

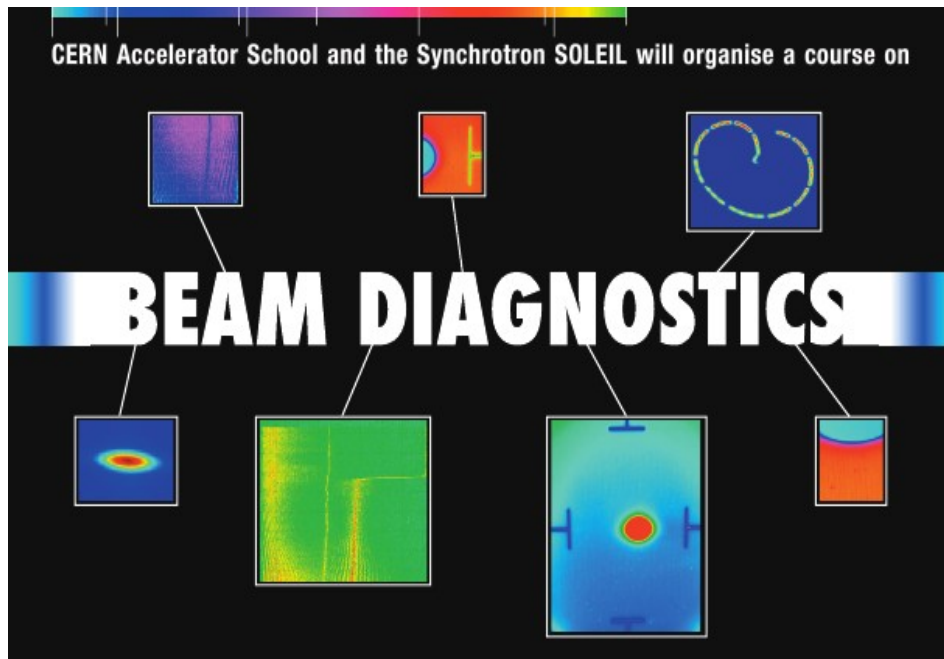
# Tune and Chromaticity Diagnostics

## Part I

**Ralph J. Steinhagen**

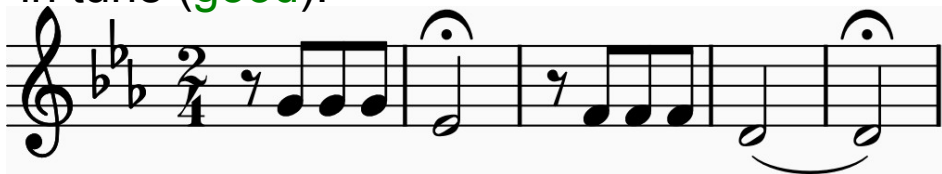
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**Acknowledgments: A. Boccardi, P. Cameron (BNL), M. Gasior,  
R. Jones, H. Schmickler, C.Y. Tan (FNAL)**

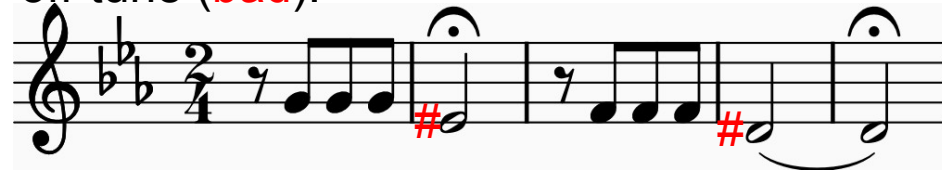


- Laymen/Musician's view (Beethoven's 5th):

in tune (good):



off-tune (bad):



- Audience will leave the concert
- ↔ Beam will leave the vacuum pipe

- Importance of tune:
  - defines beam life-time
  - strong impact on beam physics experiments:



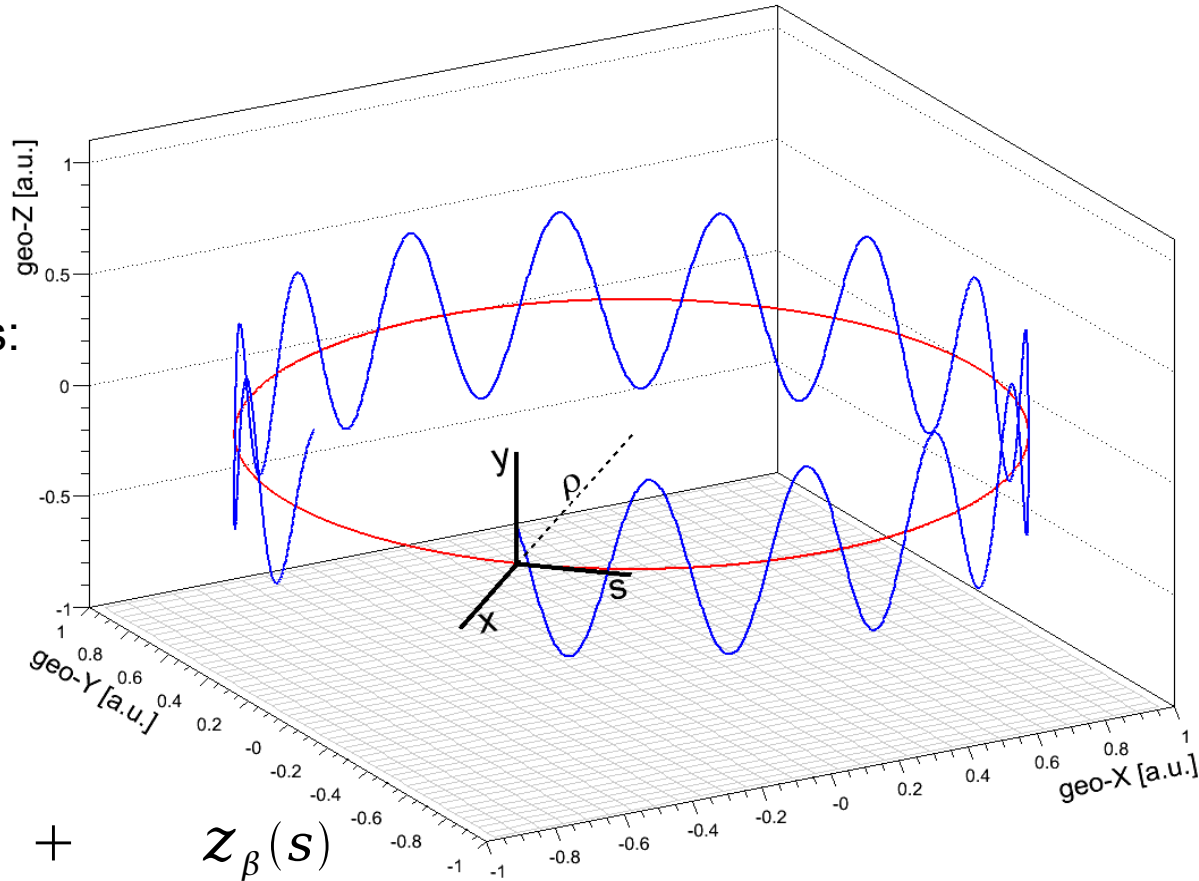
"I don't think we've quite repeated the experiment - last time we did it, the glass gave out a middle 'c'!"

- Hill's equation

... the mother of all accelerator physics:

$$z'' + k(s) \cdot z = f(s, t)$$

- $k(s)$ : focusing strength, defines:
  - phase advance  $\mu(s)$
  - betatron function  $\beta(s)$
- $f(s, t)$ : driving force



- first-order solution:

$$z(s) = \underbrace{z_{co}(s)}_{\text{closed orbit}} + \underbrace{D(s) \cdot \frac{\Delta p}{p}}_{\text{dispersion orbit}} + \underbrace{z_{\beta}(s)}_{\text{betatron oscillations}}$$

- $D(s)$ : dispersion function [m] → typically: few cm to a few meters
- $\Delta p/p$ : relative momentum offset w.r.t. c.o. → typically:  $10^{-3} \dots 10^{-4}$

- Main tune dependent part:

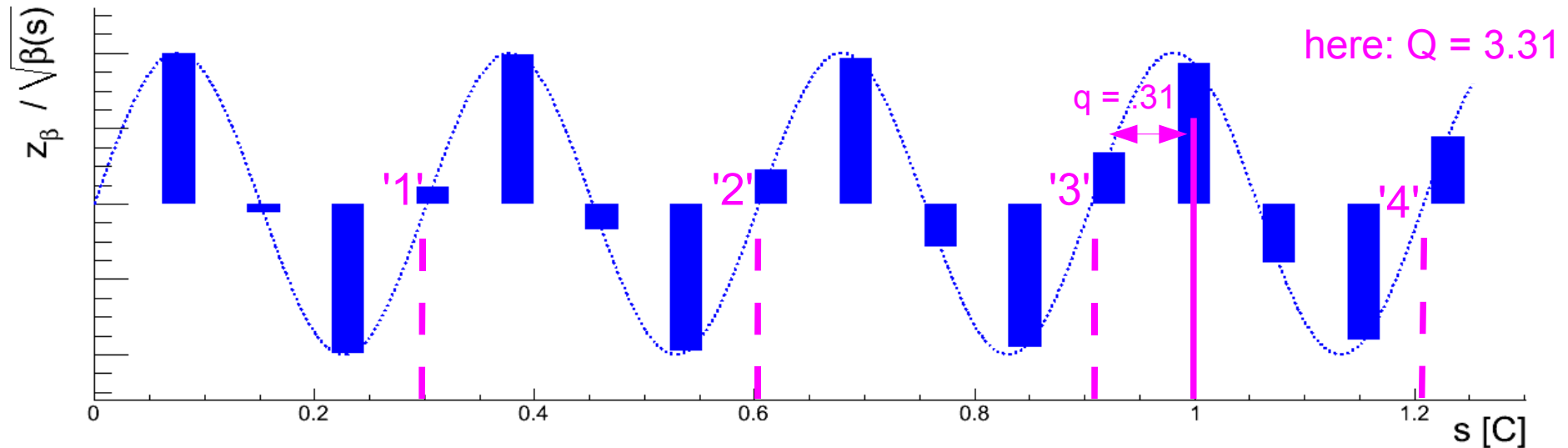
$$z_{\beta}(s) = \sqrt{\epsilon_i \beta(s)} \cdot \sin(\mu(s) + \phi_i)$$

$\epsilon_i \phi_i$  : initial particle state

→ particle describe sinusoidal oscillations in a circular accelerator

- Free Betatron Oscillations:

$$z_{\beta}(s) = \sqrt{\epsilon_i \beta(s)} \cdot \sin(\mu(s) + \phi_i)$$



- Betatron Phase Advance:  $\mu(s)$
- Tune defined as betatron phase advance over one turn:

$$Q := \frac{1}{2\pi} \oint_C \mu(s) ds$$

common:  $Q = \underbrace{Q_{int}}_{\text{integer tune}} + \underbrace{q_{frac}}_{\text{fractional tune}}$

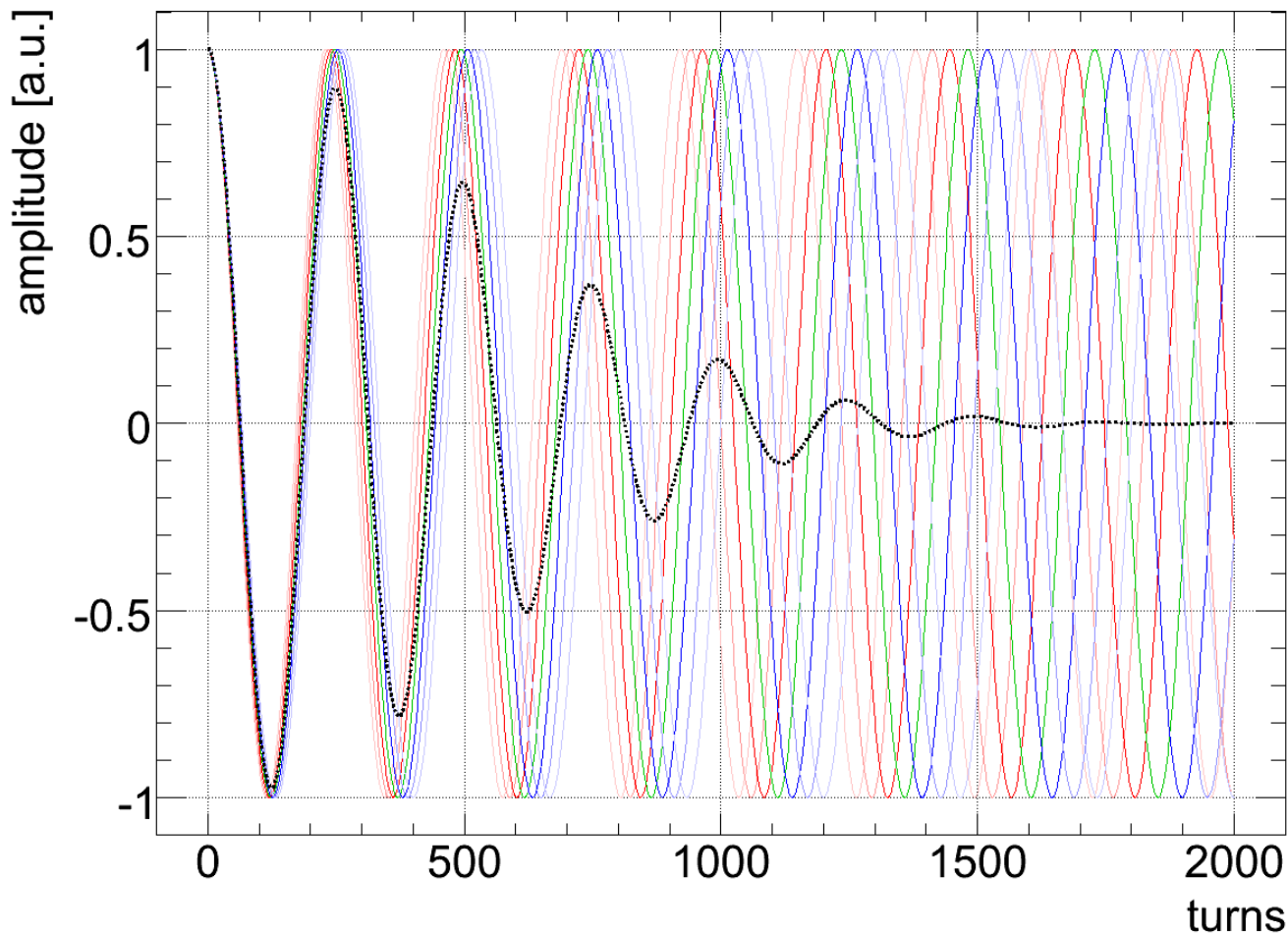
- Tune measurement options:

1. Single-turn: 'count oscillations along circumference' (usually while threading 'first turn')
2. Turn-by-turn: pick and observe the oscillation at a given single BPM

$$\Delta z_{\beta} = \sqrt{\epsilon_i \beta} \cdot \sin(\mu + \phi_i + 2\pi Q \cdot n)$$

→ FFT analysis returns  $q_{frac}$

- Individual bunch particles usually differ slightly w.r.t. their individual tune  
 → Literature: “Landau Damping” (Historic misnomer: particle energy is preserved!)



- E.g. if  $f(\Delta Q)$  is a narrow Gaussian distribution with  $\sigma_Q \ll Q$ :

$$\bar{z}(t) = \underbrace{\bar{z}_0}_{\text{dampening}} \cdot \underbrace{e^{-\frac{1}{2} \cdot \sigma_Q^2 n^2} \cdot \cos(2\pi Q \cdot n)}_{\text{tune oscillations}}$$

→ large tune spread ↔ fast damping of e.g. head-tail instabilities

→ Tune oscillations are usually damped

## Part I:

- Recap: What the .... is 'Q', Oscillations Dampening → **just done**
  - Perturbation Sources, Requirements
- Tune Diagnostics
  - Classic Fourier-Transform Based
    - Detectors: BPMs, Diode-Peak-Detection, (Schottky → **F. Casper**)
  - Phase-Locked-Loop (PLL) Systems
- Advanced Topic → **your choice**

## Part II: → **in about an hour**

- Recap: Definitions, Requirements & Constraints
- Classic Chromaticity Diagnostics
  - Momentum shift  $\Delta p/p$  based Q' tracking methods → LHC examples
- Collective Effects
  - Head-tail phase shift
  - De-coherence based methods: PLL Side-Exciter

- Why do we need to measure the tune at all? Does it change?

- Quadrupole strength (hor. focusing):  $k(s) = \frac{q}{p} \frac{\partial B}{\partial x}$

- Quadrupole gradient errors:  $k(s) \rightarrow k_0(s) + \Delta k(s)$

- saturation of iron yoke, magnet calibration errors, power converter ripple, etc.

$$\Delta Q = \frac{1}{4\pi} \beta(s) \cdot \Delta k(s)$$

→ watch out for quadrupole errors at large beta functions (e.g. final focus)!

- Energy perturbation  $p \rightarrow p_0 + \frac{\Delta p}{p_0}$

- Main dipoles vs. quadrupoles mismatch → *natural chromaticity*  $Q'_{nat}$

$$\Delta Q = -\frac{1}{4\pi} \beta(s) \cdot \left( k(s) \cdot \frac{\Delta p}{p_0} \right) \sim Q'_{nat} \cdot \frac{\Delta p}{p_0}$$

- RF frequency change (aka. radial steering)

$$\Delta Q := Q' \cdot \frac{\Delta p}{p_0}$$

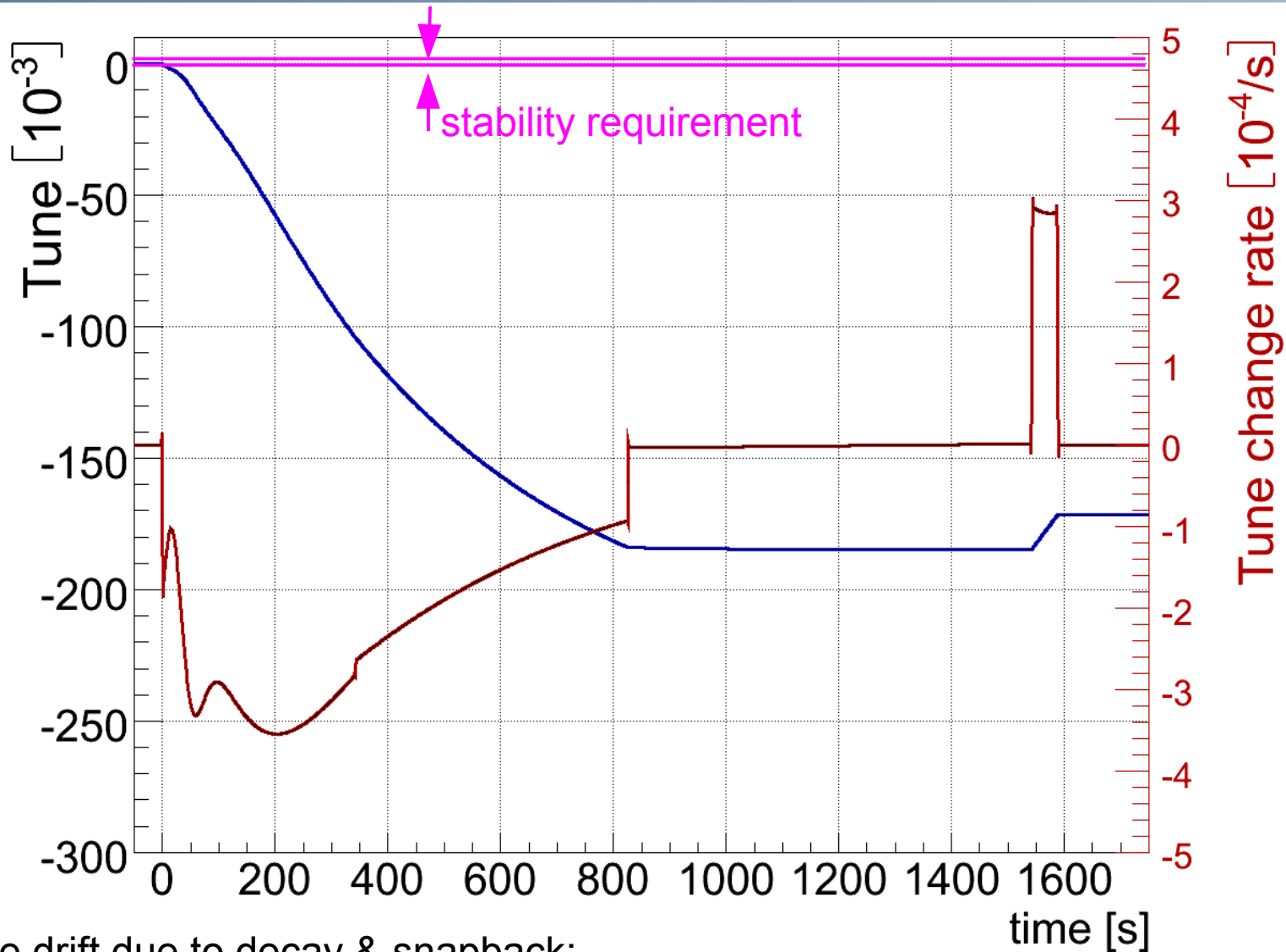
→ defines machine's *chromaticity*  $Q'$

subtle but important difference:  
LHC:  $Q'_{nat} \approx -140$  but  $Q' \approx 1$

→ next lecture

→ bottom line: tune is usually not a constant

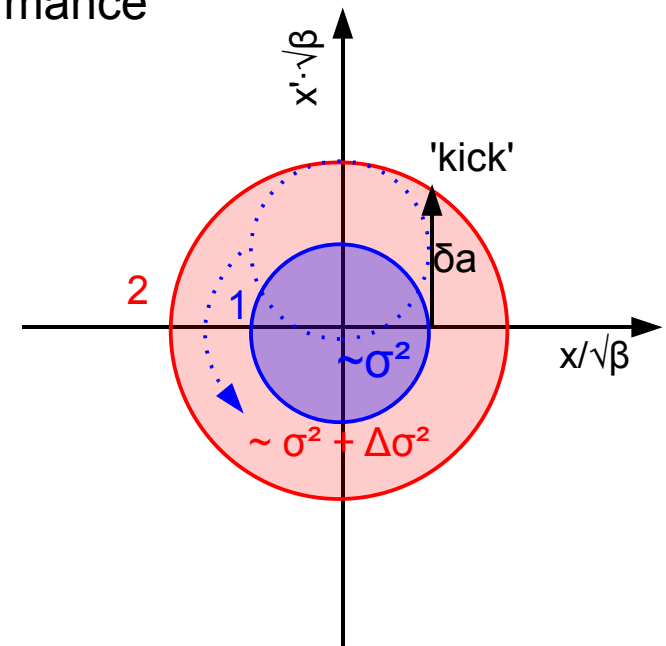




- LHC Tune drift due to decay & snapback:
  - effect intrinsic to superconducting magnets
  - Tune drift (without  $b_3$  effects):  $\Delta Q \approx 0.1$
  - Tune change rate:  $\Delta Q/\Delta t|_{\max} < 10^{-3} \text{ s}^{-1}$



- Transverse beam size as an impact on accelerator performance
  - smaller beam-sizes  $\sigma$  favourable
    - HEP colliders: higher luminosity
    - Light Sources: higher brightness



- beam size increases quadratically with angular kick  $\delta a$

$$\frac{\Delta \sigma}{\sigma} \approx \frac{1}{2} \left( \frac{\delta a}{\sigma} \right)^2$$

- N.B. for electrons, esp. synchrotron light sources, this is partially compensated by energy losses due to synchrotron light radiation.
- Protons: memory effect – the beam does not forgive...!
  - LHC limit:  $\delta a \ll 10 \mu\text{m} = \sim 1/20 \sigma$  !!
- Further constraints on kick amplitudes: aperture limitations due to functional insertion, machine protection systems, ..

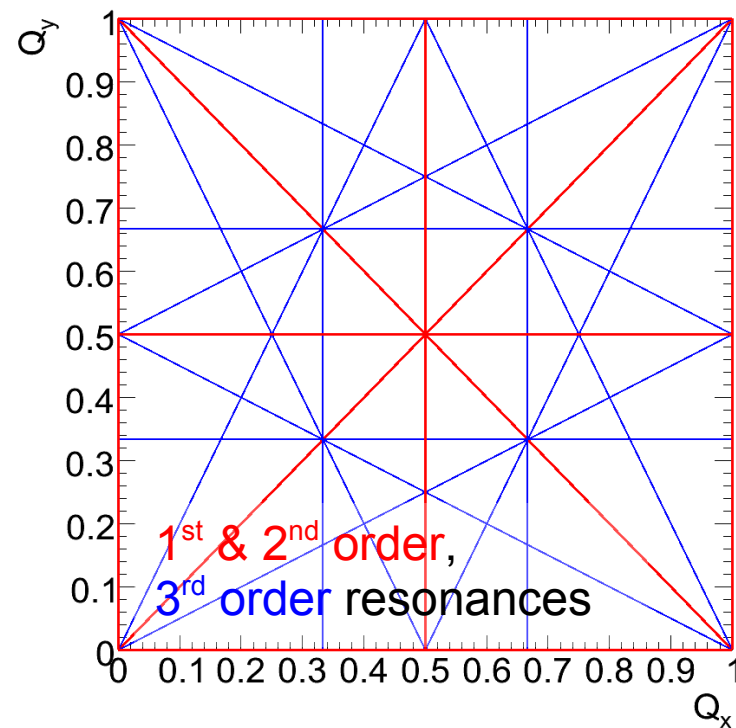
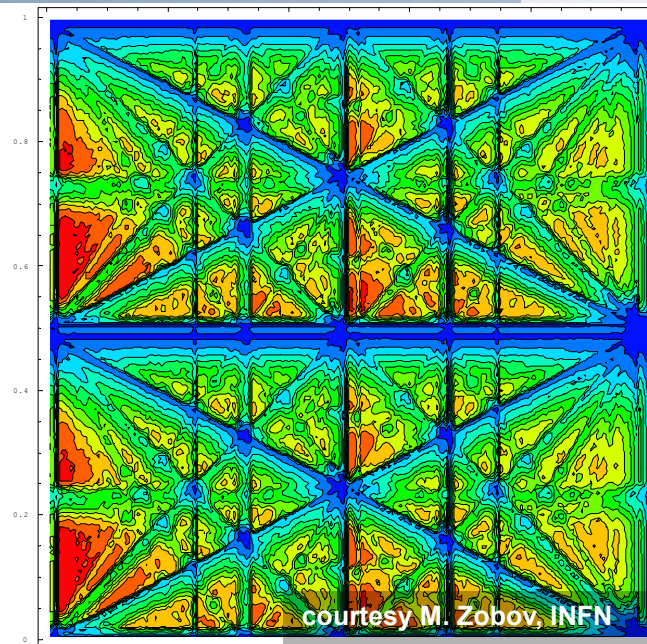
→ Limit excitation to necessary minimum, favours passive/sensitive systems

- Unstable particle motion reduces beam-lifetime (~dynamic aperture) if resonance condition is met:

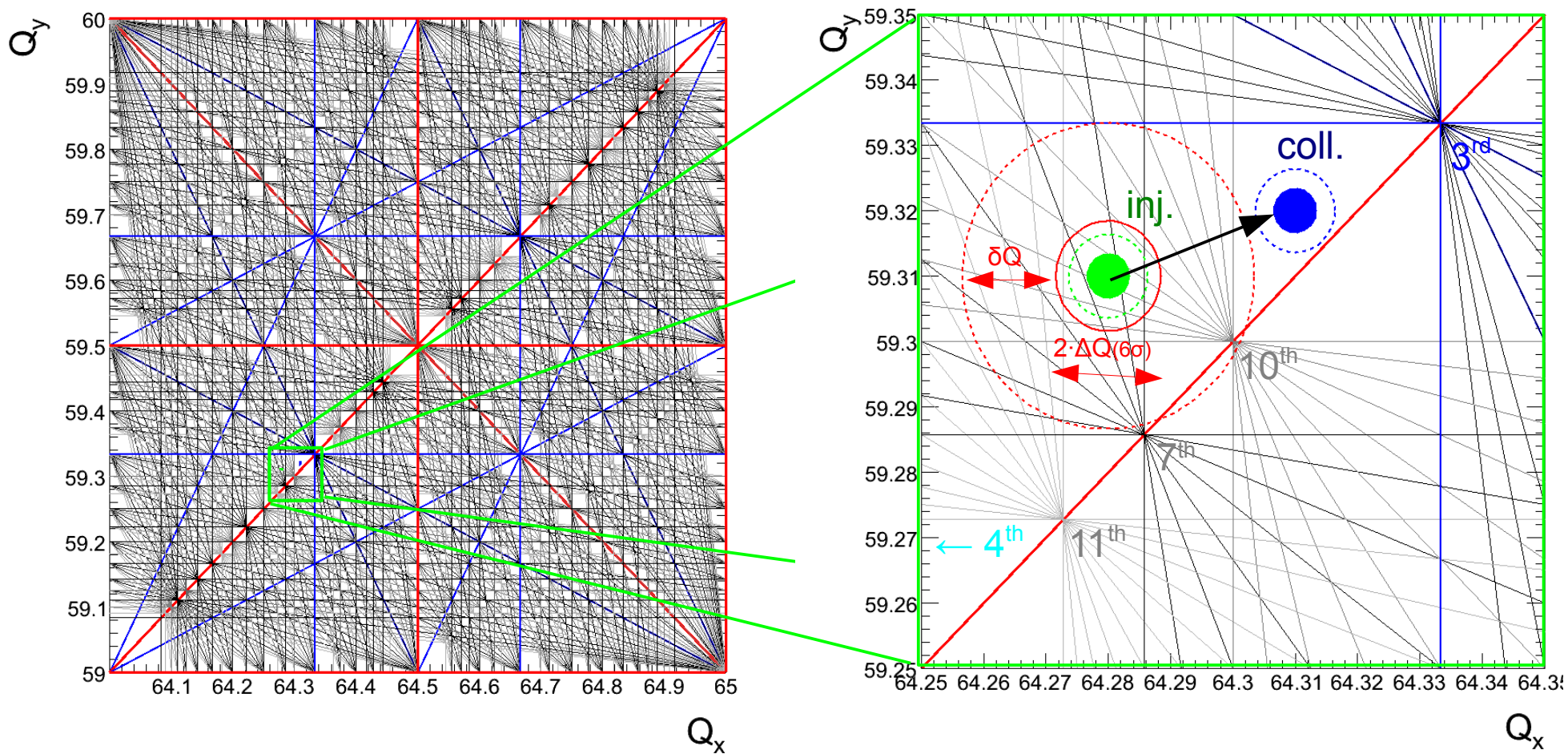
$$p = m \cdot Q_x + n \cdot Q_y \quad \wedge \quad m, n, p \in \mathbb{Z}$$

- similar relation also in between  $Q_x$  &  $Q_s$  (important for lepton accelerators)
- Resonance order:  $O = |m| + |n|$
- Lepton accelerator: avoid up to  $\sim 3^{\text{rd}}$  order
- Hadron colliders:
  - negligible synchrotron radiation damping
  - need often to avoid up to the  $12^{\text{th}}$  order

*“Hadron beams are like elephants – treat them bad and they'll never forgive you!”*



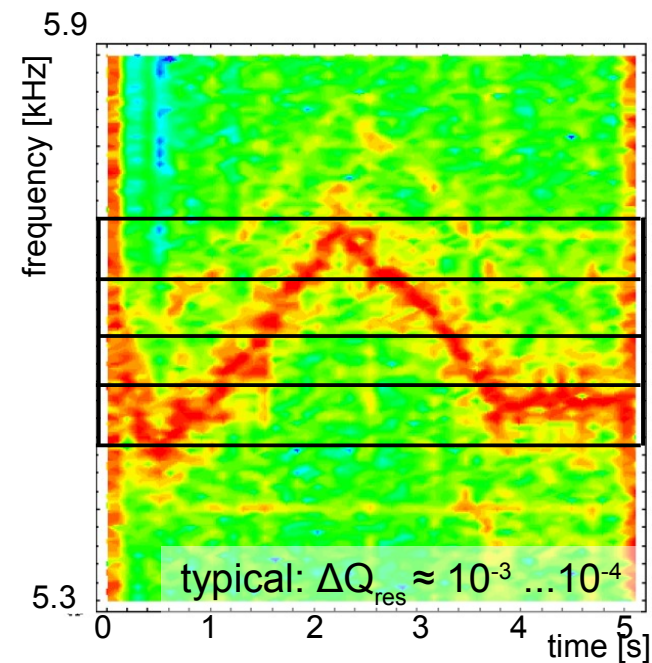
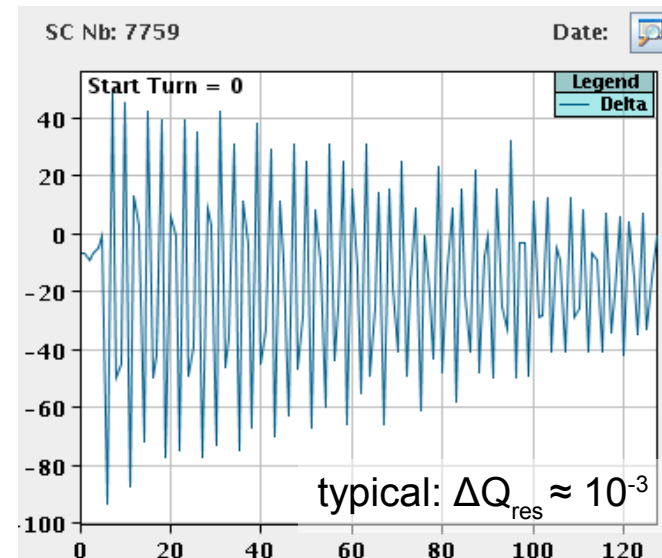
- Example LHC: Tune stability requirement:  $\Delta Q \approx 0.001$  vs. exp. drifts  $\sim 0.06$



- N.B. need to stay much further off these resonance lines due to
  - finite tune width: chromaticity, space charge, momentum spread, detuning with amplitude and resonance's stop band itself

- Classic, using BPMs with 'kick' or 'chirp' excitation
  - limited by aperture constraints
    - Performance reduction
      - typically:  $\Delta z \leq 0.1 \sigma$
    - Loss of particles & protection
      - LHC:  $\Delta z \leq 25 \mu\text{m}$  &  $\Delta p/p \leq 5 \cdot 10^{-5}$
  - limited by emittance blow-up
  
- Passive monitoring of residual oscillations:
  - Schottky monitors
  - Diode-Detection based Base-Band-Q (BBQ) meter
  
- Active Phase-Locked-Loop (PLL) systems
  - In combination with RF modulation
    - chromaticity tracking

typical:  $\Delta Q_{\text{res}} \approx 10^{-3} \dots 10^{-5}$





- Control Theory → System Identification



- Example (first order) beam response  $\approx$  damped harmonic oscillator resonance

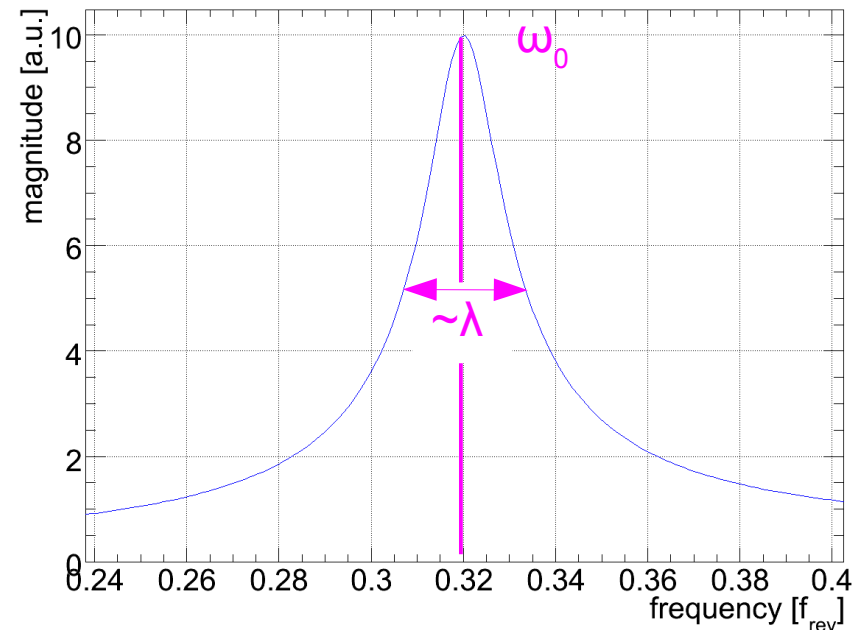
( $\omega_0$ : resonant frequency (Q),  $\lambda$ : tune resonance width ( $\sigma_Q$ ),

$\omega$ : driving frequency)

$$|G(\omega)| := \left| \frac{X(s)}{E(s)} \right| \approx \frac{\omega_0}{\sqrt{(\omega^2 - \omega_0^2)^2 + (2\lambda\omega_0\omega)^2}}$$

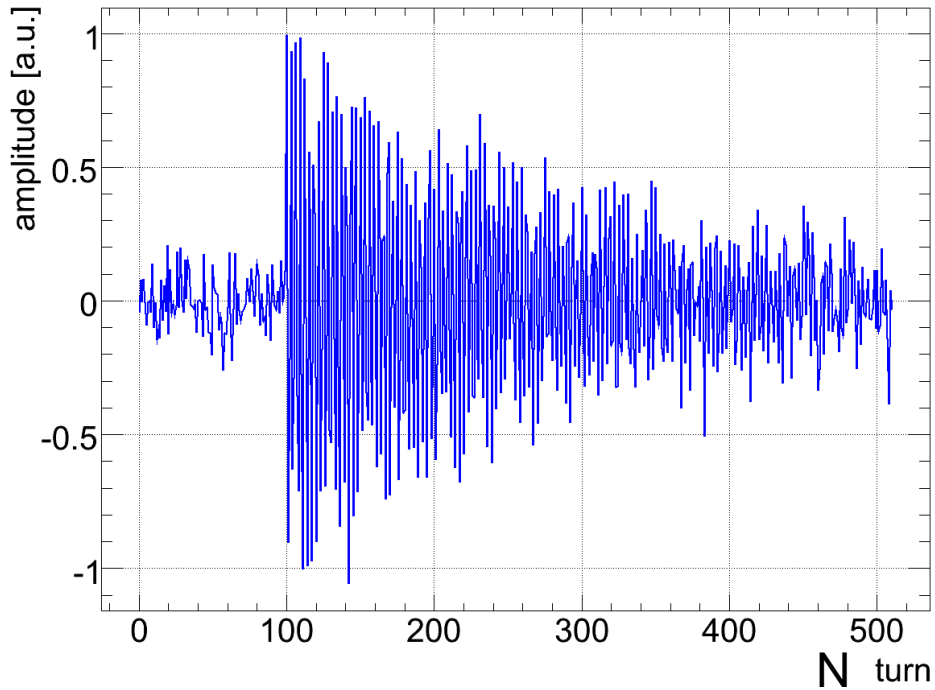
- Excitation choices:

- White or remnant noise
  - no information on signal phase
- Single-turn transverse kick (classic)
- Frequency Sweep aka. 'Chirp'
  - focuses excitation power on frequency range of interest → less  $\epsilon$ -blow-up, constant power
- Phase-Locked-Loop Systems = resonant excitation on the Tune



- Note: Exciter and pickup have additional non-beam related responses!

- .... how an kick-induced beam oscillation usually looks like (no sync. beating)

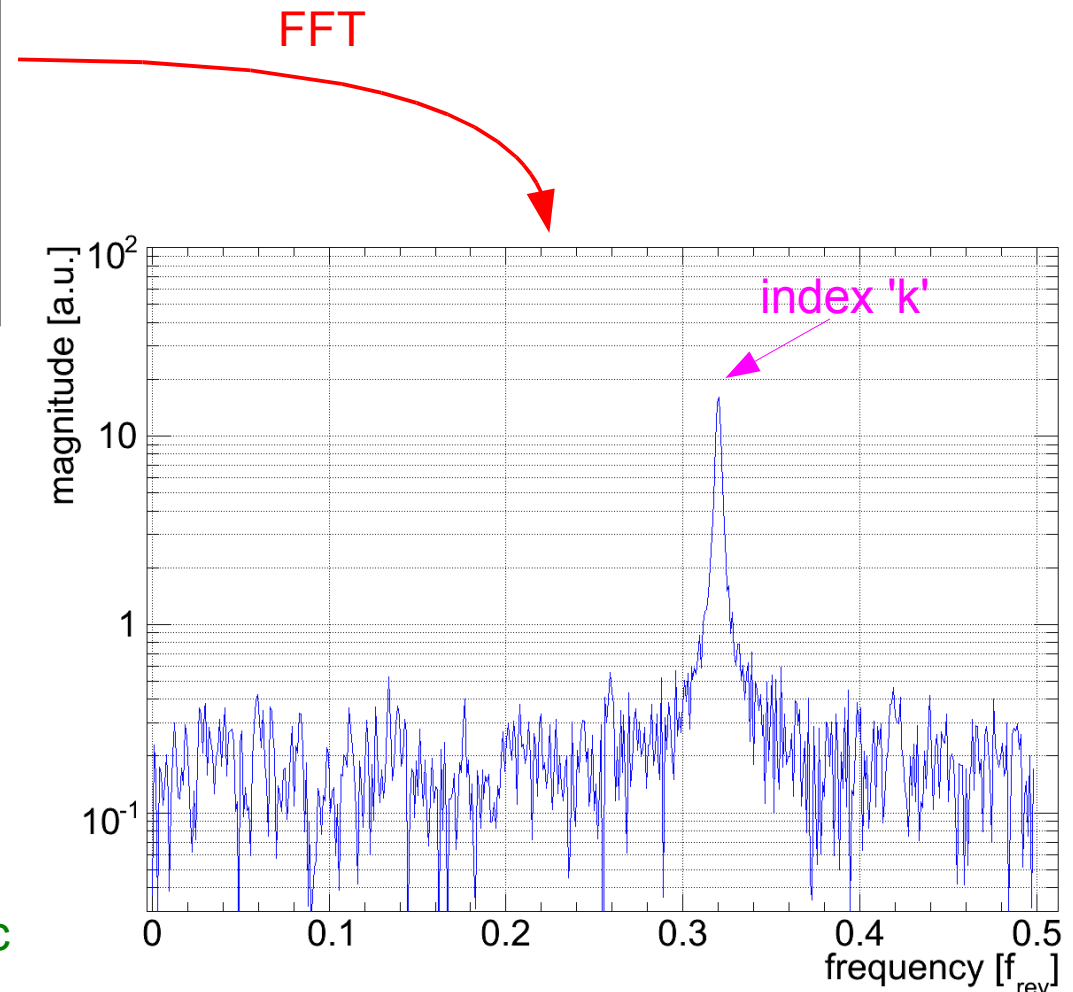


- Fourier analysis of turn-by-turn data:

- magnitude peaks at  $q_{frac}$

$$q_{frac} \approx \frac{k}{N}$$

- N.B. no information on  $Q_{int}$  !
- improve resolution by fitting central bin width → additional topic



Underlying measurement related to BPM design:

- Usual choices:
  - wall-current, button, shoebox, strip-line pickup (→ **P. Fork lecture**)
  - resonant pickups (e.g. Schottky → **F. Caspers**)

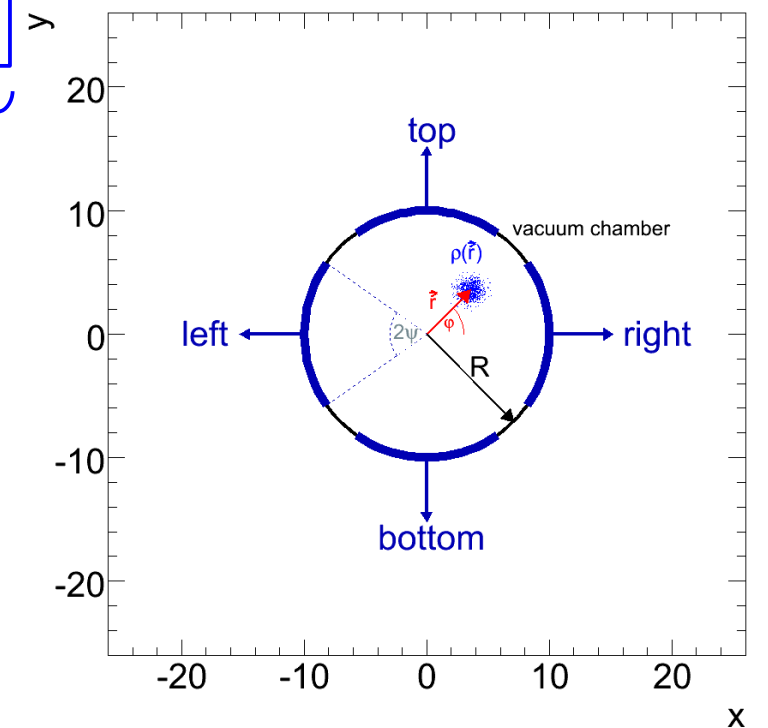
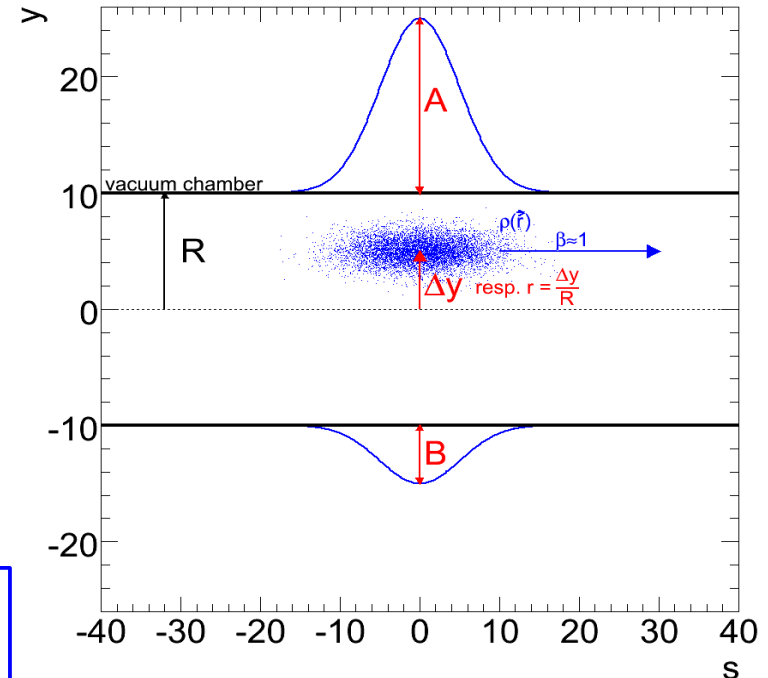
Single charge image density on pickup segment<sup>1</sup>:

$$I_{L/R}(t) = \frac{I_{\omega}(t)}{2\pi} \left[ \underbrace{2\psi}_{\text{longitudinal beam signal}} \mp \underbrace{2\frac{x}{R} \sin(\psi) + \frac{x^2 - y^2}{R^2} \sin(2\psi)}_{\text{transverse beam signal}} + h.o. \right]$$

longitudinal  
beam signal

transverse  
beam signal

- real-life signal is usually further convoluted with pickup and acquisition electronics response<sup>2,3</sup>!
- will elaborate a bit more on above equation

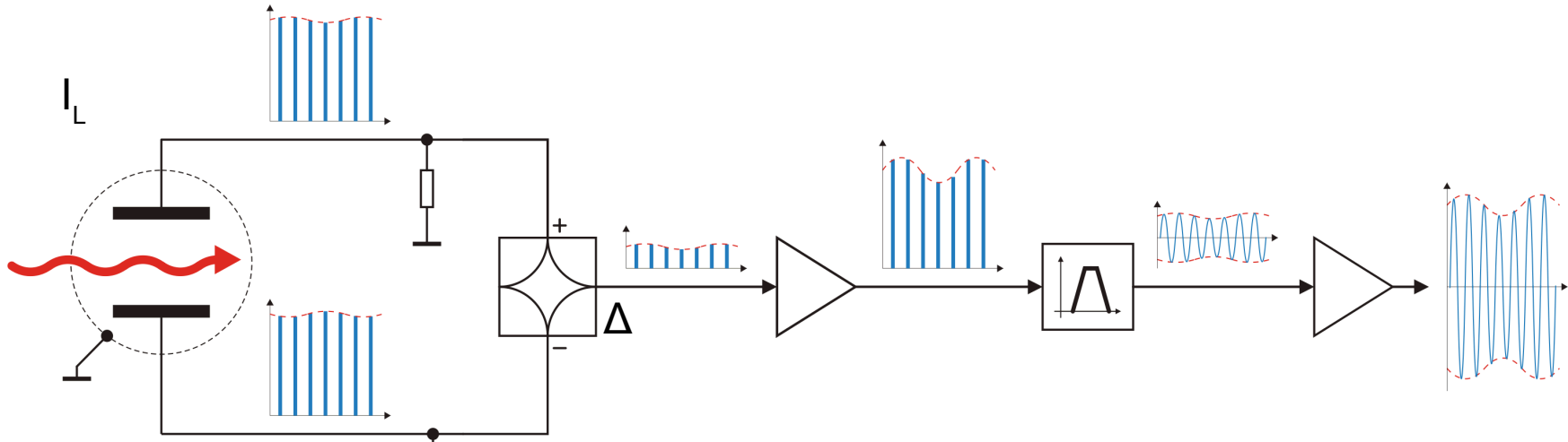


<sup>1</sup>R. Littauer, "Beam Instrumentation", SLAC Summer School, 1982. (p.902)

<sup>2</sup>D. McGinnis, "The Design of Beam Pickup and Kickers", BIW'94, 1994

<sup>3</sup>G. Vismara, "Signal Processing for Beam Position Monitors", CERN-SL-2000-056-BI





$$I_{L/R}(t) = \frac{I_{\omega}(t)}{2\pi} \left[ 2\psi \mp 2\frac{x}{R} \sin(\psi) + \frac{x^2 - y^2}{R^2} \sin(2\psi) + h.o. \right]$$

'intensity'      'position dependence'      'beam size dependence'

- Classic detection approach:  $\Sigma$ - $\Delta$  hybrid (or direct pickup signal sampling)

$$\rightarrow \frac{x}{R} \approx \frac{\Delta}{\Sigma} = \frac{I_L - I_R}{I_L + I_R} \quad R: \text{pickup half-aperture}$$

- Eliminates most 'common mode' signal (e.g. intensity),
- However ADC needs still to accommodate 'common mode' signals due to:
  - Closed orbit offset
  - 2<sup>nd</sup> order: intensity bleed-through intrinsic to any  $\Sigma$ - $\Delta$  hybrid

- A little bit in more detail:

$$I_{L/R}(t) = \frac{I_{\omega}(\sigma_s, t)}{2\pi} \cdot \left[ \underbrace{2\psi}_{\text{longitudinal beam signal (PM)}} + \underbrace{2\frac{x}{R}\sin(\psi) + \frac{x^2 - y^2}{R^2}\sin(2\psi)}_{\text{transverse beam signal (AM)}} + h.o. \right]$$

- N.B. multiplication in time-domain  $\leftrightarrow$  convolution in frequency domain

- Some important observations:

1. Transverse pickups are also sensitive to modulation of the longitudinal carrier signal

2. For tune measurement important beam-observable is  $x_{\beta}$ :

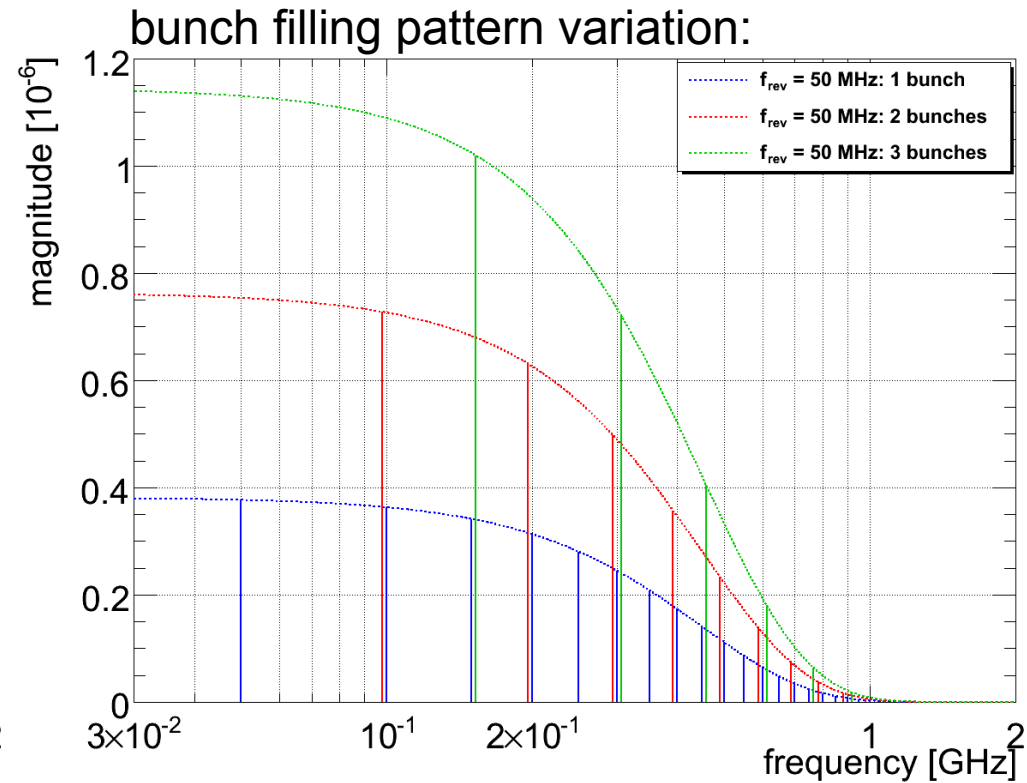
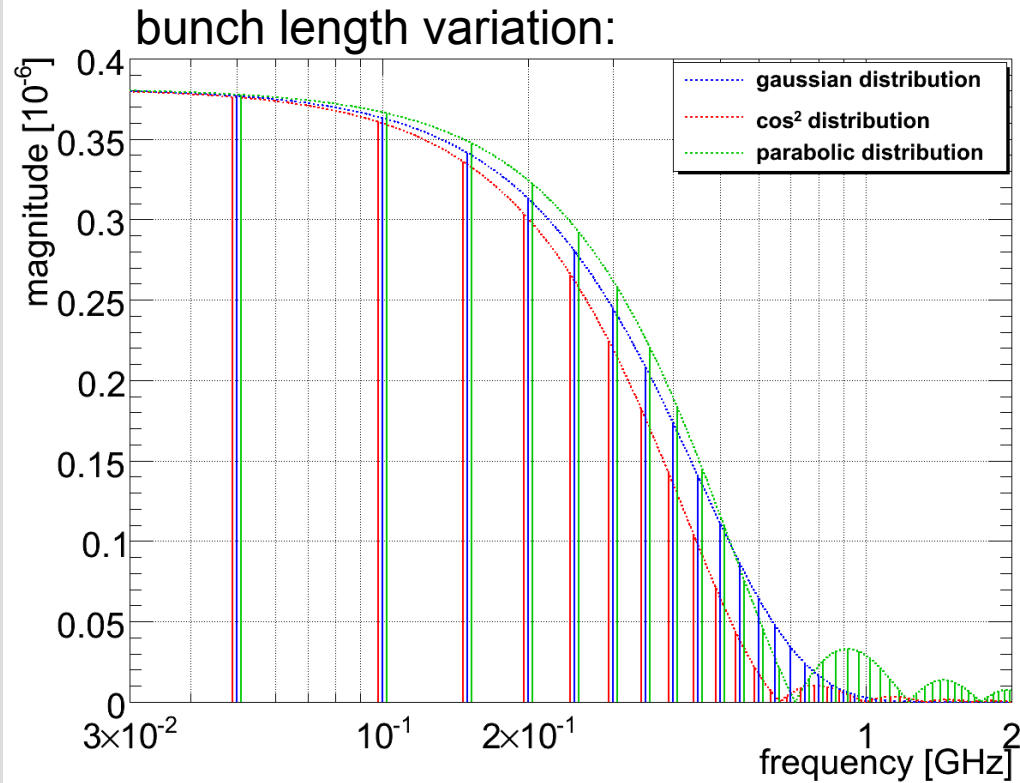
$$x \rightarrow x_{co} + D \cdot \frac{\Delta p}{p} + x_{\beta} \rightarrow I_{L/R}(t) \sim I_{CM} + \Delta I(x_{beta})$$

- 'Common-mode' signal  $I_{CM}$  limits dynamic range and ADC resolution

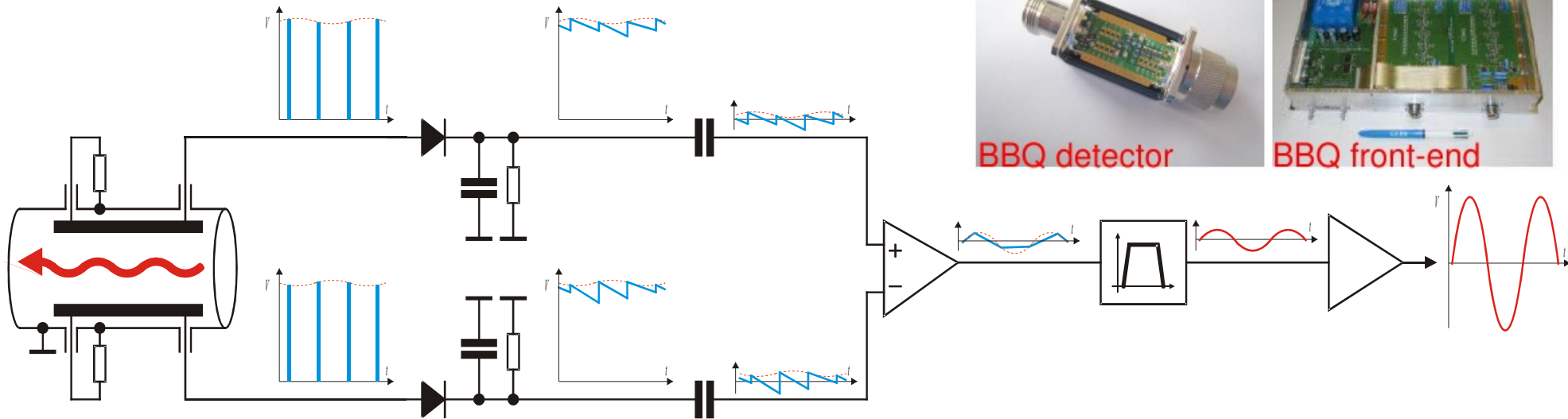
- Example:  $R \approx 44$  mm & nm resolution  $\rightarrow$  required sensitivity  $\Delta I/I_{CM} \sim 10^{-8}$

- most BPM systems:  $\Delta I/I_{CM} \sim 10^{-3} \rightarrow$  need something different for 'nm' resolution
- with e.g. good  $\Sigma$ - $\Delta$  hybrid:  $\Delta I/I_{CM} \sim 10^{-5}$

3. Higher Order term ' $x^2 - y^2$ ':  $I_{L/R}(t)$  sensitive to beam size  $\rightarrow$  a.k.a. 'quadrupolar pickup'



- Longitudinal carrier signal changes with shape, arrival time (synchrotron oscillations) and number of circulating bunches:
  - processing chain has to accommodate this through e.g. multiple gain stages
  - optimise for one bandwidth → in-/less sensitive if number of bunches change



BBQ detector



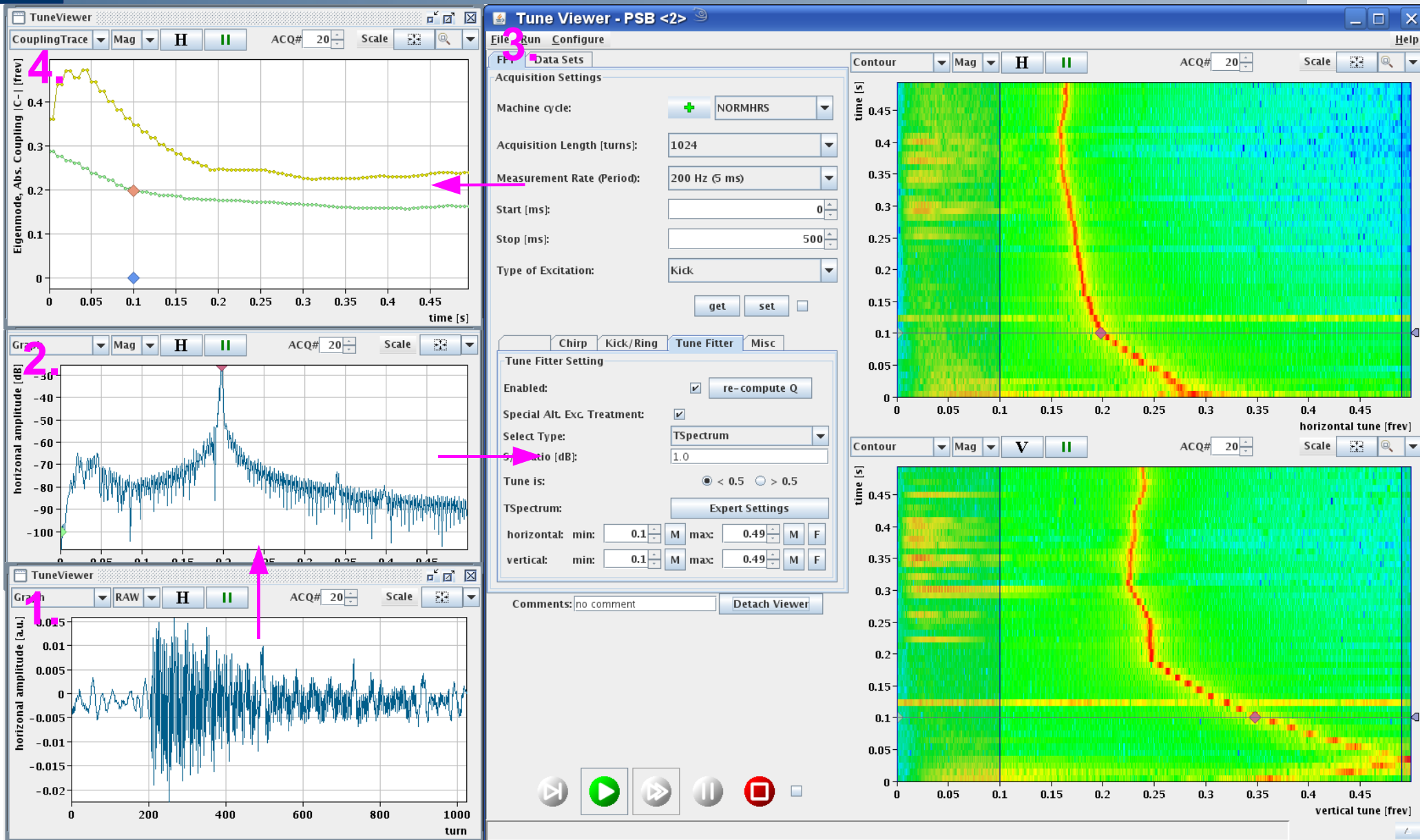
BBQ front-end

- Basic principle: AC-coupled peak detector<sup>1</sup>

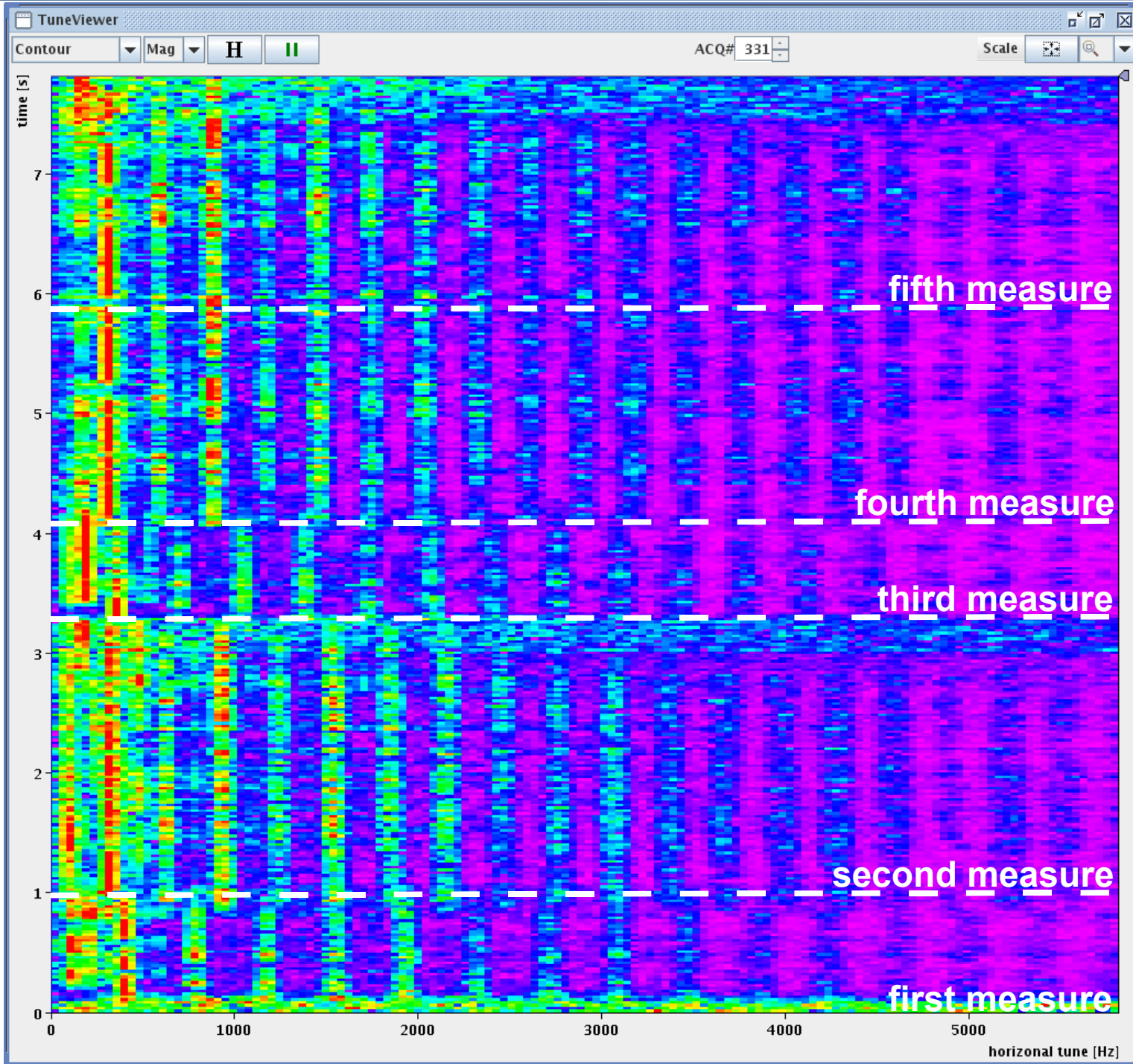
- intrinsically down samples spectra: ... GHz  $\rightarrow$  kHz (independent on filling pattern)
  - thus 'Base-Band-Tune Meter' (aka. BBQ)
  - Base-band operation: very high sensitivity/resolution ADC available
  - Measured resolution estimate:  $< 10$  nm  $\rightarrow$   $\epsilon$  blow-up is a non-issue
- AC-coupling removes common-mode  $\rightarrow$  only relative changes play a role
  - capacitance keeps the “memory” of the to be rejected signal
- no saturation, self-triggered, no gain changes to accommodate single vs. multiple bunches or low vs. high intensity beam

- However: no specific bunch-by-bunch information (unless using gating)

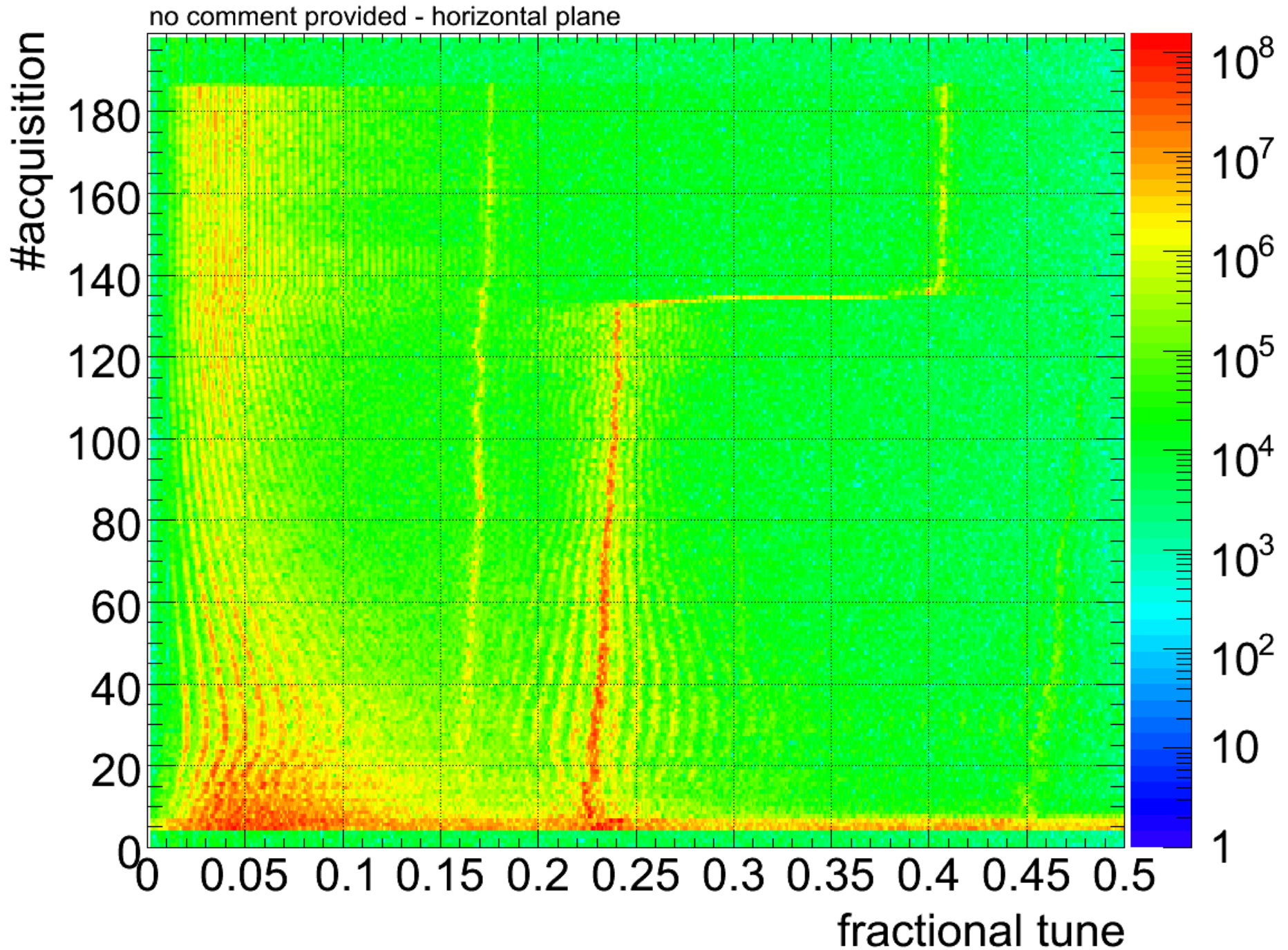
<sup>1</sup>M. Gasior, “The principle and first results of betatron tune measurement by direct diode detection”, CERN-LHC-Project-Report-853, 2005



- BBQ → fast ADC → FPGA based digital signal processing chain, FFTs @ 500 – 1 kHz!
  - provides real-time Q diagnostics for operation

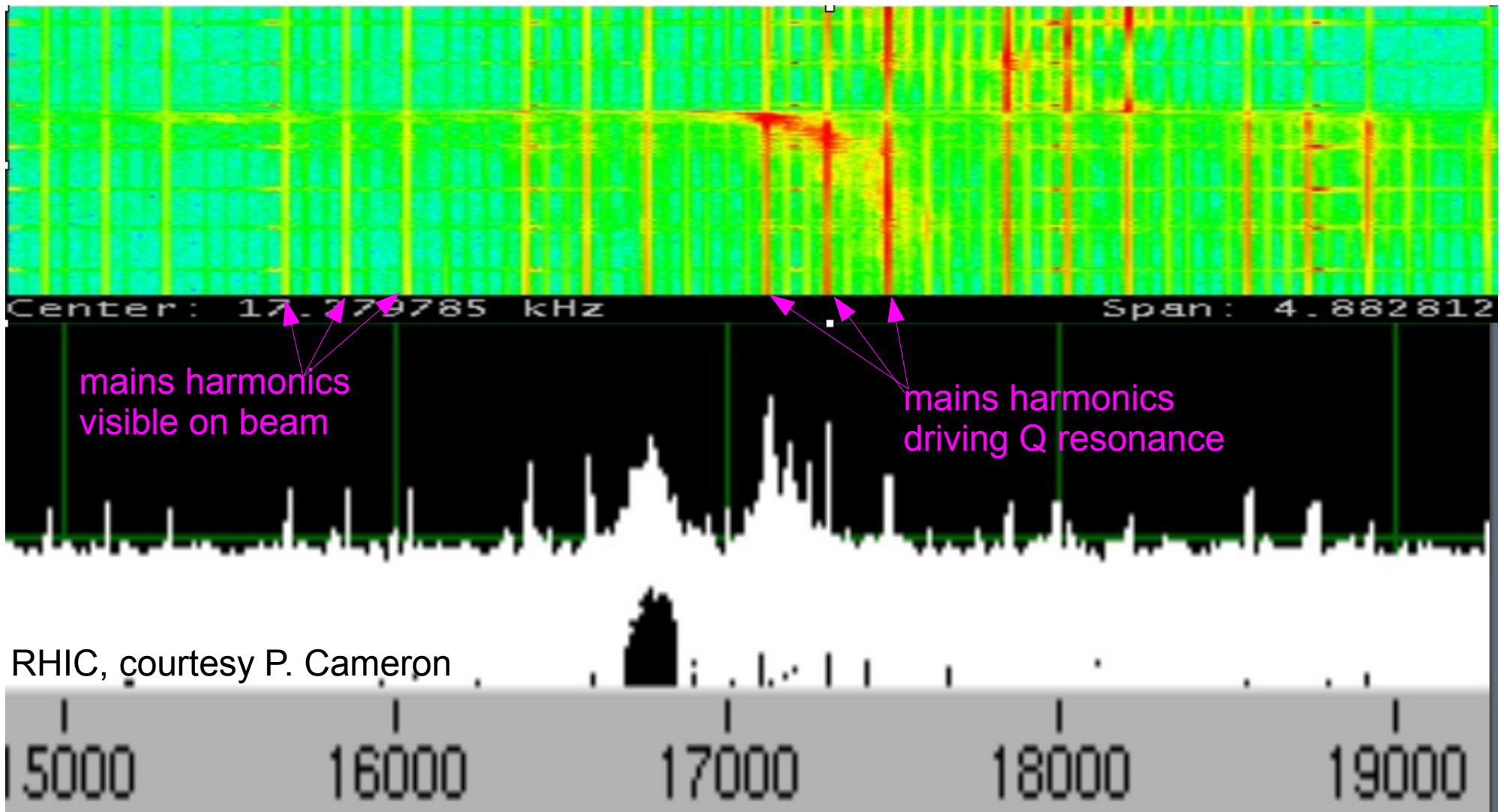


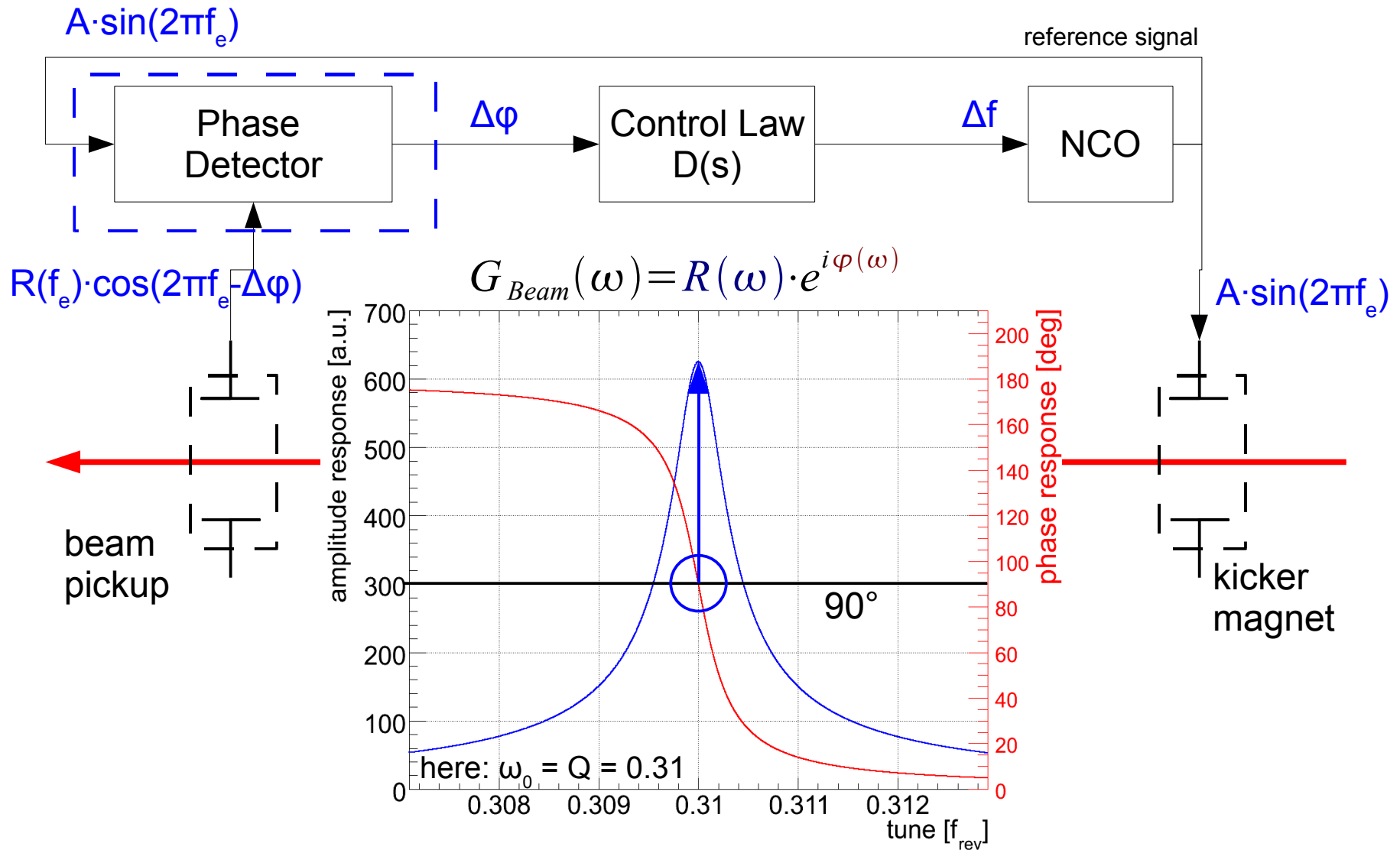




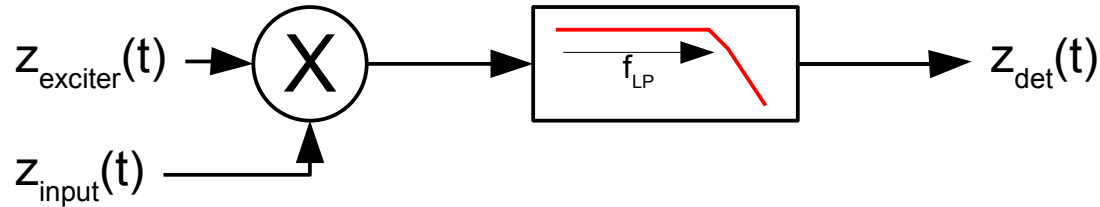


- BBQ system's high sensitivity revealed mains harmonics at RHIC and Tevatron
  - drives beam at tune resonance  $\rightarrow$  emittance blow-up, particle loss



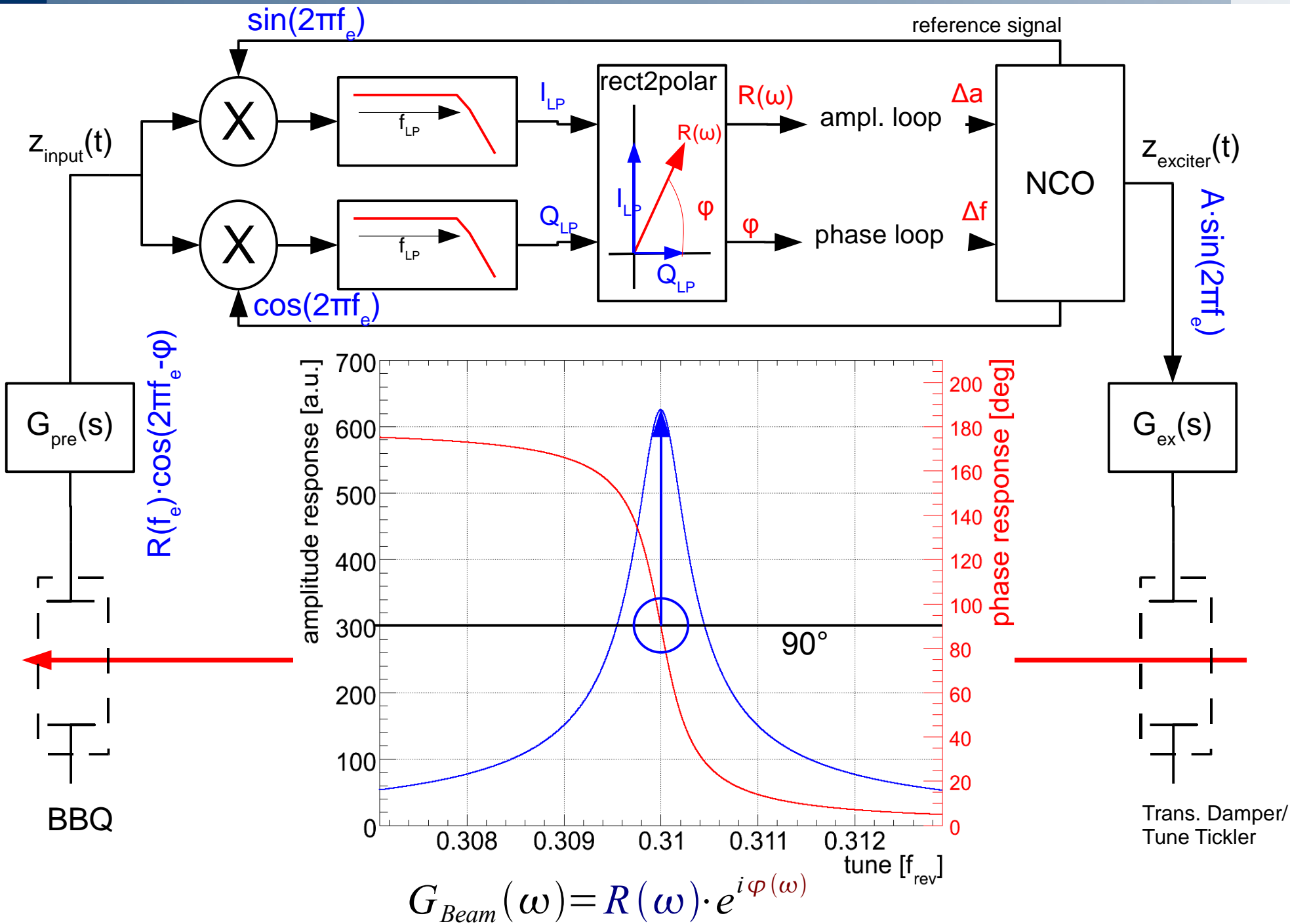


- BTF provides also information on collective effects (landau  $\rightarrow$  spread distribution):
  - impedance, stability diagram, lattice non-linearities ( $Q'$ ,  $Q''$ ), etc.

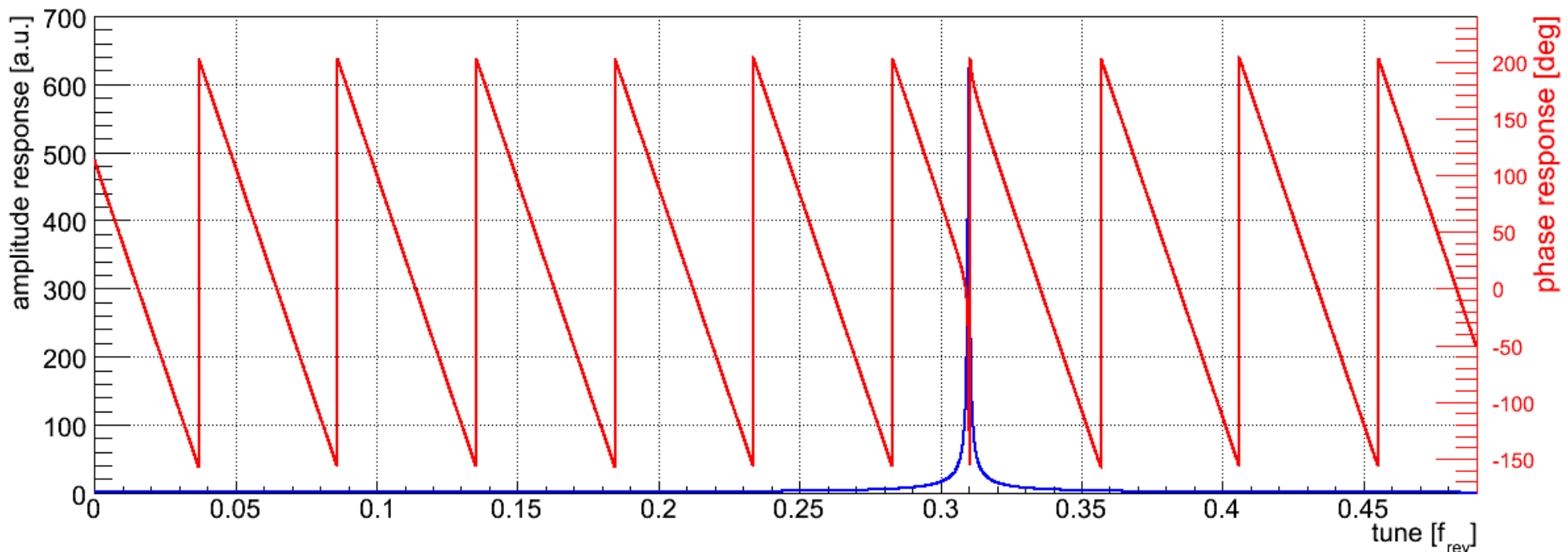


$$\begin{aligned}
 z_{det}(t) &= LP(z_{input}(t) \cdot z_{exciter}(t)) \\
 &= LP(R(f_e) \cdot \cos(2\pi f_e - \Delta\varphi(t)) \cdot A \sin(2\pi f_e)) \\
 &= \underbrace{\frac{AR}{2} \sin(\Delta\varphi(t))}_{\text{for small phases}} + \cancel{\frac{AR}{2} \sin(4\pi f_e - \Delta\varphi(t))} \\
 &\approx \Delta\varphi(t) \quad \text{removed by low-pass filter}
 \end{aligned}$$

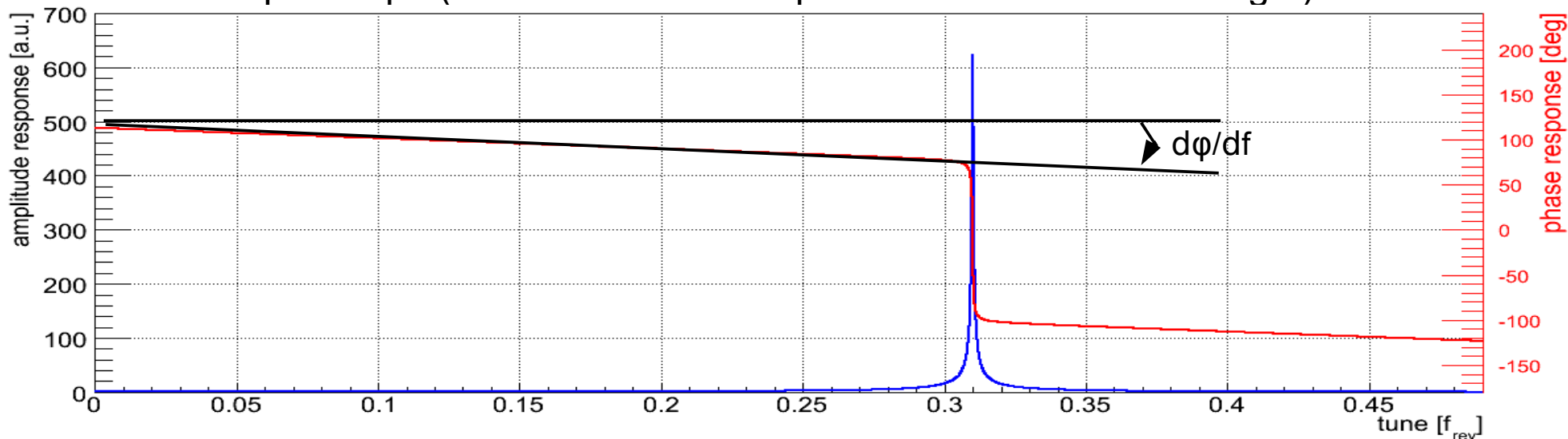
- Pro: robust analogue circuit implementation possible
- Con:
  - non-linear control signal for large phase difference  $\Delta\varphi$
  - Control signal depends on beam response's amplitude  $R(f_e)$



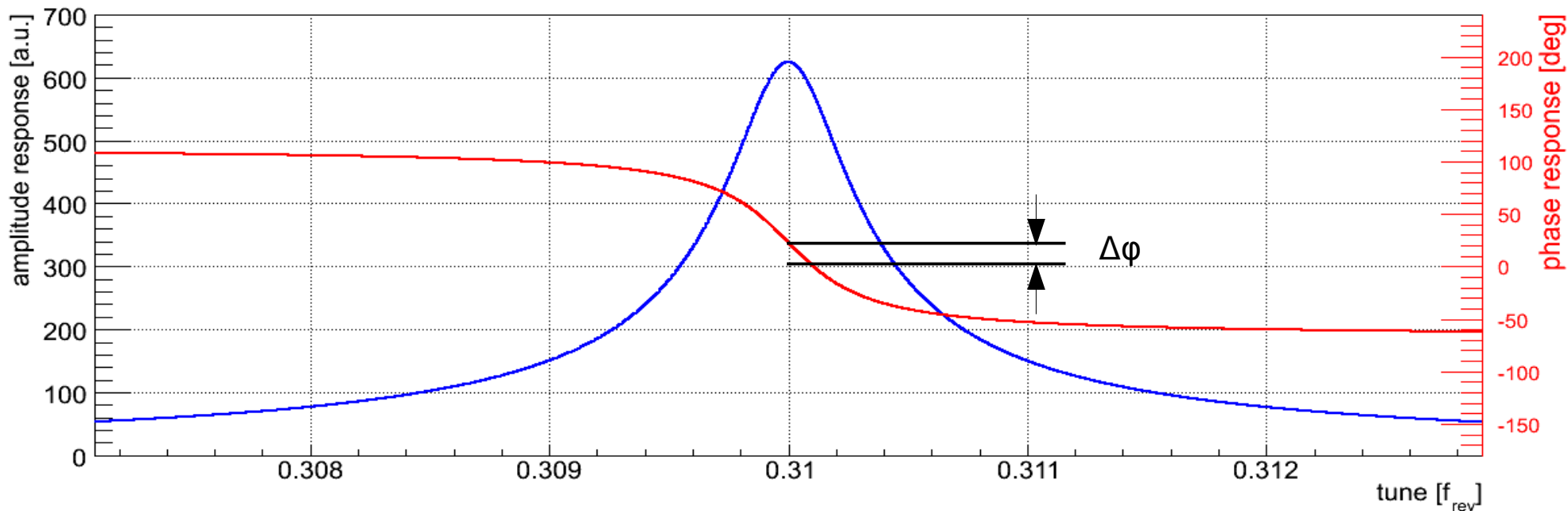
- BTF functions do not always look always as pretty as reports suggests or claim
  - an insider view on the real story:
- BTF and compensation consists of the adjustment of four parameters, preferably with stable beam condition ('chicken-egg' problem)
  - 1<sup>st</sup> step: verify necessary excitation amplitude and plane mapping (obvious?)
  - 2<sup>nd</sup> step: verify long sample delay (once per installation, constant)
    - full range BTF and count  $\pm\pi$  wrap-around  $\rightarrow$  number of delayed samples



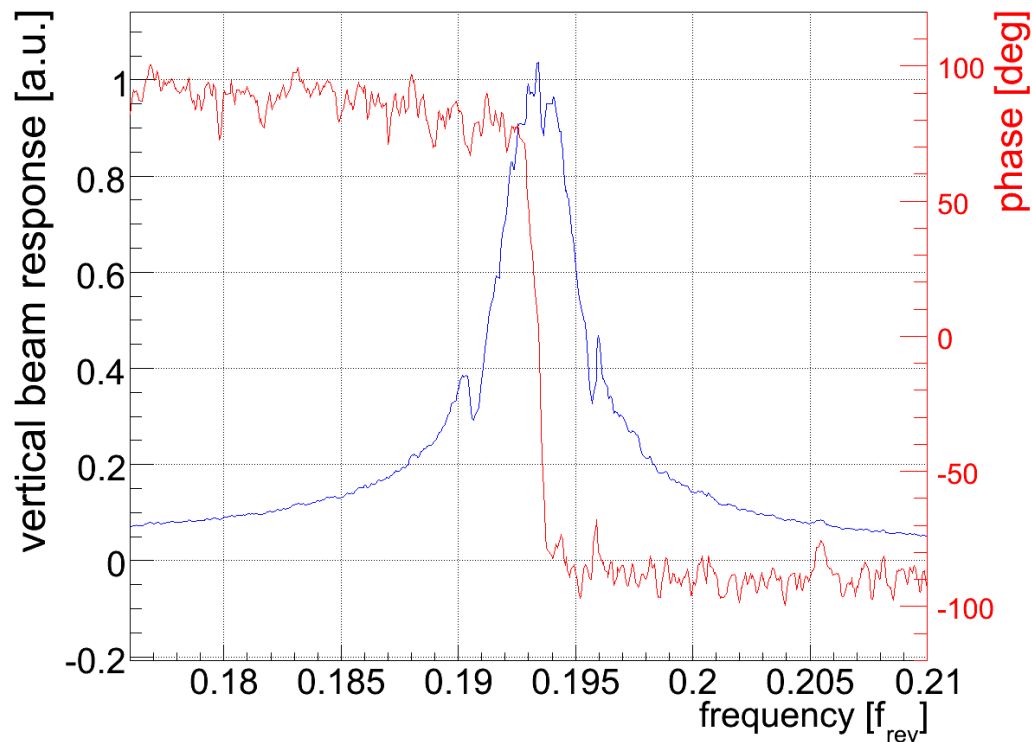
■ Measure  $d\phi/df$  slope ( ~ front-end non-lin. phase and kicker cable length)



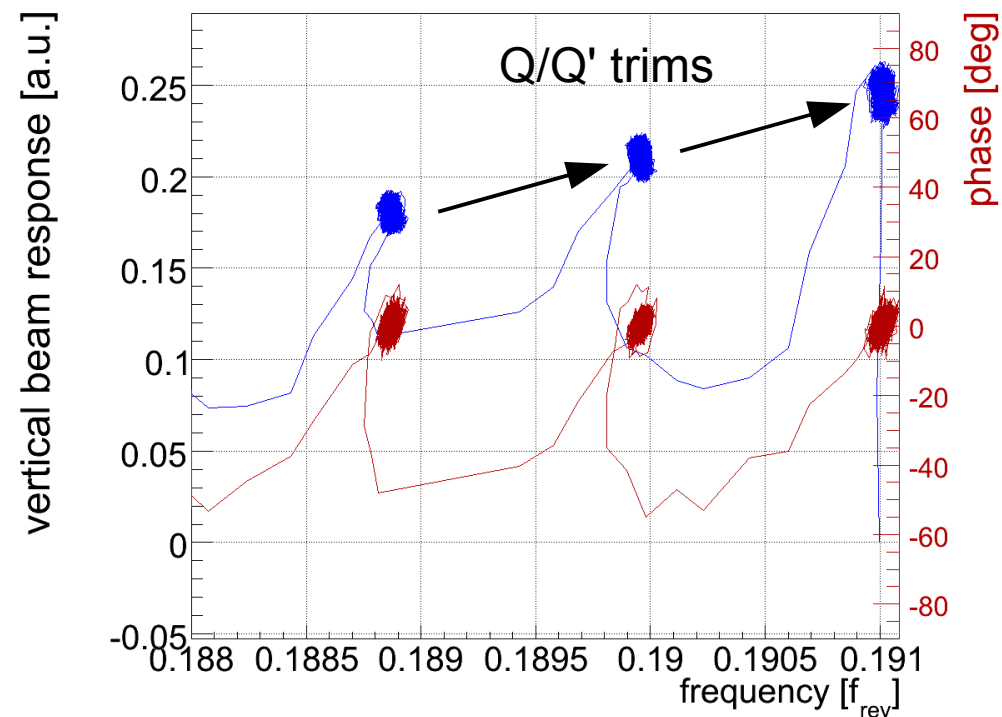
■ Adjustments of the locking phase (tune-peak – phase matching)



- What's published in papers and CAS reports:

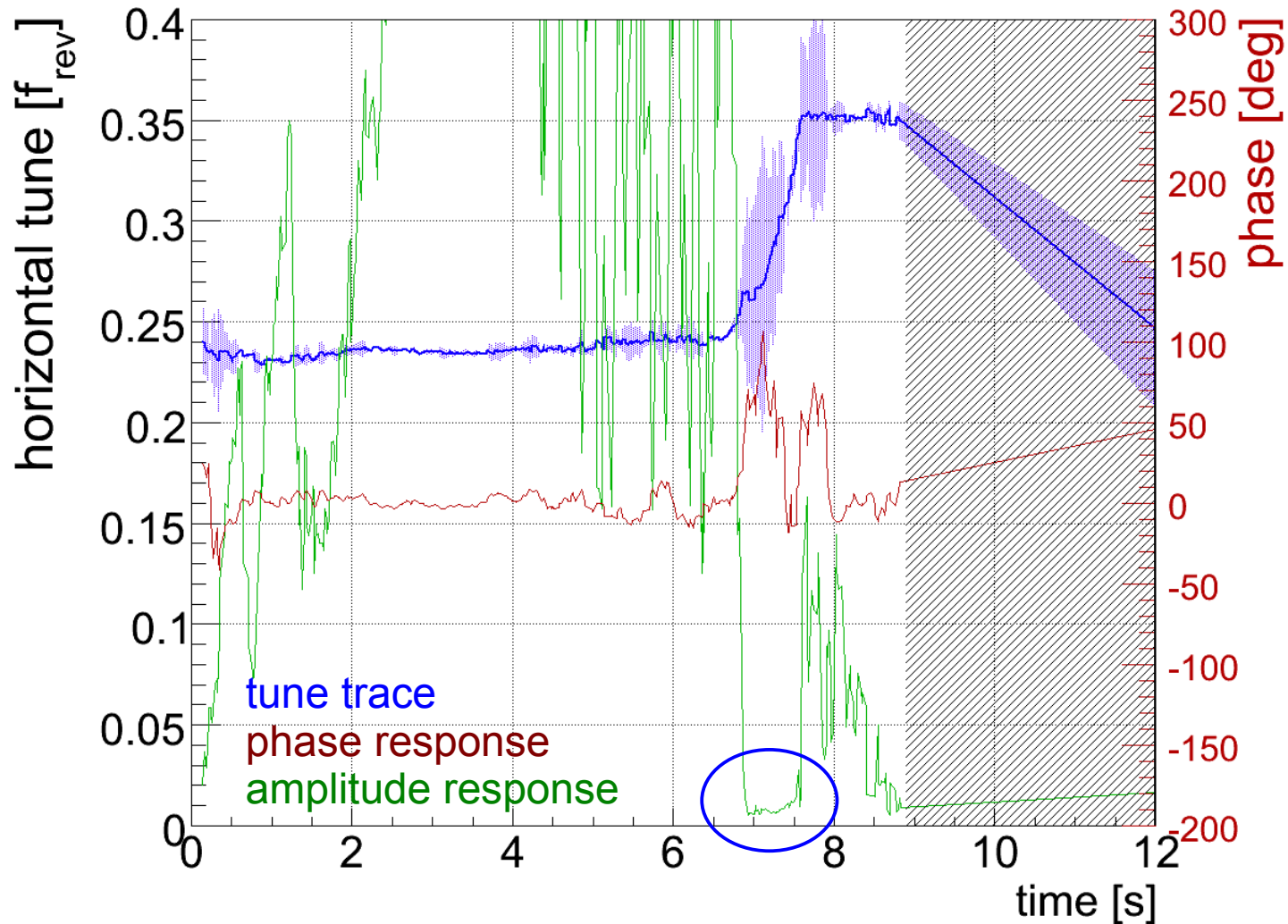


switch on PLL

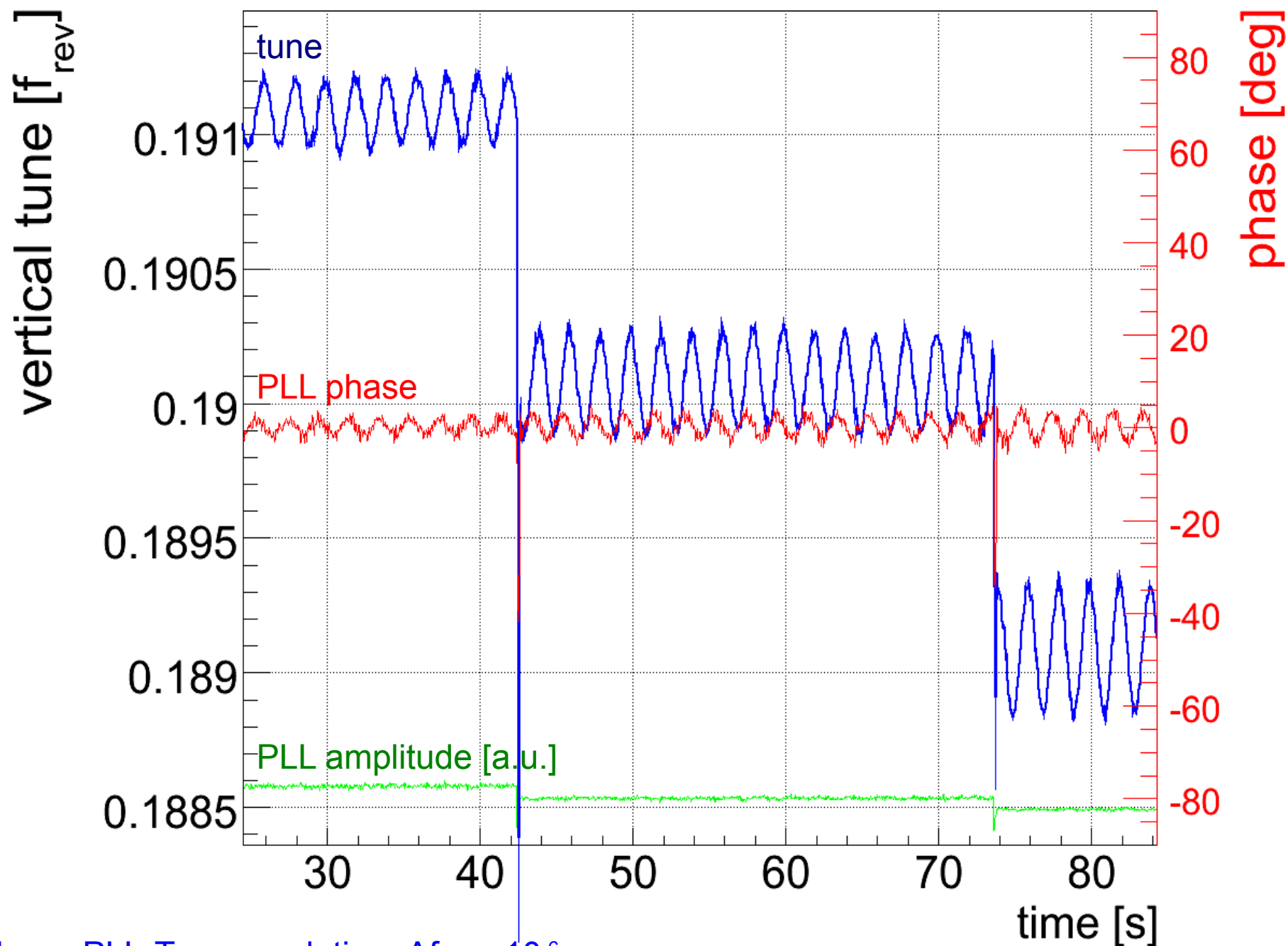




- Two domains of tracking, either slow and very precise (low loop bandwidth) or fast:



- Phase error and **non-vanishing amplitude** indicates lock
- here:  $\Delta Q/\Delta t|_{\text{max}} \approx 0.3$  within 300 ms  $f_{\text{rev}} \approx 43$  kHz

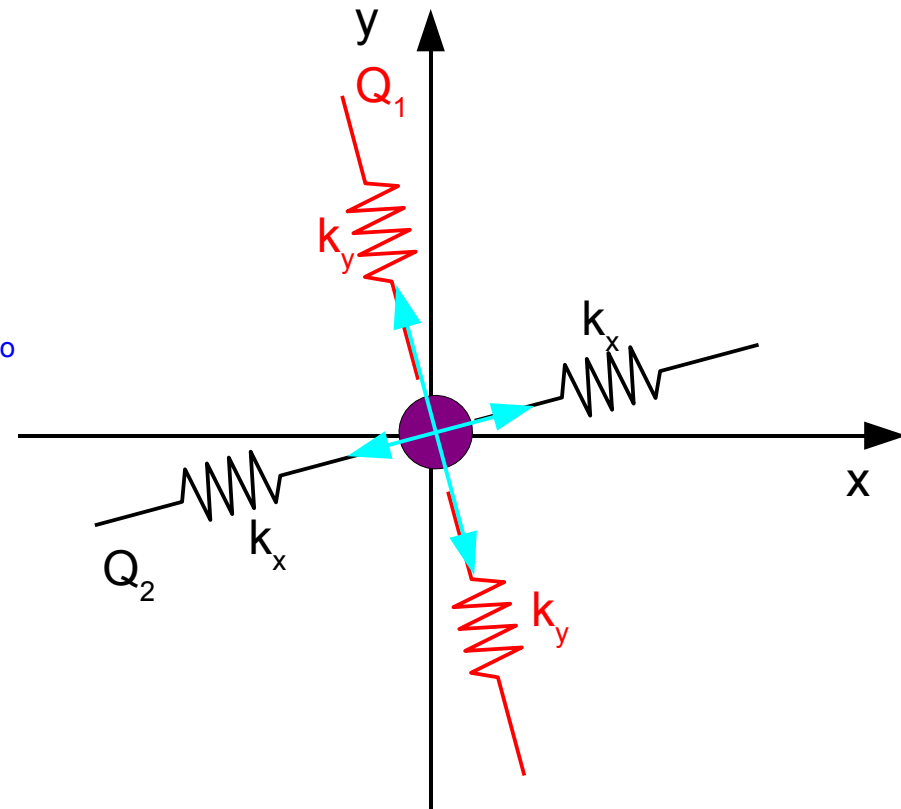


here: PLL-Tune resolution:  $\Delta f_{\text{res}} \approx 10^{-6}$

→ more during the second part

- Feed-down due to systematic closed orbit offset  $\Delta x_{co}$ :

- horizontal plane:
  - add. quadrupole → **tune shift**  $\sim \Delta x_{co}$
  - + small dipole kick  $\sim (\Delta x_{co})^2$
- vertical plane:
  - add. skew-quadrupole → **coupling**  $\sim \Delta y_{co}$
  - + small dipole kick  $\sim (\Delta y_{co})^2$ 
    - first order: rotates oscillation plane



- Feed-down due to closed orbit + change of sextupolar field:

- important for superconducting accelerators:** large changes of persistent currents (decay & snapback phenomena)
  - also visible while changing (trimming)  $Q'$
  - Higher order effects: space charge, beam-beam, ...

- In the presence of coupling (solenoids, skew-quadrupoles):

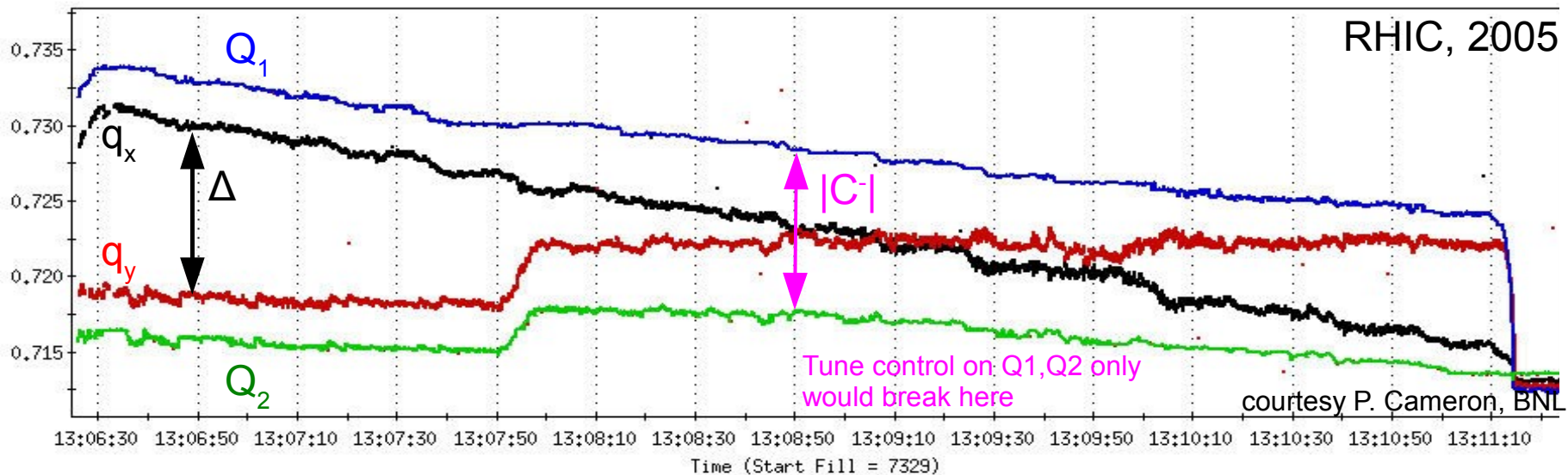
$$\begin{cases} \boxed{x}'' + k(s) \cdot \boxed{x} = \boxed{\kappa(s) \cdot y} \\ \boxed{y}'' + k(s) \cdot \boxed{y} = \boxed{\kappa(s) \cdot x} \end{cases}$$

classic harmonic oscillator, defines unperturbed tunes:  $q_x, q_y$ 
coupling terms

$$\kappa(s) = \frac{q}{2p} \left( \frac{\partial B}{\partial y} - \frac{\partial B}{\partial x} \right)$$

- assuming weak coupling, eigenmodes ( $Q_1, Q_2$ ) may be rotated w.r.t. unperturbed tunes ( $q_x, q_y, \Delta = |q_y - q_x|$ )

$$Q_{1,2} = \frac{1}{2} \left( q_x + q_y \pm \sqrt{\Delta^2 + |C^-|^2} \right)$$



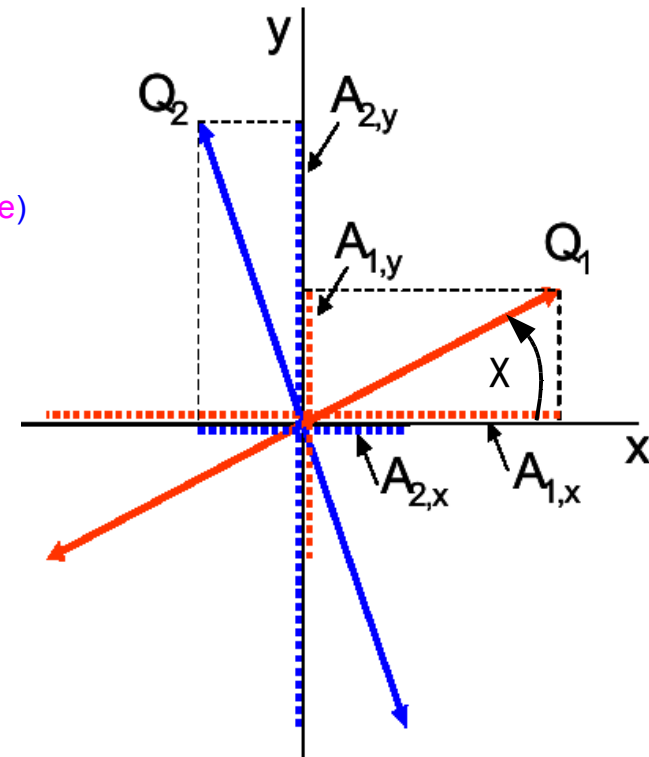
Possible improvement:

- Optimise tune working point (larger tune-split),
- Vertical orbit stabilisation in lattice sextupoles (Orbit FB → M. Böge)
- Active compensation and correction of coupling
  - ratio between regular and cross-term:
    - $A_{1,x}$ : eigenmode amplitude '1' in vert. plane
    - $A_{1,y}$ : eigenmode amplitude '1' in hor. plane

$$r_1 = \frac{A_{1,y}}{A_{1,x}} \quad \wedge \quad r_2 = \frac{A_{2,x}}{A_{2,y}}$$

$$\Rightarrow \boxed{|C^-| = |Q_1 - Q_2| \cdot \frac{2\sqrt{r_1 r_2}}{(1 + r_1 r_2)} \quad \wedge \quad \Delta = |Q_1 - Q_2| \cdot \frac{(1 - r_1 r_2)}{(1 + r_1 r_2)}}$$

- decouples beam feedback control
  - $q_x, q_y \rightarrow$  quadrupole circuits strength
  - $|C^-|, \chi \rightarrow$  skew-quadrupole circuits strength



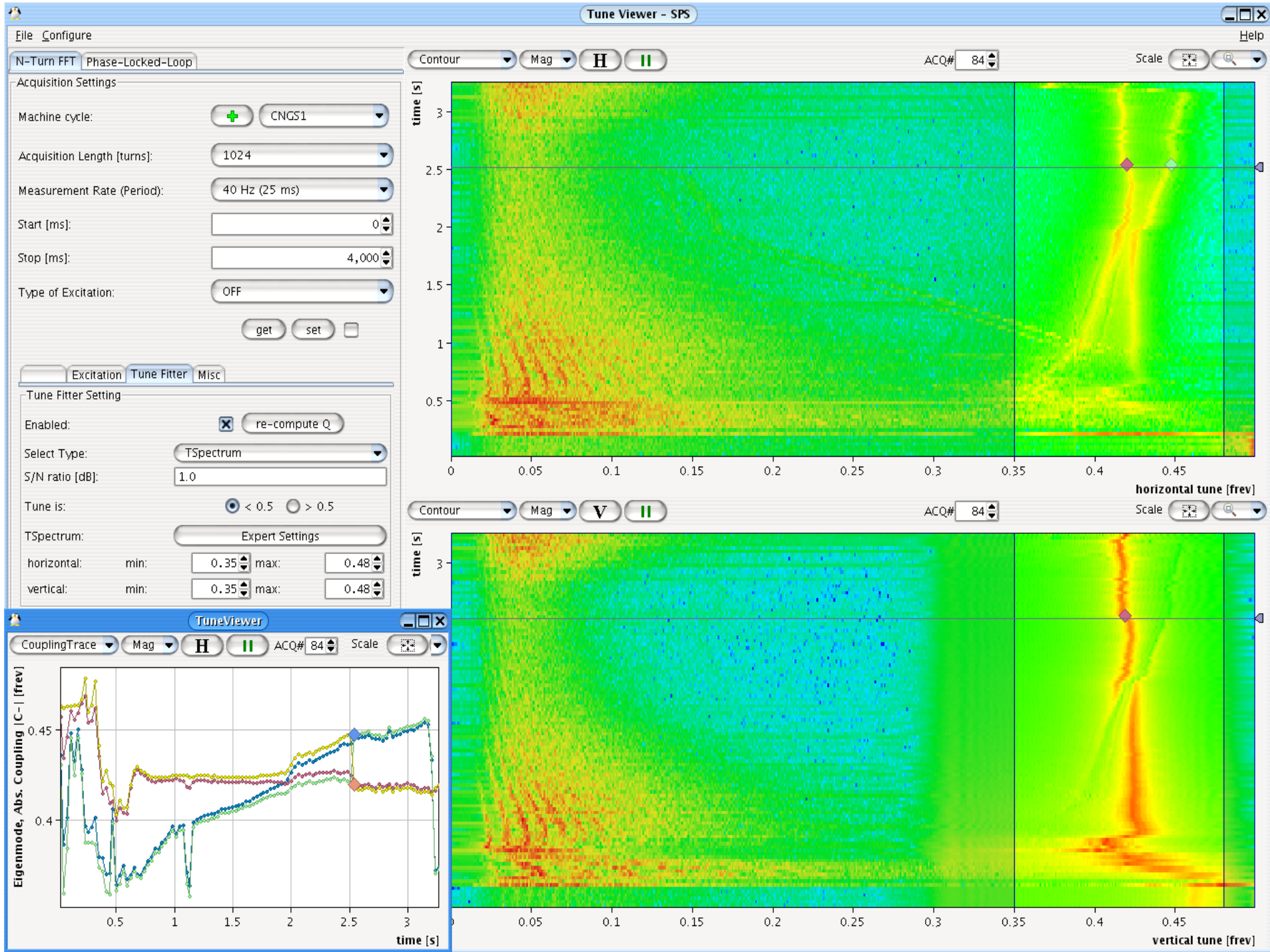




# Betatron Coupling Detection Example: CERN-SPS



Q & Q' Diagnostics, CAS Dourdan, France, Ralph.Steinhausen@CERN.ch, 2008-05-31



- That's all – questions?



- If interested: some additional advanced topics not covered so far (see Appendix):
  - Classic Tune Frequency Analysis
    - Improving Frequency Resolution of FFT based Spectra
  - Tune Phase-Locked-Loop Locking issues in the presence of:
    - Coupled Bunch Instabilities
    - Synchrotron Side-bands
    - Changing Tune Width ( $Q'$  dependence, amplitude detuning, impedance, ...)
  - Feedback on Tune, Chromaticity and Coupling



## Additional Slides

# Additional Topic I: Improving Frequency Resolution of Fast-Fourier-Transform based Spectra

- Tune frequency resolution can be improved through FFT based Interpolation algorithms  
( $k$ : index of highest bin,  $N$ : total number of turns,  $M_k$ : magnitude of bin  $k$ )

- Some common approaches:

- No interpolation:  $q \approx \frac{k}{N}$

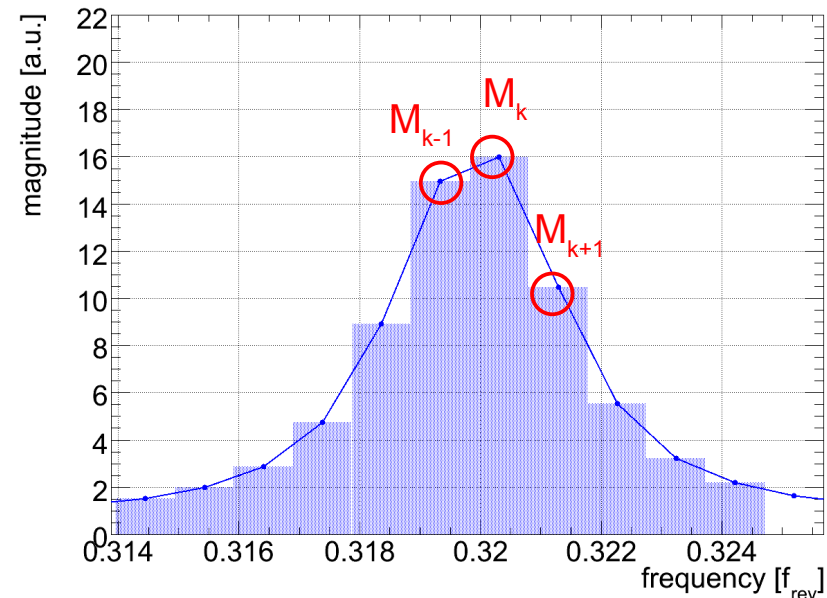
- Barycentre (n=1) & cubic (n=3) fit:  $q \approx \frac{M_{k-1}^n(k-1) + M_k^n(k) + M_{k+1}^n(k+1)}{N(M_{k-1}^n + M_k^n + M_{k+1}^n)}$

- Parabolic fit:  $q \approx \frac{k}{N} + 0.5 \cdot \frac{M_{k+1} - M_{k-1}}{2M_k - M_{k-1} - M_{k+1}}$

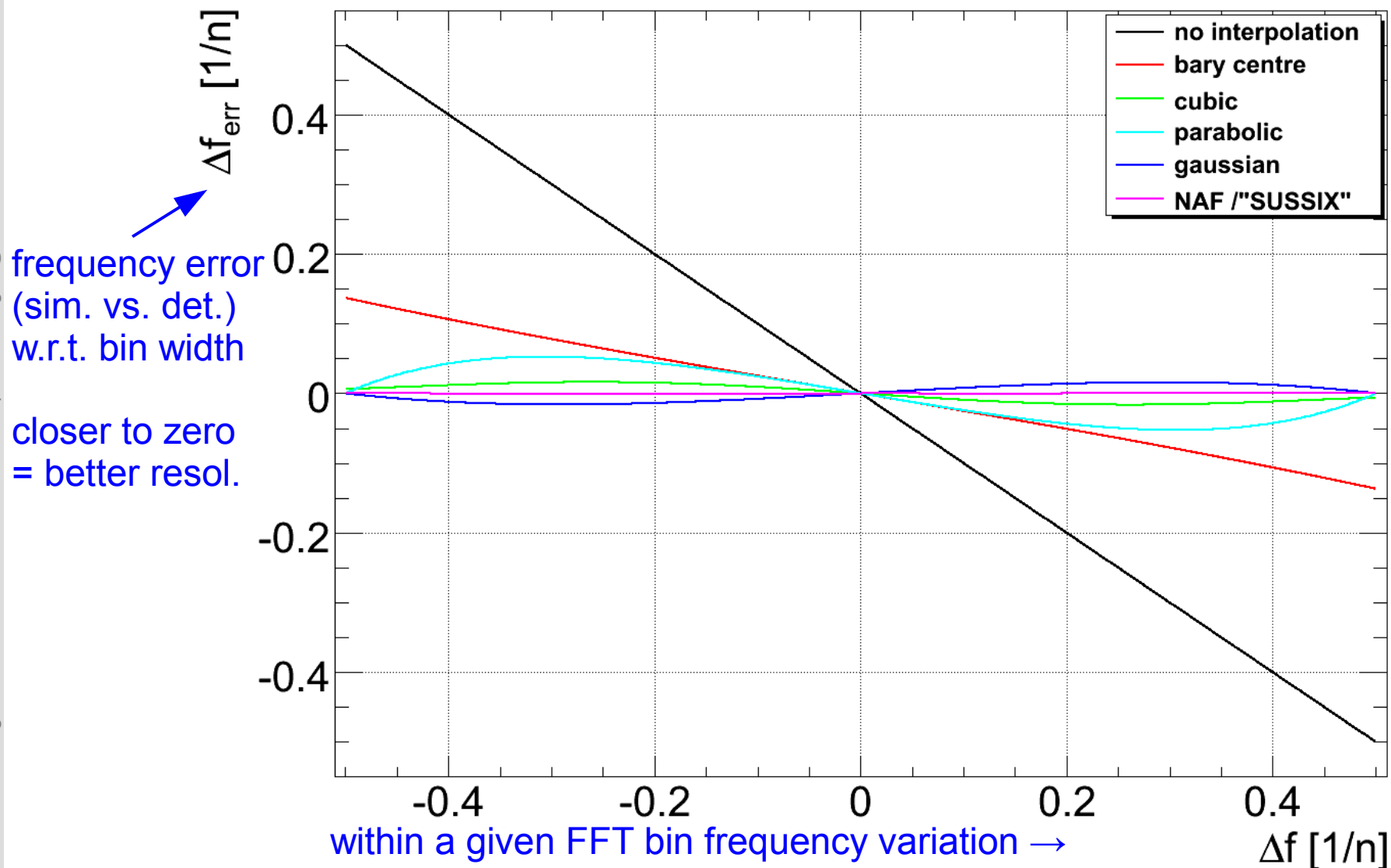
- Gaussian fit:  $q \approx \frac{k}{N} + 0.5 \cdot \frac{\log(M_{k+1}/M_{k-1})}{\log(M_k^2/(M_{k-1}M_{k+1}))}$

- NAFF/"SUSSIX":  $q \approx \frac{k}{N} \pm \frac{1}{\pi} \cdot \text{atan} \left( \frac{|M_{k\pm 1}| \sin(\frac{\pi}{N})}{|M_k| + |M_{k\pm 1}| \cos(\frac{\pi}{N})} \right)$

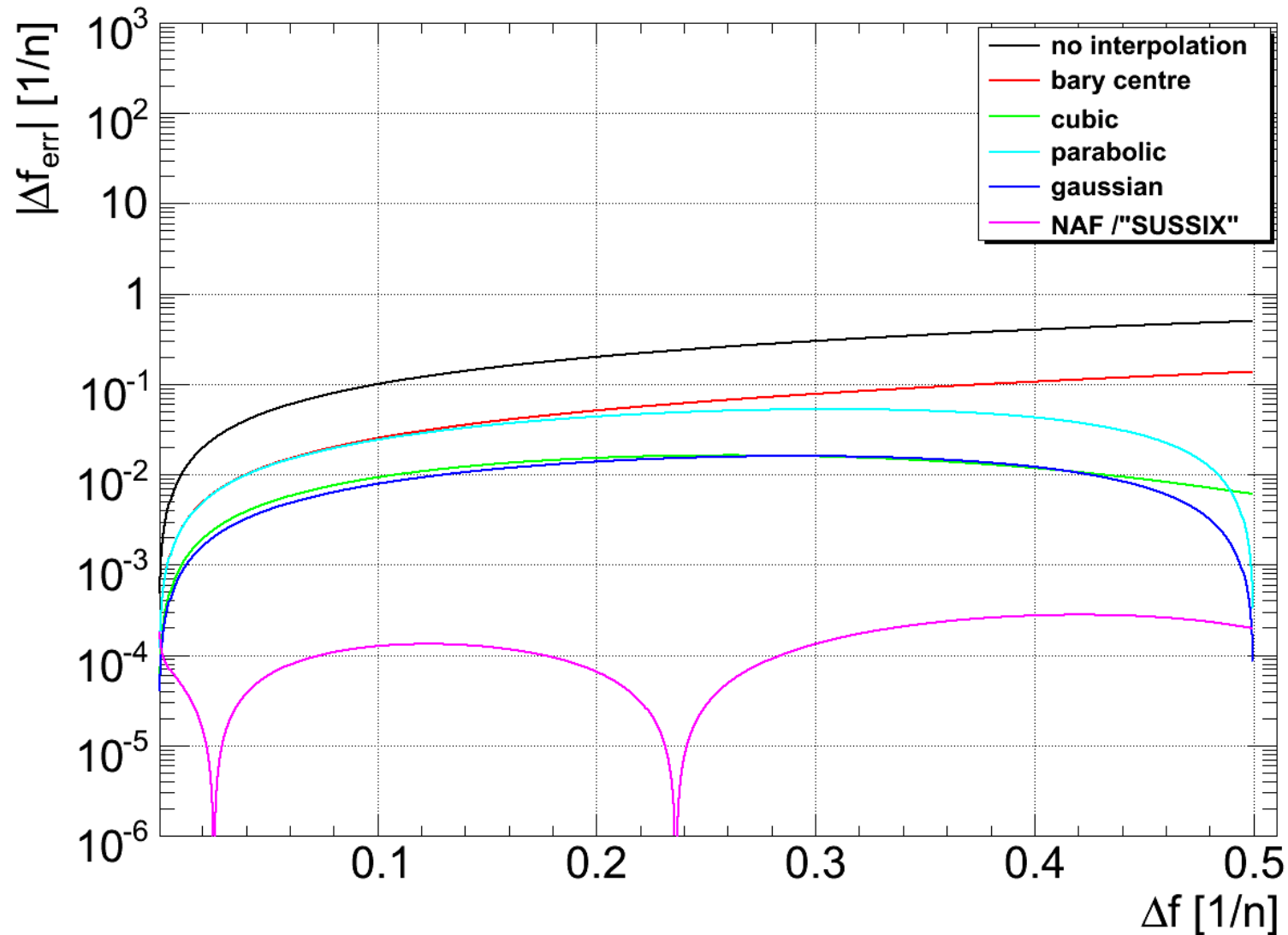
- Test case: controlled oscillation at a given frequency which is varied within one bin, normalised to sampling frequency



- 1024 turns: perfect sinusoidal oscillation & within one bin varying frequency
  - introducing some

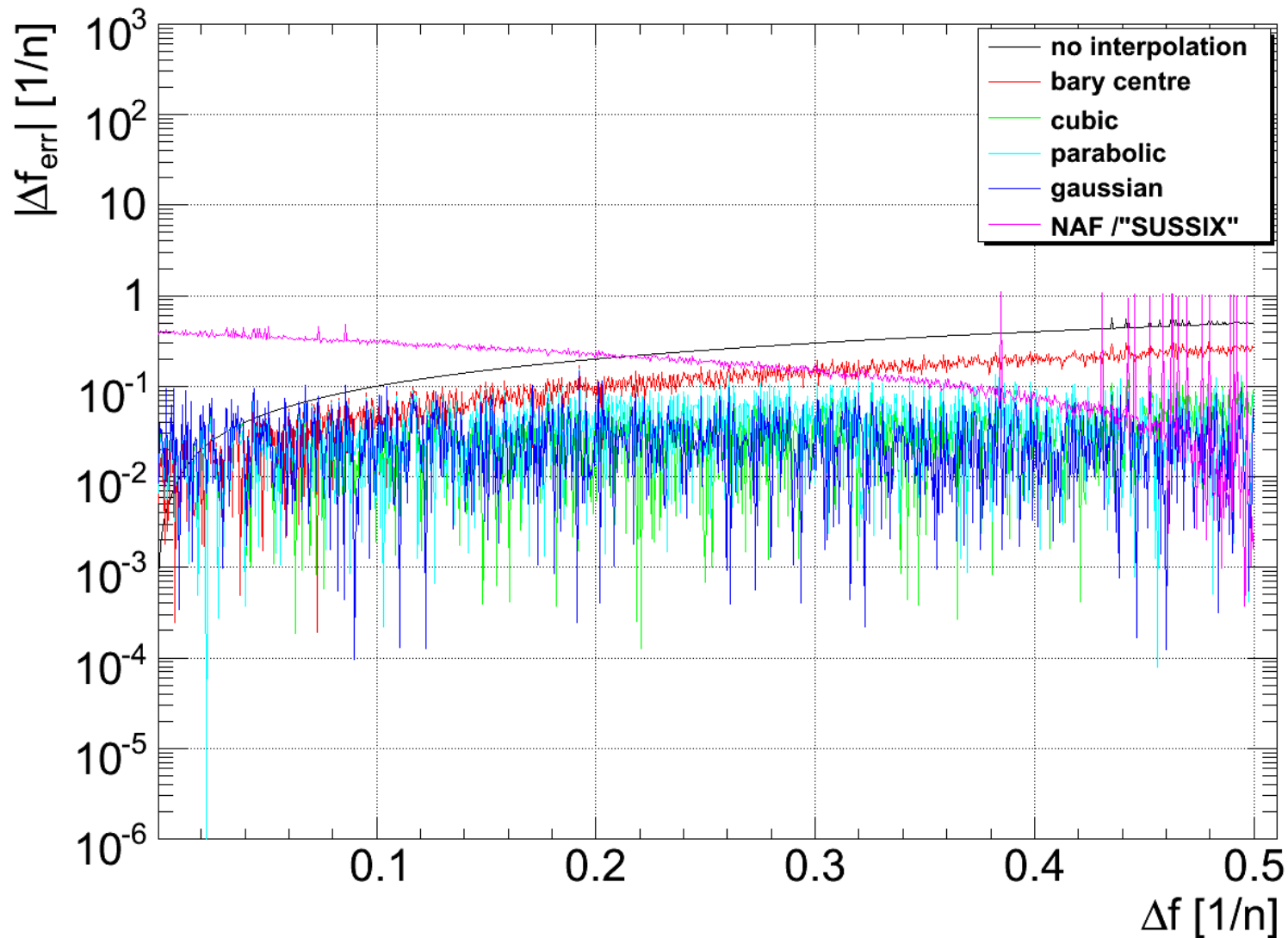


- same plot as before but: absolute error, logarithmic scale and considering frequency only within half a bin width (symmetry!)



- ... what about more realistic signals with damping, noise ...?

- same as before + 0.1 r.m.s. noise vs. kick amplitude of '1'



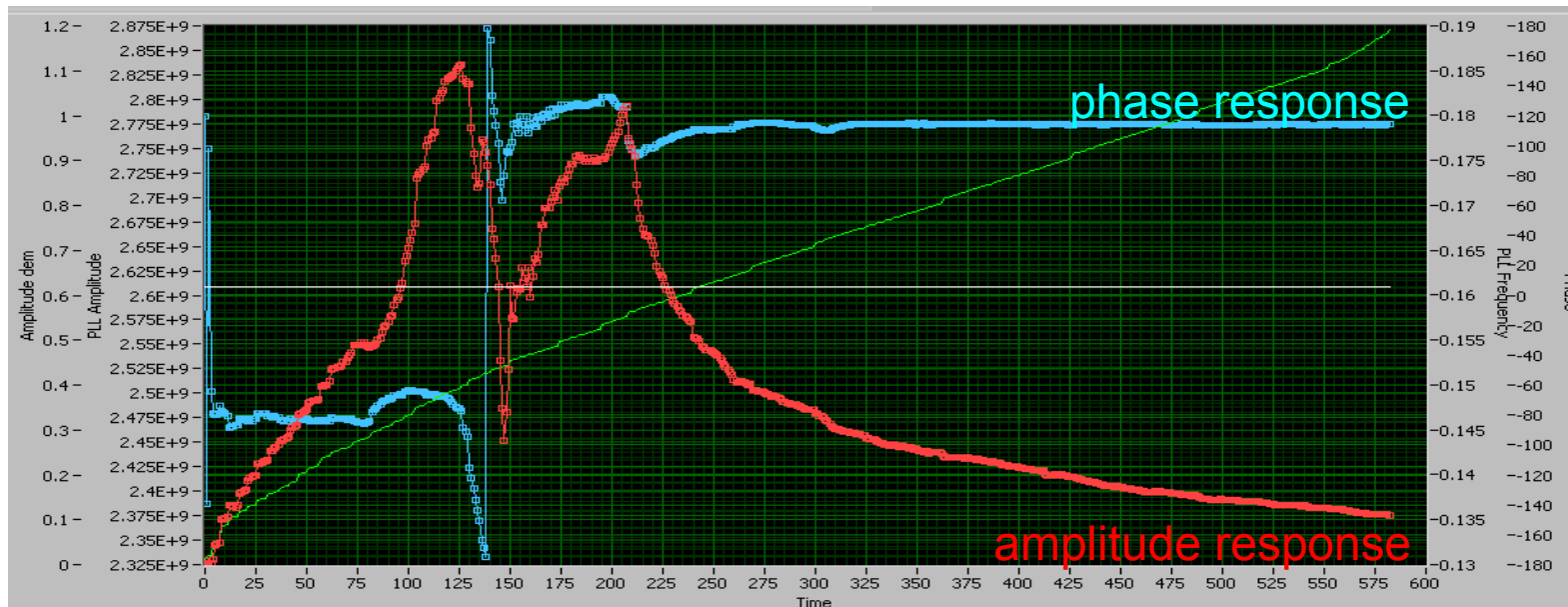
- Measurement noise is the limiting the resolution, cubic, barycentre, parabolic and Gaussian interpolation seem to yield similar performance. → **Gaussian-fit of central peak gives good results in most cases.**

# Additional Topic II:

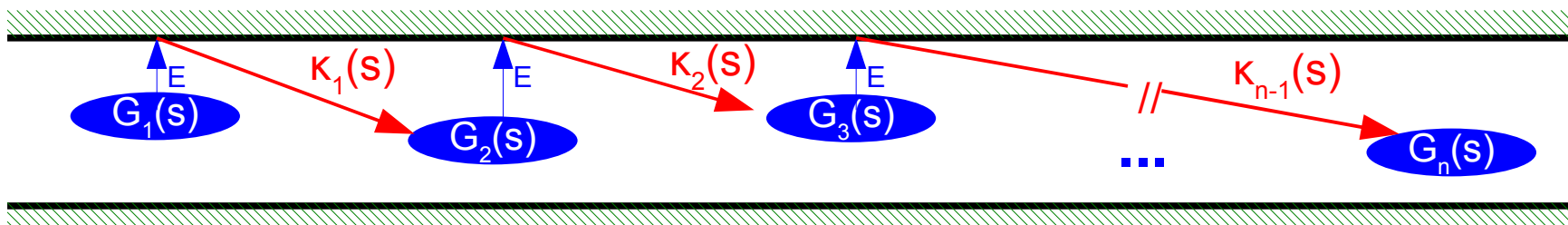
## Phase-Locked-Loop Locking in the Presence Coupled Bunch Instabilities, Synchrotron Side Bands and Tune Width Dependence



- Coupled bunch effects can hamper lock became more pronounced during later MDs
  - possible causes: impedance driven wake fields, e-cloud, beam-beam, ...

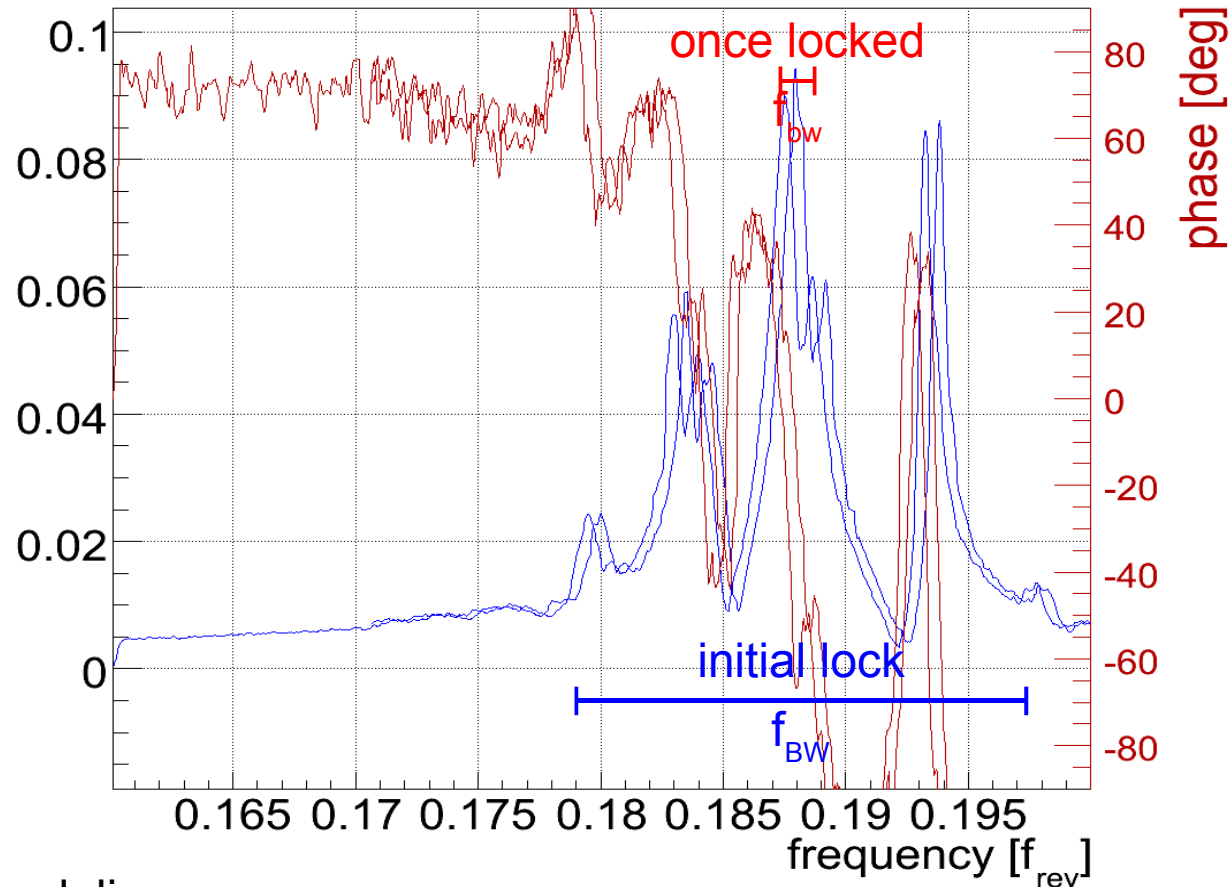


- Mechanism (impedance):



- Possible remedy:

- Detector selects and measures only one (/first) representative bunch



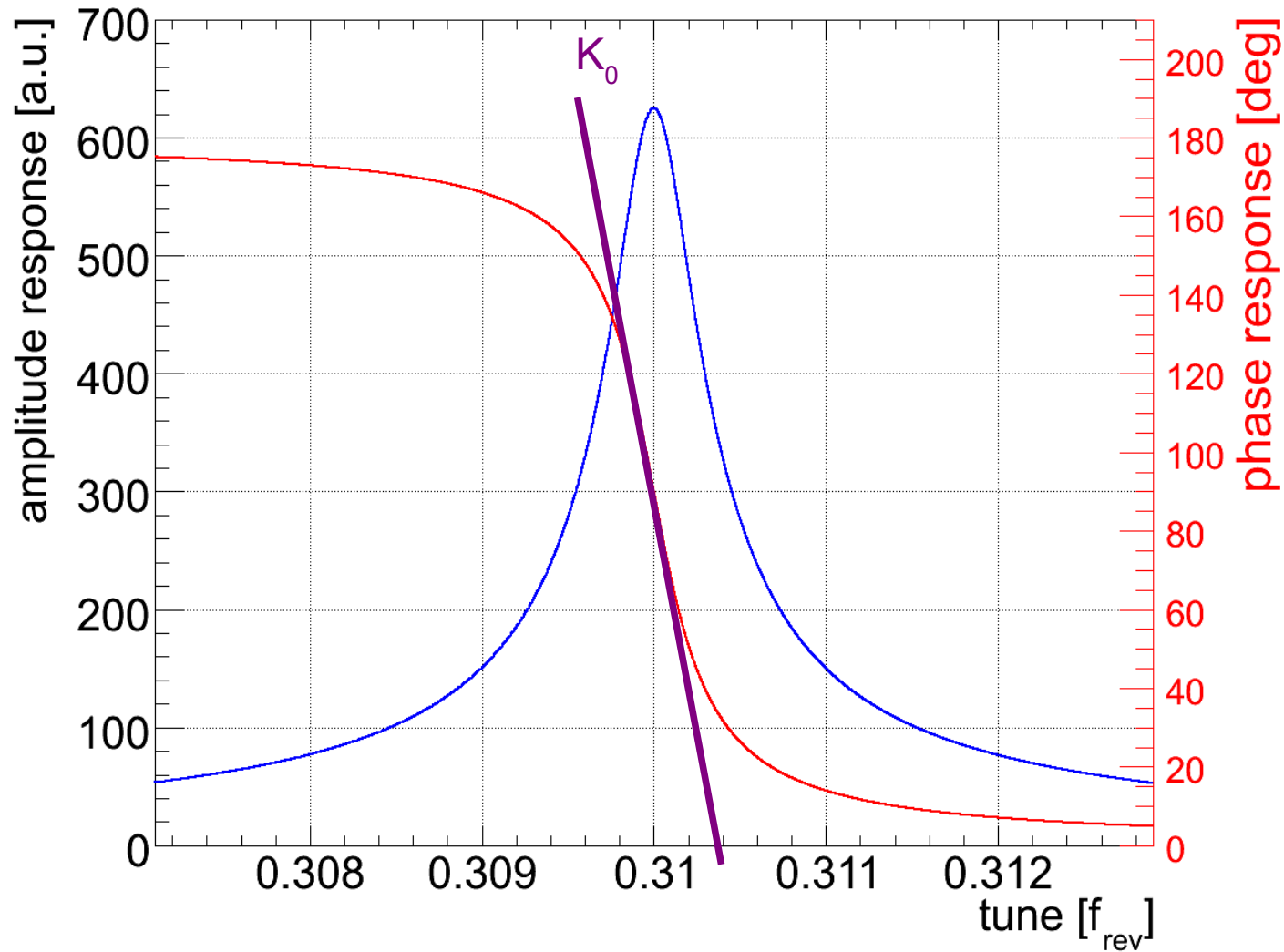
### Option I: gain scheduling

**initial lock**: open bandwidth to cover more than one side band (PLL noise ~ chirp)

- side-bands “cancel out”, strongest resonance prevails

**once locked**: reduce bandwidth for better stability/resolution

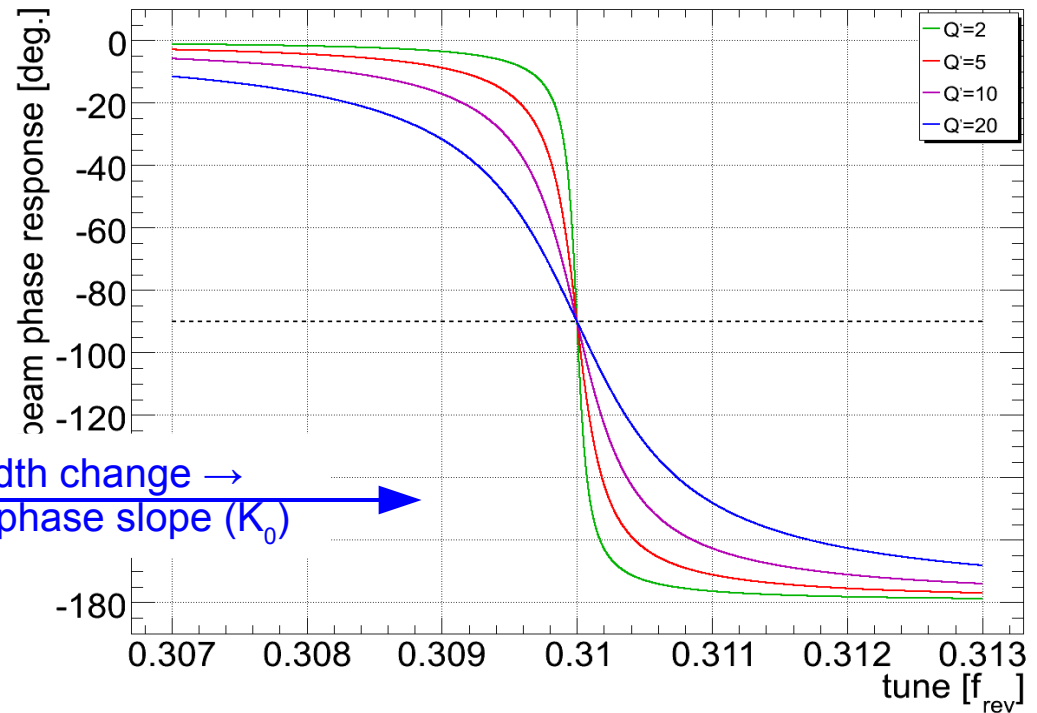
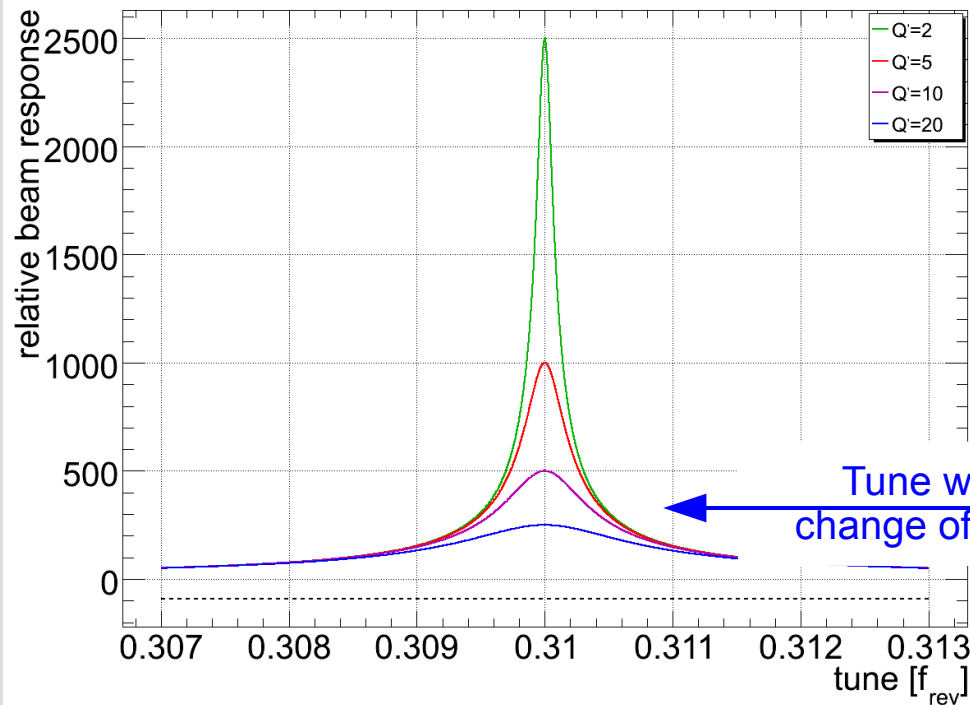
### Option II: larger excitation bandwidth, multiple exciter or broadband excitation(FNAL)



- Reminder:

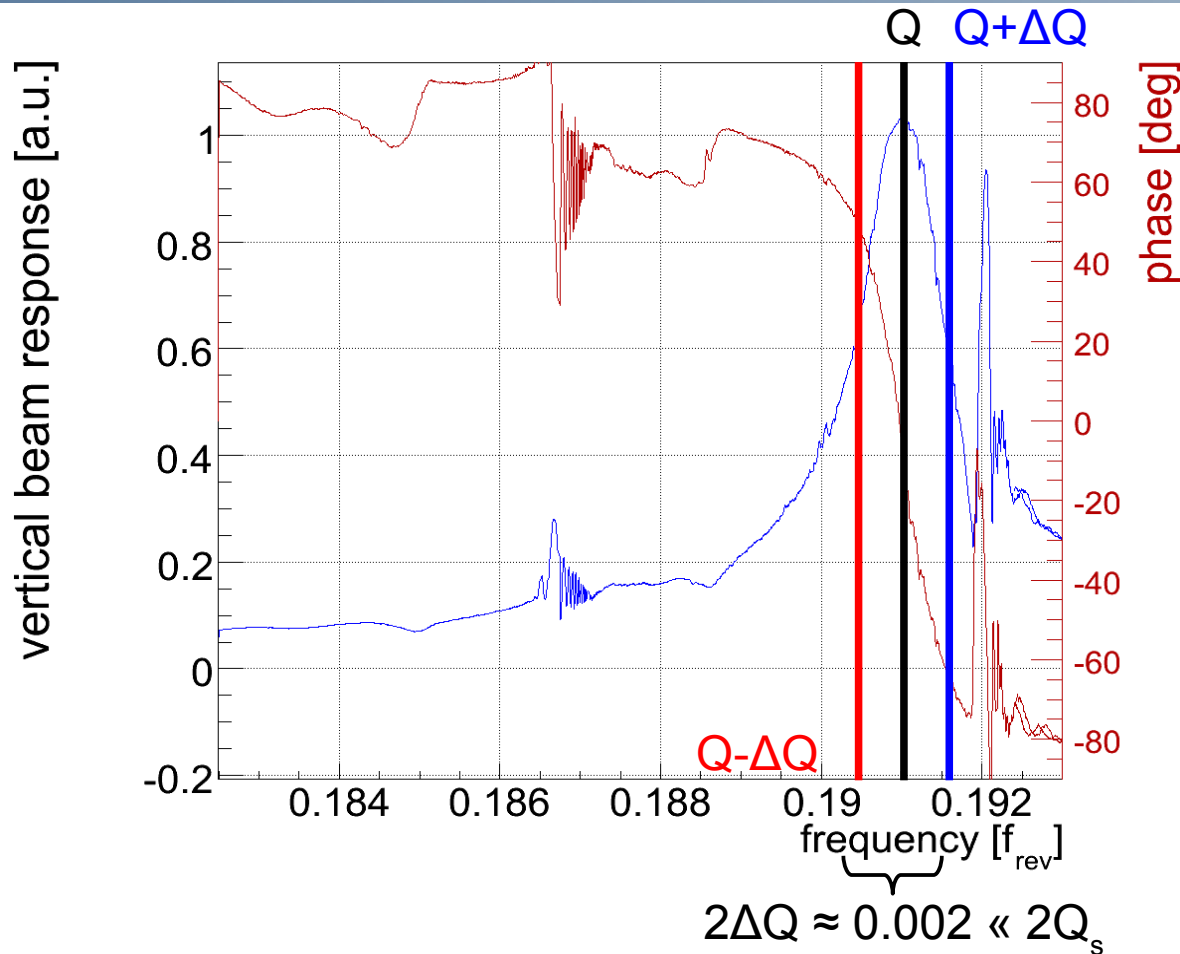
- optimal PLL Settings ( $1/\alpha \sim$  PLL bandwidth/tracking speed):

$$D(s) = K_p + K_i \frac{1}{s} \quad \text{with} \quad K_p = K_0 \frac{\tau}{\alpha} \quad \wedge \quad K_i = K_0 \frac{1}{\alpha}$$



← Tune width change →  
change of phase slope ( $K_0$ ) →

- Optimal PLL parameters (tracking speed, etc.) depend - beside measurement noise – on the effective tune width.
- Intrinsic trade-off:
  - Optimal PI for large  $\Delta Q \leftrightarrow$  sensitivity to noise (unstable loop) for small  $\Delta Q$
  - Optimal PI for small  $\Delta Q \leftrightarrow$  slow tracking speed for large  $\Delta Q$
- Can be improved by putting knowledge into the system: “gain scheduling”



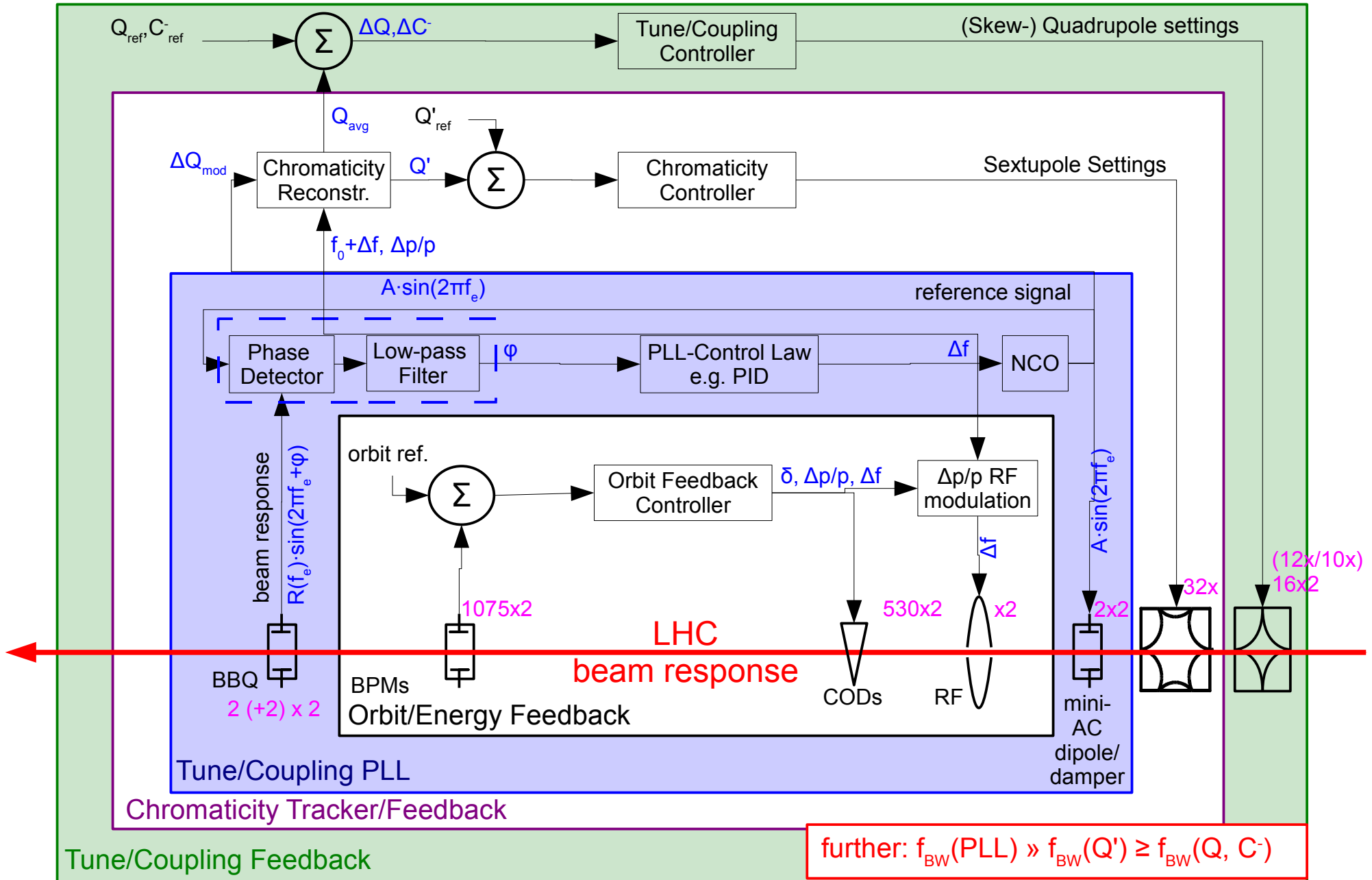
- Resonant phase change  $\leftrightarrow$  tune width change
  - “free” real-time tune footprint measurement
  - measurable dependence of  $\Delta Q \sim Q'$

driven resonance:

$$\tan(\varphi) \approx \frac{\Delta Q \cdot \omega_Q \omega_D}{\omega_Q^2 - \omega_D^2}$$

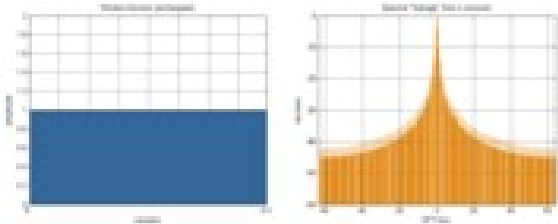
# Additional Topic III: Feed-Backs on Tune, Coupling and Chromaticity





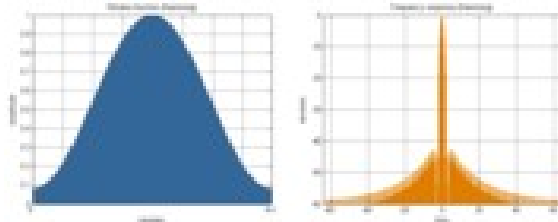
LHC FBs: 2158 input devices, 1136 output devices → total: ~3300 devices!

- rectangular, B=1.0



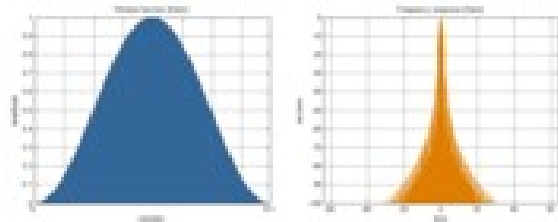
$$\omega(n) = 1$$

- Hamming, B = 1.37



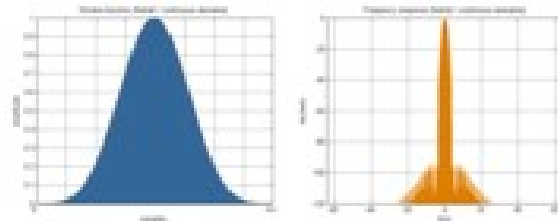
$$\omega(n) = 0.53836 - 0.46164 \cos\left(\frac{2\pi n}{N-1}\right)$$

- Von Hann, B = 1.5



$$\omega(n) = 0.5 \cdot \left[ 1 - \cos\left(\frac{2\pi n}{N-1}\right) \right]$$

- Nuttall, B = 2.01



$$\omega(n) = a_0 - a_1 \cos\left(\frac{2\pi n}{N-1}\right) + a_2 \cos\left(\frac{4\pi n}{N-1}\right) - a_3 \cos\left(\frac{6\pi n}{N-1}\right)$$

$$a_0=0.35875, \quad a_1=0.48829, \quad a_2=0.14128, \quad a_3=0.01168$$

- See wikipedia article [http://en.wikipedia.org/wiki/Window\\_function](http://en.wikipedia.org/wiki/Window_function) for details