



Introduction to Accelerators

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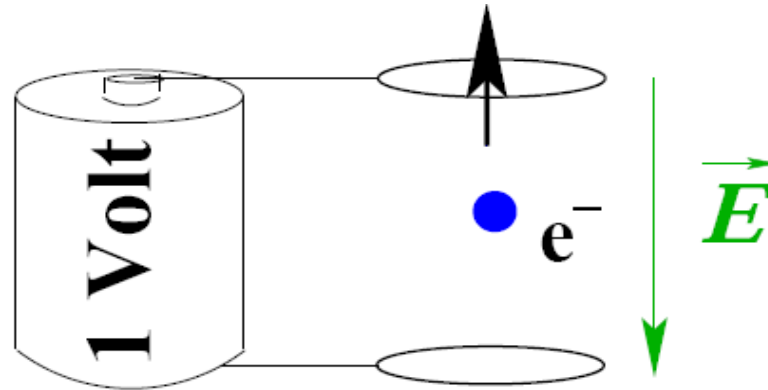


Why this Introduction?

- During this school, you will learn about **beam dynamics** in a rigorous way...
- but some of you are completely new to the field of accelerator physics.
- It seemed therefore justified to start with the introduction of a few very **basic concepts**, which will be used throughout the course.

This is a completely **intuitive approach** (no mathematics) aimed at highlighting the physical concepts, without any attempt to achieve any scientific derivation.

Units: the electronvolt (eV)



The **electronvolt (eV)** is the energy gained by an electron travelling, in vacuum, between two points with a voltage difference of 1 Volt. **$1 \text{ eV} = 1.602 \cdot 10^{-19} \text{ Joule}$**

We also frequently use the electronvolt to express masses from $E=mc^2$: **$1 \text{ eV}/c^2 = 1.783 \cdot 10^{-36} \text{ kg}$**



Beam Dynamics (1)

In order to describe the motion of the particles, each particle is characterised by:

- Its azimuthal position along the machine: s
- Its Energy: E
- Its horizontal position: x
- Its horizontal slope: x'
- Its vertical position: y
- Its vertical slope: y'

i.e. a sixth dimensional vector

(s, E, x, x', y, y')



Beam Dynamics (2)

- In an accelerator designed to operate at the energy E_{nom} , all particles having $(s, E_{nom}, 0, 0, 0, 0)$ will happily fly through the center of the vacuum chamber without any problem. These are “ideal particles”.
- The difficulties start when:
 - one introduces **dipole magnets**
 - the energy $E \neq E_{nom}$ or $(p-p_{nom}/p_{nom}) = \Delta p/p_{nom} \neq 0$
 - either of $x, x', y, y' \neq 0$



Basic problem:

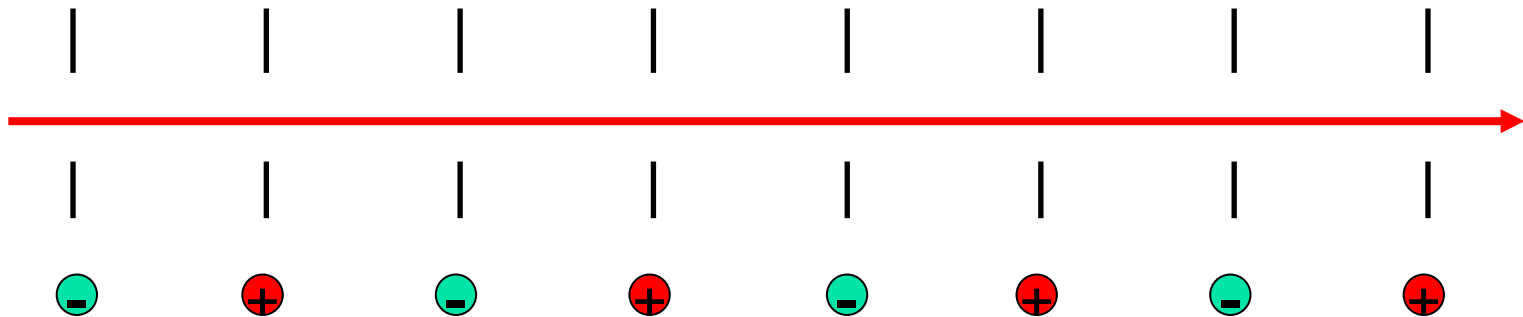
With more than 10^{10} particles per bunch, most of them will **not** be **ideal particles**, i.e. they are going to be lost !

Purpose of this lecture: how can we keep the particles in the machine ?

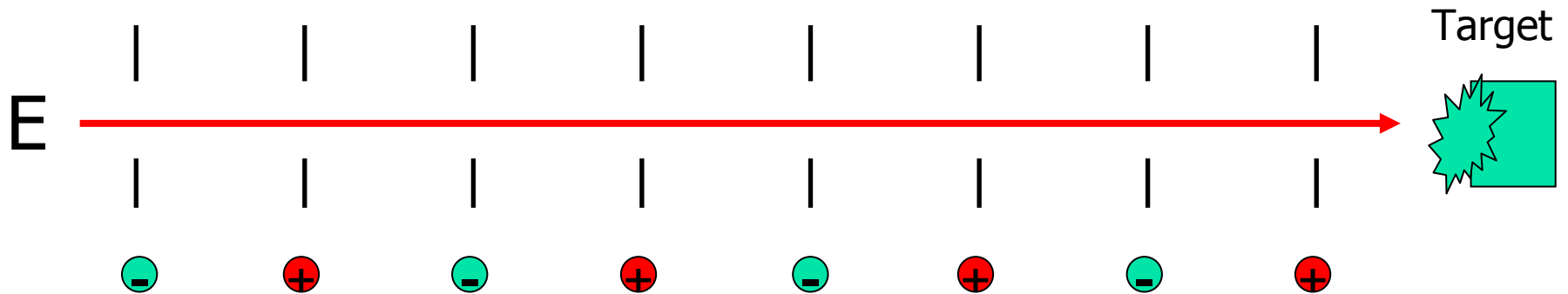
What is a Particle Accelerator?

➤ a machine to accelerate some particles ! **How is it done ?**

➤ Many different possibilities, but rather easy from the general principle:



Ideal linear machines (linacs)

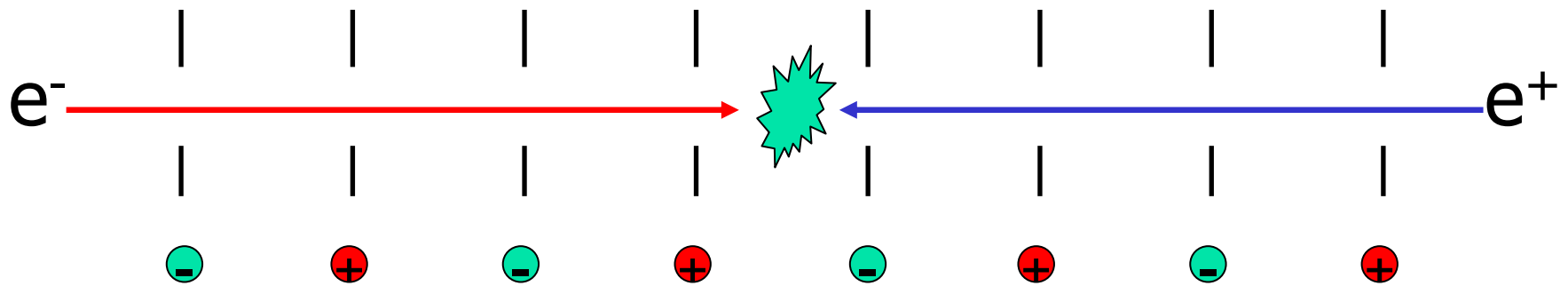


Available Energy : $E_{c.m.} = m \cdot (2+2\gamma)^{1/2} = (2m \cdot (m+E))^{1/2}$
with $\gamma = E/E_0$

Advantages: Single pass
High intensity

Drawbacks: Single pass
Available Energy

Improved solution for $E_{c.m.}$



Available Energy : $E_{c.m.} = 2m\gamma = 2E$

with $\gamma = E/E_0$

Advantages: High intensity

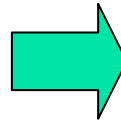
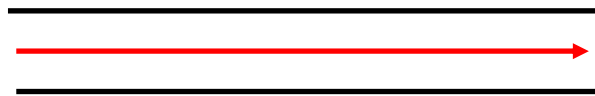
Drawbacks: Single pass

Space required

Keep particles: circular machines

Basic idea is to keep the particles in the machine for many turns.

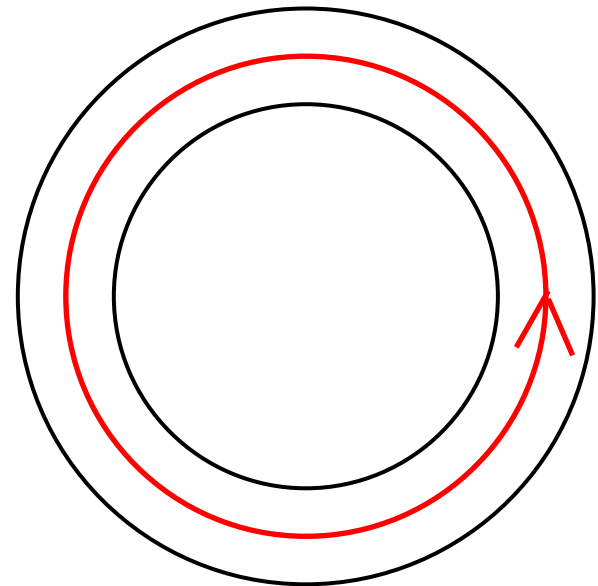
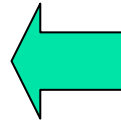
Move from the linear design



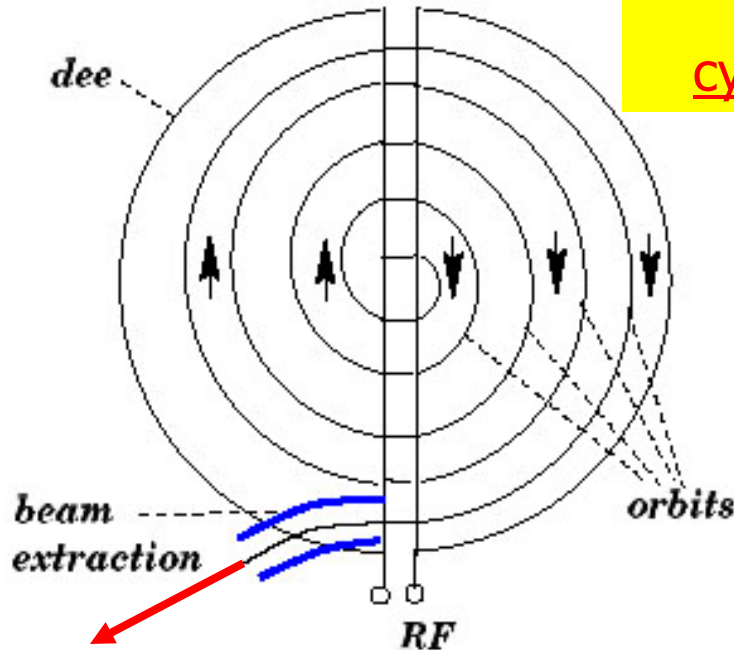
To a circular one:

➤ Need Bending

➤ Need **Dipoles!**



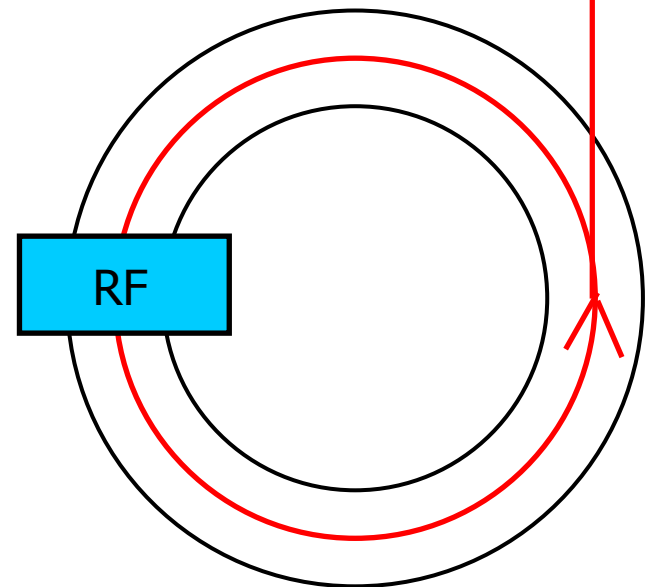
Circular machines ($E_{c.m.} \sim (mE)^{1/2}$)



fixed target:
cyclotron

huge dipole, compact design,
B = constant
low energy, single pass.

fixed target:
synchrotron

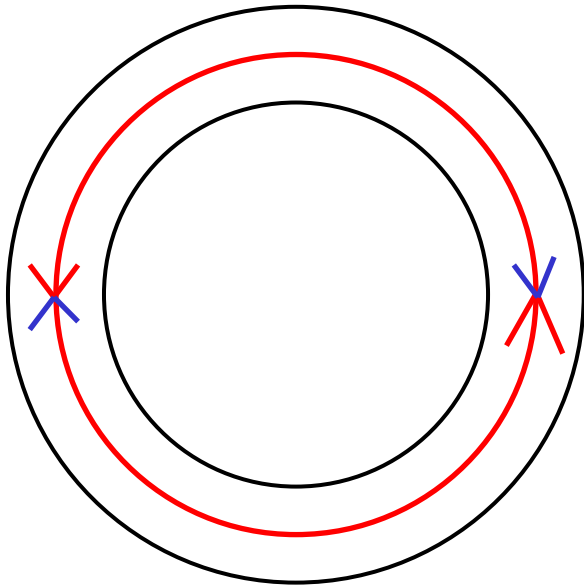


varying B, small magnets, high energy

Colliders ($E_{\text{c.m.}} = 2E$)

Colliders:

electron – positron
proton - antiproton

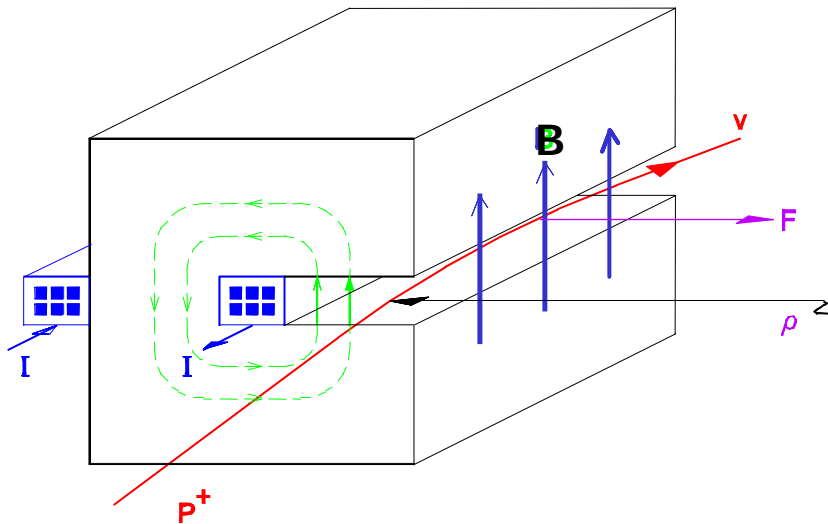


Colliders with the same type of particles (e.g. p-p) require two separate chambers. The beam are brought into a common chamber around the interaction regions

Ex: LHC

8 possible interaction regions
4 experiments collecting data

Circular machines: Dipoles



Classical mechanics:

Equilibrium between two forces

Lorentz force

Centrifugal force

$$F = e.(\underline{v} \times \underline{B})$$

$$F = mv^2/\rho$$

$$evB = mv^2/\rho$$

$$p = m_0.c.(\beta\gamma)$$

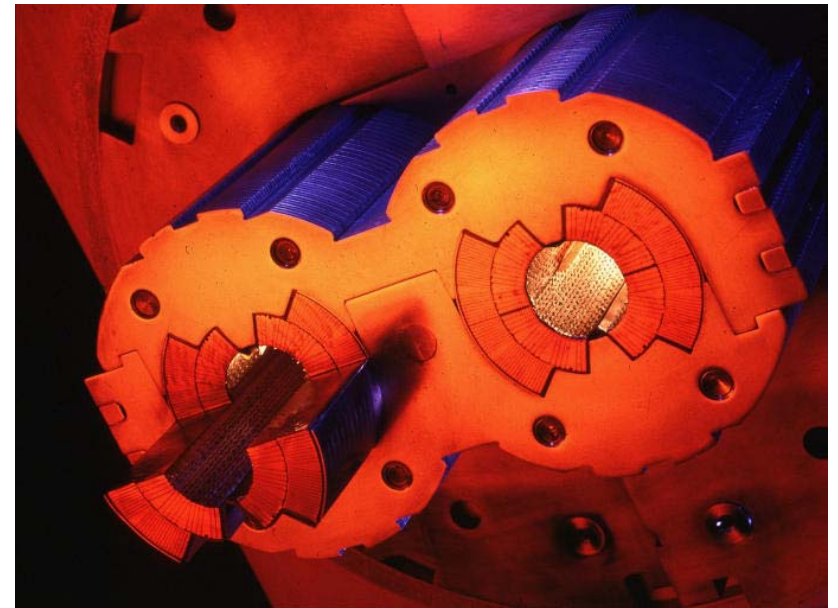


Magnetic rigidity:

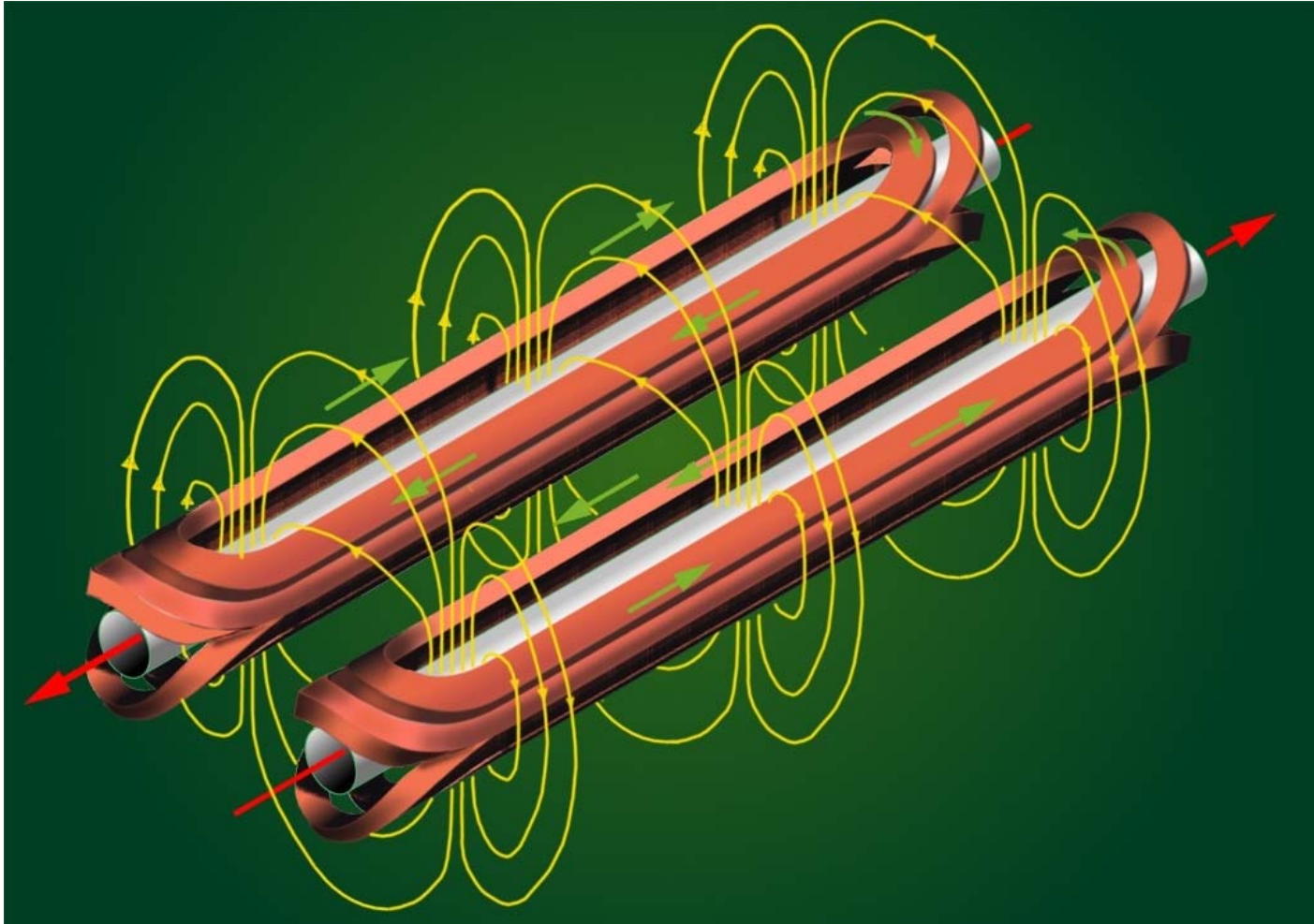
$$B\rho = mv/e = p/e$$

Relation also holds for relativistic case provided the classical momentum mv is replaced by the relativistic momentum p

Dipoles (1):



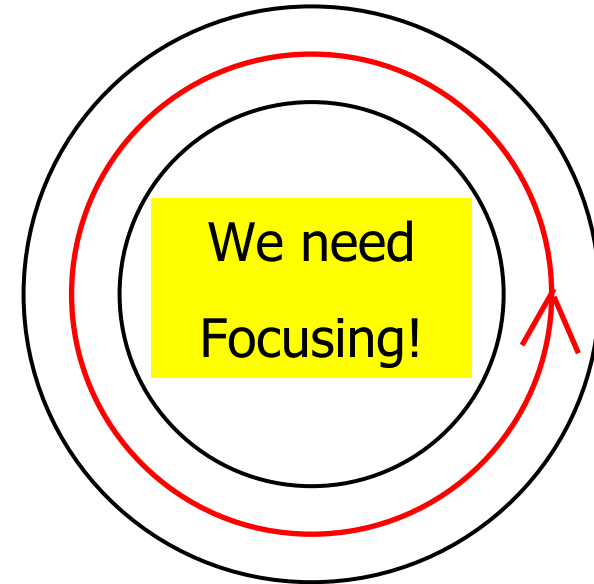
Dipoles (2):



Ideal circular machine:

- Neglecting radiation losses in the dipoles
- Neglecting gravitation

ideal particle would happily circulate on axis in the machine for ever!



Unfortunately: real life is different!

Gravitation: $\Delta y = 20$ mm in 64 msec!

Alignment of the machine

Limited physical aperture

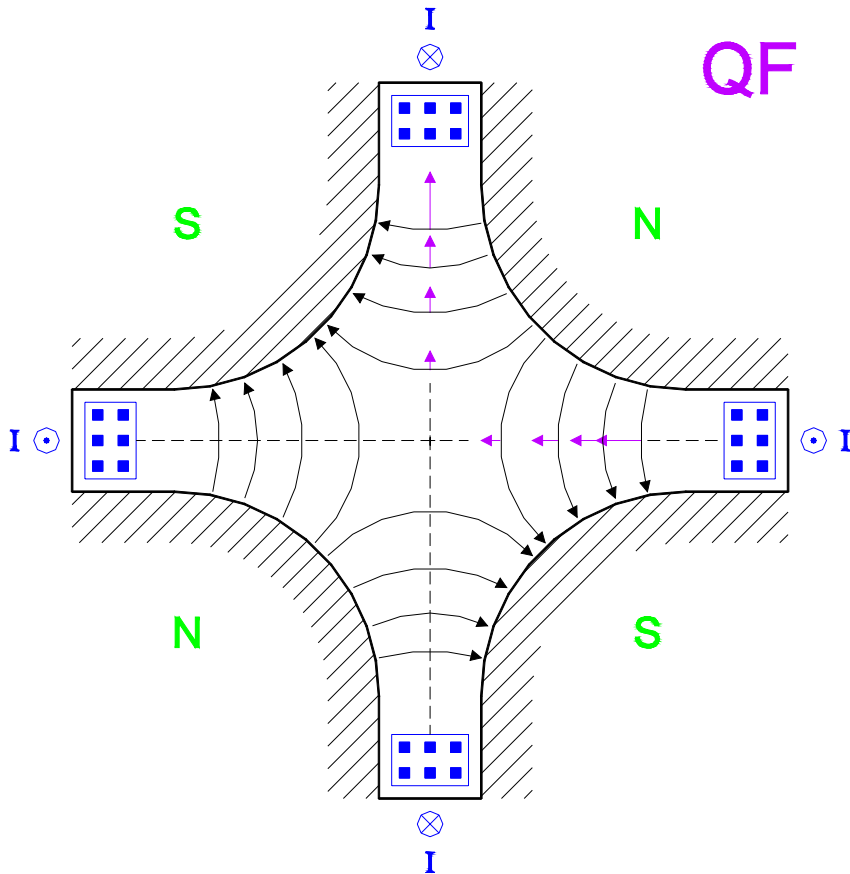
Ground motion

Field imperfections

Energy error of particles **and/or** $(x, x')_{inj} \neq (x, x')_{nominal}$

Error in magnet strength (power supplies and calibration)

Focusing with quadrupoles



$$F_x = -g \cdot x$$

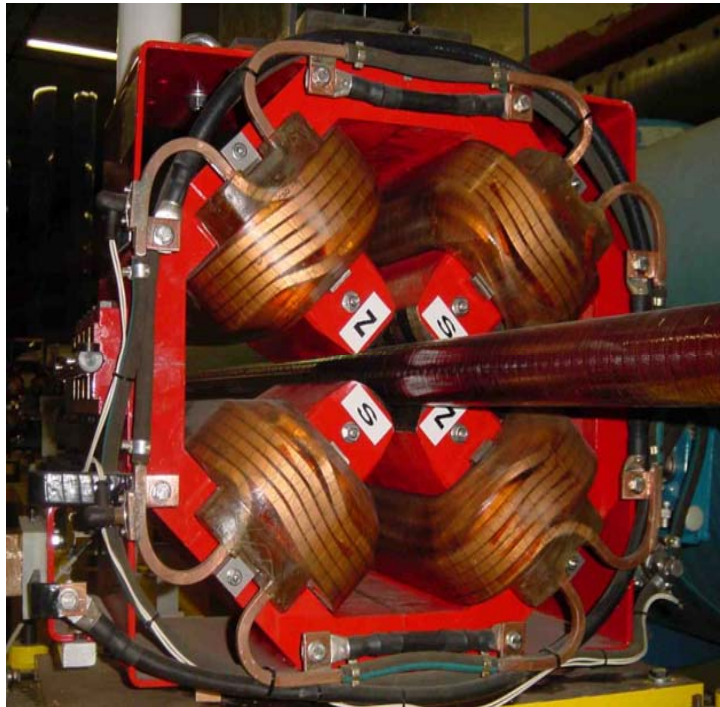
$$F_y = g \cdot y$$

Force increases **linearly** with displacement.

Unfortunately, effect is **opposite** in the two planes (H and V).

Remember: **this** quadrupole is **focusing** in the **horizontal** plane but **defocusing** in the **vertical** plane!

Quadrupoles:





Focusing properties ...

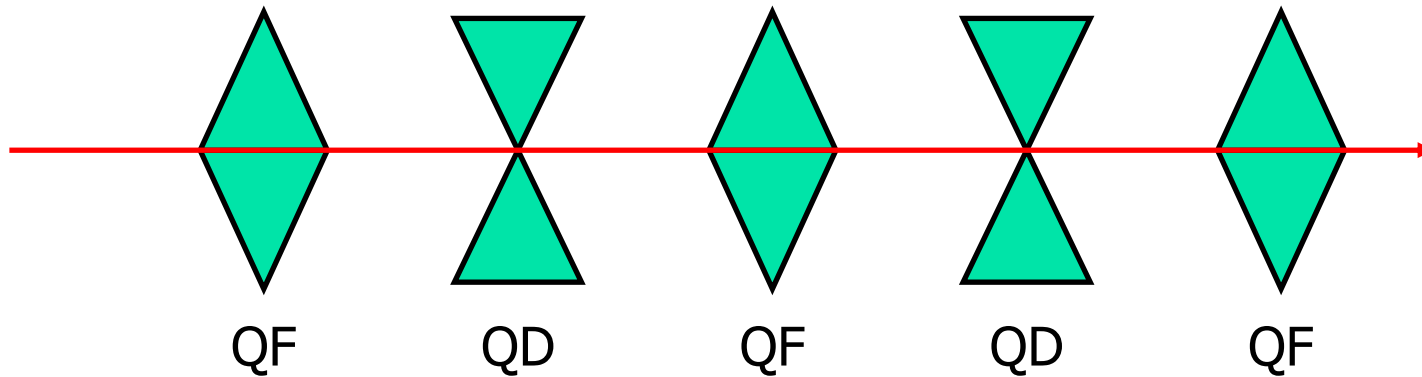
A quadrupole provides the required effect in one plane...

but the opposite effect in the other plane!

Is it really interesting ?

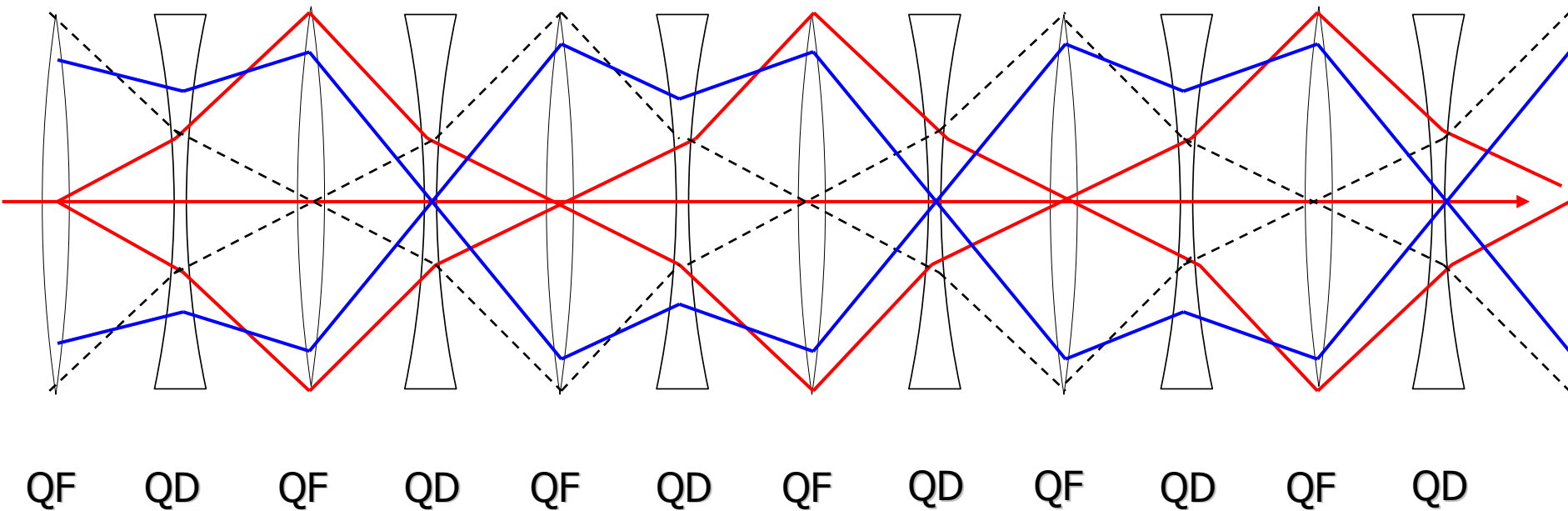
Alternating gradient focusing

Basic new idea:
Alternate QF and QD



valid for one plane only (H or V) !

Alternating gradient focusing





Alternating gradient focusing:

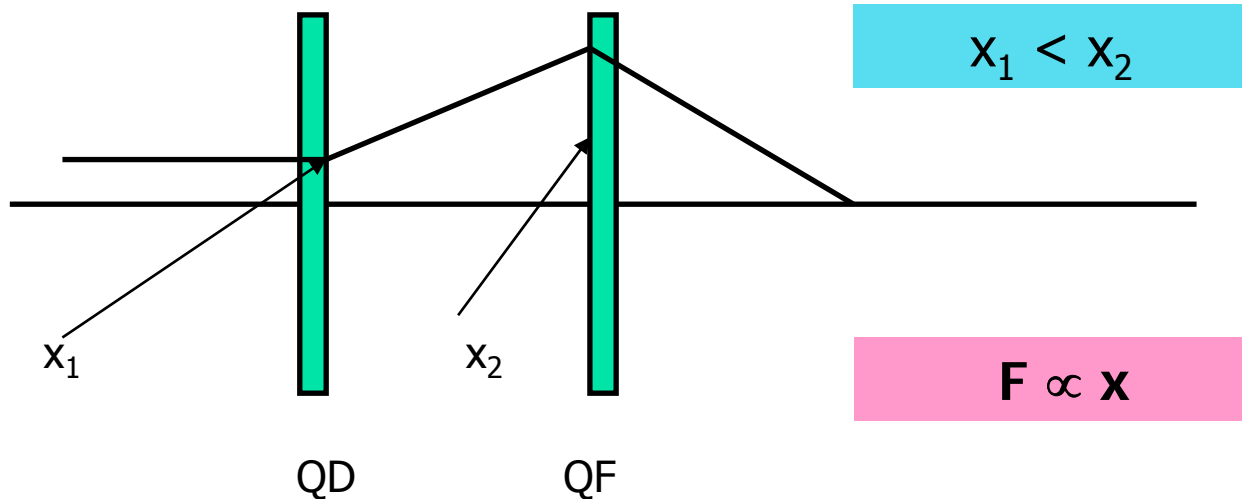
It can be shown that a section composed of **alternating focusing** and **defocusing** elements has a **net focusing effect**, provided the quadrupoles are correctly placed.

Particles for which **$x, x', y, y' \neq 0$** thus **oscillate around** the ideal particle ...

but the trajectories remain inside the vacuum chamber !

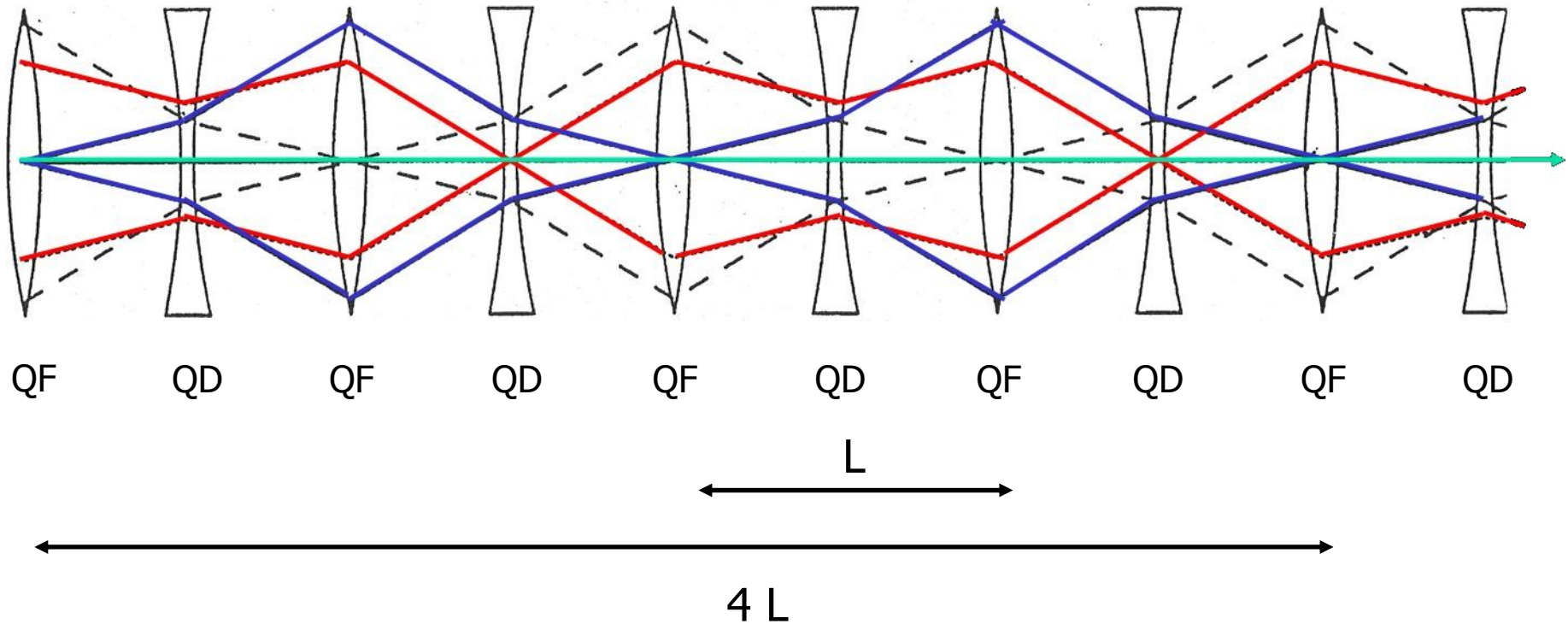
Why net focusing effect?

Purely intuitively:



Rigorous treatment rather straightforward !

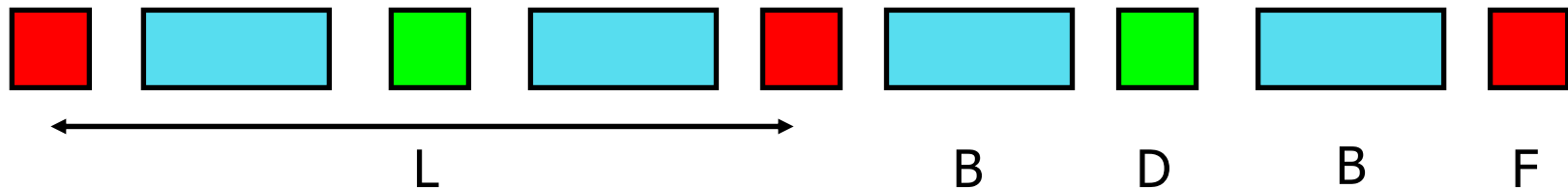
The concept of the « FODO cell »



One complete oscillation in 4 cells $\Rightarrow 90^\circ / \text{cell} \Rightarrow \mu = 90^\circ$

Real circular machines

The accelerator is composed of a **periodic** repetition of **cells**:



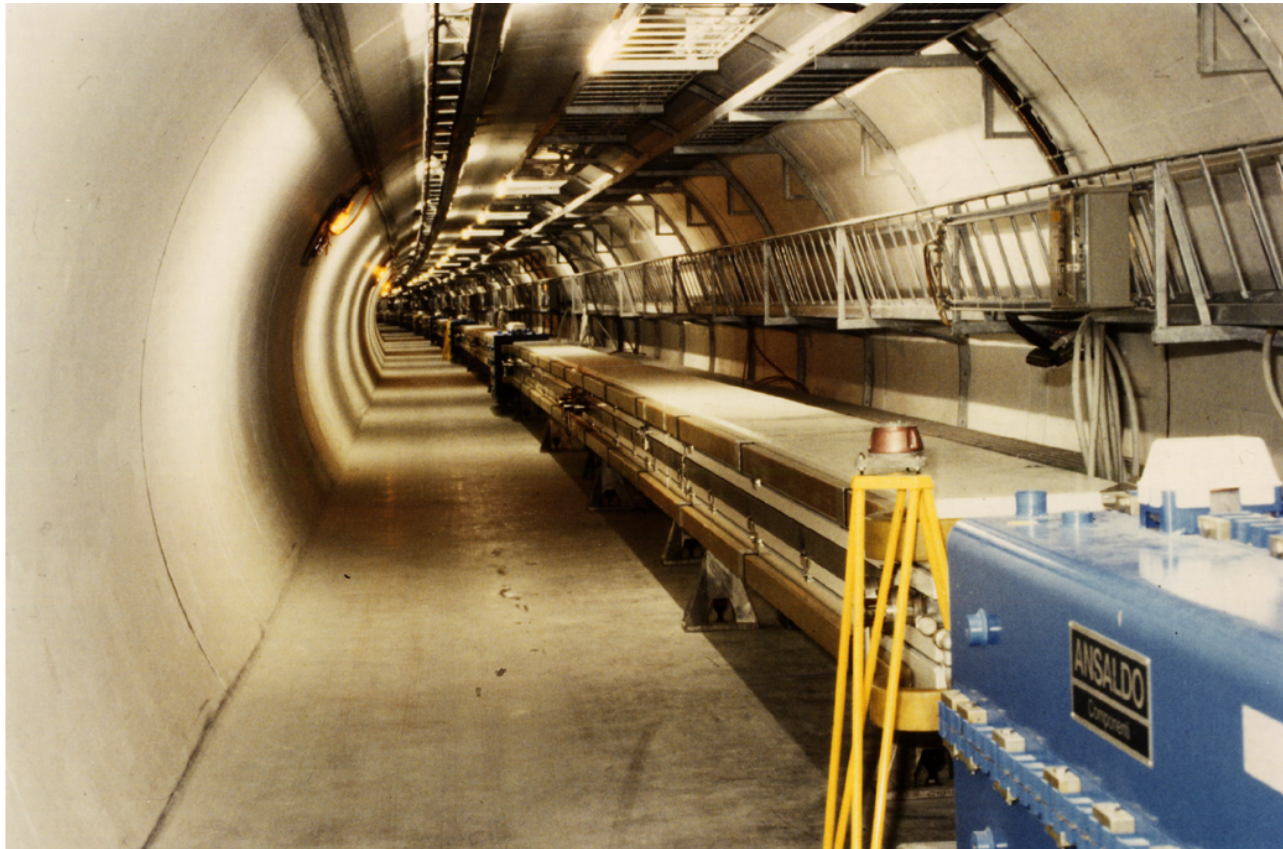
➤ The phase advance per cell μ can be modified, in each plane, by varying the strength of the quadrupoles.

➤ The ideal particle will follow a **particular** trajectory, which **closes on itself** after one revolution: **the closed orbit**.

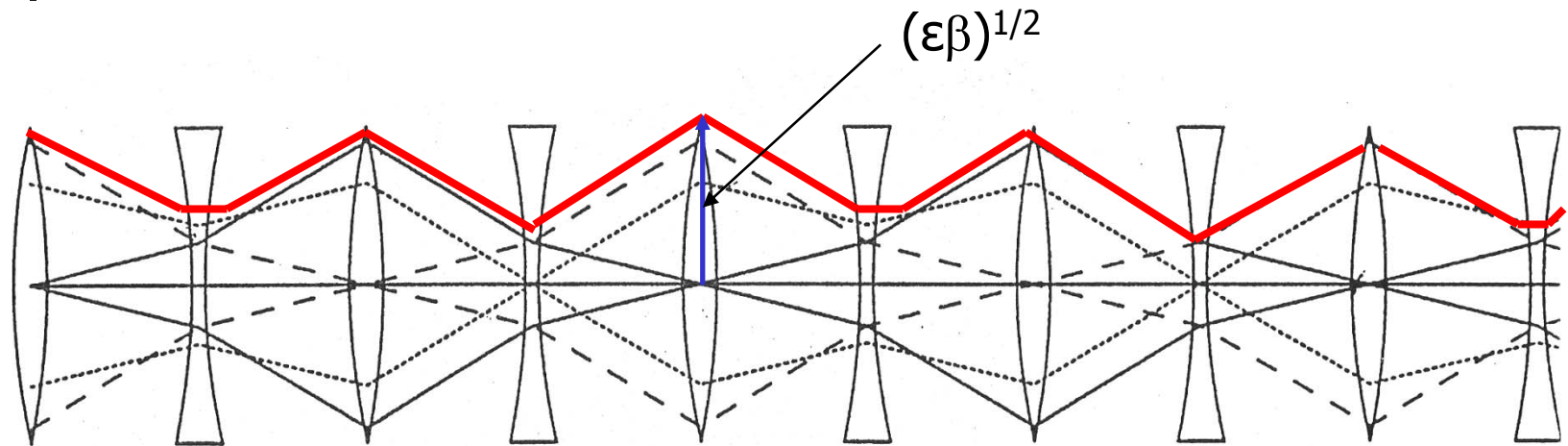
➤ The real particles will perform oscillations **around the closed orbit**.

➤ The number of **oscillations for a complete revolution** is called the **Tune Q** of the machine (Q_x and Q_y).

Regular periodic lattice: The Arc



The beta function $\beta(s)$



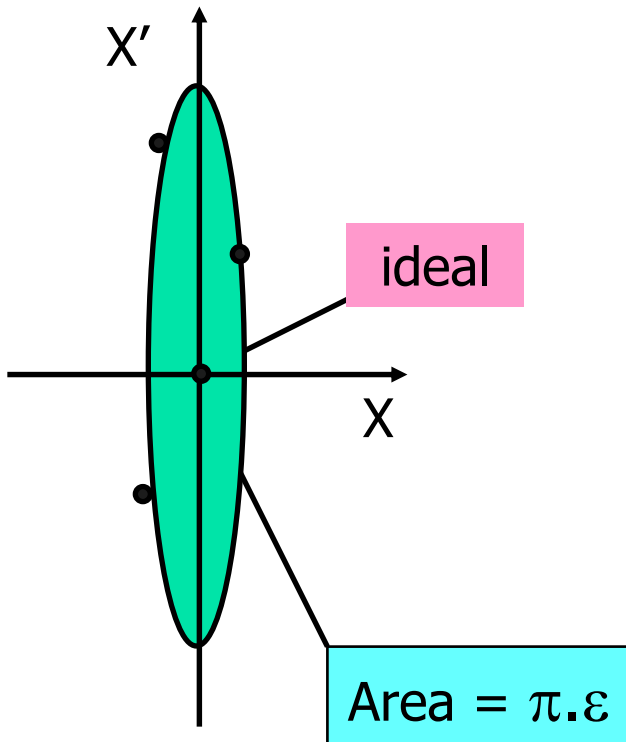
The β -function is the **envelope** around all the trajectories of the particles circulating in the machine.

The β -function has a **minimum at the QD** and a **maximum at the QF**, ensuring the net focusing effect of the lattice.

It is a **periodic function** (repetition of cells). The oscillations of the particles are called **betatron motion** or **betatron oscillations**.

Phase space at some position (s)

- Select the particle for which 65% of the particles (1σ) have a **smaller betatron motion** and plot its **position vs. its phase** (x vs. x') at some location in the machine for many turns.

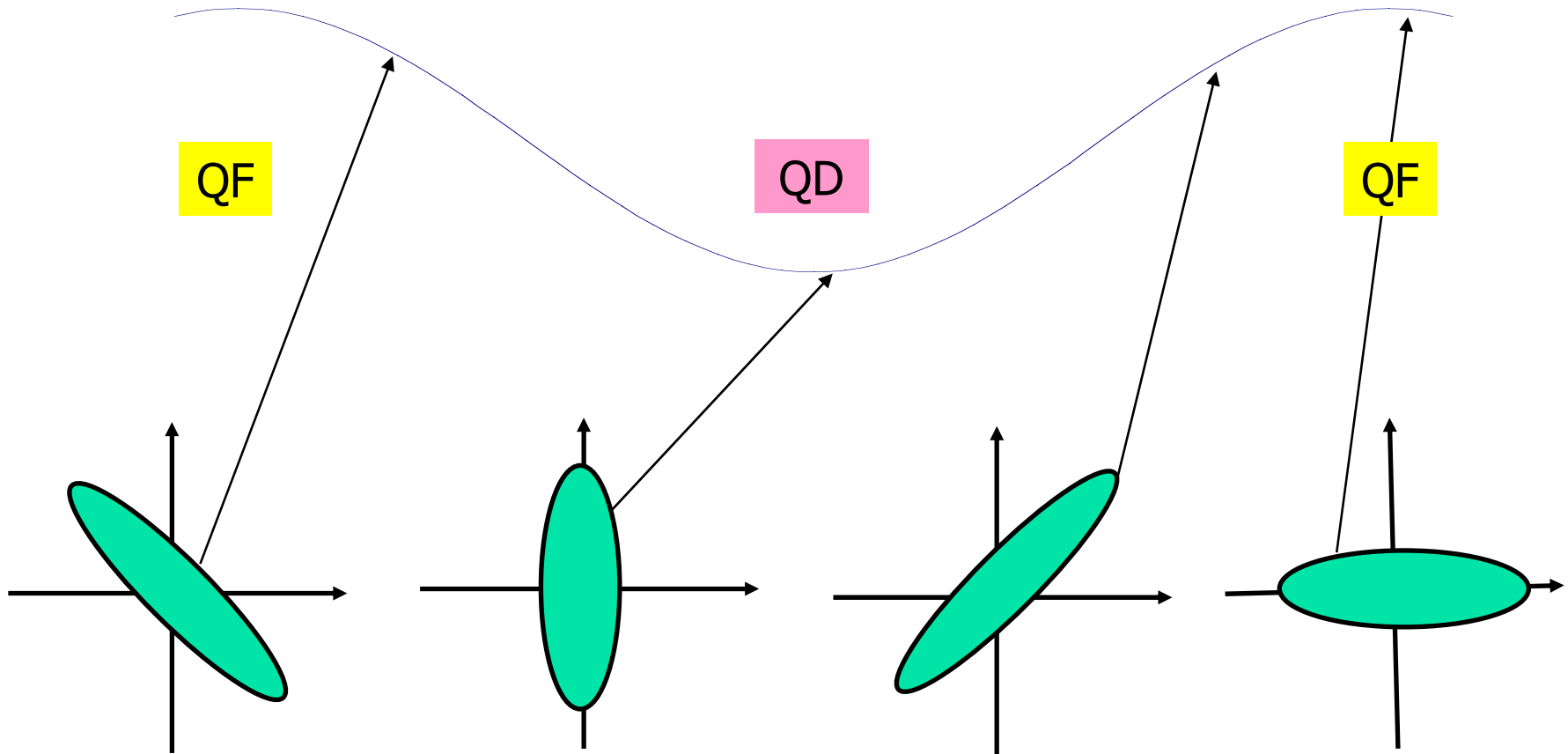


- ϵ Is the emittance of the beam [mm mrad]
- ϵ is a **property of the beam** (quality)
- Measure of how much particle depart from ideal trajectory.
- β is a **property of the machine** (quadrupoles).

Beam size [m]

$$\sigma(s) = (\epsilon \cdot \beta(s))^{1/2}$$

Emittance conservation



The shape of the ellipse varies along the machine, but its area (the emittance ε) remains constant for a given energy.

Recapitulation 1

- The fraction of the oscillation performed in a periodic cell is called the phase advance μ per cell (x or y).
- The total number of oscillations over one full turn of the machine is called the betatron tune Q (x or y).
- The envelope of the betatron oscillations is characterised by the beta function $\beta(s)$. This is a property of the quadrupole settings.
- The quality of the (injected) beam is characterised by the emittance ϵ . This is a property of the beam and is invariant around the machine.
- The r.m.s. beam size (measurable quantity) is $\sigma = (\beta \cdot \epsilon)^{1/2}$.



Off momentum particles:

- These are “non-ideal” particles, in the sense that they do not have the right energy, i.e. all particles with $\Delta p/p \neq 0$

What happens to these particles when traversing the magnets ?

Off momentum particles ($\Delta p/p \neq 0$)

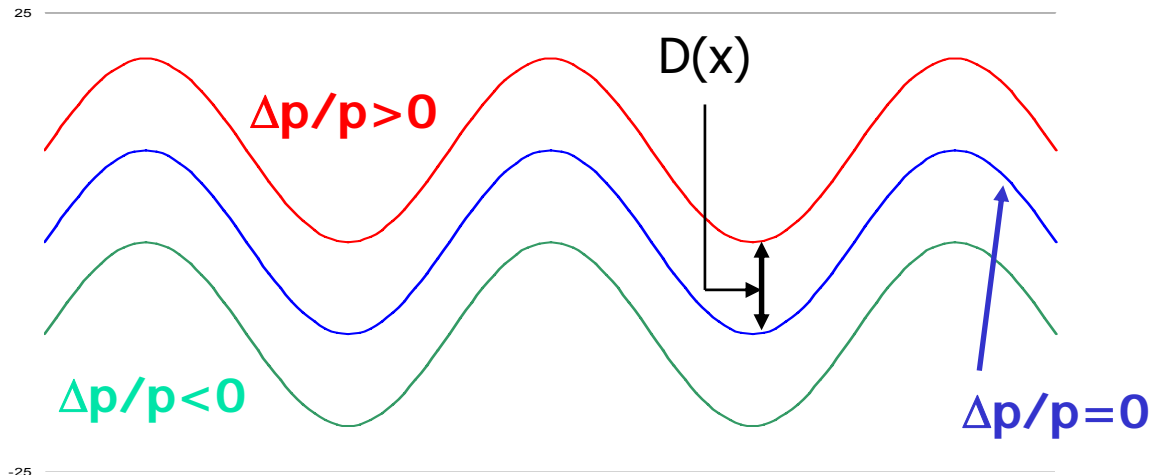
Effect from Dipoles

➤ If $\Delta p/p > 0$, particles are **less** bent in the dipoles → should spiral out !

➤ If $\Delta p/p < 0$, particles are **more** bent in the dipoles → should spiral in !

No!

There is an equilibrium with the restoring force of the quadrupoles

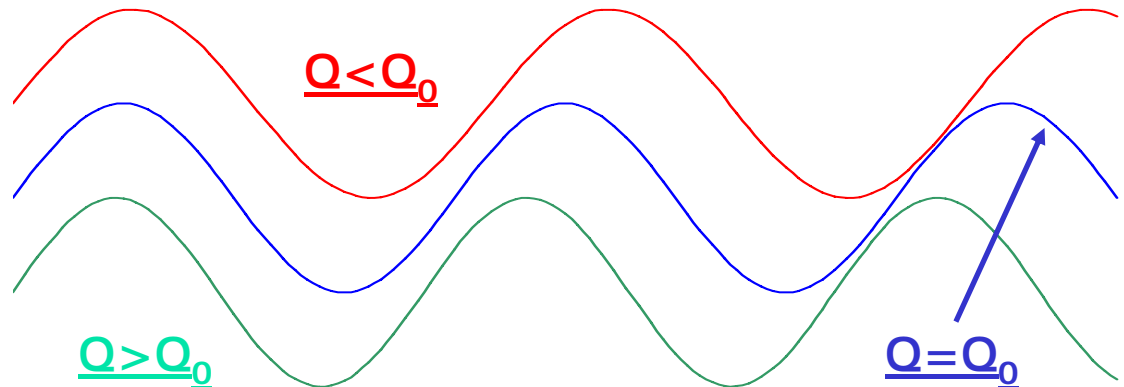


Off momentum particles ($\Delta p/p \neq 0$)

Effect from Quadrupoles

- If $\Delta p/p > 0$, particles are **less** focused in the quadrupoles → **lower Q !**
- If $\Delta p/p < 0$, particles are **more** focused in the quadrupoles → **higher Q !**

Particles with different momenta would have a different **betatron tune** $Q=f(\Delta p/p)$!





The chromaticity Q'

Particles with different momenta ($\Delta p/p$) would thus have different tunes Q .
So what ?

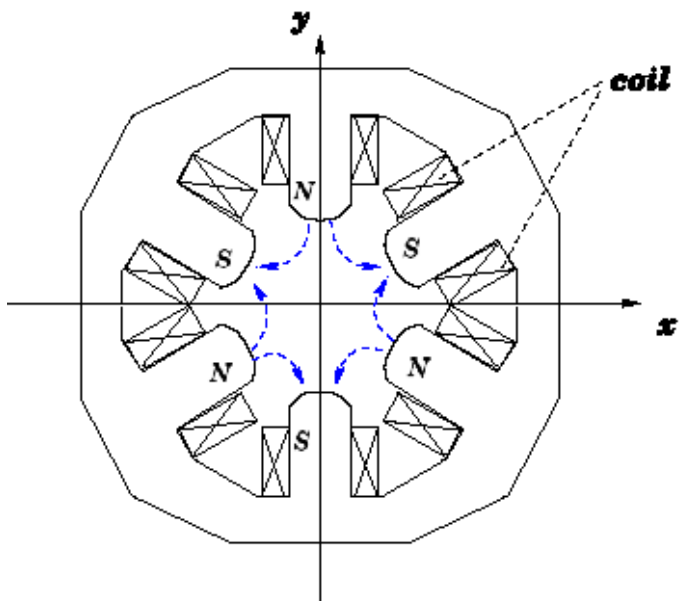
unfortunately

- The tune dependence on momentum is of **fundamental** importance for the **stability** of the machine. It is described by the **chromaticity** of the machine Q' :

$$Q' = \Delta Q / (\Delta p/p)$$

The chromaticity has to be carefully **controlled and corrected** for stability reasons. This is achieved by means of **sextupoles**.

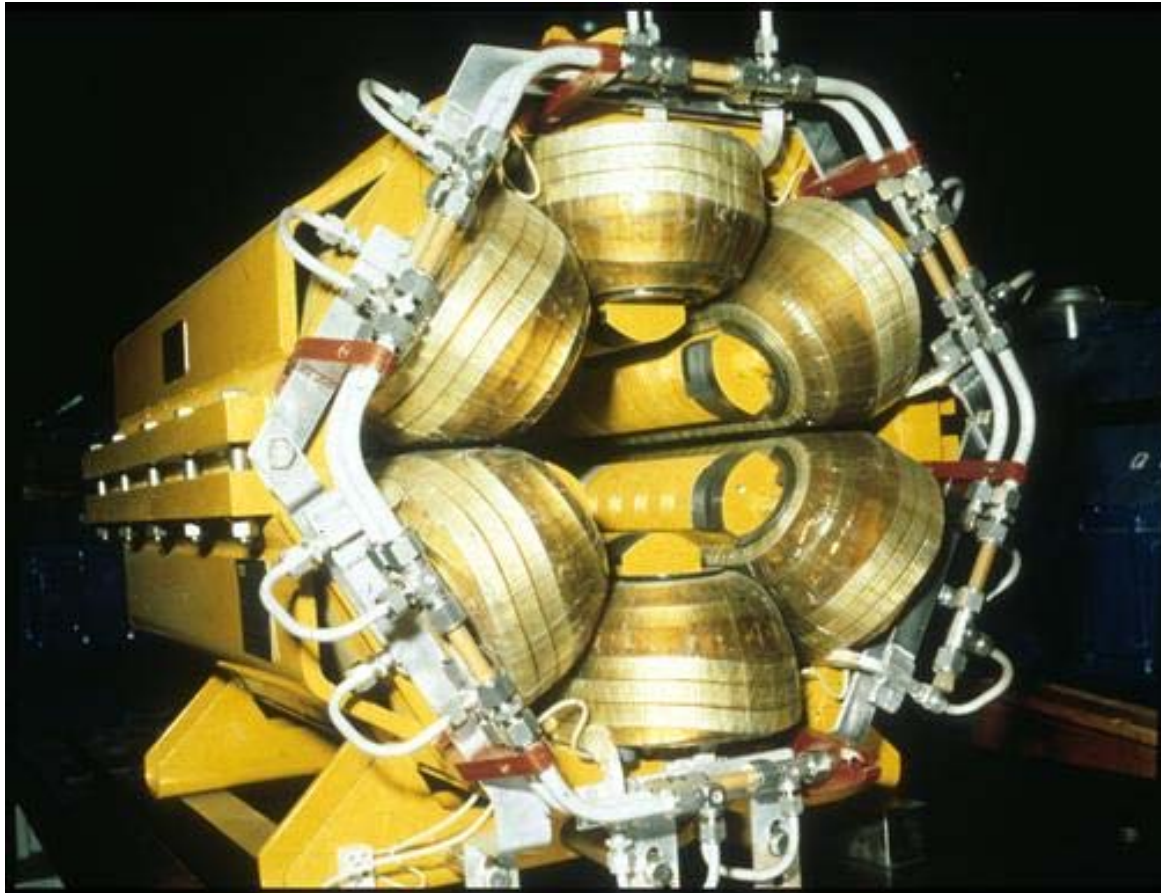
The sextupoles (SF and SD)



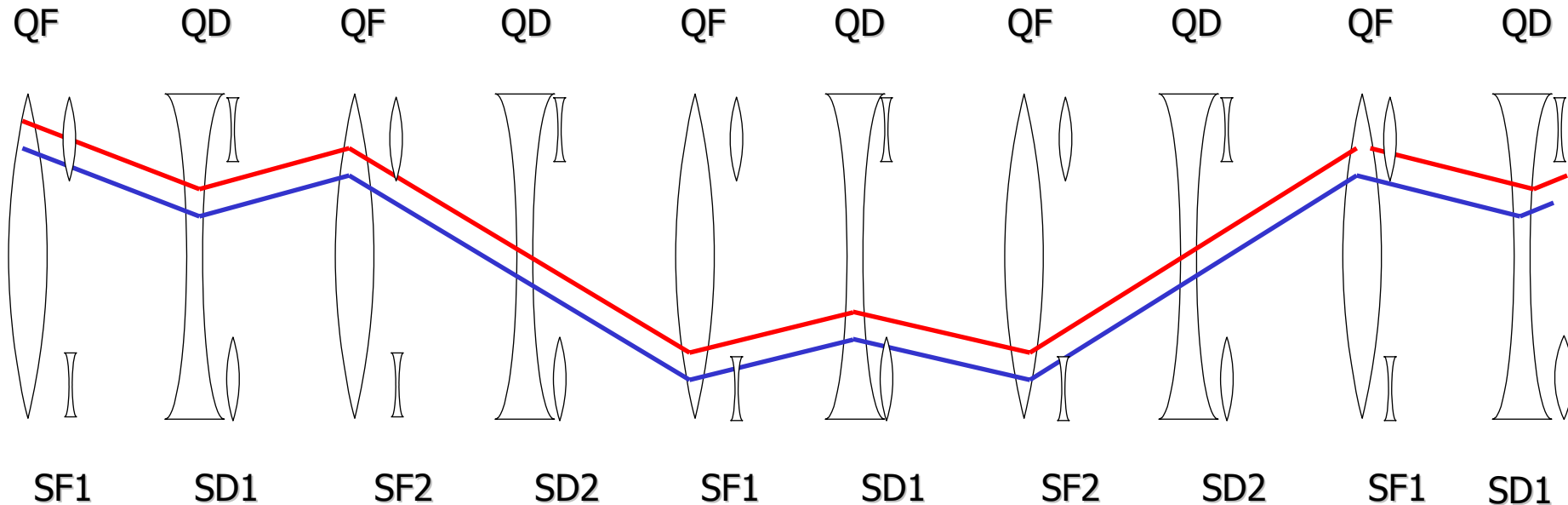
- $\Delta x' \propto x^2$
- A SF sextupole basically « **adds** » focusing for the particles with $\Delta p/p > 0$, and « **reduces** » it for $\Delta p/p < 0$.
- The chromaticity is corrected by adding a sextupole after each quadrupole of the FODO lattice.

Sextupoles:

SPS



Chromaticity correction



The undesired effect of sextupoles on particles with the **nominal energy** can be avoided by grouping the sextupoles into « families ».

Nr. of families:

$$N = (k * 180^\circ) / \mu = \text{Integer}$$

$$\text{e.g. } 180^\circ / 90^\circ = 2$$



Recapitulation 2

- For off momentum particles ($\Delta p/p \neq 0$), the magnets induce other important effects, namely:
 - The dispersion (dipoles)
 - The chromaticity (quadrupoles)



Longitudinal plane

➤ So far, we considered only the motion in the transverse planes from an intuitive point of view. The corresponding rigorous treatment will be given in the lectures on “Transverse Beam Dynamics”.

➤ The lectures on “Longitudinal Beam Dynamics” will explain the details of the corresponding longitudinal motion as well as the RF acceleration of the particles.



Natural chromaticity...

- Take a particle and slightly **increase** its momentum:

$$\square \Delta p/p > 0 \quad \square \Delta Q < 0 \quad \square Q' < 0$$

- Take a particle and slightly **decrease** its momentum:

$$\square \Delta p/p < 0 \quad \square \Delta Q > 0 \quad \square Q' < 0$$

Natural Q' is always negative !



Tunes of the machine

Why do we have to control the tunes (Q_x and Q_y) so accurately?

Because there are some (many !) forbidden values!

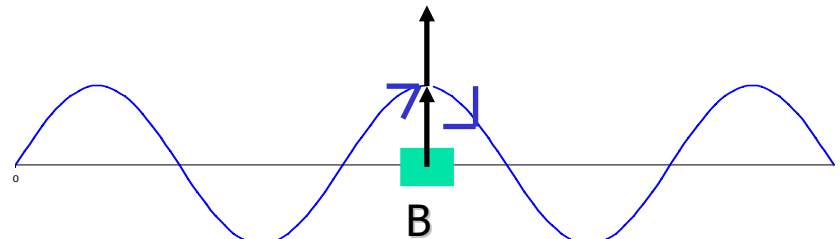
Forbidden values for Q

- An error in a **dipole** gives a kick which has always the same sign!

Integer Tune $Q = N$

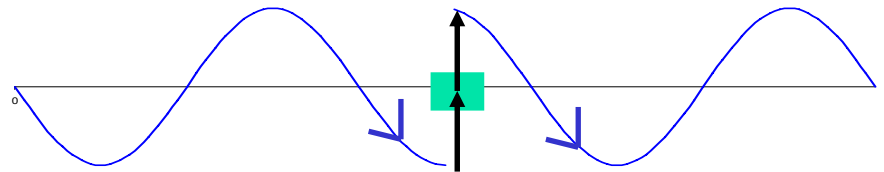
Forbidden!

The perturbation adds up!



Half-integer Tune $Q = N + 0.5$

O.K. for an error in a dipole!
The perturbation cancels
after each turn!

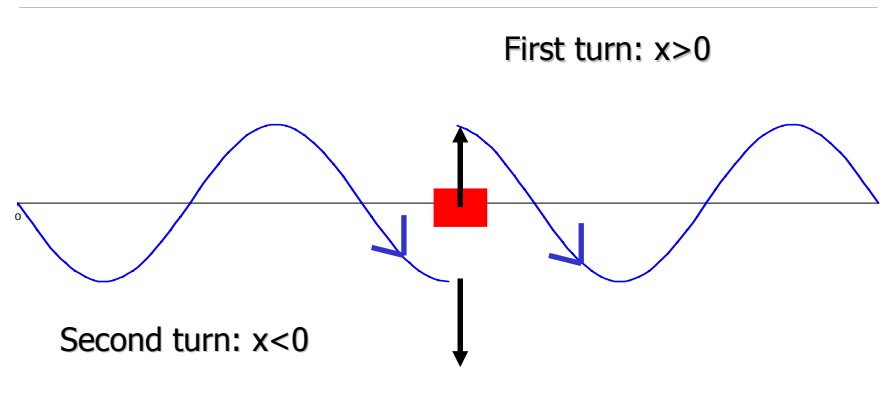


Forbidden values for Q

- An error in a **quadrupole** gives a kick whose sign depends on x

Half-integer Tune $Q = N + 0.5$

Forbidden !
The amplitude of the oscillation is steadily increasing!

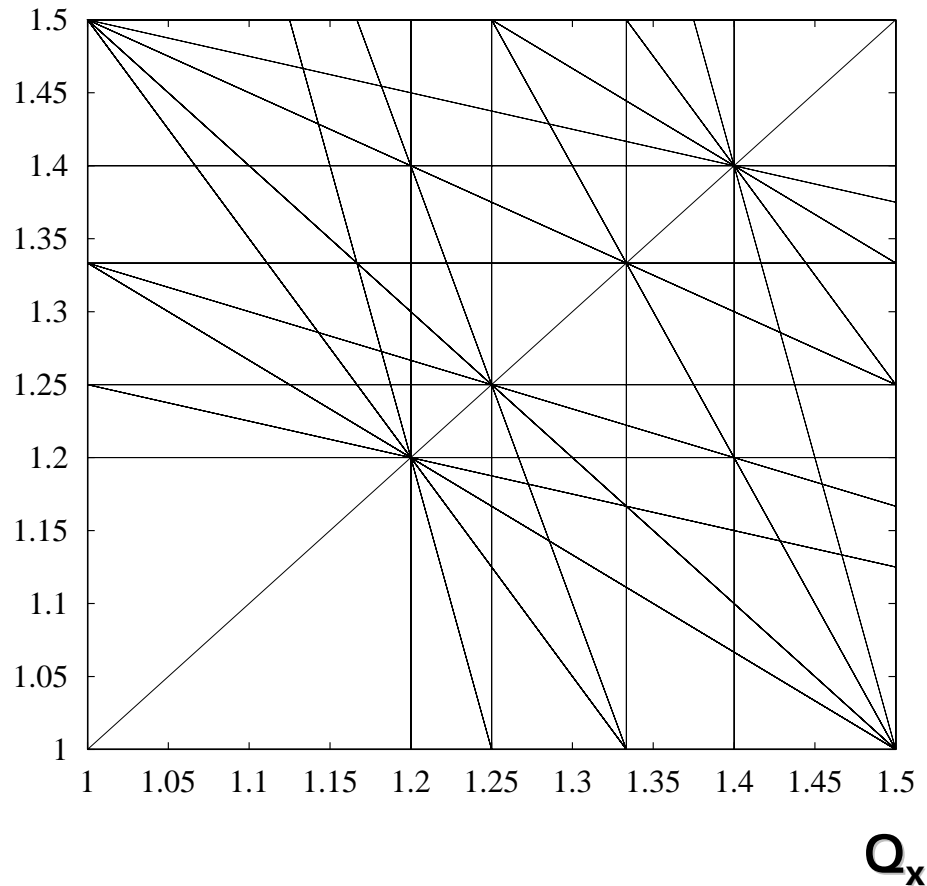


Similar conclusions for $1/3, 1/4, 1/5, \dots$

Tune diagram

an illustration for a lepton machine:

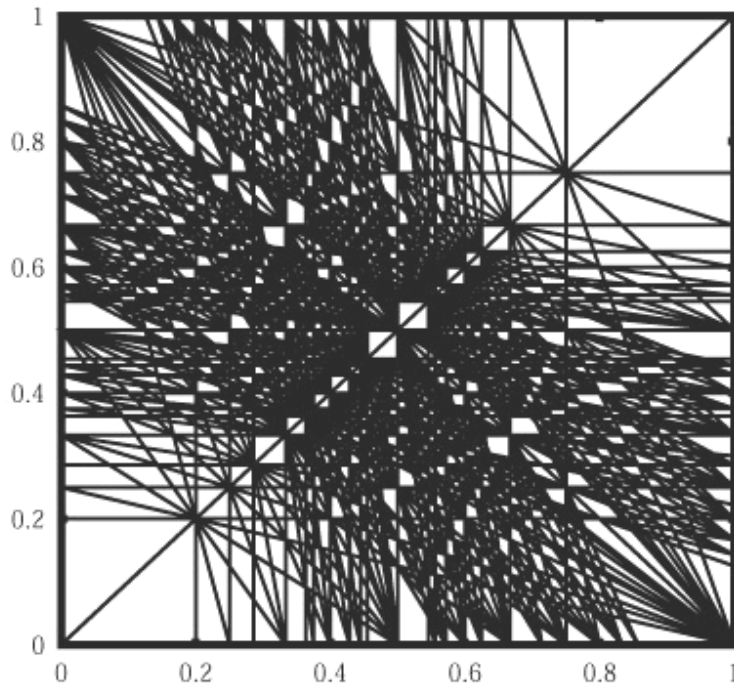
Q_y



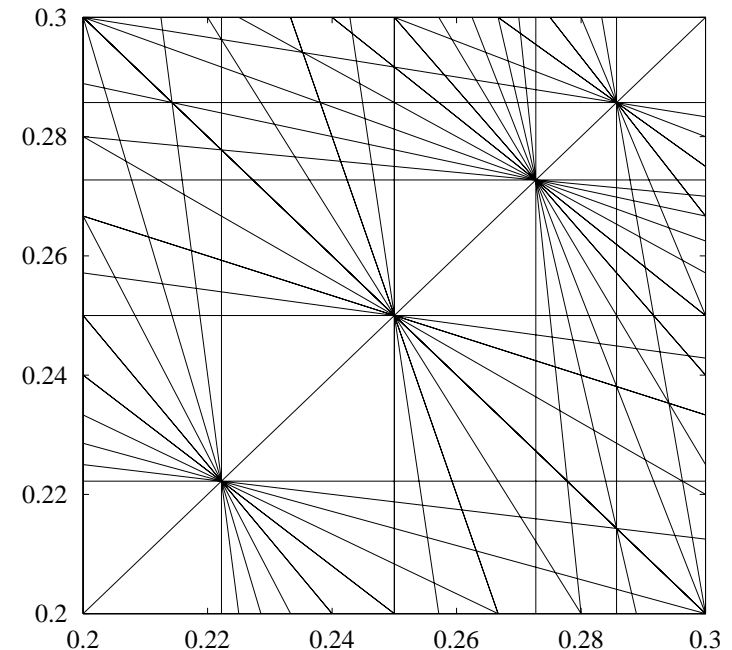
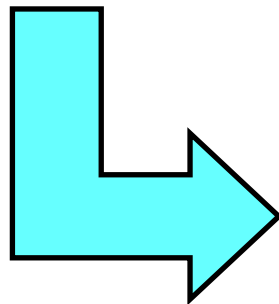
Tune values
(Q_x and/or Q_y)
which are forbidden
in order to avoid
resonances

The lowest the order of
the resonance, the most
dangerous it is.

Tune diagram for protons



The particles have a certain tune spread, the bunch thus represents a small **area** rather than a **point** in the tune diagram.





That's it for the Introduction...

Thank you very much for your attention !!!