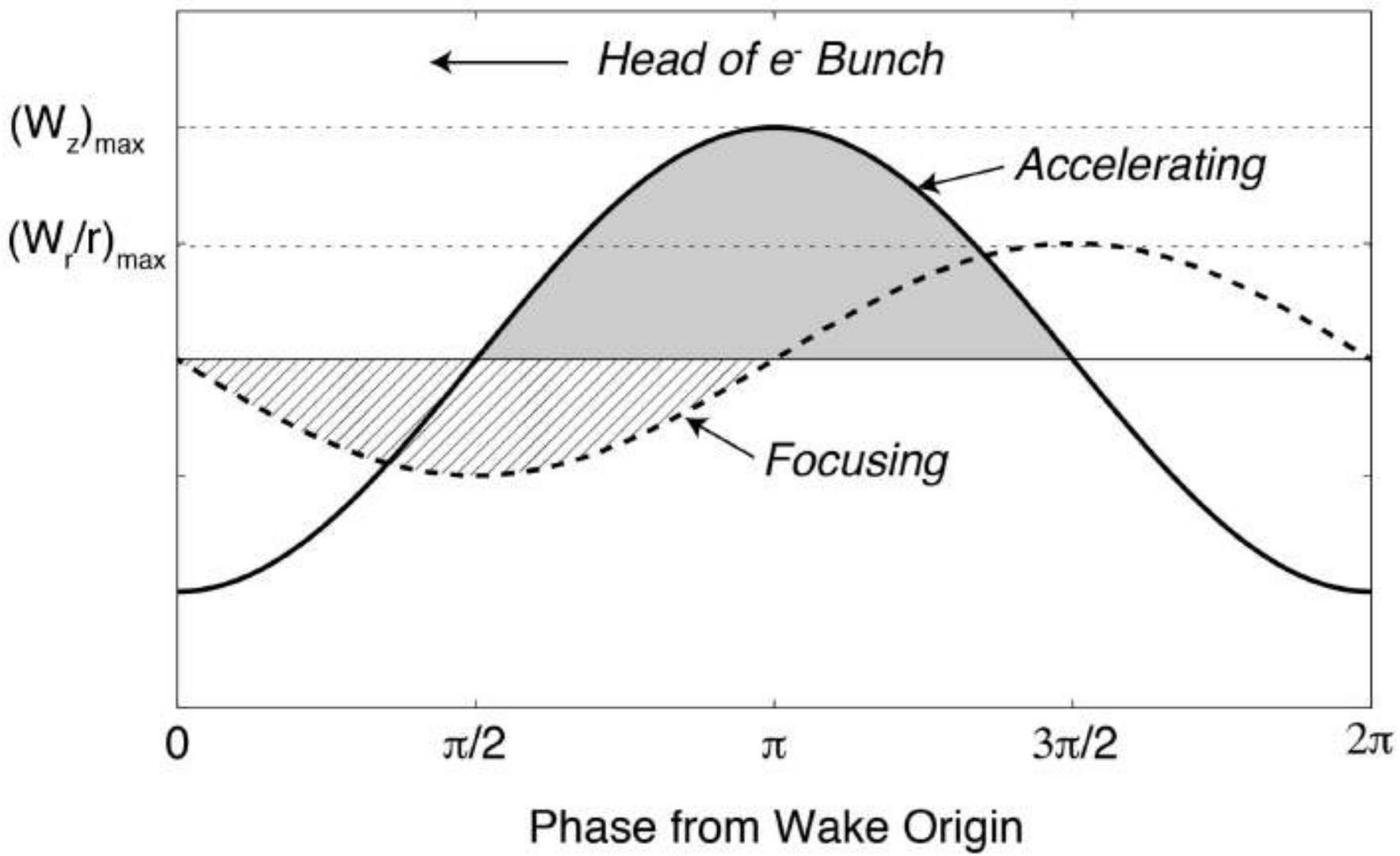


- M. Ferrario et al., Nucl. Inst. and Meth, A 637 (2011)
- A. Mostacci et al., Proc. of IPAC 2011, Spain

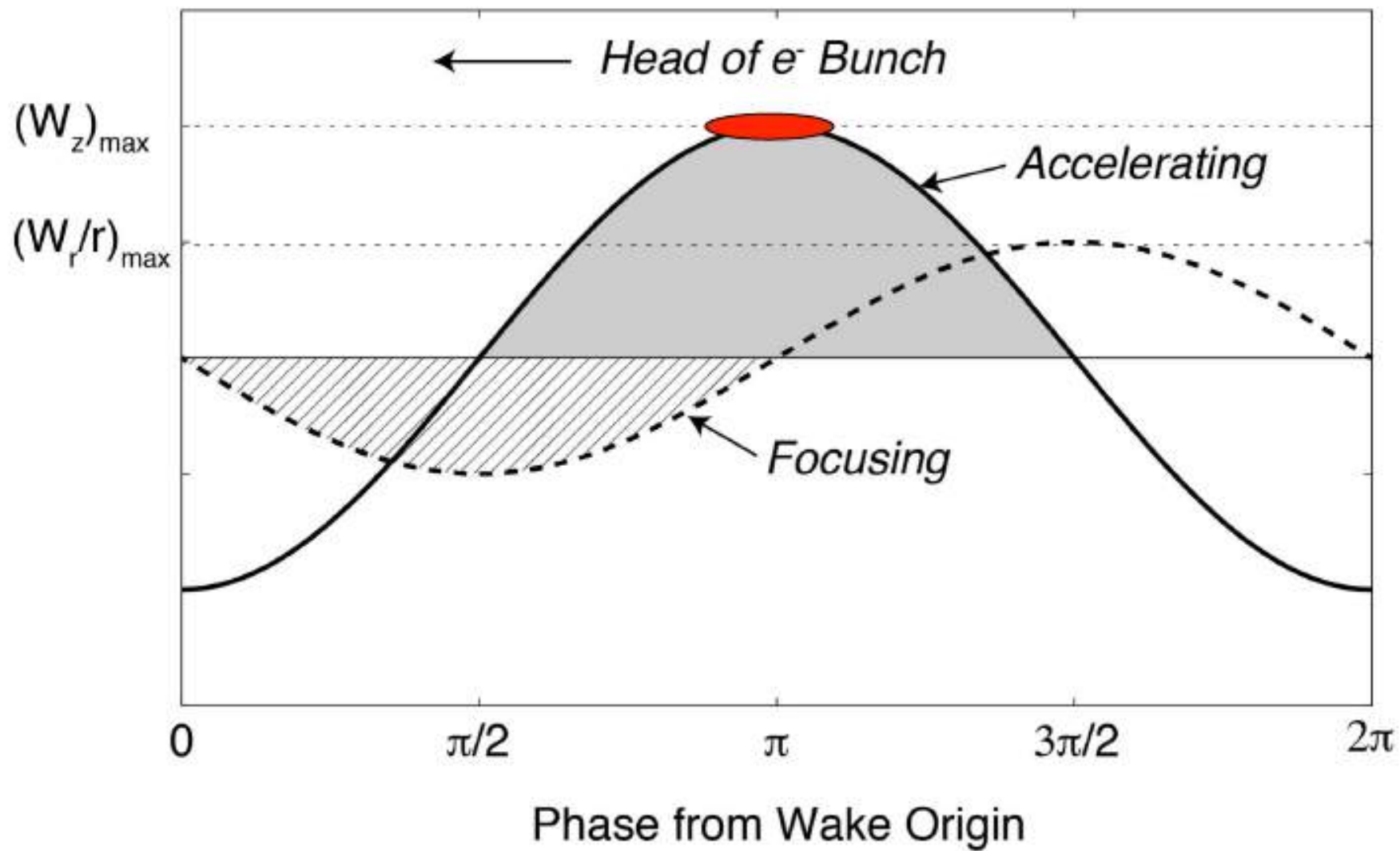
# PWFA transverse field

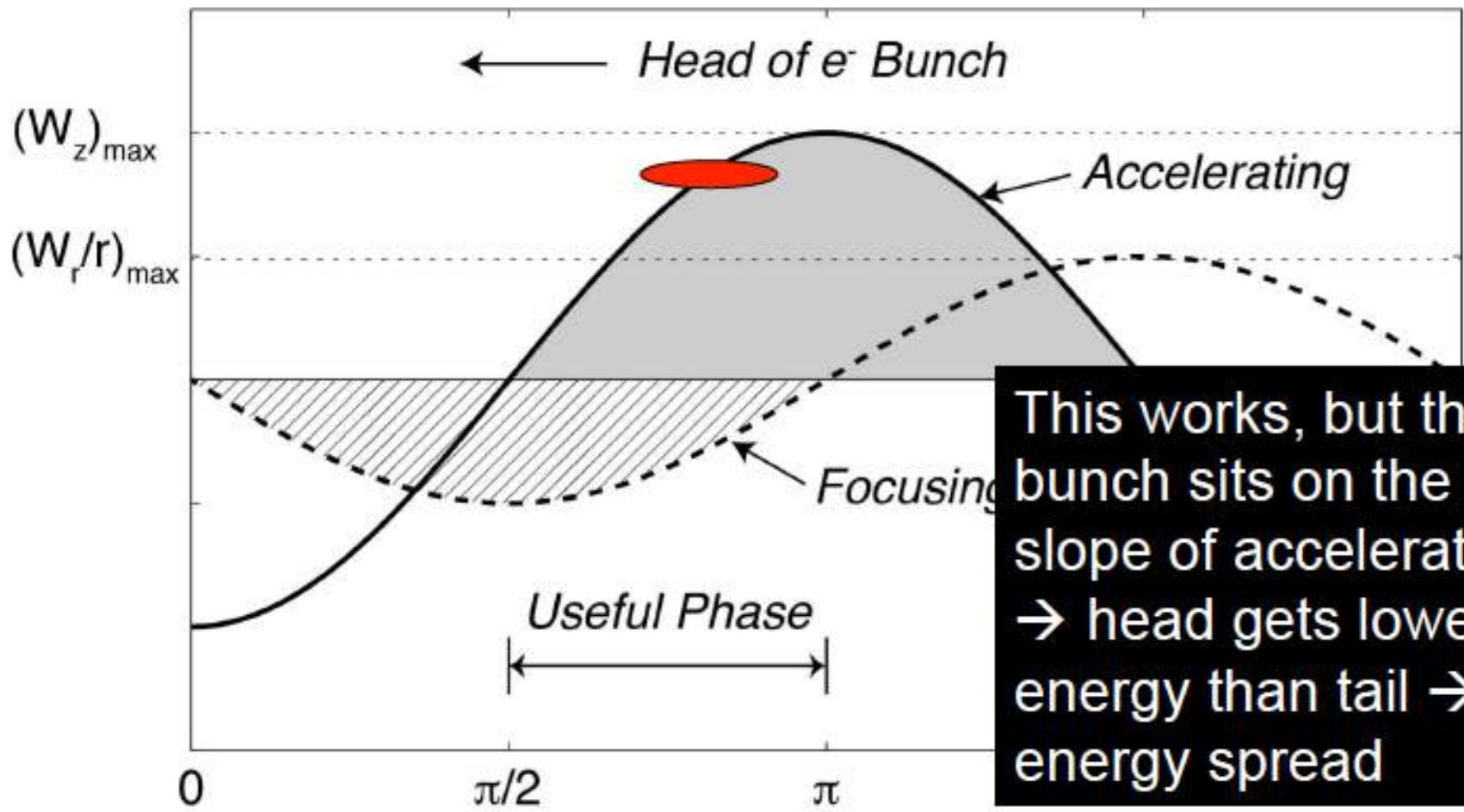


$$E_r(r, \xi) \approx (\alpha)(k_p \sigma_z) e^{-k_p^2 \sigma_z^2 / 2} \sin k_p \xi R'(r)$$







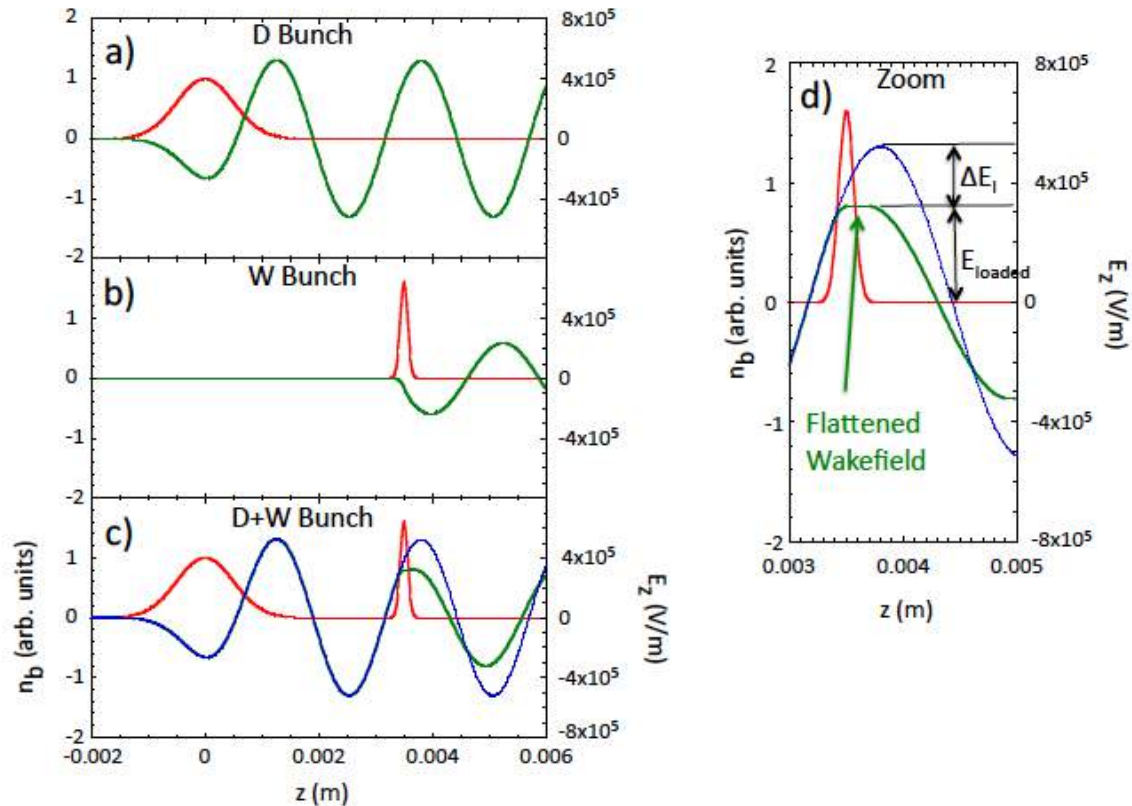


This works, but the bunch sits on the slope of acceleration → head gets lower energy than tail → energy spread

Phase from Wake Origin



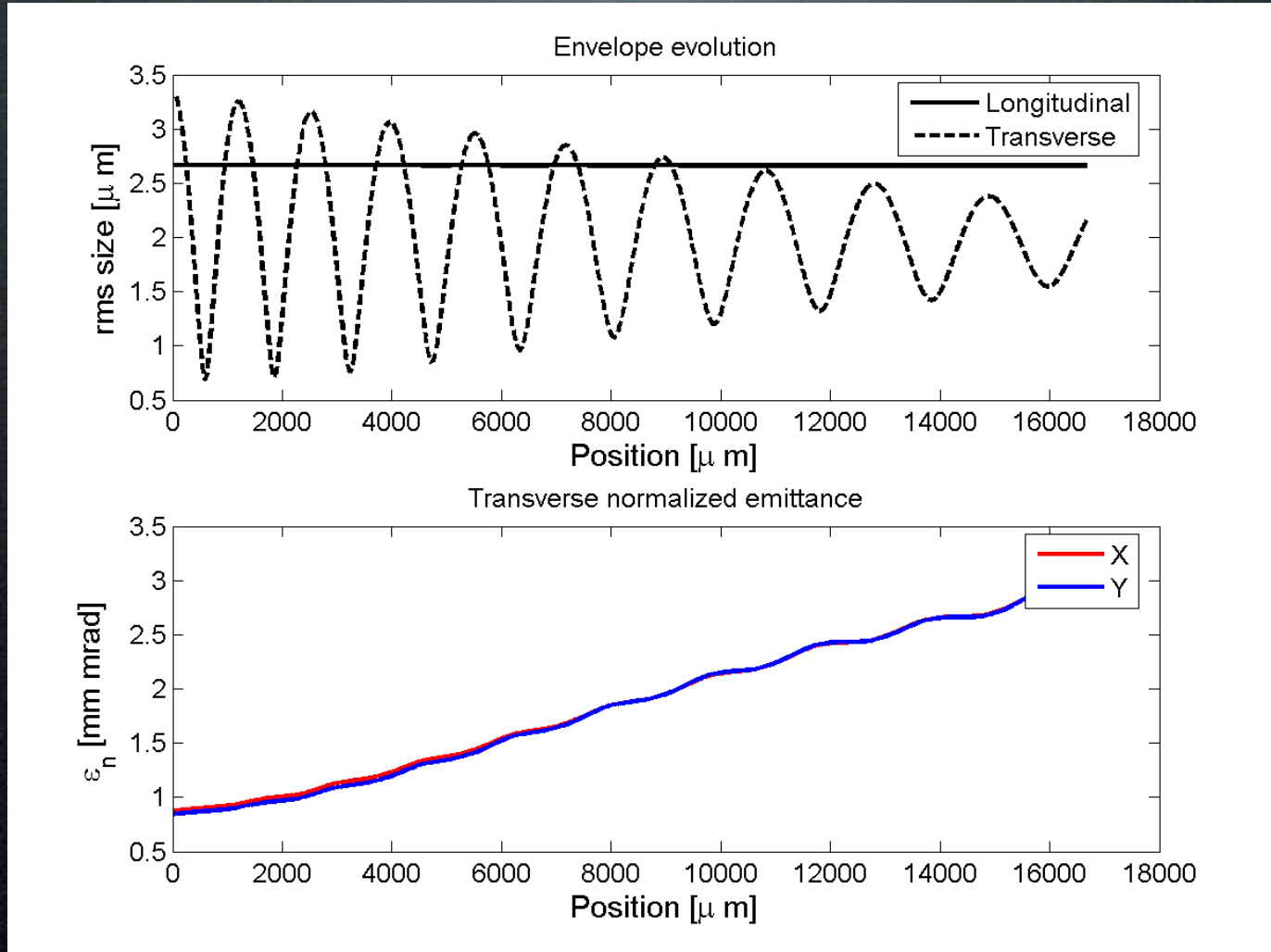
# Energy spread compensation with beam loading



**Fig. 5:** Linear beam loading example: (a) drive bunch density profile (red line) and longitudinal wakefield  $E_z$  (green line), (b) same for the witness bunch, (c) same for the drive and witness bunches together. The field of the drive bunch only is shown as the blue line in panel (c). A zoom around the witness bunch is shown in panel (d). The bunches move to the left.



# Transverse beam dynamics inside the plasma



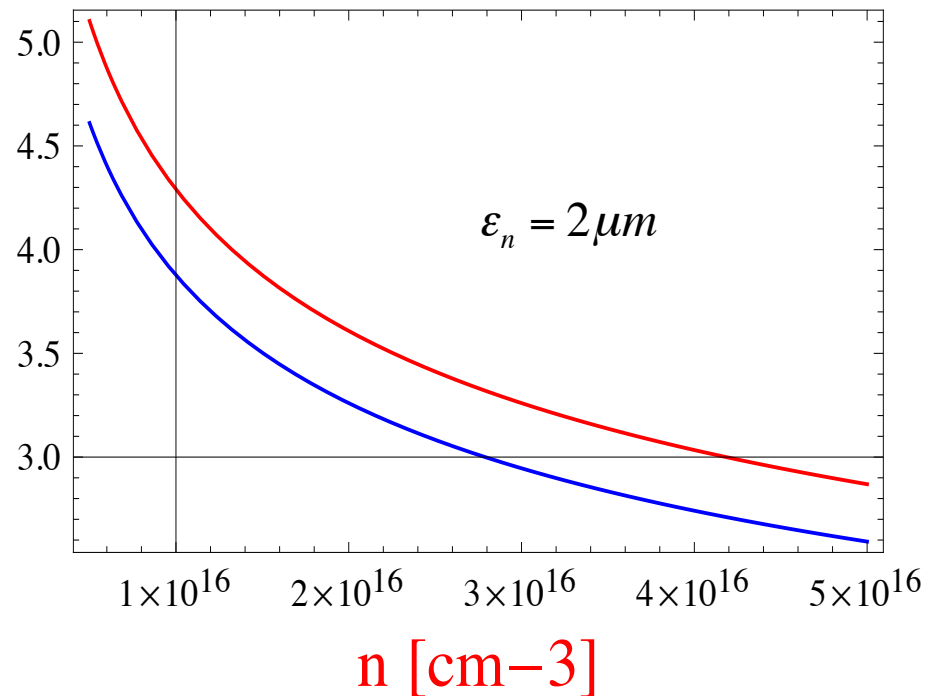
matching condition with the plasma:

$$\sigma_x'' + \frac{k_p^2}{3\gamma} \sigma_x = \frac{\epsilon_n^2}{\gamma^2 \sigma_x^3}$$

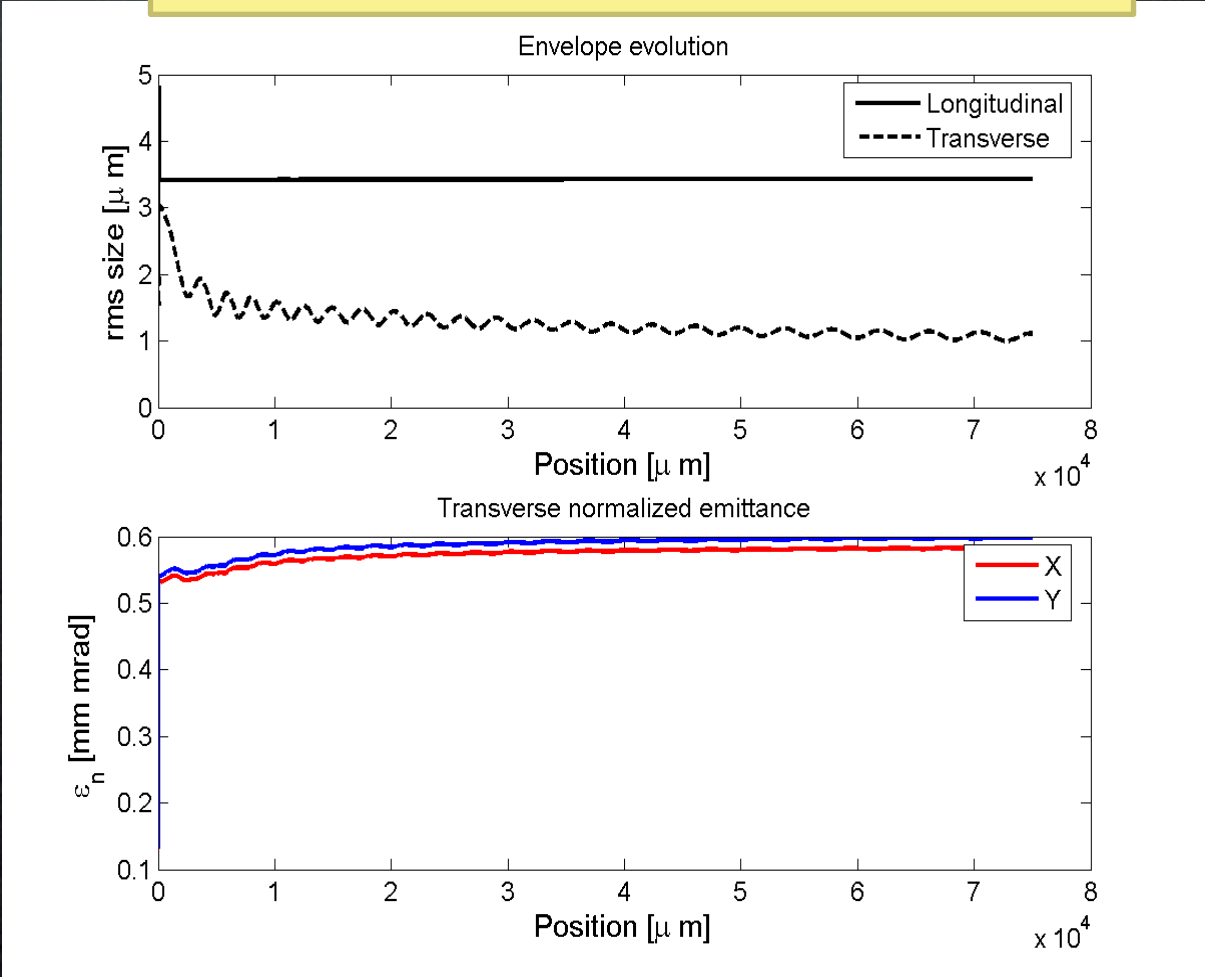
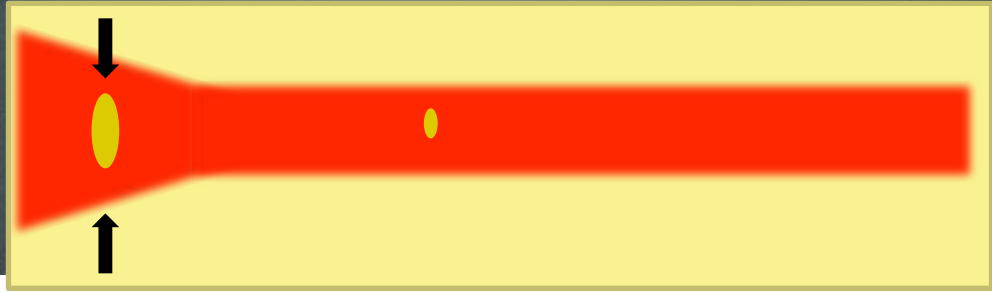
$$\sigma_\epsilon = \sqrt[4]{\frac{3}{\gamma}} \sqrt{\frac{\epsilon_n}{k_p}}$$

$$k_p^2 = \frac{e^2 n_1}{\epsilon_0 m c^2}$$

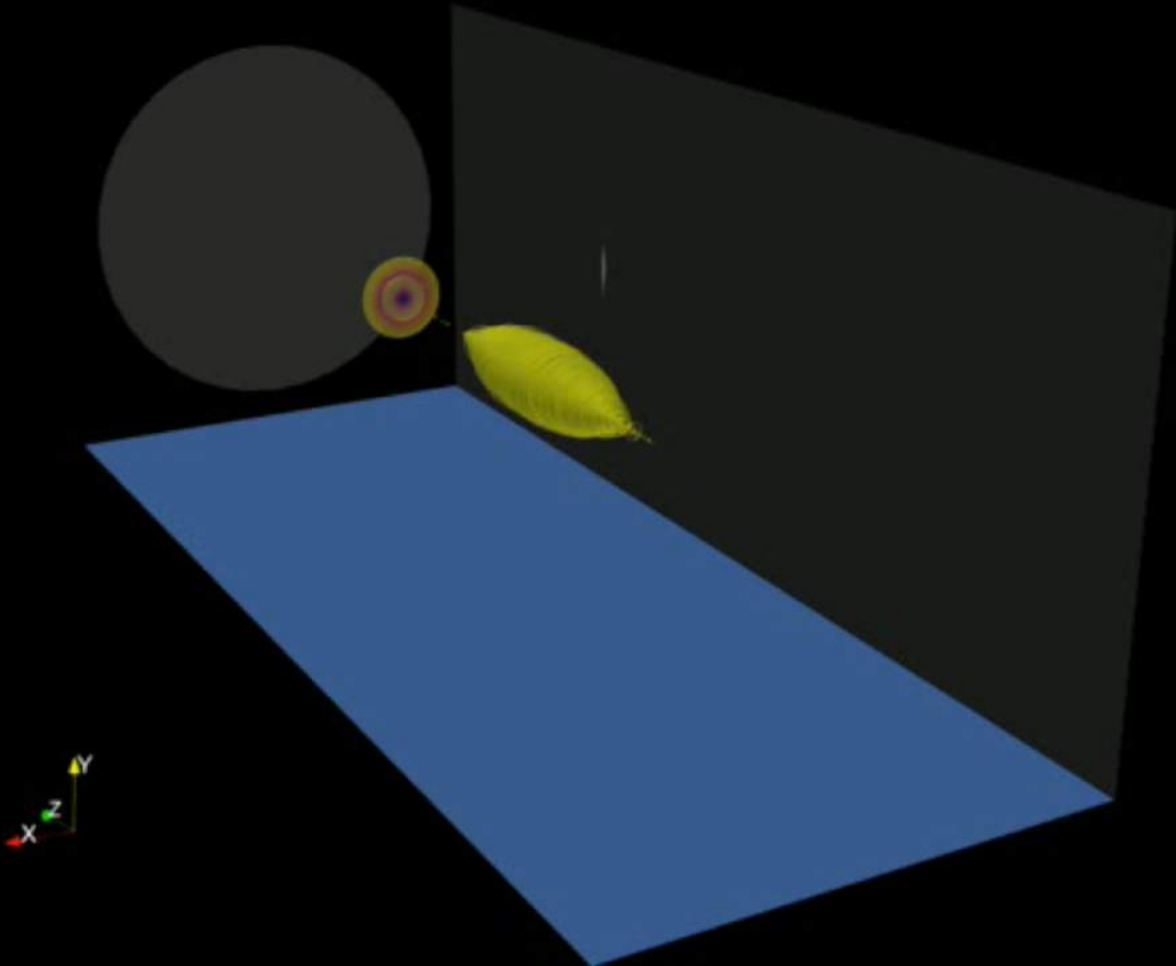
sigma\_r [um]



$$\sigma_\varepsilon = \sqrt[4]{\frac{3}{\gamma}} \sqrt{\frac{\varepsilon_n}{k_p}}$$

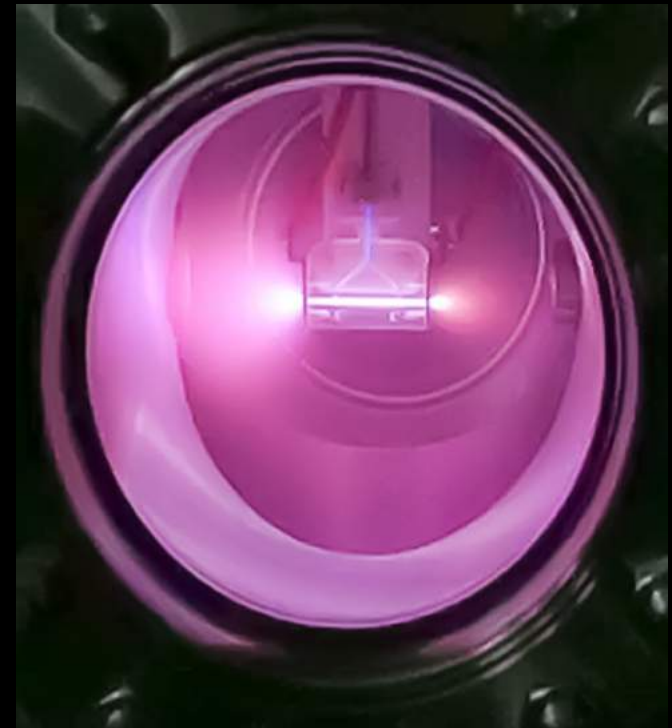








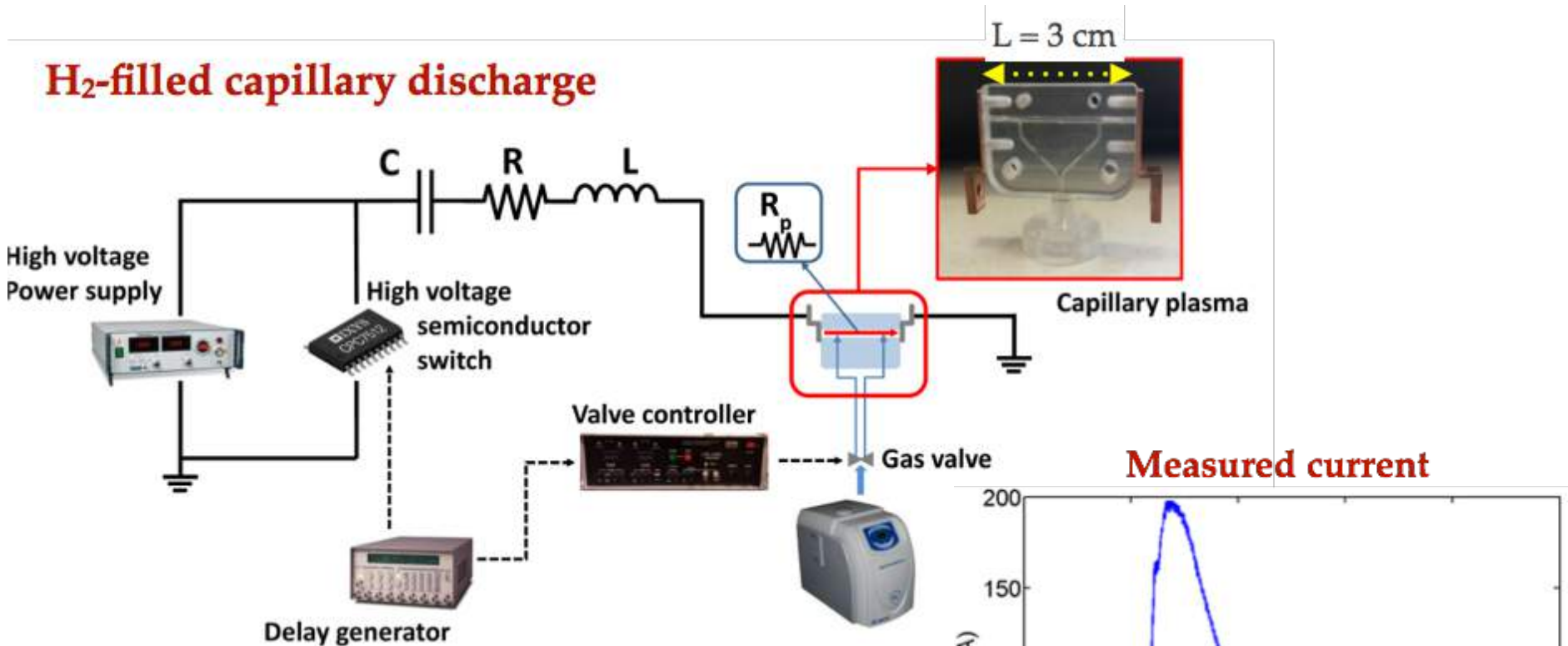
# Capillary Discharge at SPARC\_LAB



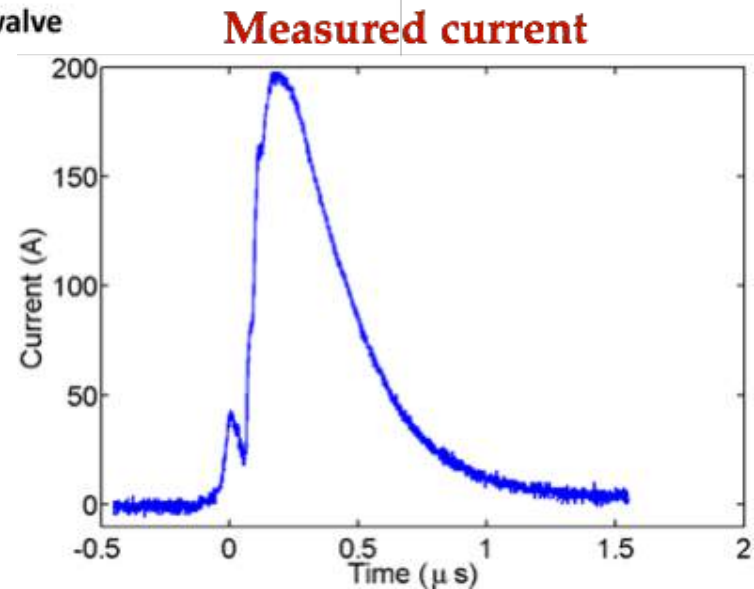


# Plasma Source

## H<sub>2</sub>-filled capillary discharge

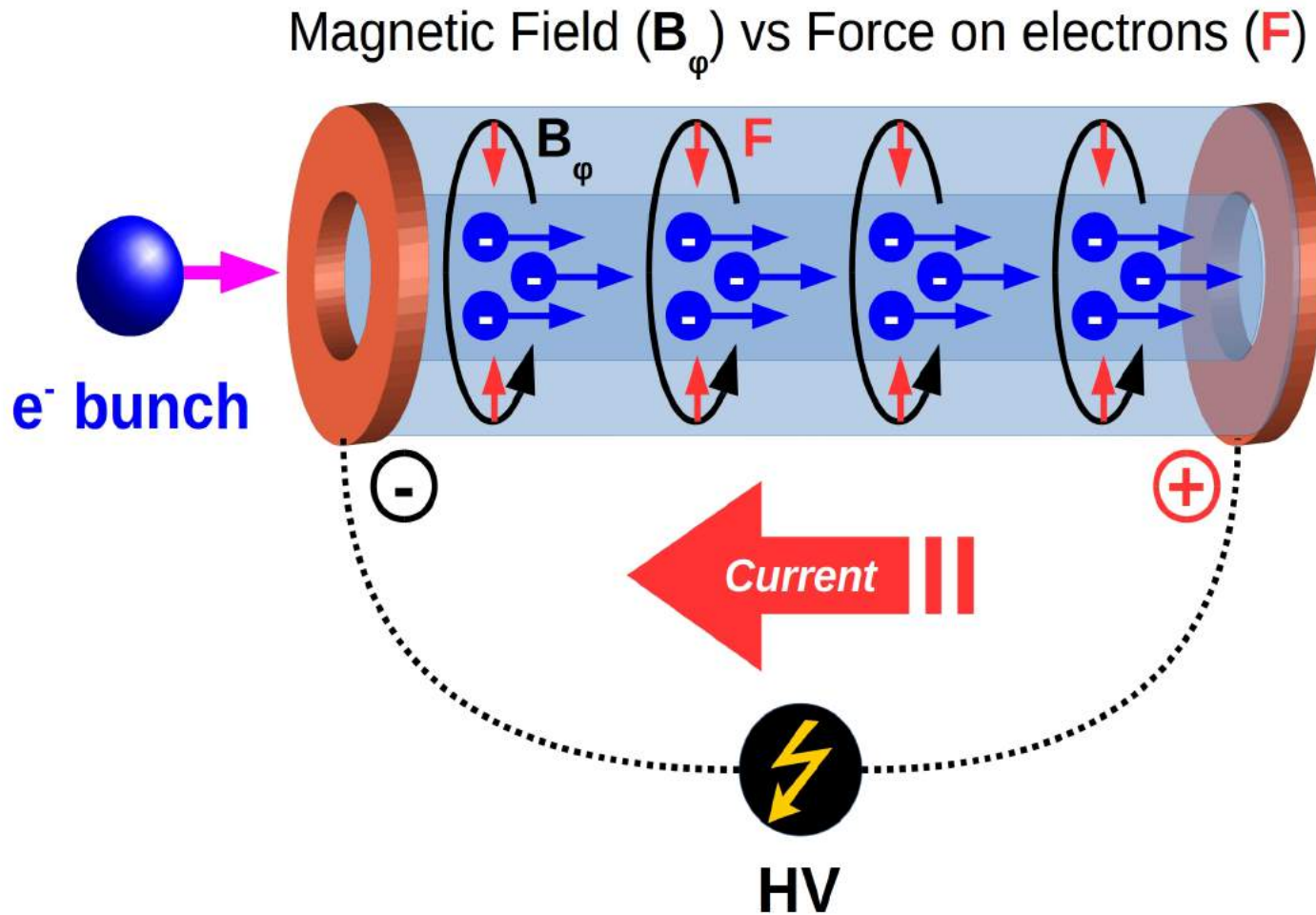


$P_{H_2} = 10$  mbar  
Total discharge duration: 800 ns  
Voltage: 20 kV  
Peak current: 200 A  
Capacitor: 6 nF

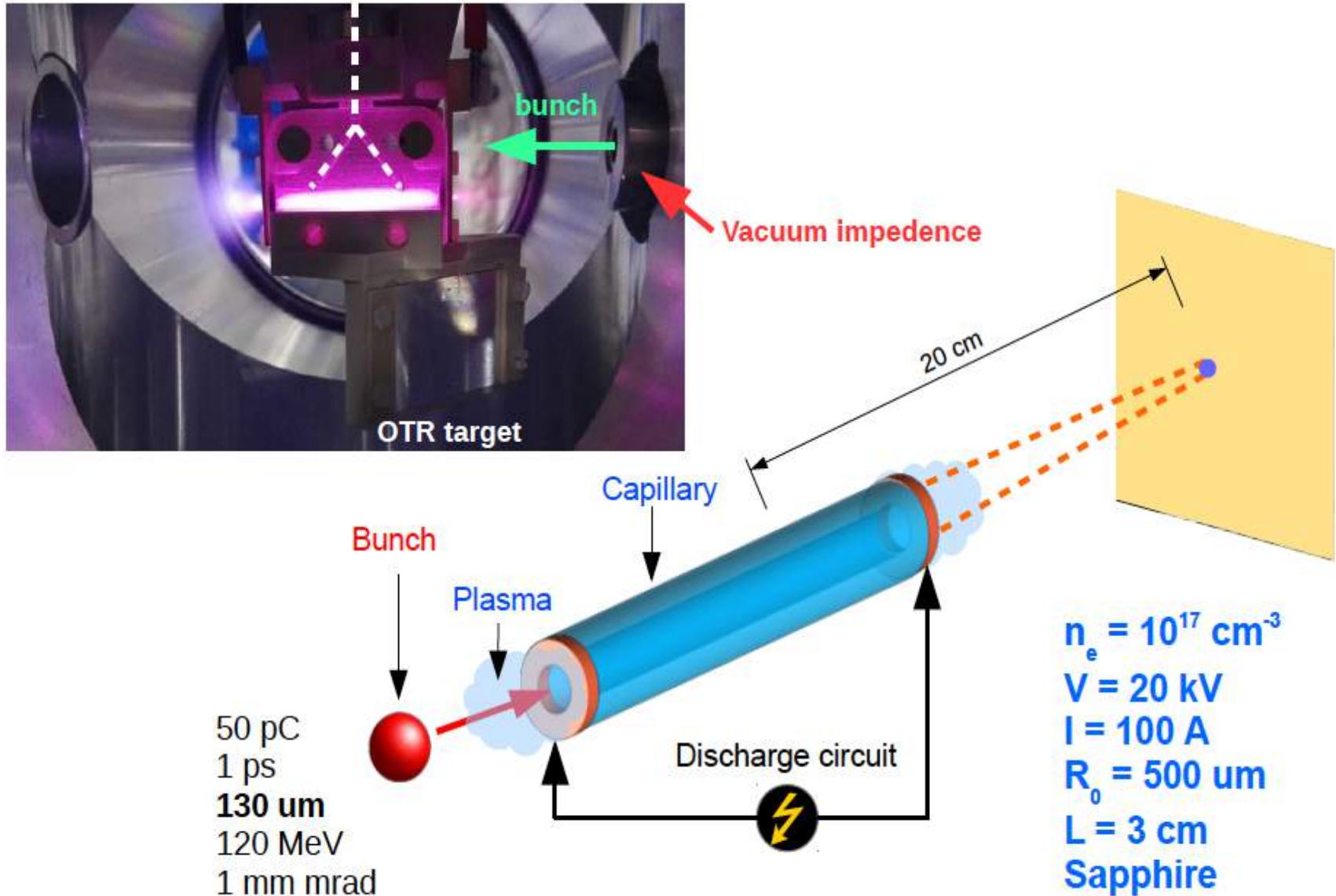


Courtesy of M. P. Anania, A. Biagioni, D. Di Giovenale, F. Filippi, S. Pella

# Active Plasma Lens

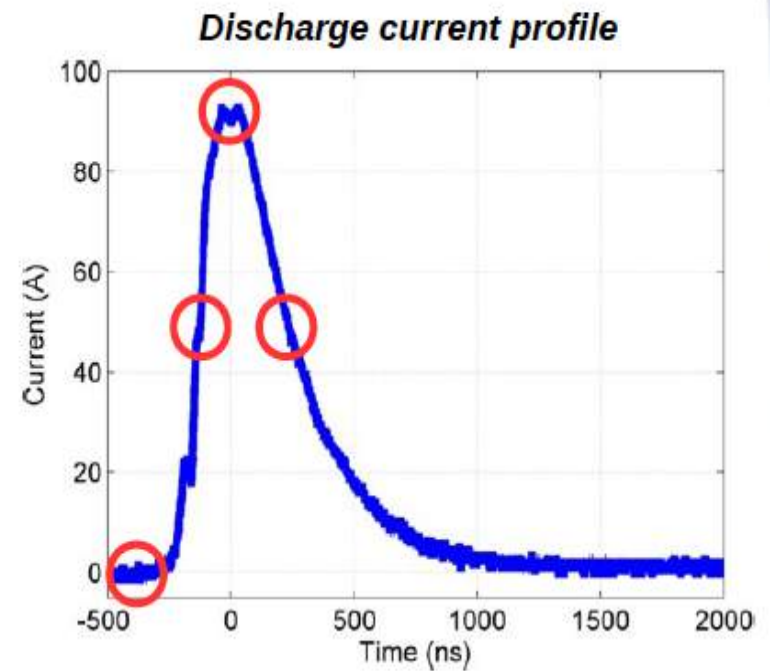
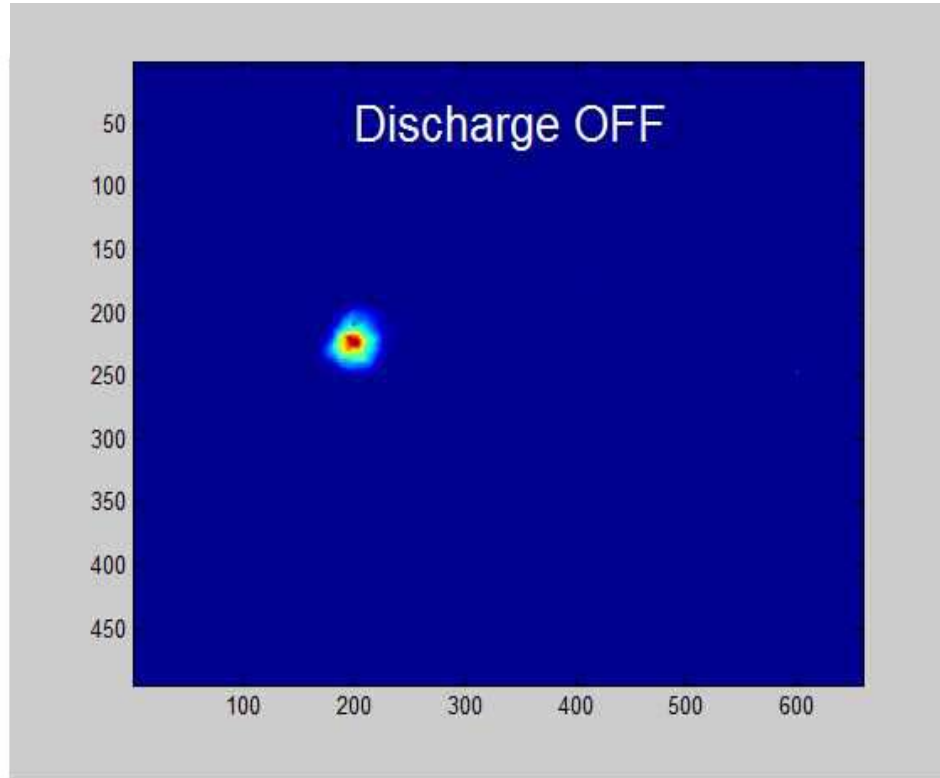


# Experimental layout



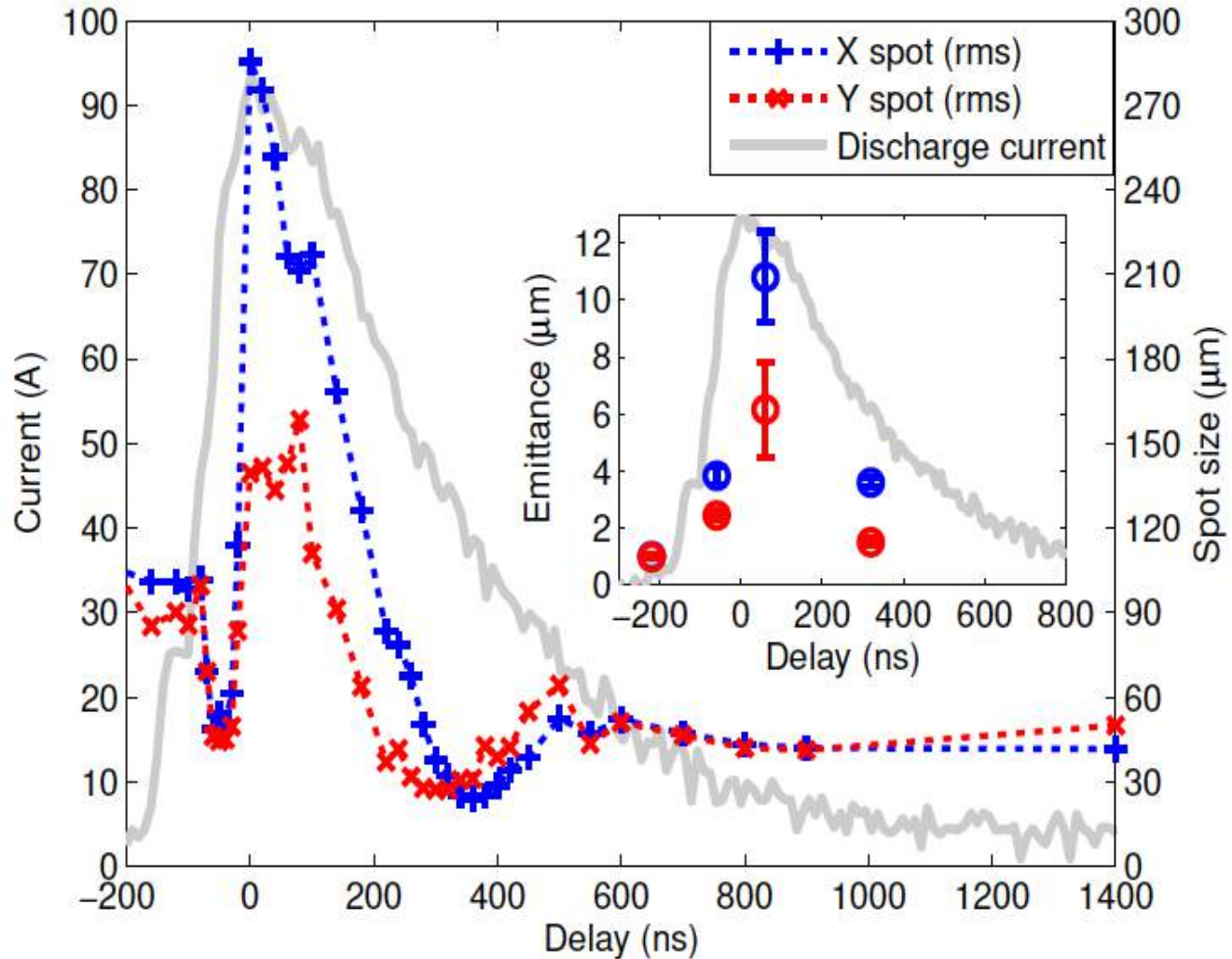


# Preliminary results

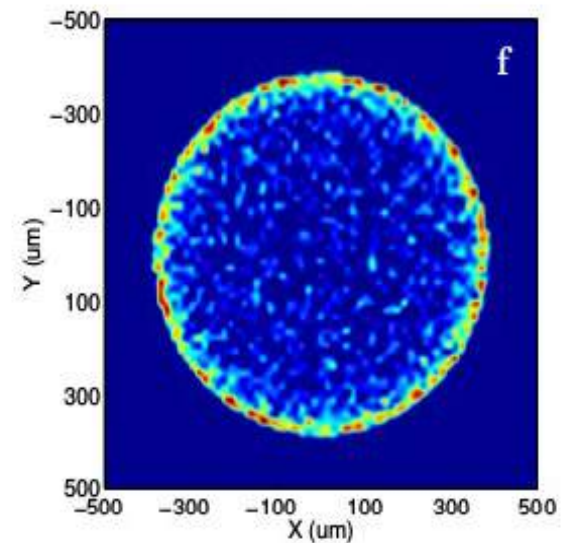
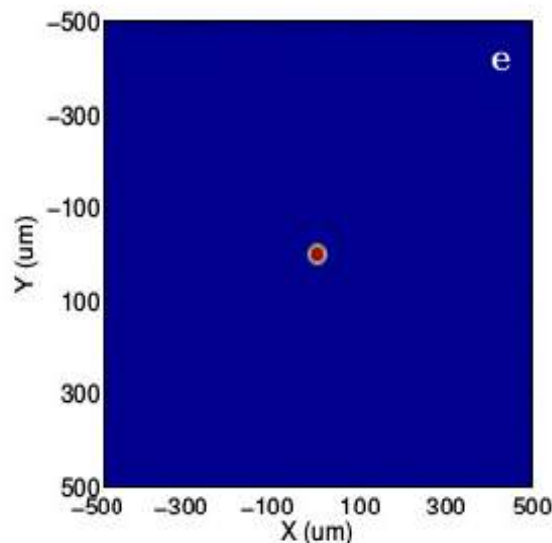
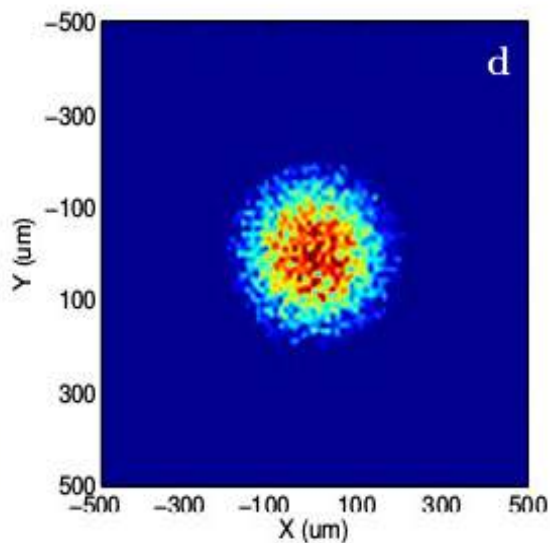
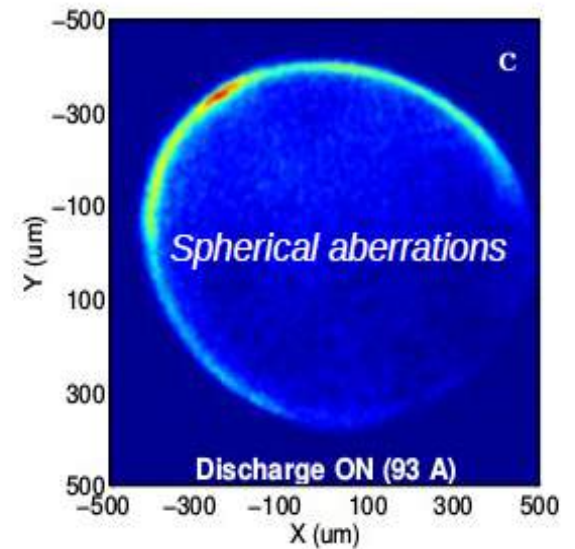
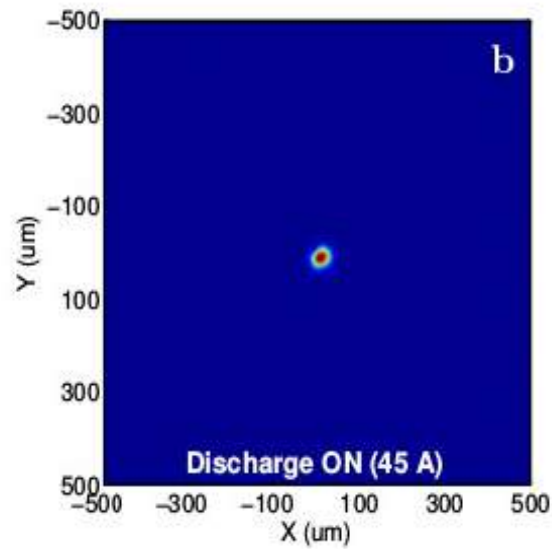
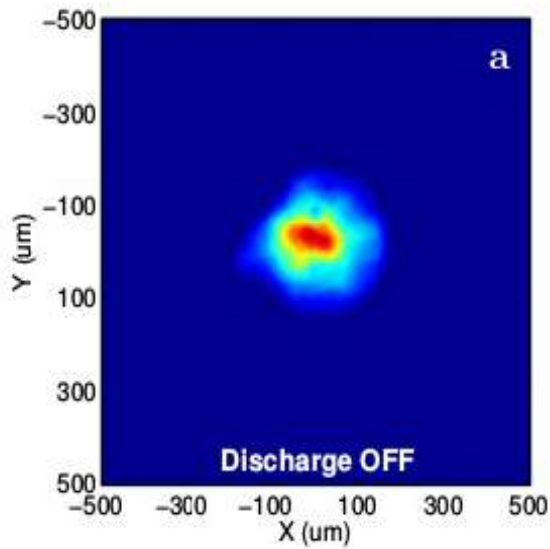


## Experimental characterization of active plasma lensing for electron beams

R. Pompili,<sup>1,a)</sup> M. P. Anania,<sup>1</sup> M. Bellaveglia,<sup>1</sup> A. Biagioni,<sup>1</sup> S. Bini,<sup>1</sup> F. Bisesto,<sup>1</sup> E. Brentegani,<sup>1</sup> G. Castorina,<sup>1,2</sup> E. Chiadroni,<sup>1</sup> A. Cianchi,<sup>3</sup> M. Croia,<sup>1</sup> D. Di Giovenale,<sup>1</sup> M. Ferrario,<sup>1</sup> F. Filippi,<sup>1</sup> A. Giribono,<sup>4</sup> V. Lollo,<sup>1</sup> A. Marocchino,<sup>1</sup> M. Marongiu,<sup>4</sup> A. Mostacci,<sup>4</sup> G. Di Pirro,<sup>1</sup> S. Romeo,<sup>1</sup> A. R. Rossi,<sup>5</sup> J. Scifo,<sup>1</sup> V. Shpakov,<sup>1</sup> C. Vaccarezza,<sup>1</sup> F. Villa,<sup>1</sup> and A. Zigler<sup>6</sup>

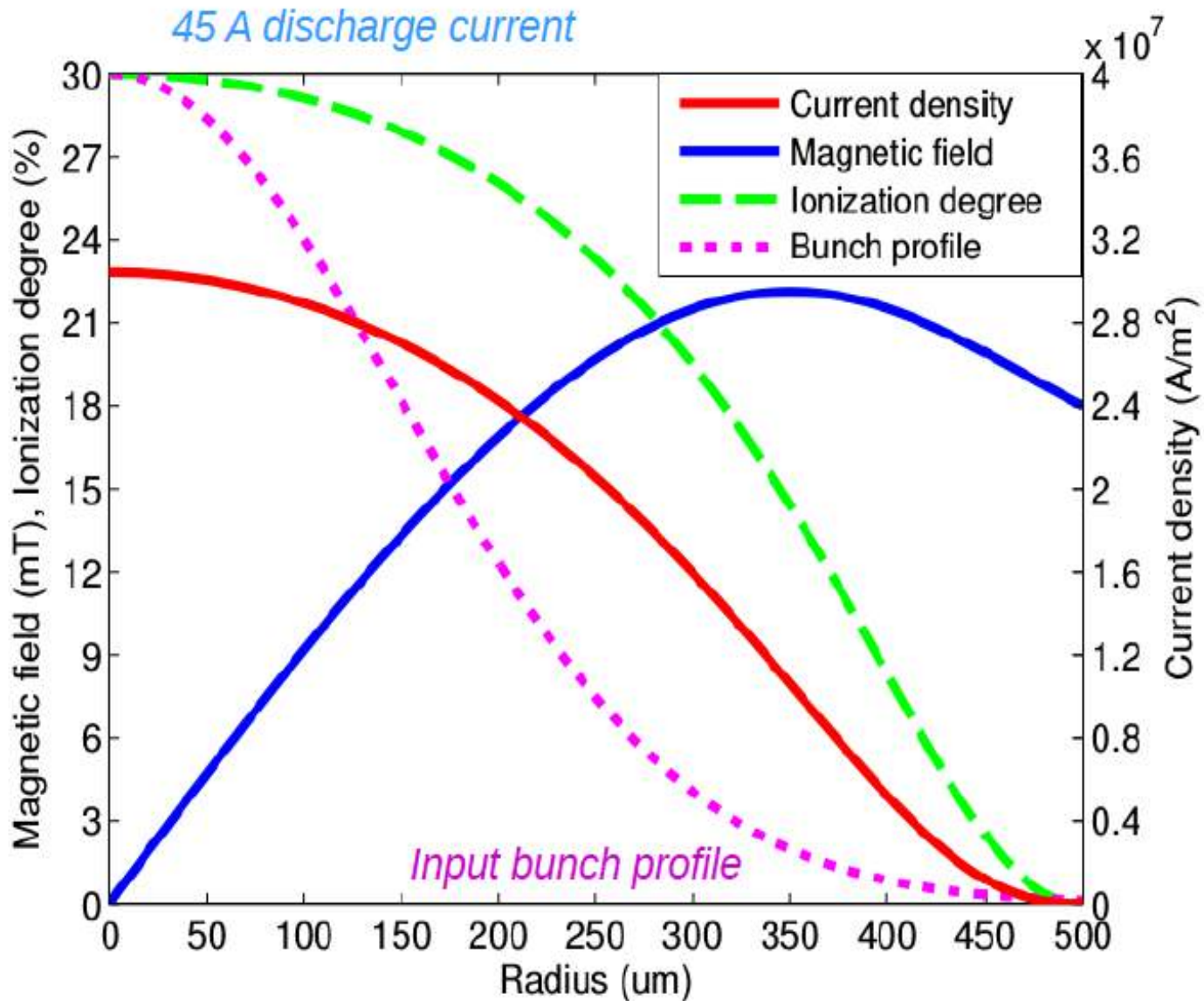


# Results vs simulations

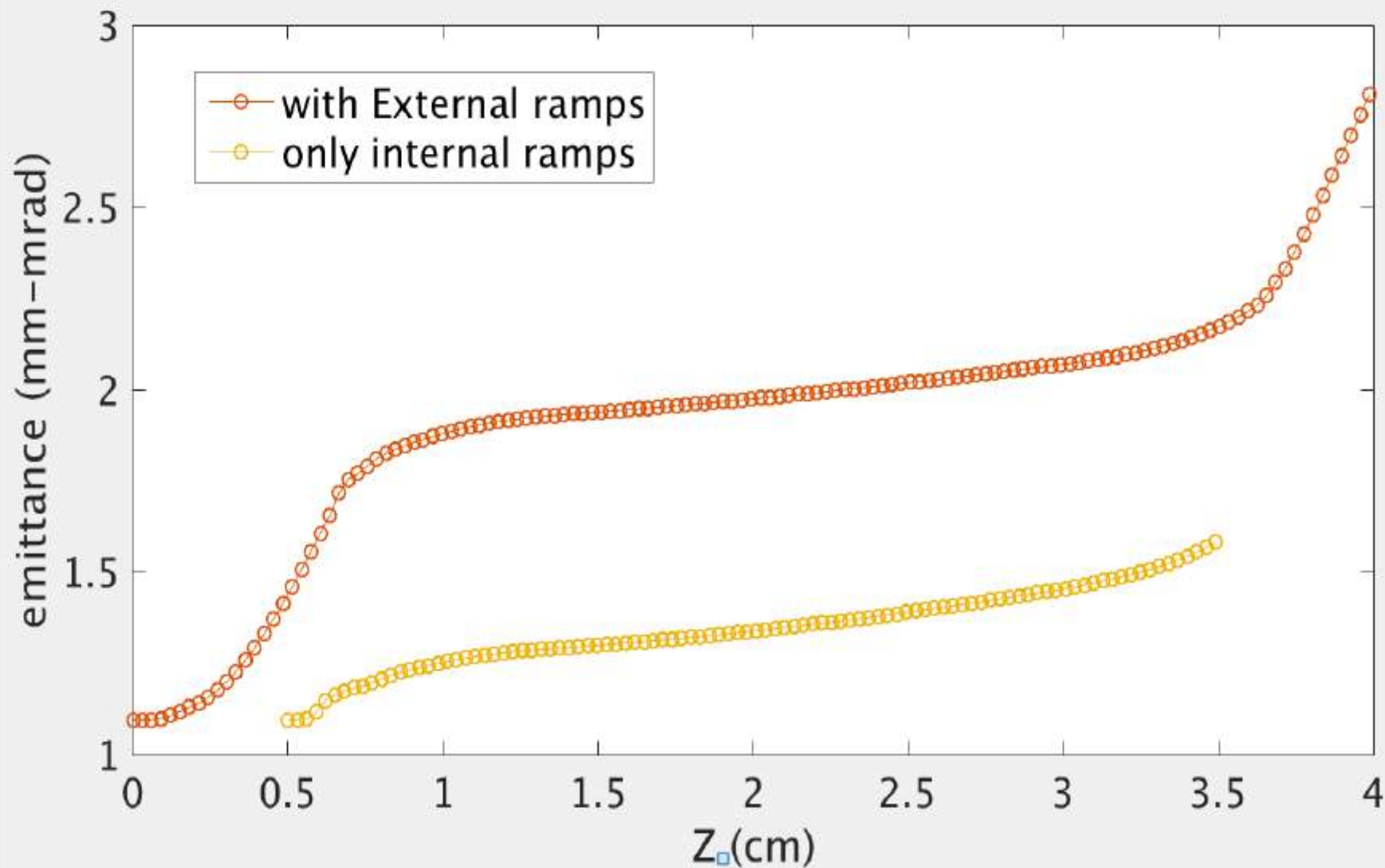


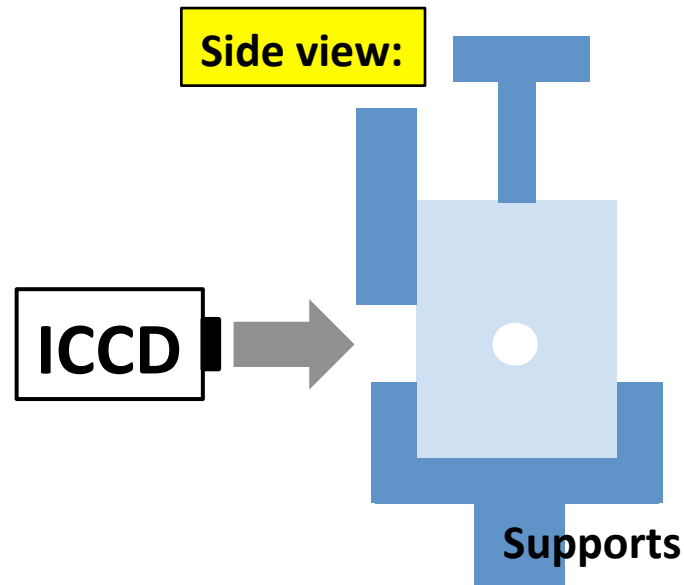
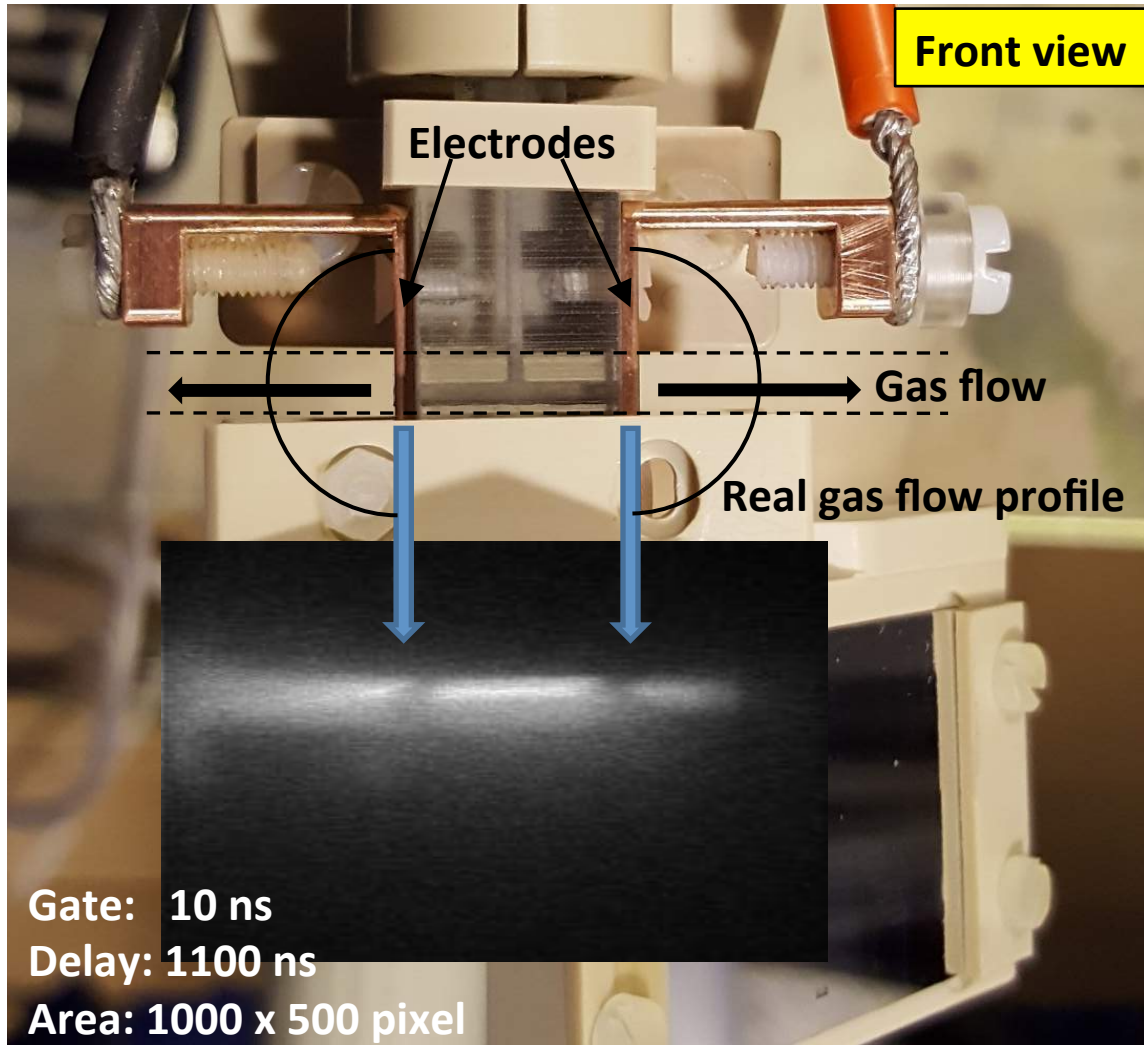


# Nonlinear focusing field



# Ramps effects on emittance without discharge





In order to see the real expansion of the plasma we have to mount the capillary of 3 cm length so that we will not see the cutting due to the supports



# velocity of plasma



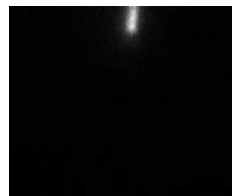
**Delay: 20 images separated  
by 100 ns = 2  $\mu$ s**

**Gate: 10 ns**

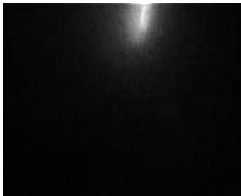
**Area: 1000 x 500 pixel**



0 ns



200 ns



400 ns



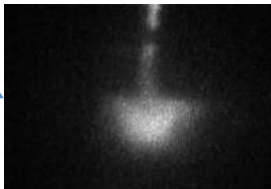
700 ns



1000 ns



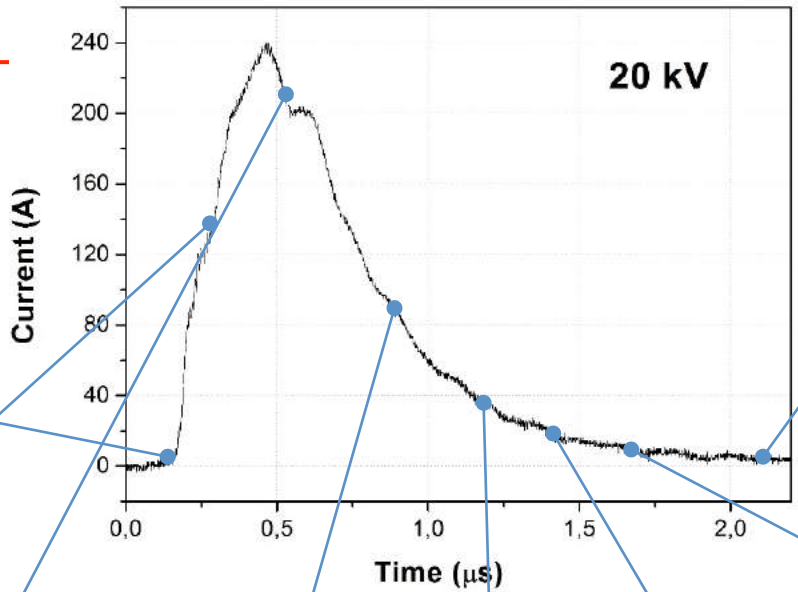
1300 ns



1600 ns



2000 ns



20 kV

Current (A)

Time ( $\mu\text{s}$ )

The CERN Accelerator School  
is organizing a course on

# PLASMA WAKE ACCELERATION

23-29 November, 2014

CERN, Geneva, Switzerland

The course will be of interest to staff and students in accelerator laboratories, university departments and companies working in or having an interest in the field of new acceleration techniques. Following introductory lectures on plasma and laser physics, the course will cover the different components

of a plasma wake accelerator and plasma beam systems. An overview of the experimental studies, diagnostic tools and state of the art wake acceleration facilities, both present and planned, will complement the theoretical part. Topical seminars and a visit of CERN will complete the programme.



Beatwave

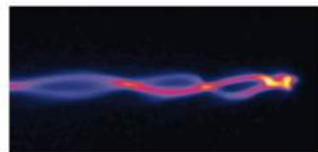


Diagram illustrating the interaction of a laser pulse with a plasma, showing the resulting wakefield structure.

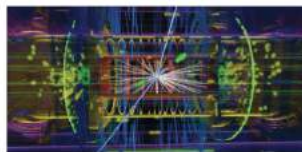


Diagram illustrating the interaction of a laser pulse with a plasma, showing the resulting wakefield structure.



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email: barbara.stasser@cern.ch  
http://cas.web.cern.ch/cas



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Supported by EU/ARIES via EuroNNAc3  
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Dielectric structures and other novel technologies  
Advanced and novel accelerators for high energy physics  
High gradient and multibunch acceleration in metallic structures  
(C-X-band and beyond) with innovative power generation schemes  
Plasma accelerators driven by: modern lasers, electron beams, proton beams  
Computations for accelerator physics advanced beam diagnostics for beams and plasma  
Novel schemes using advanced technologies (table-top FEL, medical imaging ...)



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