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# *Small Emittance Sources/Guns*

*Tsumoru Shintake*

*Spring-8/RIKEN Hyogo Japan*

*Electron source for SASE-FELs, ERLs*

# Topics

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- *This is **not** a review talk on rf photo-injector.*  
*About rf photo-injectors there are ton's of papers.*
  - *LCLS design report Chap. 6*  
*available from <http://www-ssrl.slac.stanford.edu/lcls/CDR/>*
- *In our X-ray FEL project at RIKEN/SPring-8, we choose HV pulse gun using **thermionic cathode** (CeB<sub>6</sub> single crystal).*
- *How did I make this decision.....*
  - *Basic physics around the low emittance electron source.*
  - *Practical things..*

# *before the “emittance” Such Beam Source Should be*

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- **Stable** ..... *charge, energy, pointing,*
  - *important for the machine tuning and performance.*
  - *stable X-ray beam delivery to the user.*  
(SASE fluctuation can be reduced by seeding )
- **Clean** ..... *no halo, no dark current*  
(*high field on cathode such as 100 MV/m will not be good.*)
  - *important for*
    - *undulator protection*
    - *accurate beam position monitoring*
- **Uniform** ..... *current density*
  - *eliminate gain variation due to local lasing.*
  - *deliver an uniform and stable X-ray beam to the user.*
- **Maintenance Free**

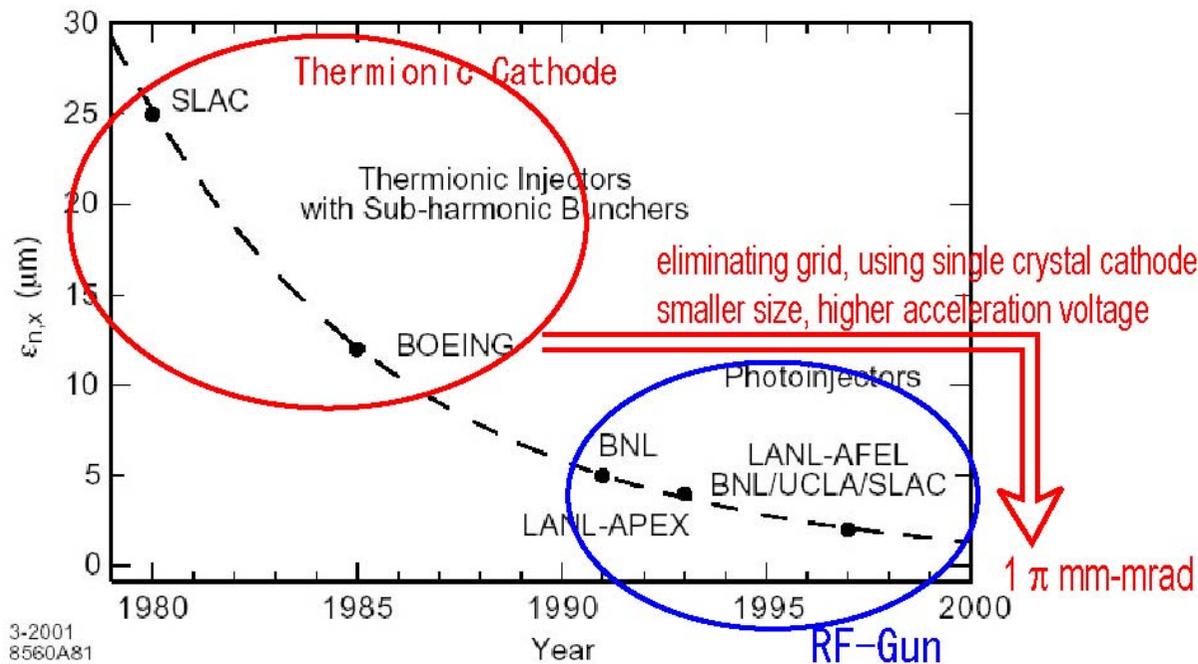
*Like a Full Moon  
(stable and clean)*



「月待図」『江戸年中行事図聚』より

# Technical Challenge of Thermionic Injectors to $1 \pi \text{ mm-mrad}$ emittance

this picture is copied from LCLS "Conceptual Design Report", SLAC-R-593 UC-414



3-2001  
8560A&1

Normalized rms transverse emittance measured by the leading thermoionic (SLAC, BOEING) and rf photocathode injectors [16]. All data are for bunched beams with approximately 1 nC of charge.

- *Eliminating control grid from cathode.*
- *Smaller size cathode, from 8 mm to 3 mm diameter.*
- *Higher gun voltage, 150 kV to 500 kV.*
- *Using single crystal CeB6 cathode.*

# SCSS & X-ray FEL Beam Parameter

X-ray FEL

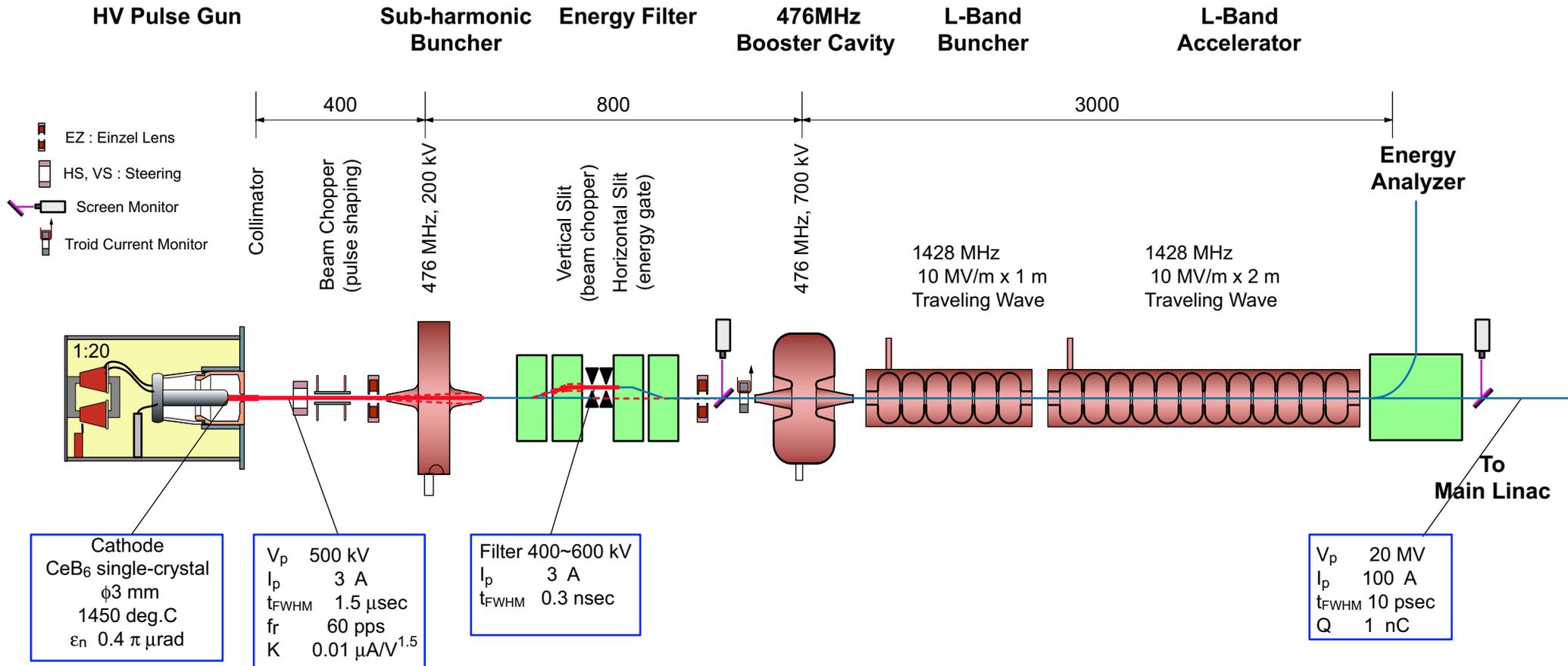
at undulator section

		SCSS	X-ray FEL	
Beam Energy	$E$	1.0	6.0	GeV
X-ray Wavelength	$\lambda$	3.6	0.1	nm
Beam Emittance	$\varepsilon_n$	2	1	$\pi$ mm.mrad
Bunch Length	$\Delta z$	150	75	$\mu$ m
	FWHM	0.5	0.25	psec
Transverse Beam Size	$\sigma_{x,y}$	100	25	$\mu$ m
Peak Current	$I_p$	2	4	kA
Charge per bunch	$q$	1	1	nC
Undulator Parameter	$\lambda_u$	15	15	mm
	$K$	1.3	1.3	
	Length	$L$	22.5	30
FEL Saturation Length	$L_{sat}$	20	30	m

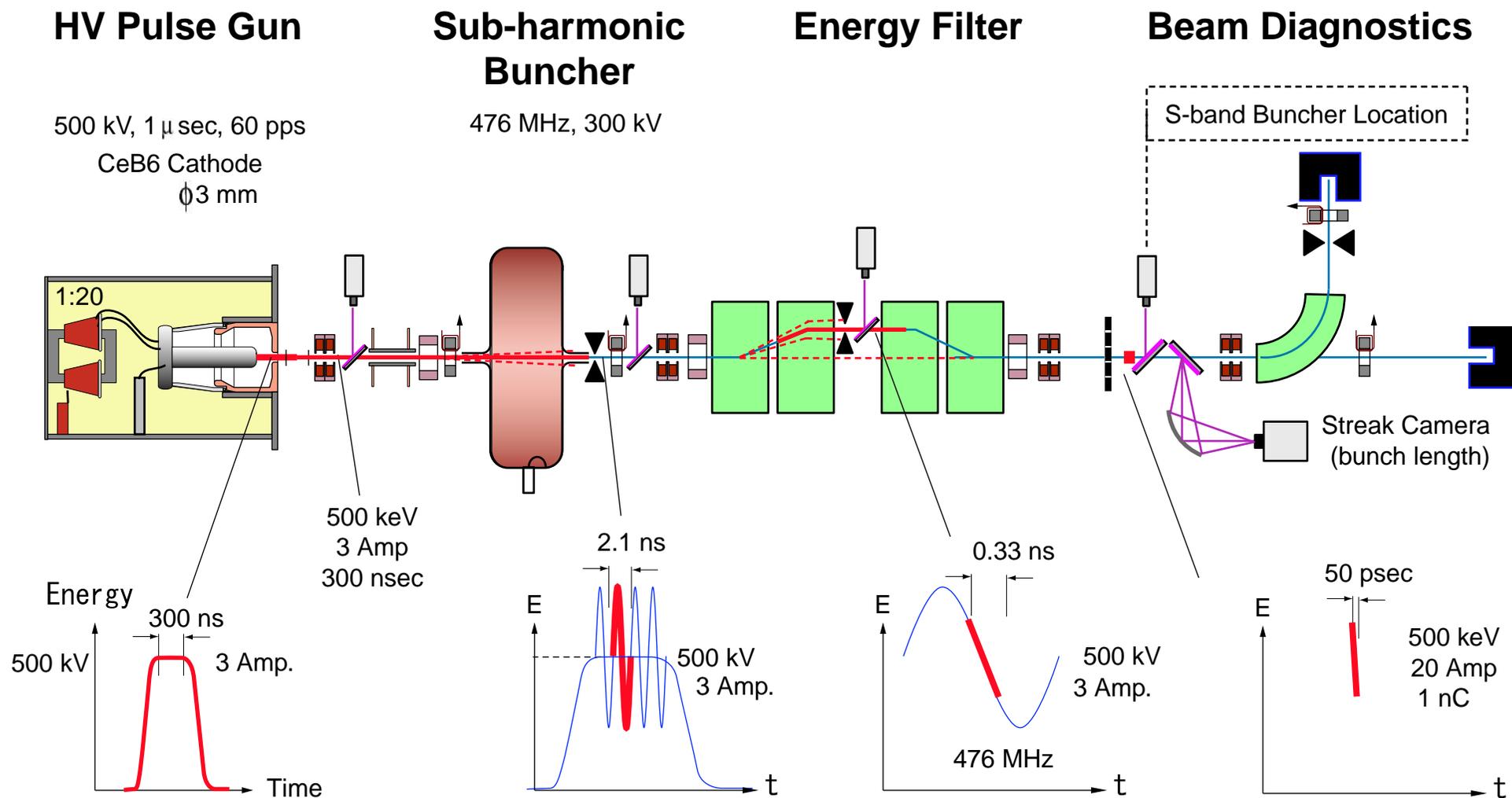
# Low Emittance Injector for SASE-FEL

2002 July

X-ray FEL



# Beam Formation in 500 kV Electron Gun



# CeB<sub>6</sub> Cathode & Heater Assembly

X-ray FEL



- CeB<sub>6</sub> Cathode 3 mm Diameter
- Emittance 0.4  $\pi$ .mm.mrad (thermal emittance, theoretical)
- Beam Current 3 Amp. at 1450 deg.C (using graphite heater)
- Current Density > 40 A/cm<sup>2</sup>

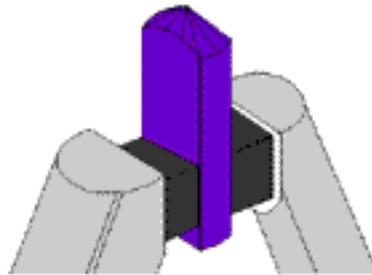




## LaB<sub>6</sub> / Cebix Cathodes



FEI Beam Technology Division is the leading manufacturer of LaB<sub>6</sub> and CeBix cathodes and maintains a close partnership with each electron beam instrument manufacturer to ensure our cathodes meet the requirements of each piece of equipment and application. Included with each cathode is a detailed handling and operating guideline for the specific instrument. FEI maintains a large inventory of cathodes, and orders are typically shipped the next day.

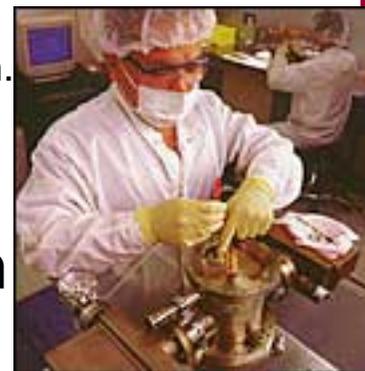


FEI **Mini Vogel Mount (MVM)** cathodes provide stable, high-brightness, long-lifetime operation for all major electron beam instruments. View FEI's standard parts for all major electron beam instruments in use worldwide on our [Applications Chart](#).

The superiority of the MVM comes from the simplicity of its design.

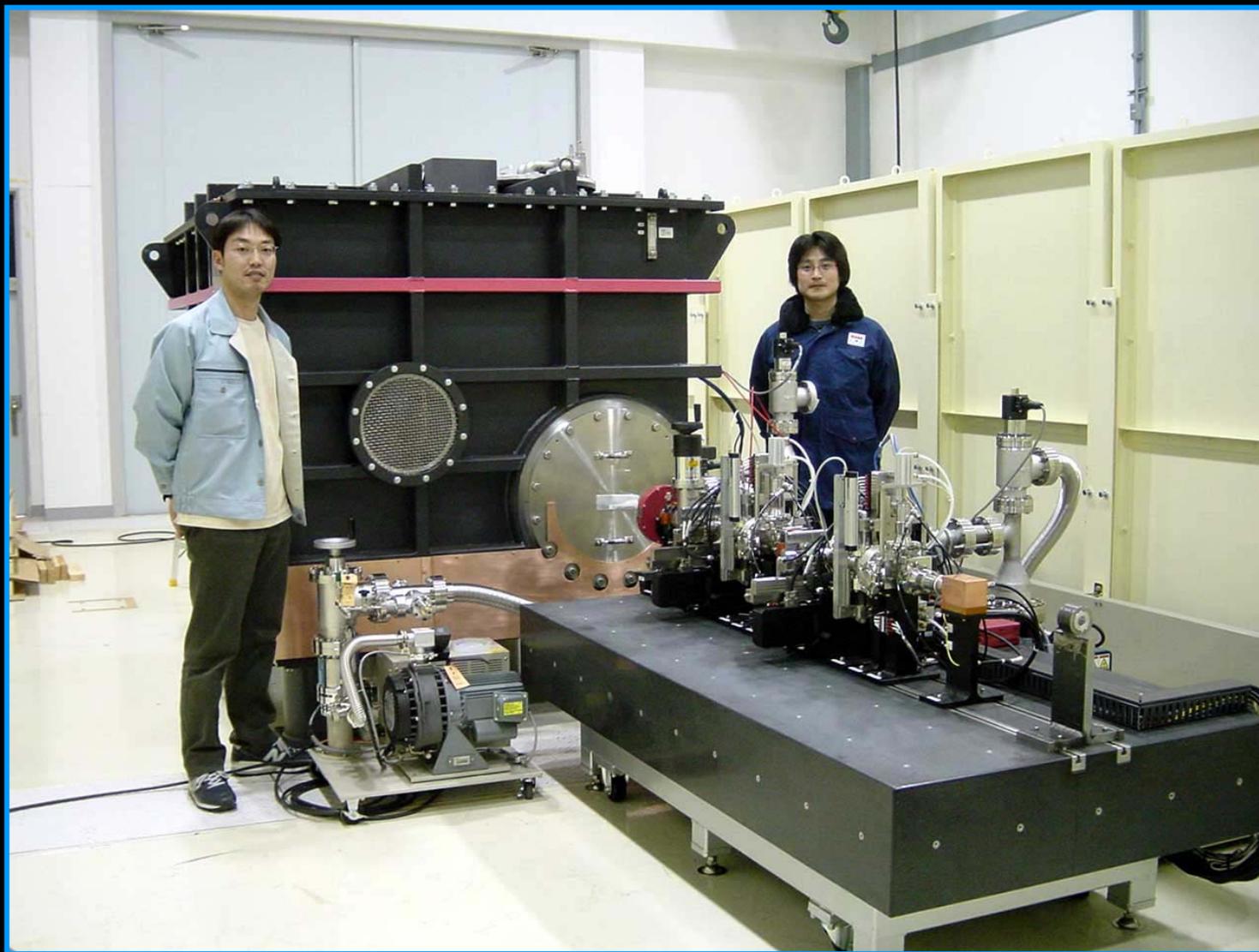


<http://www.feibeamtech.com>

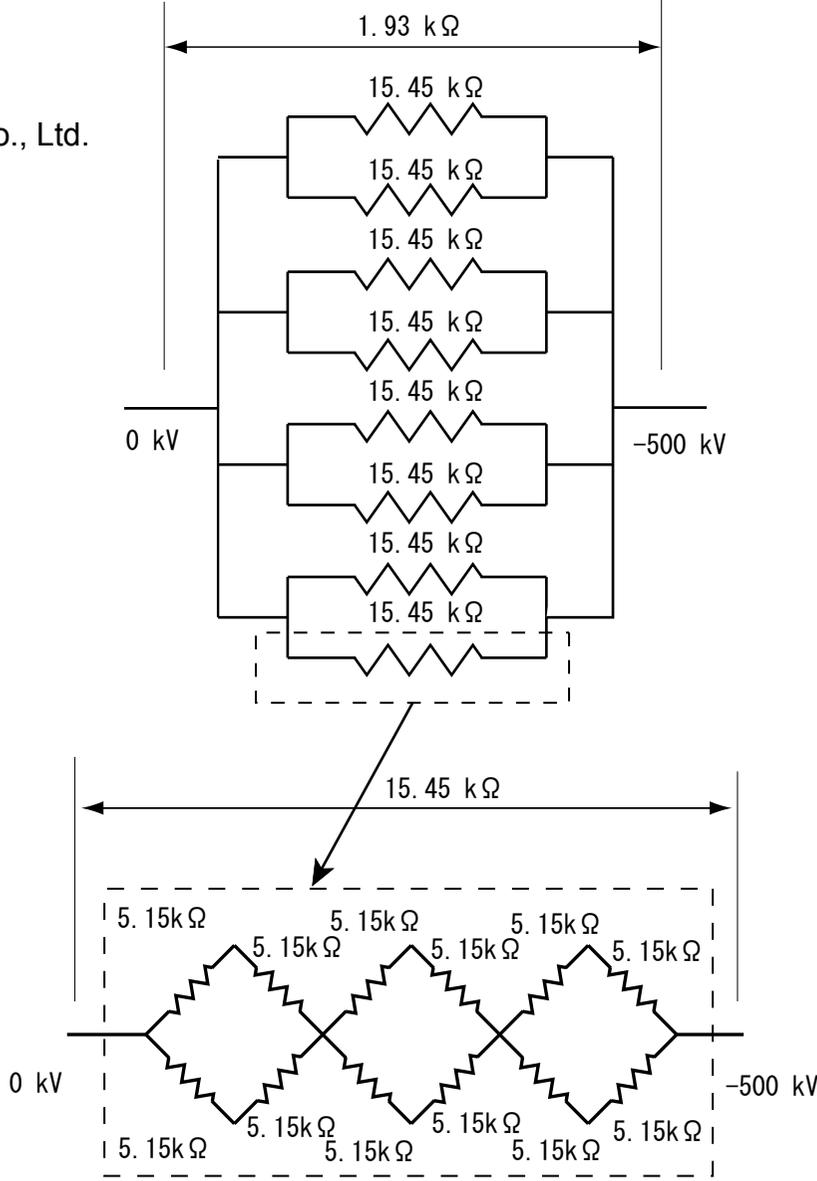


# Low-Emittance 500kV Electron Gun R&D

X-ray FEL

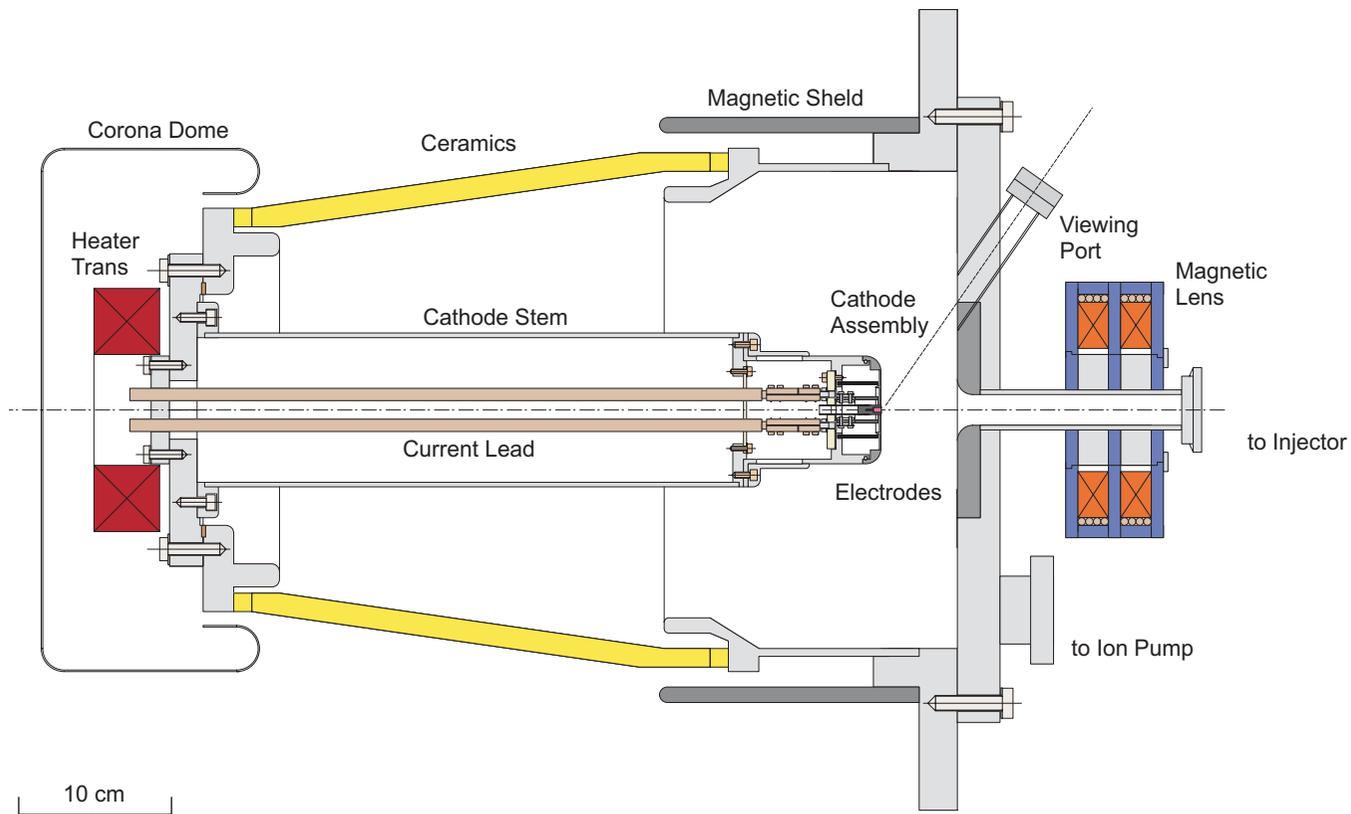


# 500kV Dummy Load



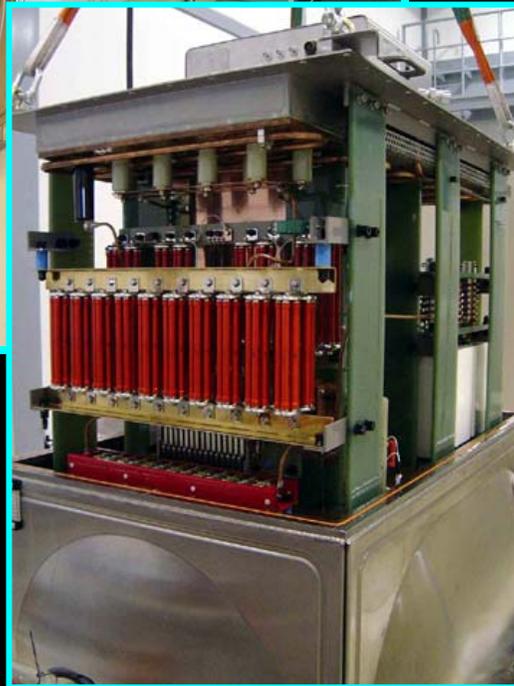
Pulse High Voltage	500 kV
Peak Current	259 A
Pulse Width	1.6 micro-s
Dummy Load Impedance	1.93 kΩ
Average Power	12.2 kW

# 500kV Electron Gun



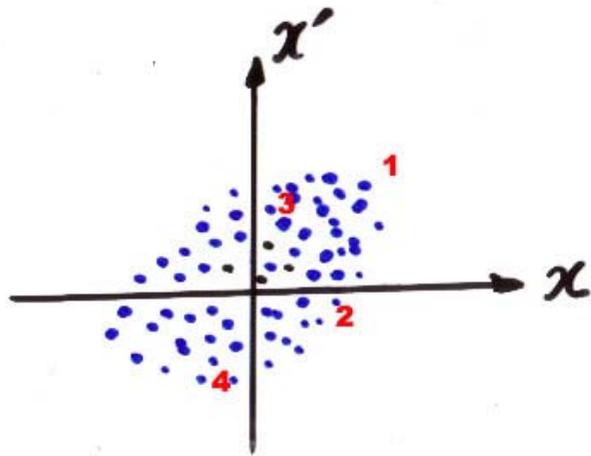
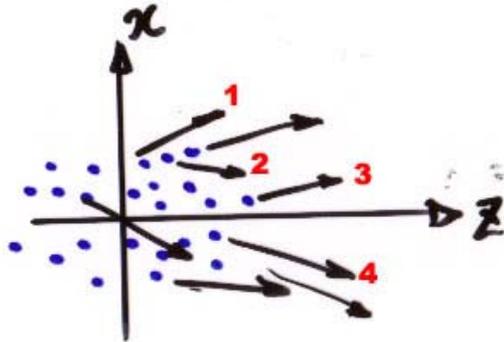
# Closed Compact Modulator for C-band Accelerator

X-ray FEL



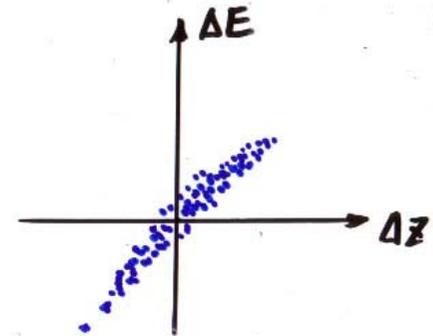
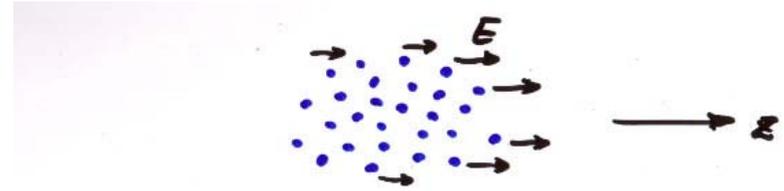
- Compact.  
W 1.7 m x D 1.2 m x H 1m.
- Good EMI shield.
- Better cooling for HV component.
- Eliminating cooling air fan.
- No dust accumulation due to high voltage in air.
- No environmental effects: moisture and temperature variation.

# Transverse & Longitudinal Emittance



Transverse  $(x, x')$ ,  $(y, y')$

**R&D Challenge**



$$\gamma \ell_z = \gamma \sigma_z^2 \sigma_\delta^2 - \langle z\delta \rangle^2$$

no correlation  $\langle z\delta \rangle = 0$

$$= \gamma \sigma_z \cdot \sigma_\delta = \sigma_z \cdot \frac{\sigma_E}{mc^2}$$

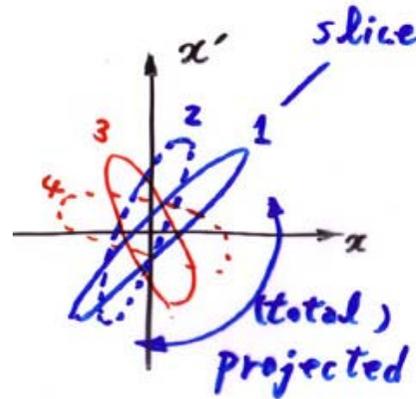
Longitudinal  $(z, E)$

**Easily obtainable.**

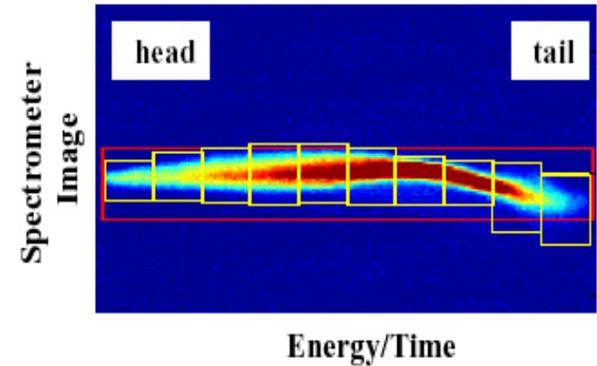
# All Three Emittance have to be Small

- **Slice Emittance**

$\mathcal{E}_{slice}$  (un co-related)  
*thermal emittance*  
 define SASE FEL gain



courtesy D. Dowell

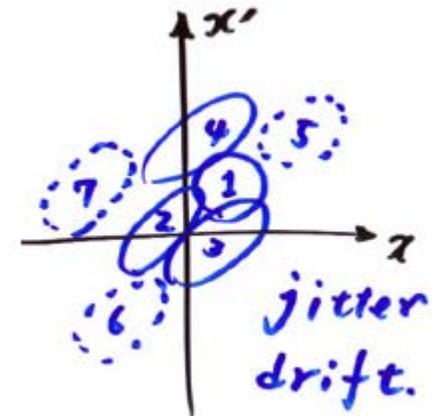


- **Projected Emittance**  
 (bunch emittance)

$\mathcal{E}_{proj}$  (co-related)  
*measure of wakefield, space-charge, etc.*

- **Multiple Bunch Emittance**

$\mathcal{E}_{multi}$  (include jitter & drift)  
*stability measure, affect X-ray beam delivery efficiency*

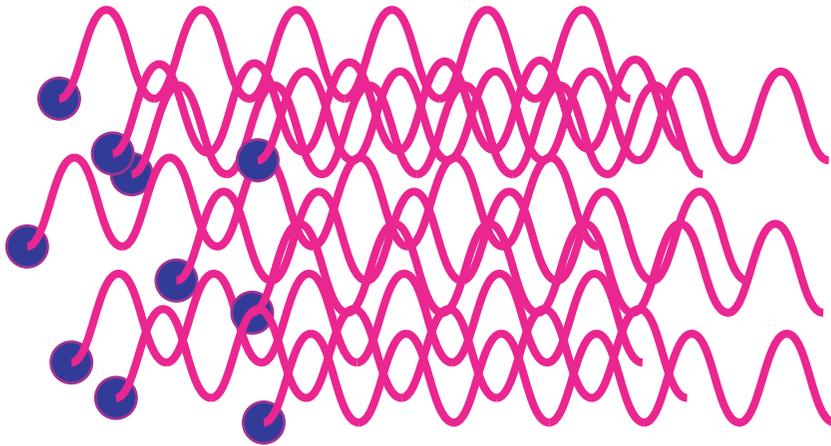


$\mathcal{E}_{slice}$	<	$\mathcal{E}_{proj}$	<	$\mathcal{E}_{multi}$
(fs)		(ps)		(sec, minute,)

# From SR to FEL

SR or ERL

Spontaneous Radiation



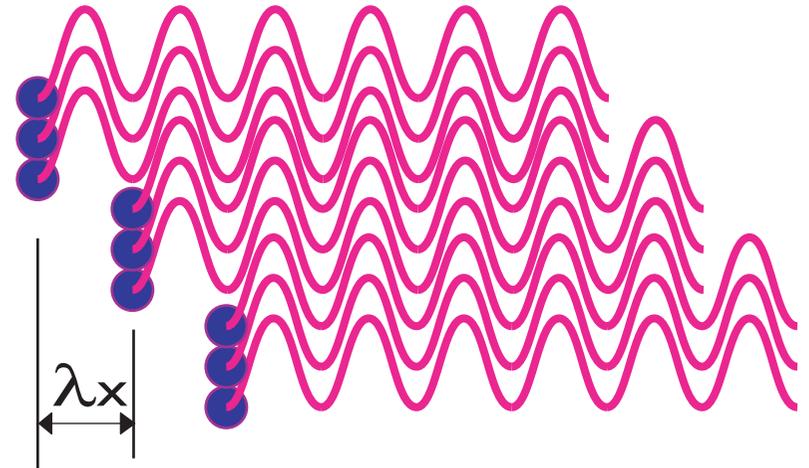
***N***-electrons  
random distribution

$$E_{spt} \sim \sqrt{N} E_1$$

$$P_{spt} \sim N P_1$$

FEL: Free Electron Laser

Coherent Radiation

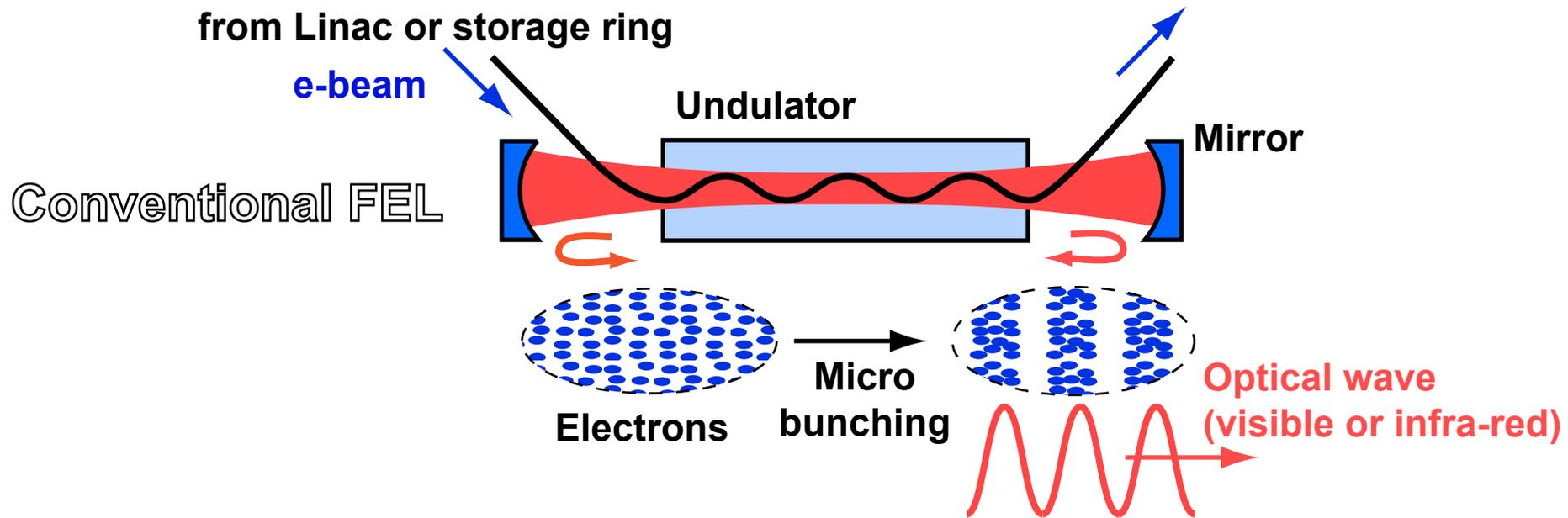


***N***-electrons  
micro-bunched

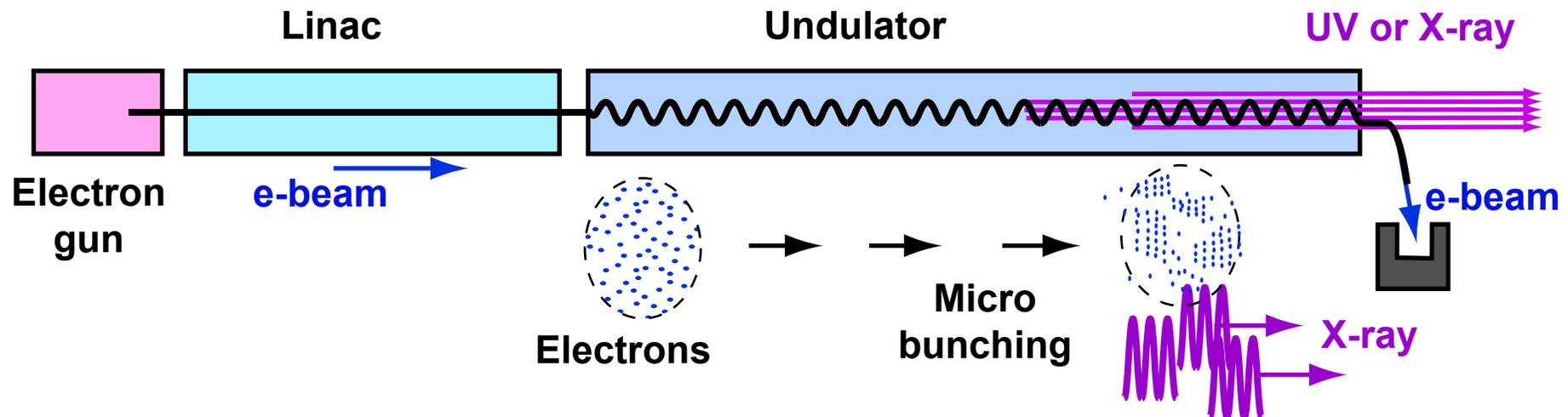
$$E_{coherent} \sim N E_1$$

$$P_{coherent} \sim N^2 P_1$$

# From Cavity Type FEL to SASE-FEL



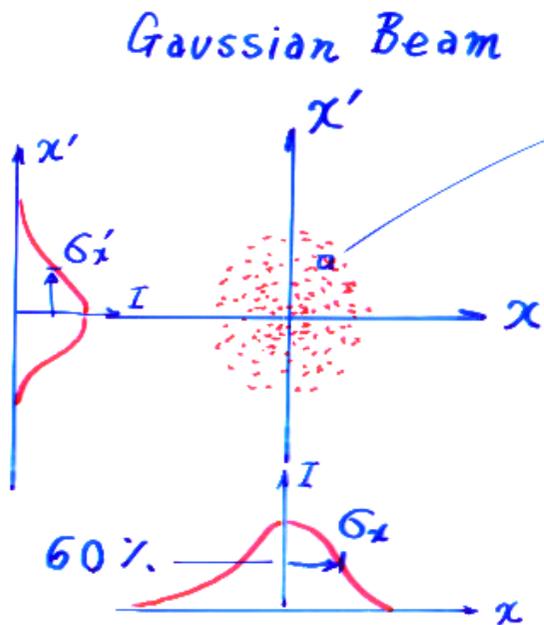
## SASE-FEL



# Emittance Definition

We use Normalized r.m.s. emittance

$$\begin{aligned} \epsilon_n &= \gamma\beta \sqrt{\langle x^2 \rangle \langle x'^2 \rangle - \langle xx' \rangle^2} \\ &= \frac{1}{m_0 c} \sqrt{\langle x^2 \rangle \langle P_x^2 \rangle - \langle x P_x \rangle^2} \end{aligned}$$



$$dN = \frac{Ne}{2\pi\sigma_x\sigma_{x'}} e^{-\left(\frac{x^2}{2\sigma_x^2} + \frac{x'^2}{2\sigma_{x'}^2}\right)} dx dx'$$

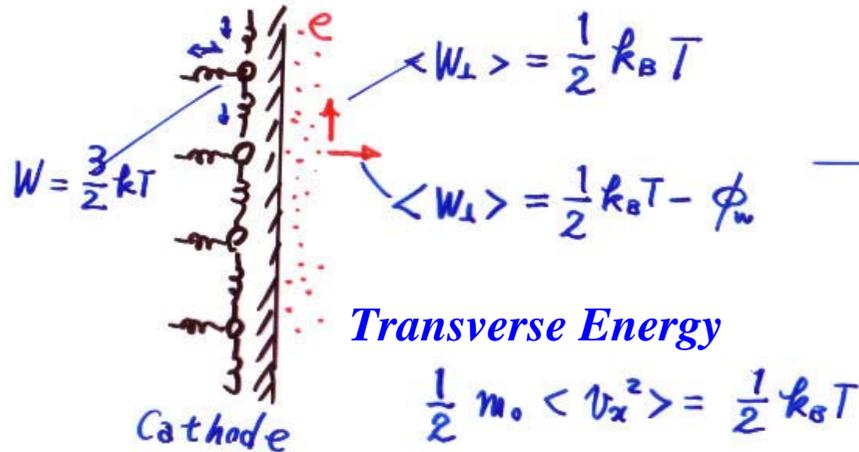
$\langle xx' \rangle^2 = 0$

$$\langle x^2 \rangle = \frac{\iint x^2 dN}{\iint dN} = \dots = \sigma_x^2$$

$$\langle x'^2 \rangle = \frac{\iint x'^2 dN}{\iint dN} = \dots = \sigma_{x'}^2$$

$$\epsilon_n = \gamma\beta \sigma_x \sigma_{x'} \iff \text{Area: } A = \pi \sigma_x \sigma_{x'}$$

# Thermal Emittance (lower limit)



Transverse Energy

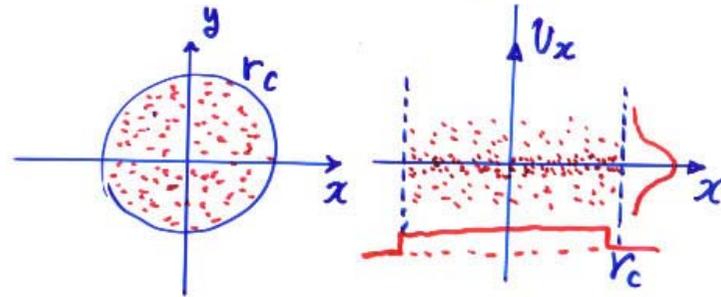
$$\frac{1}{2} m_0 \langle v_x^2 \rangle = \frac{1}{2} k_B T$$

Transverse Momentum

$$\begin{aligned} \sqrt{\langle p_x^2 \rangle} &= \gamma m_0 \sqrt{\langle v_x^2 \rangle} \\ &= \gamma \sqrt{m_0 k_B T} \end{aligned}$$

Thermal Emittance

$$\begin{aligned} \epsilon_{xN} &= \frac{1}{m_0 c} \sqrt{\langle x^2 \rangle \langle p_x^2 \rangle - \langle x p_x \rangle^2} \\ &= \frac{\gamma r_c}{2} \sqrt{\frac{k_B T}{m_0 c^2}} \end{aligned}$$



r.m.s Cathode Size

$$\begin{aligned} \langle x^2 \rangle &= \frac{\int x^2 dS}{\int dS} \\ &= \frac{r_c^2}{4} \\ \sqrt{\langle x^2 \rangle} &= r_c / 2 \end{aligned}$$

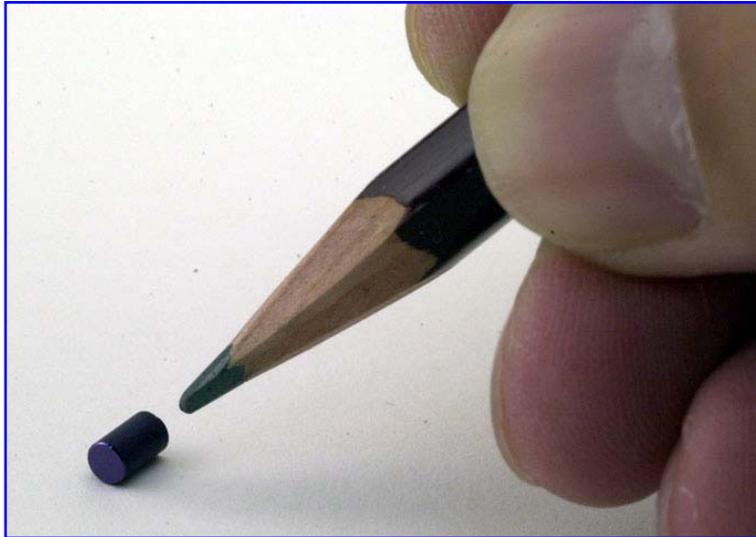
Boltzmann constant

$$k_B = 1.38 \times 10^{-23} \text{ (J/deg)}$$

# Use Small Size Cathode

... First Strategy for smaller thermal emittance

- *Thermionic cathode*



**3mm** diameter cathode (CeB6)  
is used in a low emittance injector.  
(SCSS SPring-8/RIKEN)

Operating Temperature 1450°C

$$w_e = \frac{3}{2} k_B T = 223 \text{ meV}$$

Thermal Emittance

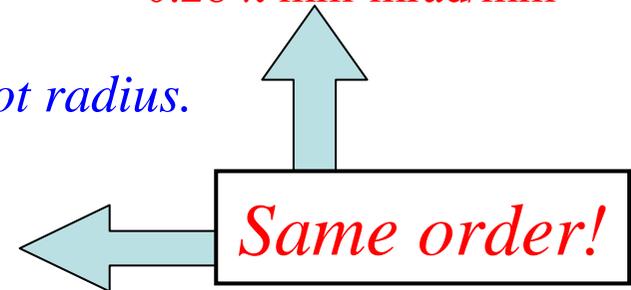
$$\varepsilon_{xN} = \frac{\gamma r_c}{2} \sqrt{\frac{k_B T}{m_0 c^2}} = 0.4 \pi \text{ mm-mrad}$$

0.28  $\pi$  mm-mrad/mm

- *RF photo-cathode injector.*

Today's RF photo injectors use  $\sim 1$  mm spot radius.

$$\varepsilon_{xN} = \frac{\gamma r_c}{2} \sqrt{\frac{k_B T_e}{m_0 c^2}} = 0.35 \pi \text{ mm-mrad}$$



$T_e$  is “measured” effective electron temperature of copper cathode using 266 nm laser (ref. 2).  $k T_e = 0.27 \text{ eV}$  (2360°C).

# How “meV” can be important on “GeV” beam?

- *Transverse energy along x-direction.*

$$\frac{1}{2} m_e \langle v_x^2 \rangle = \frac{1}{2} k_B T = 74 \text{ meV}$$

*10<sup>-12</sup> difference!*

←→ **GeV**  
*wrong feeling*

- *Transverse momentum*

$$p_x = \gamma m_0 v_x = \beta \frac{\gamma m_0 c^2}{c} = \beta \frac{E}{c} = 274 \text{ eV/c}$$

*Momentum ratio*

*~10<sup>-7</sup>*

←→ **GeV/c**

$$w_k = 74 \text{ meV}, v = 161 \text{ km/sec}, \quad \varepsilon_{xN} = \gamma \sigma_x \sigma_x' = 1\pi \text{ mm-mrad}$$

$$\beta = 0.00054, \gamma = 1$$

$$\gamma = 2 \times 10^3, \sigma_x = 1 \text{ mm}, \sigma_x' = 5 \times 10^{-7}$$

***Emittance is a measure of the transverse momentum.***

## Use Small Size Cathode....2

to reduce transverse kick by rf acceleration field.

*TM<sub>0</sub>-mode axial-symmetric field in the accelerating rf cavity.*

Using Maxwell's equation,  $\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \rightarrow \frac{1}{r} \frac{\partial}{\partial r} (r H_\theta) = j\omega \epsilon_0 E_z$

Integrating near the axis,  $H_\theta = \frac{r}{2} \cdot j\omega \epsilon_0 E_{z0} \propto \frac{r}{\lambda}$

From the Gauss's law,

$$\nabla \cdot \mathbf{E} = 0 \rightarrow \frac{1}{r} \frac{\partial}{\partial r} (r E_r) + \frac{\partial E_z}{\partial z} = 0 \rightarrow E_r = -\frac{r}{2} \left( \frac{\partial E_z}{\partial z} \right)_{r=0} \propto \frac{r}{\lambda}$$

Transverse Kick

$$P_r = q(E_r + u_z \times B_\theta) \propto \frac{r}{\lambda}$$

For low emittance,

- smaller cathode size and
- lower rf frequency is desirable.

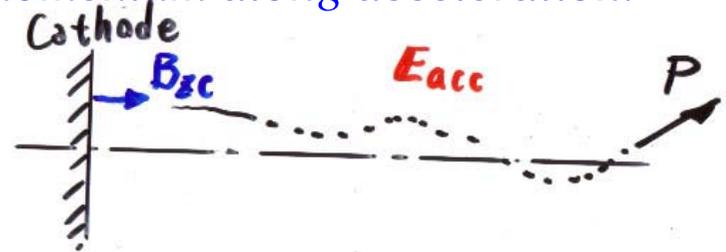
use 1 mm laser spot than 10 mm spot.

use 476 MHz, 714 MHz or 1.3 GHz, than S-band, C-band or X-band. Note on available field gradient.

# Magnetic Field on the Cathode Must be Zero

- Conservation of transverse generalized momentum along acceleration.

$$\mathbf{p}_{\perp} = \gamma m_0 \mathbf{u}_{\perp} + q \mathbf{A}_{\perp} = \text{constant}$$



- Initial condition on the cathode.

If magnetic field is not zero,

$$\mathbf{p}_{\perp}^0 = q \mathbf{A}_{\perp}^0 \quad u_{\perp} = 0, \quad A_{\perp} \neq 0$$

- At high energy after acceleration, the transverse momentum in free space is given by  $\mathbf{p}_{\perp} = \gamma m_0 \mathbf{u}_{\perp} = \mathbf{p}_{\perp}^0 = q \mathbf{A}_{\perp}^0$  Bush's Theorem

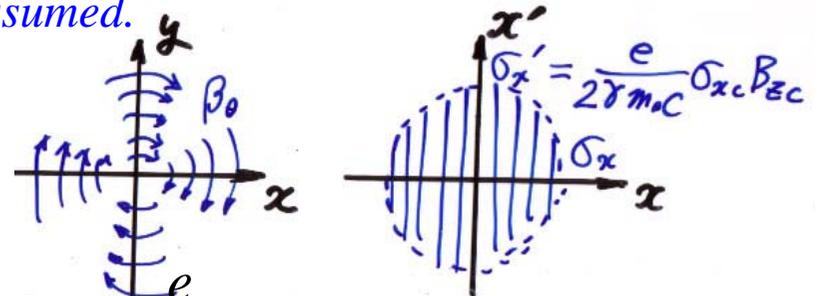
- Axial symmetric case. Field leakage on the cathode :  $B_{zc}$   $\mathbf{A}_{\theta} = \frac{1}{2} r_c B_{zc}$

$$\varepsilon_{xN} = \frac{e}{2m_0 c} \sigma_{xc}^2 B_{zc}$$

*laminar flow is assumed.*

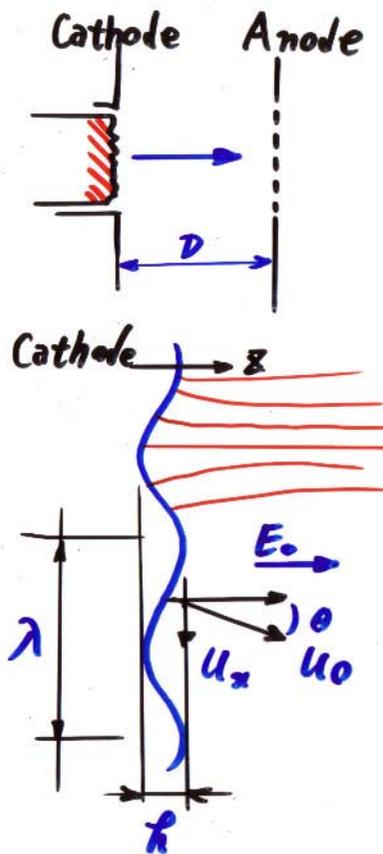
$$\sigma_x = \sigma_{xc} = 1 \text{ mm}, B_{zc} = 100 \text{ Gauss}$$

$$\varepsilon_{xN} = 2.9 \pi \text{ mm-mrad}$$



$$\gamma \beta_{\theta} = \frac{e}{2m_0 c} r_c B_{zc}$$

# Emittance Dilution due to Surface Roughness of Cathode



Simple roughness model

Simplified surface curve:  $z = \frac{h}{2} \cos\left(\frac{2\pi z}{\lambda}\right)$

Tangential velocity:  $u_x = \frac{dz}{dx} \cdot u_0 = \frac{\pi h}{\lambda} \cdot u_0 \sin\left(\frac{2\pi z}{\lambda}\right)$

Average velocity:  $\langle u_x^2 \rangle = \frac{1}{2} \left(\frac{\pi h}{\lambda}\right)^2 \langle u_0^2 \rangle = \frac{ehE_0}{2m_0} \left(\frac{\pi h}{\lambda}\right)^2$

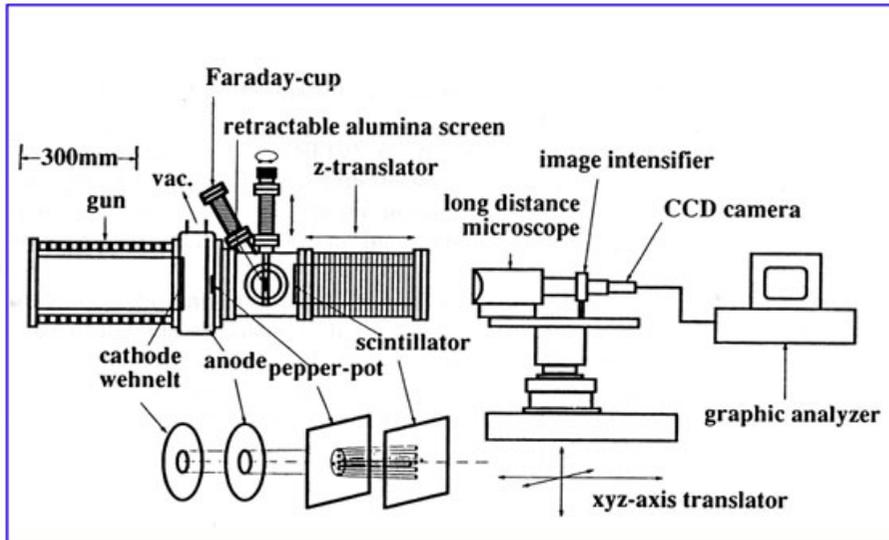
Emittance:  $\varepsilon_{nN} = \frac{1}{m_0 c} \sqrt{\langle P_x^2 \rangle} \cdot \sigma_x = \frac{\gamma r_c}{2} \sqrt{\frac{eE_0}{2m_0 c^2}} \cdot \frac{\pi h^{1.5}}{\lambda}$

Required Smoothness

Accelerating Field: $E_0$ (MV/m)	1	10	<b>100</b>
Roughness: $h$ ( $\mu\text{m}$ )	3	1	<b>0.3</b>

$\varepsilon_{nN} < 0.3 \pi \text{ mm-mrad}$ ,  $\lambda = 10 \mu\text{m}$ ,  $r_c = 1 \text{ mm}$

# Emittance Measurement on Thermionic Cathode



H. Kobayashi, et.al, "Emittance Measurement for High Brightness Electron Guns", Proc. 1992 Linear Accelerator Conference, Ottawa, Aug.23-28, 1992

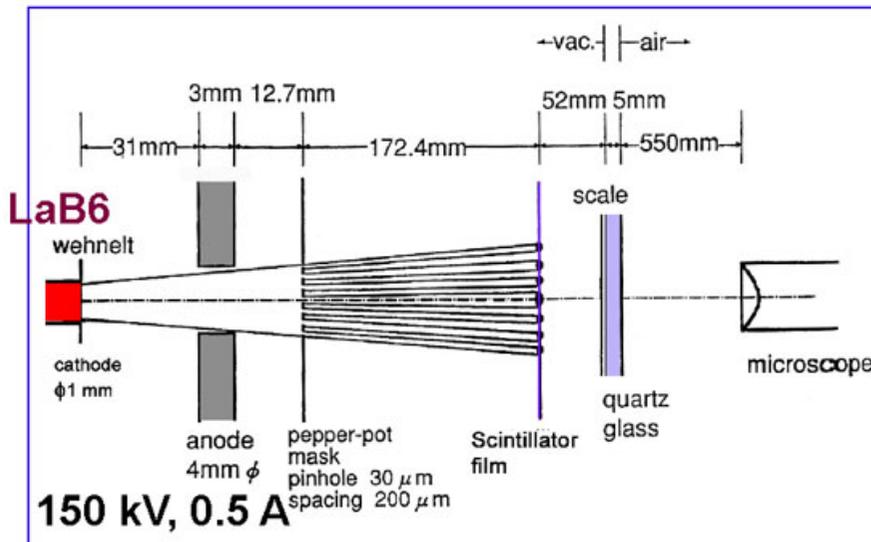
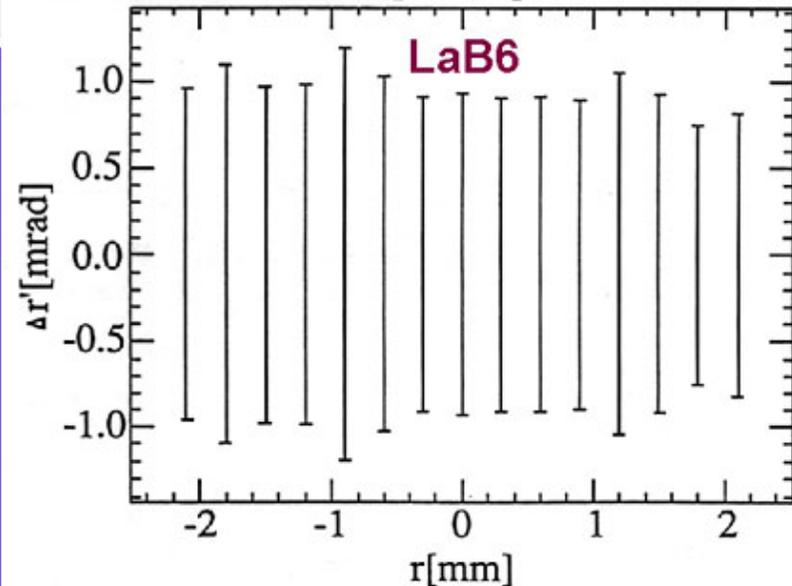
Y. Yamazaki, T. Kurihara, H. Kobayashi, I Sato and A. Asami @ KEK "High-precision pepper-pot technique for a low-emittance electron beam" Nucl. Instrum. and Methods in Phys. Research A322(1992) 139-145

Measured Emittance =  $0.34 \pi \text{ mm-mrad}$

Theoretical Emittance =  $0.13 \pi \text{ mm-mrad}$

resolution limited.

Area(95%) of the phase space  $1.37 [ \pi \text{ mm.mrad} ]$



# Beam Quality from Three Type Cathode

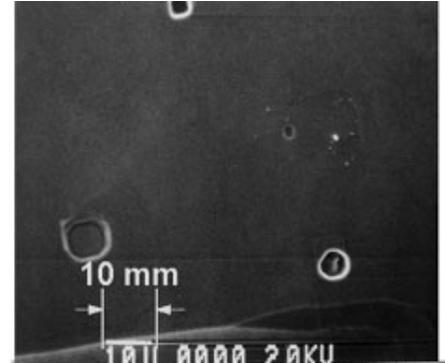
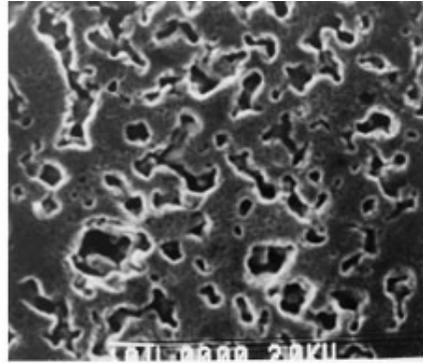
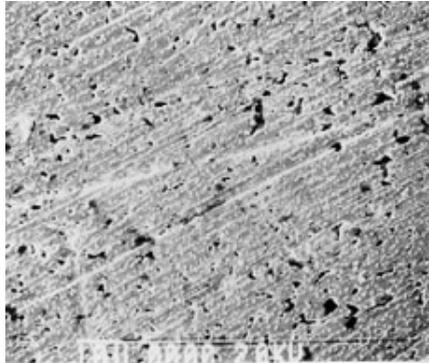
Hitoshi Kobayashi KEK

Dispenser Cathode

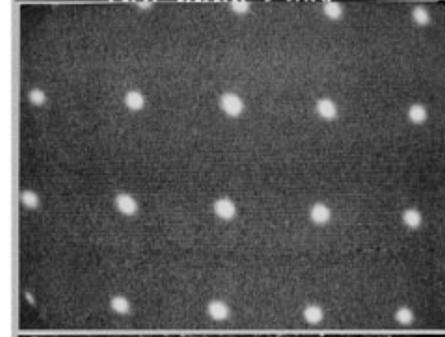
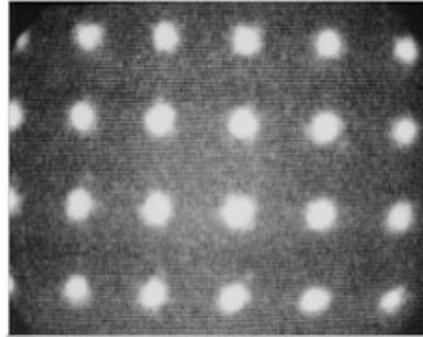
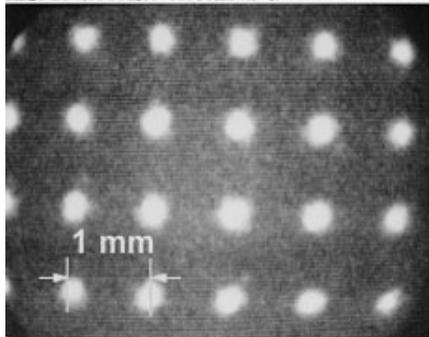
LaB6 Poly-crystal

LaB6 Single Crystal

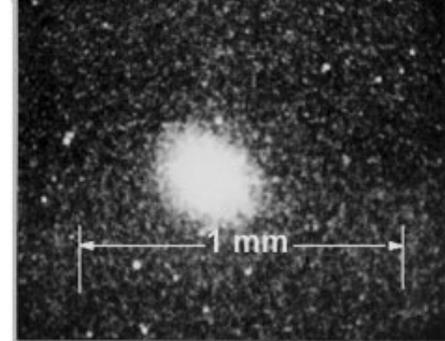
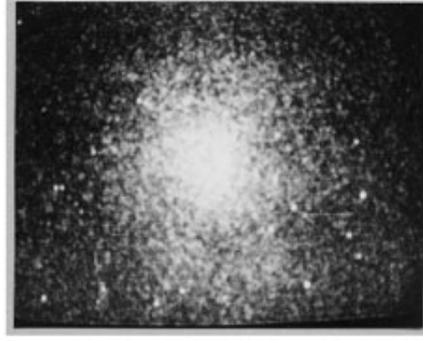
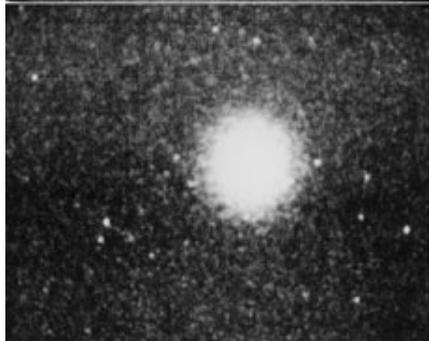
SEM Image



Pepper-pot Image

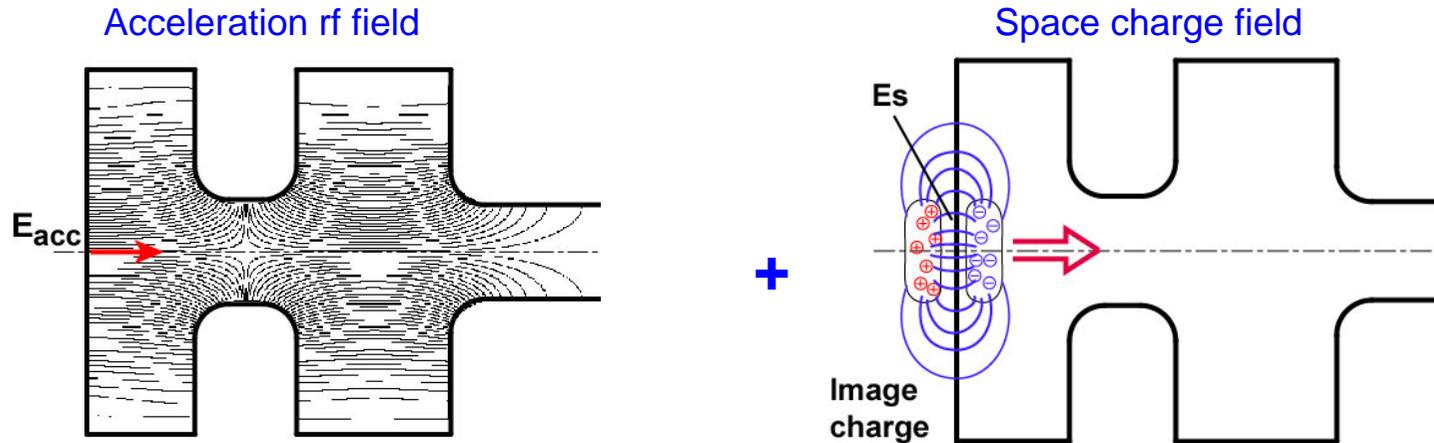


Magnified Spot



# Space Charge Effect...in rf gun

- Maxwell's equation is linear, thus any field can be analyzed by linear superposition. (no  $E^2$ ,  $H^2$ ,  $EH$  terms)
- RF-acceleration field in the cavity + space charge.



Maximum space charge field at cathode

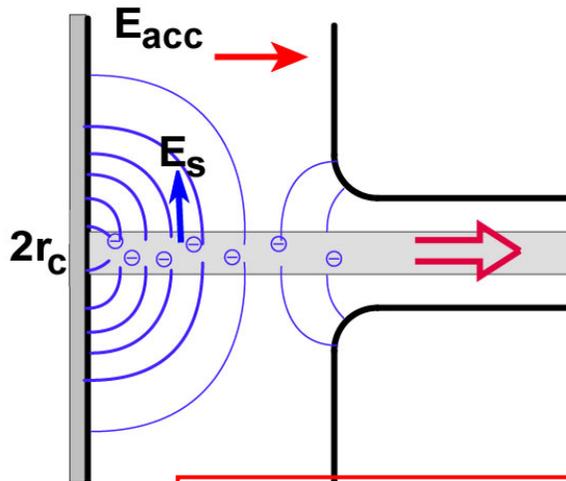
$$E_{spc.max} = \frac{Q}{\pi\epsilon_0 r_c^2}$$

$$Q = 1 \text{ nC}, r_c = 1 \text{ mm} \rightarrow E_{spc.max} = 36 \text{ MV/m}$$

*The non-linear space charge field runs with bunch.*

*The tail electron feels lower acceleration field. To overcome this energy difference, we need very high field gradient: 100 MV/m ~ 140 MV/m.*

# Space Charge Effect...in HV gun.



- DC or long pulse beam.
- Beam becomes a column shape.
- The nonlinear space charge field is localized just near to the cathode (distance  $\sim r_c$ ), thus space-charge field contribution to emittance is small.

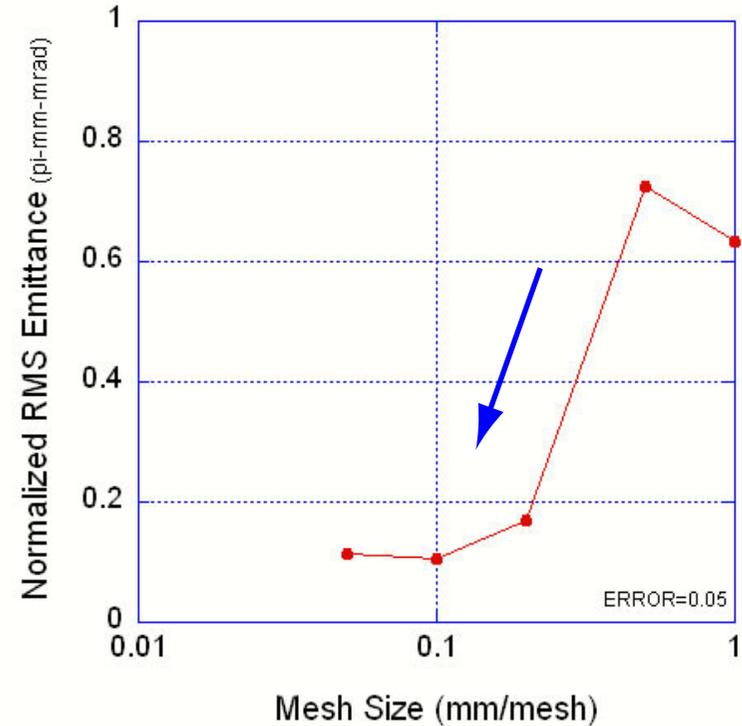
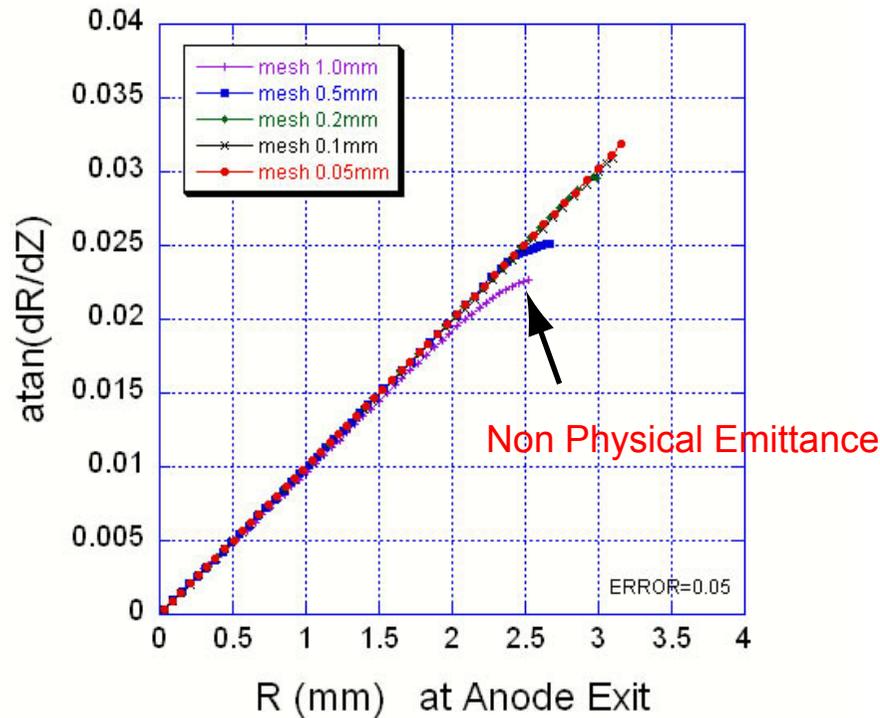
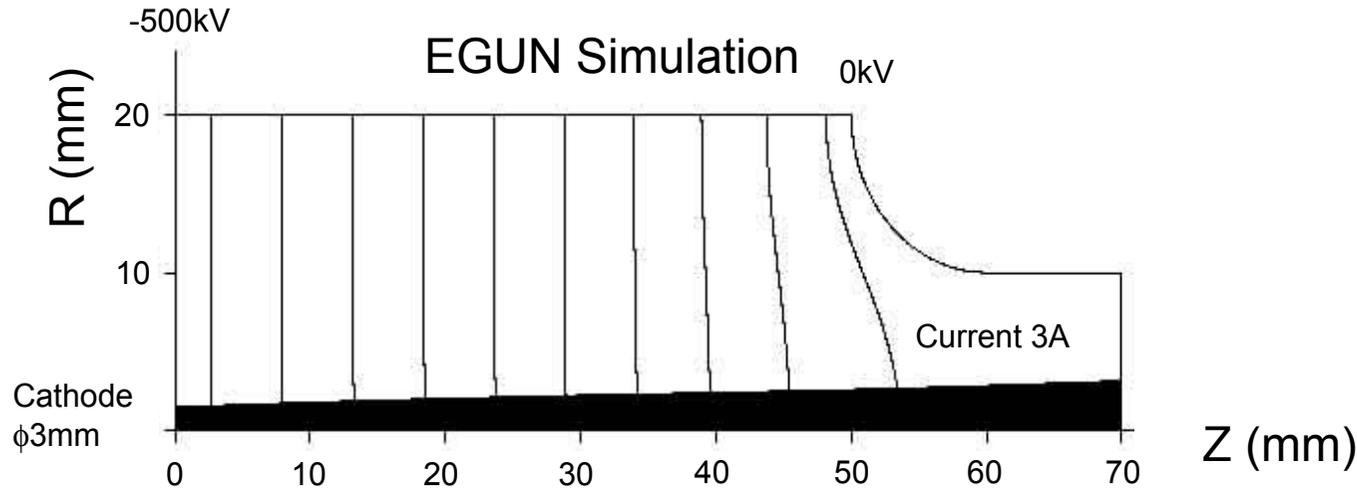
*Specially in “temperature limited cathode condition”, beam trajectory becomes straight, and space charge field is very linear, thus the emittance becomes negligibly small.*

$$E_s \simeq E_r = \frac{I_z}{\pi \epsilon_0 R^2 u_z} \cdot r$$

*To do this, we apply 500 kV on 5 cm gap for 3 A beam in SCSS gun.*

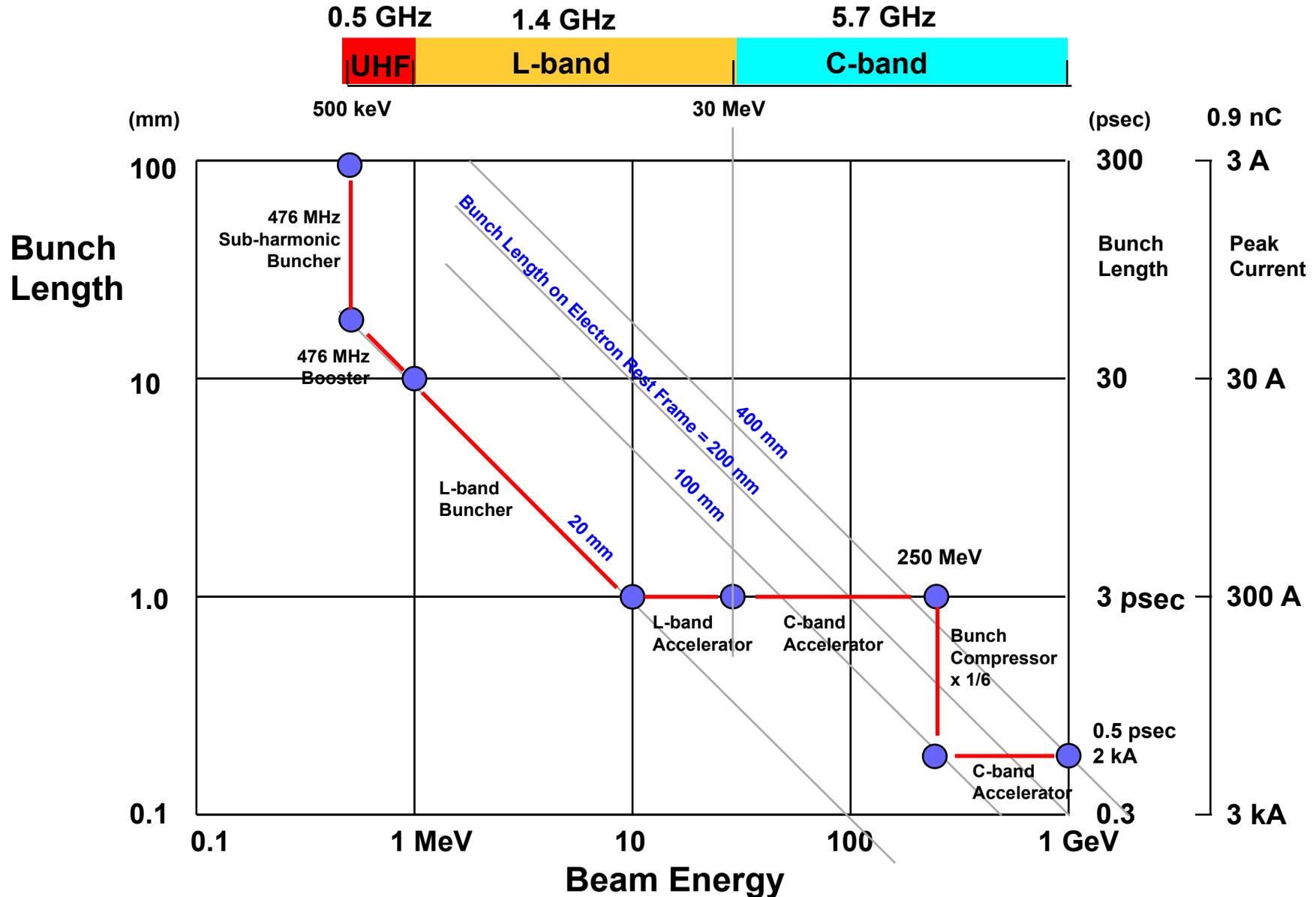
*E-Gun simulation says:  $\epsilon_n < 0.1 \pi$  mm-mrad thermal emittance is not included.*

# Emittance at Parallel Electrodes



# Bunch Compression after the Gun

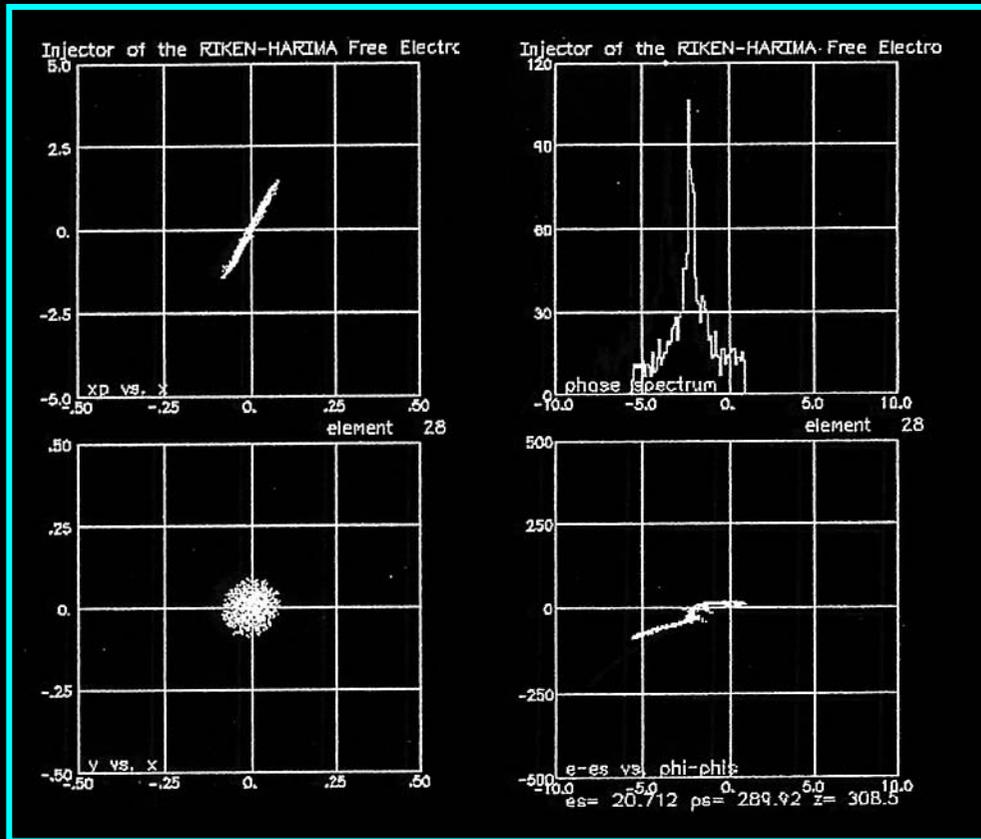
Maintain bunch-length on the electron rest-frame being constant.



# Injector Beam Simulation Result

At exit of L-band accelerator.

By K. Sawada, SUMITOMO Heavy Ind.



- $Q$  0.92 nC
- $\epsilon_n$  1.9  $\pi$ mm.mrad
- $\Delta t$  12.7 ps
- $\Delta\phi$  6.5 deg.

Parmella Simulation

# Dark Current from the Gun

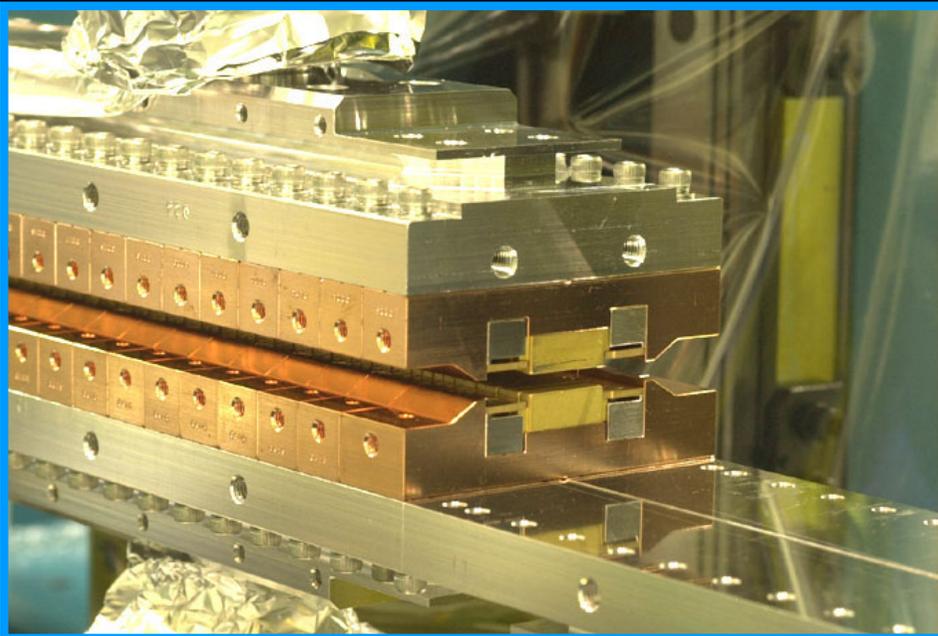
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- *The dark current will potentially **kill** the permanent magnet in the undulator.*
- *The dark current will create an **off-set** information on the beam-position monitor.*
- *Dark current generated from the accelerating structure can be easily cleaned by Q-magnet focusing because of the energy difference.*
- *But, the dark current from the gun will remain with core beam.*
- ***RF-gun is potential source of the dark current because of its high field 100 MV/m.***
- *SCSS high voltage gun runs at 10 MV/m, dark current generation will be small.*

# SCSS Undulator In-Vacuum Type

X-ray FEL

- Segment length 4.5 m
- $N=300/\text{segment}$ ,  $\lambda_u=15\text{mm}$
- Mechanical minimum gap = 2mm
- Nominal gap  $\sim 3.5\text{mm}$  ( $K\sim 1.3$ )
- 45-deg. tilted Halbach type
- More compact than ordinary ones



First prototype model arrived, Sep. 2002

# *Photocathode rf-injector v.s. Thermionic HV gun injector*

	<i>Photocathode rf-injector ( Cu, Mg )</i>	<i>Thermionic-cathode injector (LaB6, CeB6)</i>
<i>Thermal Emittance</i>	<i>0.4 <math>\pi</math> mm-mrad/mm</i>	<i>0.3 <math>\pi</math> mm-mrad/mm</i>
<i>Emittance at 20~100 MeV</i>	<i>best record 2 <math>\pi</math> mm-mrad/nC/3psec on the simulation &lt; 1 is possible.</i>	<i>2 <math>\pi</math> mm-mrad/nC/3psec (simulation only, no experimental data)</i>
<i>Cathode Lifetime</i>	<i>OK for copper cathode, but QE is low (&lt; QE10<sup>-4</sup>). Mg cathode has high QE (0.2%), but needs cleaning once per week.</i>	<i>OK &gt;3000 hour lifetime is reported in electron microscope application.</i>
<i>Surface Roughness</i>	<i>h &lt; 300 nm/10<math>\mu</math>m Need to maintain a flat surface under HV processing and discharge at 100 MV/m. It can limit actual lifetime of the cathode.</i>	<i>OK LaB6 and CeB6 keep is flat surface in nm range after usage at high temperature (auto cleaning by evaporation of cathode material)</i>

# *Photocathode rf-injector v.s. Thermionic HV gun injector*

Continued

	<i>Photocathode rf-injector ( Cu, Mg )</i>	<i>Thermionic-cathode injector (LaB6, CeB6)</i>
<i>Timing</i>	<i>Laser light defines timing of the electron emission from the cathode, and it also carries timing to the seeding, and also pump-probe experiment in sub-pico-sec or femto-sec range.</i>	<i>Need careful design on phase-stable rf drive line. Timing of electron bunching is defined by phase in the rf-power system for the buncher, which is basically same order as the mode-locking in pulsed laser.</i>
<i>Cost &amp; Manpower</i>	<i>To maintain the laser system, we need manpower. We have to pay for the laser system and accessories (clean room, temperature control system).</i>	<i>More costly: The high voltage gun and the buncher system will cost more than rf-gun cavity. But it will be comparable to total system cost of photo-injector including laser system and its accessories.</i>
<i>Who is doing?</i>	<i>BNL, ANL, SLAC GTF, DESY TTF, DESY Zeuthen, CERN CTF, Frascati SPARC, Tokyo University, and many others.</i>	<i>SCSS(RIKEN/SPring-8) Small emittance gun using BaO cathode: MIT 17GHz injector</i>